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Observations of the Cabibbo-Suppressed decays $\Lambda_c^+ \to n\pi^+\pi^0$, $n\pi^+\pi^-\pi^+$ and the Cabibbo-Favored decay $\Lambda_c^+ \to nK^-\pi^+\pi^+$*

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Abstract: Using electron-positron annihilation data samples corresponding to an integrated luminosity of 4.5 fb^{-1}, collected by the BESIII detector in the energy region between 4599.53 MeV and 4698.82 MeV, we report the first observations of the Cabibbo-suppressed decays $\Lambda_c^+ \to n\pi^+\pi^0$, $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$, and the Cabibbo-favored decay $\Lambda_c^+ \to nK^-\pi^+\pi^+$ with statistical significances of 7.9σ, 7.8σ, and $> 10\sigma$, respectively. The branching fractions of these decays are measured to be $\mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^0) = (0.64 \pm 0.09 \pm 0.02)\%$, $\mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^-\pi^+) = (0.45 \pm 0.07 \pm 0.03)\%$, and $\mathcal{B}(\Lambda_c^+ \to nK^-\pi^+\pi^+) = (1.90 \pm 0.08 \pm 0.09)\%$, where the first uncertainties are statistical and the second are systematic. We find that the branching fraction of the decay $\Lambda_c^+ \to n\pi^+\pi^0$ is about one order of magnitude higher than that of $\Lambda_c^+ \to n\pi^+$.

Keywords: Λ_c^+ baryon, Branching fraction, BESIII detector

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I. INTRODUCTION

Charmed baryons provide an excellent laboratory for studying the properties of QCD in the case where a heavy quark couples to two light quarks. Currently, we do not have reliable phenomenological models to describe these complex baryon decays. More than two decades ago, a general formulation of a topological-diagram scheme for the nonleptonic weak decays of baryons was proposed to calculate the amplitudes of different topological diagrams [1]. The factorizable external W-emission amplitude T and internal W-emission amplitude C, as well as the non-factorizable inner W-emission amplitude C' and W-exchange amplitude E are introduced in the representation of these diagrams. In this scheme, one amplitude can be factorized into two parts, the decay constant of the emitted meson and the heavy-to-light transition form factor, which can be directly calculated. By way of example, the topological diagrams for $\Lambda_c^+ \to n\pi^+\pi^0$ can be seen in Fig. 1. We know that here the non-factorizable contribution is essential for understanding the weak decays of charmed baryons, in contrast to the negligible effect they contribute in heavy-meson decays [2]. Hence, studies of charmed baryon decays involving non-factorizable components are critical for understanding the underlying dynamics of charmed-baryon decays in general.

In addition, an alternative and model-inpendent approach based on quark flavor $SU(3)$ symmetry has been proposed to describe charmed-baryon decays. Even though it is an approximate method, it is a powerful and reliable tool to extract useful information about these transitions [3–10]. Improved measurements of charmed-baryon decays will be important in further testing the validity of this method.

Many experimental studies of various Λ_c^+ decays have been reported, with a focus on decays to two-body hadronic final states [11–21]. However, more studies of multi-body hadronic decays are required to gain deeper insight into the nature of non-perturbative QCD, as these involve rich intermediate processes and the branching fractions (BF) are in general higher. There is a significant discrepancy between the predicted BFs of $\Lambda_c^+ \to N\bar{K}\pi\pi$ decays reported in Ref. [22] and the measured values for $\Lambda_c^+ \to pK^-\pi^+\pi^0$ and $\Lambda_c^+ \to p\bar{K}^0\pi^+\pi^-$ [23]. Clarifying this tension requires knowledge of the BFs of hadronic Λ_c^+ decays to final states involving a neutron. Fewer studies have been performed for these decays because of the difficulty of direct neutron detection.

The BESIII collaboration has measured the BFs of the Cabibbo-favored (CF) decay $\Lambda_c^+ \to nK_S^0\pi^+$ [14] and the Cabibbo-suppressed (CS) decay $\Lambda_c^+ \to n\pi^+$ [18]. These studies provide critical tests for isospin and $SU(3)$ symmetries in charmed-baryon decays, the violation of which could lead to enhanced CP-violation effects [24, 25]. Taking the BFs of $\Lambda_c^+ \to nK_S^0\pi^+$, $\Lambda_c^+ \to pK_S^0\pi^0$ and $\Lambda_c^+ \to pK^-\pi^+$ from the Particle Data Group (PDG) [23], the amplitudes of these three decays satisfy the triangle relation and obey isospin symmetry [3]. However, the ratio of the BF of $\Lambda_c^+ \to n\pi^+$ to that of $\Lambda_c^+ \to p\pi^0$ is measured to be larger than 7.2 which disagrees with the $SU(3)$ flavor symmetry model prediction [3, 6, 26]. Precise measurements of the BFs of additional Λ_c^+ decays to final states involving a neutron, such as $\Lambda_c^+ \to n\pi^+\pi^0$, $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$, and $\Lambda_c^+ \to nK^-\pi^+\pi^+$, are highly desirable to allow for futher tests of isospin and $SU(3)$ symmetries in charmed-baryon decays. Throughout this paper, charge-conjugate modes are implicitly assumed.

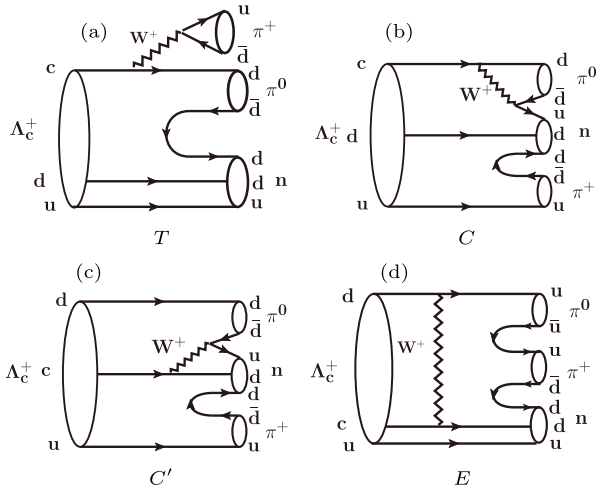

Fig. 1. Topological diagrams of the decay $\Lambda_c^+ \to n\pi^+\pi^0$ via (a) external W-emission T, (b) internal W-emission C, (c) inner W-emission C', and (d) W-exchange diagram E.

In this paper, two CS decays $\Lambda_c^+ \to n\pi^+\pi^0$, $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$, and the CF decay $\Lambda_c^+ \to nK^-\pi^+\pi^+$ are studied by employing the double-tag (DT) method [27], where the neutron is reconstructed via the missing-mass technique. Our analysis is based on electron-positron annihilation data samples collected at seven center-of-mass (c.m.) energies from $\sqrt{s} = 4599.53$ MeV to 4698.82 MeV by the BESIII detector. These data samples correspond to an integrated luminosity of 4.5 fb^{-1} [28–31], as listed in Table 1. Within this energy range, the electron-positron collisions provide a clean environment for the production $\Lambda_c^+\bar{\Lambda}_c^-$ pairs, which offers a unique opportunity to carry out model-independent measurements of the BFs of various Λ_c^+ decays involving neutrons. The DT method allows us to measure the BFs without any theoretical input or external information on the cross section of Λ_c^+ production.

Table 1. The c.m. energies and integrated luminosities for the data samples.

$\sqrt{s}$ /MeV	Int. luminosity/pb^{-1}
4599.53	$586.90 \pm 0.10 \pm 3.90$
4611.86	$103.65 \pm 0.05 \pm 0.55$
4628.00	$521.53 \pm 0.11 \pm 2.76$
4640.91	$551.65 \pm 0.12 \pm 2.92$
4661.24	$529.43 \pm 0.12 \pm 2.81$
4681.92	$1667.39 \pm 0.21 \pm 8.84$
4698.82	$535.54 \pm 0.12 \pm 2.84$

The $\bar{\Lambda}_c^-$ baryons are reconstructed with twelve exclusive hadronic decay modes, as listed in Table 2. This data set is referred to as the single-tag (ST) sample. The π^0, K_S^0, $\bar{\Lambda}$, $\bar{\Sigma}^0$, and $\bar{\Sigma}^-$ particles are reconstructed via individual dominant decay modes. Those events in which any of the signal decays $\Lambda_c^+ \to n\pi^+\pi^0$, $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$, and $\Lambda_c^+ \to nK^-\pi^+\pi^+$ are reconstructed in the system recoiling against the ST $\bar{\Lambda}_c^-$ candidates are denoted as DT candidates.

II. BESIII DETECTOR AND MONTE CARLO SIMULATION

The BESIII detector [32] records symmetric e^+e^- collisions provided by the BEPCII storage ring [33], which operates in the center-of-mass energy range from 2.0 GeV to 4.95 GeV. BESIII has collected large data

Table 2. Requirements on ΔE, ST yields, ST and DT efficiencies for the signal decays of $\Lambda_c^+ \to n\pi^+\pi^0$ ($\mathcal{E}$), $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ ($\mathcal{D}$) and $\Lambda_c^+ \to nK^-\pi^+\pi^+$ ($\mathcal{F}$) at $\sqrt{s} = 4681.92$ MeV. The uncertainties are statistical only. The quoted efficiencies do not include any subdecay BFs. Entries of "…" are for the cases where knowledge of the DT efficiencies is not required in the analysis.

Tag mode	ΔE /MeV	$N_i^{\rm ST}$	$\varepsilon_i^{\rm ST}(\%)$	$\varepsilon_i^{\rm DT}(\mathcal{E})(\%)$	$\varepsilon_i^{\rm DT}(\mathcal{D})(\%)$	$\varepsilon_i^{\rm DT}(\mathcal{F})(\%)$
$\bar{p}K_S^0$	$(-21, 18)$	3376±61	48.6	12.58	11.55	18.50
$\bar{p}K^+\pi^-$	$(-29, 26)$	17508±147	47.1	10.77	10.59	16.61
$\bar{p}K_S^0\pi^0$	$(-49, 34)$	1785±63	19.7	4.74	4.65	7.15
$\bar{p}K_S^0\pi^-\pi^+$	$(-34, 31)$	1511±57	20.7	4.22	4.13	5.99
$\bar{p}K^+\pi^-\pi^0$	$(-60, 41)$	5111±128	17.8	…	4.13	6.60
$\bar{\Lambda}\pi^-$	$(-23, 21)$	2074±48	41.3	9.95	9.12	15.08
$\bar{\Lambda}\pi^-\pi^0$	$(-50, 41)$	4380±88	18.2	4.05	3.85	6.00
$\bar{\Lambda}\pi^-\pi^+\pi^-$	$(-40, 36)$	2059±61	13.9	2.80	2.67	3.93
$\bar{\Sigma}^0\pi^-$	$(-33, 31)$	1398±42	25.9	6.31	5.89	8.46
$\bar{\Sigma}^-\pi^0$	$(-67, 32)$	879±43	22.5	5.78	5.21	8.22
$\bar{\Sigma}^-\pi^-\pi^+$	$(-40, 32)$	3027±88	22.1	5.40	5.21	7.77
$\bar{p}\pi^-\pi^+$	$(-26, 20)$	1596±80	53.6	…	…	17.35

samples in this energy region [34]. The cylindrical core of the BESIII detector covers 93% of the full solid angle and consists of a helium-based multilayer drift chamber (MDC), a plastic scintillator time-of-flight system (TOF), and a CsI(Tl) electromagnetic calorimeter (EMC), which are all enclosed in a superconducting solenoidal magnet providing a 1.0 T magnetic field. The solenoid is supported by an octagonal flux-return yoke with resistive plate counter muon identification modules interleaved with steel. The charged-particle momentum resolution at 1 GeV/c is 0.5%, and the specific ionization energy loss (dE/dx) resolution is 6% for electrons from Bhabha scattering. The EMC measures photon energies with a resolution of 2.5% (5%) at 1 GeV in the barrel (end-cap) region. The time resolution in the TOF barrel region is 68 ps, while that in the end-cap region is 60 ps [35–37]. High-statistics Monte Carlo (MC) simulation samples for the process $e^+e^- \to$ inclusive are produced with the KKMC generator [38] by incorporating the initial-state radiation (ISR) effects and the beam-energy spread. The inclusive MC simulation sample, which consists of $\Lambda_c^+\bar{\Lambda}_c^-$ events, $D_{(s)}^{(*)}$ production, ISR return to lower-mass ψ states, and continuum processes $e^+e^- \to q\bar{q}$ ($q = u,d,s$), is generated to determine the ST detection efficiencies and estimate the potential background. For this MC simulation sample, all the known decay modes of charmed hadrons and charmonia are modeled with evtgen [39, 40] using BFs taken from the PDG [23], and the remaining unknown decays are modeled with LUNDCHARM [41]. Furthermore, exclusive DT signal MC simulation events, where the $\bar{\Lambda}_c^-$ decays into any of the tag modes and the Λ_c^+ decays into any of the signal modes of $n\pi^+\pi^0$, $n\pi^+\pi^-\pi^+$ or $nK^-\pi^+\pi^+$, are used to determine the DT detection efficiencies. The Born cross sections are taken into account when producing the MC simulation sample of $\Lambda_c^+\bar{\Lambda}_c^-$ pairs. The $\Lambda_c^+ \to n\pi^+\pi^0$ and $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ signal MC samples are simulated evenly distributed in phase space since the angular distributions, momentum distributions and the two-body invariant mass distributions of the final state particles of the signal MC simulation samples are in good agreement with data with the current size of data set. For the signal MC simulation sample of $\Lambda_c^+ \to nK^-\pi^+\pi^+$, the key kinematic distributions mentioned above have been reweighted to agree with those of data. All final tracks and photons are fed into a GEANT4-based [42, 43] detector simulation package, which includes the geometric description of the BESIII detector.

III. ANALYSIS

Charged tracks detected in the MDC must satisfy $|\cos\theta| < 0.93$ (where θ is defined with respect to the z-axis, which is the symmetry axis of the MDC) and have a distance of closest approach to the interaction point (IP) of less than 10 cm along the beam axis and less than 1 cm in the perpendicular plane, except for those used for reconstructing K_S^0 and $\bar{\Lambda}$ decays. Particle identification (PID) for charged tracks combines measurements of the energy deposited in the MDC (dE/dx) and the flight time in the TOF to form likelihoods $\mathcal{L}(h)(h = p,K,\pi)$ for each hadron h hypothesis. Tracks are identified as protons when their likelihoods satisfy $\mathcal{L}(p) > \mathcal{L}(K)$ and $\mathcal{L}(p) > \mathcal{L}(\pi)$, while charged kaons and pions are identified by comparing the likelihoods for the kaon and pion hypotheses, $\mathcal{L}(K) > \mathcal{L}(\pi)$ or $\mathcal{L}(\pi) > \mathcal{L}(K)$, respectively.

Neutral showers are reconstructed in the EMC. Showers not associated with any charged track are identified as photon candidates. The deposited energy of each shower in the EMC must be greater than 25 MeV in the barrel region, corresponding to the polar angle $|\cos\theta| < 0.80$, and greater than 50 MeV in the end-cap region, corresponding to $0.86 < |\cos\theta| < 0.92$. The EMC time difference from the event start time is required to be less than 700 ns, to suppress electronic noise and showers unrelated to the event. The π^0 candidates are reconstructed from photon pairs with invariant masses within 115 MeV/c^2 < $M(\gamma\gamma)$ < 150 MeV/c^2. To improve momentum resolution, a kinematic fit constraining the photon pairs to the π^0 known mass is performed and the resulting four-momentum of the π^0 candidate is used for further analysis.

Candidates for K_S^0 and $\bar{\Lambda}$ decays are formed by $\pi^+\pi^-$ and $\bar{p}\pi^+$ combinations, respectively. For these tracks, their distances of closest approaches to the IP must be within ±20 cm along the beam direction. No distance constraint in the transverse plane is required. The charged pion is not subjected to the PID requirement described above, while the proton PID is imposed. The two final-state tracks are constrained to originate from a common decay vertex by requiring the χ^2 of the vertex fit to be less than 100. Furthermore, the decay vertex is required to be separated from the IP by a distance of at least twice the fitted vertex resolution. The fitted momenta of the $\pi^+\pi^-$ and $\bar{p}\pi^+$ combinations are used in the subsequent analysis. We require 487 MeV/c^2 < $M(\pi^+\pi^-)$ < 511 MeV/c^2 and 1111 MeV/c^2 < $M(\bar{p}\pi^+)$ < 1121 MeV/c^2 to select K_S^0 and $\bar{\Lambda}$ candidates, respectively, which corresponds to a window of about three times the standard deviations either side of the known masses. The $\bar{\Sigma}^0$ and $\bar{\Sigma}^-$ candidates are reconstructed with the $\gamma\bar{\Lambda}$ and $\bar{p}\pi^0$ combinations with invariant masses being in the intervals of 1179 MeV/c^2 < $M(\gamma\bar{\Lambda})$ < 1203 MeV/c^2 and 1176 MeV/c^2 < $M(\bar{p}\pi^0)$ < 1200 MeV/c^2, respectively.

For the tag modes $\bar{\Lambda}_c^- \to \bar{p}K_S^0\pi^0$, $\bar{\Lambda}_c^- \to \bar{p}K_S^0\pi^-\pi^+$, $\bar{\Lambda}_c^- \to \bar{\Sigma}^-\pi^-\pi^+$, and $\bar{\Lambda}_c^- \to \bar{p}\pi^-\pi^+$, possible backgrounds with $\bar{\Lambda} \to \bar{p}\pi^+$ are rejected by requiring $M(\bar{p}\pi^+)$ to be outside the range (1110, 1120) MeV/c^2. In the selection of $\bar{\Lambda}_c^- \to \bar{p}K_S^0\pi^0$ decays, candidate events within the range 1170 MeV/c^2 < $M(\bar{p}\pi^0)$ < 1200 MeV/c^2 are excluded to

suppress background from $\bar{\Sigma}^-$ decays. To remove K^0_S decays in the selection of $\bar{\Lambda}_c^- \to \bar{\Lambda}\pi^-\pi^+\pi^-$, $\bar{\Lambda}_c^- \to \bar{\Sigma}^-\pi^0$, $\bar{\Lambda}_c^- \to \bar{\Sigma}^-\pi^-\pi^+$, and $\bar{\Lambda}_c^- \to \bar{p}\pi^-\pi^+$ candidates, the invariant masses of any $\pi^+\pi^-$ and $\pi^0\pi^0$ pairs are required to lie outside of the range $(480,520)$ MeV/c^2. MC simulation studies indicate that peaking backgrounds and cross feeds among the twelve ST modes are negligible after applying the above veto procedures.

The ST $\bar{\Lambda}_c^-$ baryons are identified using the beam-constrained mass $M_{\rm BC} \equiv \sqrt{E^2_{\rm beam}/c^4 - p^2/c^2}$, where $E_{\rm beam}$ is the average value of the e^+ and e^- beam energies and p is the measured $\bar{\Lambda}_c^-$ momentum in the c.m. system of the e^+e^- collision. To improve the signal purity, the energy difference $\Delta E \equiv E - E_{\rm beam}$ for the $\bar{\Lambda}_c^-$ candidate is required to fulfil a mode-dependent ΔE requirement shown in Table 2, corresponding to approximately three times the resolutions. Here, E is the total reconstructed energy of the $\bar{\Lambda}_c^-$ candidate. For each ST mode, if more than one candidate satisfies the above requirements, we select the one with the minimal $|\Delta E|$. Figure 2 shows the $M_{\rm BC}$ distributions of various ST modes for the data sample at $\sqrt{s} = 4681.92$ MeV, where a clear $\bar{\Lambda}_c^-$ signal peak can be seen in each mode.

To obtain the ST yields, unbinned maximum-likelihood fits are performed to the $M_{\rm BC}$ distributions, where the signal shapes are modeled with the MC-simulated shape convolved with a Gaussian function representing the resolution difference between data and MC simulation, and the background shapes are described by an ARGUS function [44]. The fit results for the data sample at $\sqrt{s}$ =4681.92 MeV are shown in Fig. 2. The fits to the $M_{\rm BC}$ distributions for the other six data samples at different c.m. energies are shown in Appendix B. Candidates with $M_{\rm BC} \in (2275,2300)$ MeV/c^2 for the data sample at $\sqrt{s} = 4599.53$ MeV, $M_{\rm BC} \in (2275,2306)$ MeV/c^2 for the data samples at $\sqrt{s} =$ 4611.86 MeV, 4628.00 MeV, 4640.91 MeV, and $M_{\rm BC} \in (2275,2310)$ MeV/c^2 for the data samples at $\sqrt{s} =$ 4661.24 MeV, 4681.92 MeV, 4698.82 MeV are retained for further analysis. The differences in selection requirements between data sets are necessary as the resolution and effects of ISR vary with collision energy. The fitted ST yields, ST and DT efficiencies for each ST mode at $\sqrt{s} = 4681.92$ MeV are summarized in Table 2; those for the other c.m. energy points can be found in Appendix B.

Searches are performed for the decays $\Lambda_c^+ \to n\pi^+\pi^0$, $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ and $\Lambda_c^+ \to nK^-\pi^+\pi^+$ among the remaining tracks and showers recoiling against the ST $\bar{\Lambda}_c^-$ candidates. In the case of $\Lambda_c^+ \to n\pi^+\pi^0$, the event is allowed to contain only one pion with opposite charge to the tagged $\bar{\Lambda}_c^-$ satisfying the same selection criteria as described above. The π^0 candidate giving rise to the smallest χ^2 for the mass-constrained kinematic fit is retained. When searching for $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ and $\Lambda_c^+ \to nK^-\pi^+\pi^+$ decays,

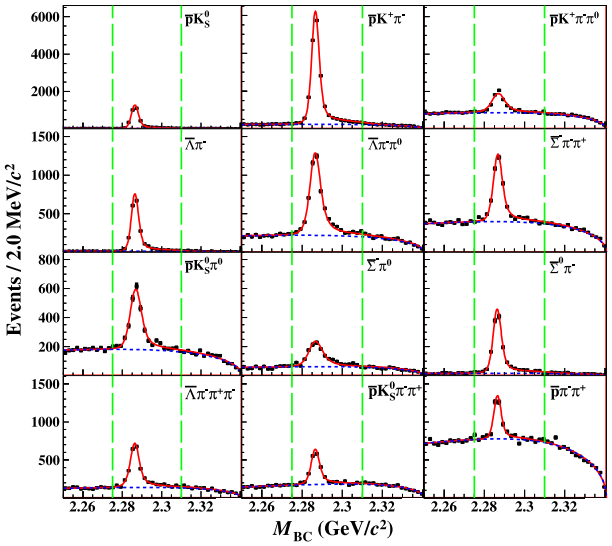

Fig. 2. (color online) The $M_{\rm BC}$ distributions of various ST modes for the data sample at $\sqrt{s}$=4681.92 MeV. The points with error bars represent data. The (red) solid curves indicate the fit results and the (blue) dashed curves describe the fitted background shapes. The ranges between green dashed lines are the signal regions.

events are selected with only three remaining charged tracks, satisfying the desired charge and PID criteria. In each of the decays, the kinematic variable $M_{\rm miss} \equiv \sqrt{E^2_{\rm miss}/c^4 - |\vec{p}_{\rm miss}|^2/c^2}$ is used to infer the presence of the undetected neutron. Here, $E_{\rm miss}$ and $\vec{p}_{\rm miss}$ are calculated by $E_{\rm miss} \equiv E_{\rm beam} - E_{\rm rec}$ and $\vec{p}_{\rm miss} \equiv \vec{p}_{\Lambda_c^+} - \vec{p}_{\rm rec}$, where $E_{\rm rec}(\vec{p}_{\rm rec})$ is the energy (momentum) of the reconstructed final-state particles in the e^+e^- c.m. system. The momentum of the Λ_c^+ baryon $\vec{p}_{\Lambda_c^+}$ is calculated by $\vec{p}_{\Lambda_c^+} \equiv -\hat{p}_{\rm tag}\sqrt{E^2_{\rm beam}/c^2 - m^2_{\Lambda_c^+}c^2}$, where $\hat{p}_{\rm tag}$ is the momentum direction of the ST $\bar{\Lambda}_c^-$ and $m_{\Lambda_c^+}$ is the known mass of the Λ_c^+ baryon [23]. In the case of $\Lambda_c^+ \to n\pi^+\pi^0$, $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$, and $\Lambda_c^+ \to nK^-\pi^+\pi^+$ decays, the $M_{\rm miss}(\pi^+\pi^0)$, $M_{\rm miss}(\pi^+\pi^-\pi^+)$ and $M_{\rm miss}(K^-\pi^+\pi^+)$ spectra are expected to peak around the known mass of the neutron *i.e.* at 939.6 MeV/c^2 [23]. A study of the inclusive MC simulation sample reveals that the dominant background events for the signal mode $\Lambda_c^+ \to n\pi^+\pi^0$ are from the processes $\Lambda_c^+ \to \Lambda\pi^+$ with $\Lambda \to n\pi^0$, $\Lambda_c^+ \to \Sigma^+\pi^0$ with $\Sigma^+ \to n\pi^+$ and $\Lambda_c^+ \to \Sigma^0\pi^+$ with $\Sigma^0 \to \gamma\Lambda(\to n\pi^0)$. In order to reject these background events for $\Lambda_c^+ \to n\pi^+\pi^0$, the following selection criteria are applied: $M_{\rm miss}(\pi^+) > 1300$ MeV/c^2 and $M_{\rm miss}(\pi^0) > 1370$ MeV/c^2. For the signal mode $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$, the peaking backgrounds from the decays $\Lambda_c^+ \to nK^0_S(\to \pi^+\pi^-)\pi^+$, $\Lambda_c^+ \to \Sigma^+(\to n\pi^+)\pi^+\pi^-$, and $\Lambda_c^+ \to \Sigma^-(\to n\pi^-)\pi^+\pi^+$ are suppressed by requiring $M_{\pi^+\pi^-} \notin (487,511)$ MeV/c^2, $M_{\rm miss}(\pi^+\pi^-) \notin (1150,1250)$ MeV/c^2 and $M_{\rm miss}(\pi^+\pi^+) \notin (1150,1250)$ MeV/c^2. Note that, based on the study of the inclusive MC simulation

sample, ten ST modes (after excluding $\bar{\Lambda}_c^- \to \bar{p}K^+\pi^-\pi^0$ and $\bar{\Lambda}_c^- \to \bar{p}\pi^-\pi^+$) are used in the analysis of the decay $\Lambda_c^+ \to n\pi^+\pi^0$, and eleven ST modes (after excluding $\bar{\Lambda}_c^- \to \bar{p}\pi^-\pi^+$) are used in the analysis of the decay $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ to improve the background level in the $M_{\rm miss}$ spectra.

The resulting $M_{\rm miss}$ distributions of the DT candidate events summed over all data samples at seven c.m. energies are shown in Fig. 3. A peak around the neutron mass is observed in Fig. 3(a) representing the $\Lambda_c^+ \to n\pi^+\pi^0$ signal. Moreover, there are two prominent structures peaking around the Λ and Σ^0 mass regions, which correspond to the CF decays $\Lambda_c^+ \to \Lambda\pi^+\pi^0$ and $\Lambda_c^+ \to \Sigma^0\pi^+\pi^0$, respectively. Significant signal peaks around the neutron mass are also observed for $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ and $\Lambda_c^+ \to nK^-\pi^+\pi^+$ in Fig. 3(b) and Fig. 3(c), respectively.

The total DT signal yield is obtained by performing an unbinned maximum-likelihood fit on the $M_{\rm miss}$ distribution. The neutron, Λ and Σ^0 signals are modeled by individual MC-derived shapes convolved with Gaussian functions that account for the shift and resolution difference between data and MC simulation. The Gaussian-function parameters are left free and are shared by the three signal processes in the fit to the $M_{\rm miss}(\pi^+\pi^0)$ spectra. The potential background events in the $M_{\rm miss}$ distributions are classified into two categories. Those directly originating from continuum hadron production in the e^+e^- annihilation are denoted as $q\bar{q}$ background. Those from $e^+e^- \to \Lambda_c^+\bar{\Lambda}_c^-$ events excluding the contributions from the corresponding signal process are referred to as $\Lambda_c^+\bar{\Lambda}_c^-$ background. In the fit to the $M_{\rm miss}(\pi^+\pi^0)$ distribution, the $q\bar{q}$ background is described by a second-order Chebyshev polynomial with free parameters and the $\Lambda_c^+\bar{\Lambda}_c^-$ backgroud shape is taken from the inclusive MC simulation sample. For the fit to the $M_{\rm miss}(\pi^+\pi^-\pi^+)$ spectrum, only one first-order Chebyshev polynomial with free parameters is used to model all the background. Similarly, for the $M_{\rm miss}(K^-\pi^+\pi^+)$ distribution, one second-order Chebyshev function with fixed parameters is taken as the background shape where the parameters are derived from the fit to the inclusive MC simulation sample. The DT signal and background yields are left free in the fit to

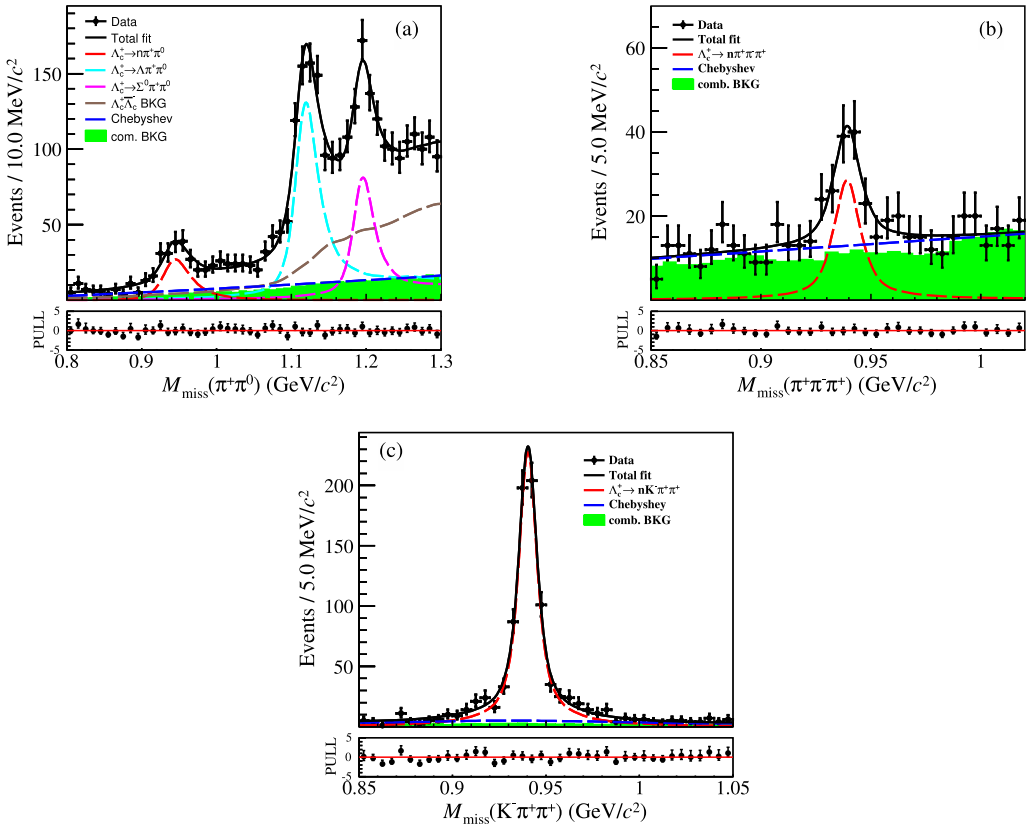

Fig. 3. (color online) The $M_{\rm miss}$ distributions of the surviving DT candidate events for (a) $\Lambda_c^+ \to n\pi^+\pi^0$, (b) $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ and (c) $\Lambda_c^+ \to nK^-\pi^+\pi^+$ decays with fit results overlaid. The points with error bars are data combined from seven c.m. energy points. The solid black curves are the fit results. The red, cyan, and pink dashed curves indicate the neutron, Λ, and Σ signal shapes, respectively. The brown dashed curve is the $\Lambda_c^+\bar{\Lambda}_c^-$ background shape for $\Lambda_c^+ \to n\pi^+\pi^0$. The blue dashed curves represent the fitted combinatorial background shape. The green histograms are the simulated combinatorial background shapes from the inclusive MC sample. The plots in the bottom of each graphs show the pull value of each bin, in which the values are expected to fluctuate around 0.

the $M_{\rm miss}(\pi^+\pi^0)$, $M_{\rm miss}(\pi^+\pi^-\pi^+)$ and $M_{\rm miss}(K^-\pi^+\pi^+)$ spectra. Figure 3 shows the results of the fits to the $M_{\rm miss}$ distributions. From these fits, we determine the DT signal yields of $\Lambda_c^+\to n\pi^+\pi^0$, $\Lambda_c^+\to n\pi^+\pi^-\pi^+$, and $\Lambda_c^+\to nK^-\pi^+\pi^+$ to be $N^{\rm DT}_{n\pi^+\pi^0}=150.9\pm21.4$, $N^{\rm DT}_{n\pi^+\pi^-\pi^+}=120.6\pm17.9$ and $N^{\rm DT}_{nK^-\pi^+\pi^+}=805.8\pm33.1$, respectively, where the uncertainties are statistical only. The statistical significances of the $\Lambda_c^+\to n\pi^+\pi^0$, $\Lambda_c^+\to n\pi^+\pi^-\pi^+$ and $\Lambda_c^+\to nK^-\pi^+\pi^+$ signals are 7.9σ, 7.8σ, and $>10\sigma$, respectively, which are evaluated by the changes in the likelihoods between the nominal fit and the fit with the signal yield set to zero, and accounting for the change in the number of degrees of freedom.

The BFs of the decays $\Lambda_c^+\to n\pi^+\pi^0$, $\Lambda_c^+\to n\pi^+\pi^-\pi^+$, and $\Lambda_c^+\to nK^-\pi^+\pi^+$ are determined as

$$\mathcal{B}=\frac{N^{\rm DT}}{\sum_{ij}N^{\rm ST}_{ij}\cdot(\epsilon^{\rm DT}_{ij}/\epsilon^{\rm ST}_{ij})\cdot\mathcal{B}_{\rm int}}, \tag{1}$$

where i and j represent the ST modes and the data samples at different c.m. energies, respectively. The factor $\mathcal{B}_{\rm int}$ is $(98.823\pm0.034)\,\%$, which is the BF of $\pi^0\to\gamma\gamma$ [23], is only present for $\Lambda_c^+\to n\pi^+\pi^0$. $N^{\rm ST}_{ij}$, $\epsilon^{\rm ST}_{ij}$, and $\epsilon^{\rm DT}_{ij}$ are the ST yields, ST efficiencies, and DT efficiencies, respectively. The detection efficiencies $\epsilon^{\rm ST}_{ij}$ and $\epsilon^{\rm DT}_{ij}$ are estimated from the inclusive MC simulation sample and exclusive DT signal MC simulation samples, respectively. The ST and DT efficiencies for the data sample at $\sqrt{s}=4681.92$ MeV are summarized in Table 2. The detection efficiencies for the other data samples are summarized in Appendix B. The obtained BFs are summaried in Table 3.

Most systematic uncertainties from the ST side cancel in the determination of the BFs, as is clear from Eq. (1). However, effects from the signal side can lead to systematic bias, for example, the requirement of no extra charged track, tracking efficiency, PID efficiency, π^0 reconstruction, peaking background veto, $M_{\rm miss}$ fit, ST $\bar{\Lambda}_c^-$ yield, MC modeling, and MC sample size. The systematic uncertainty due to the requirement of no extra charged track is assigned as 1.1% from the study of a control sample of $e^+e^-\to\Lambda_c^+\bar{\Lambda}_c^-$ with $\Lambda_c^+\to nK^-\pi^+\pi^+$ and $\bar{\Lambda}_c^-$ decays to tag modes. The systematic uncertainties associated with the efficiencies of the tracking and PID of charged particles are estimated to be 1% by using control samples of $e^+e^-\to\pi^+\pi^+\pi^-\pi^-$ and $e^+e^-\to K^+K^-\pi^+\pi^-$ events collected at c.m. energies above $\sqrt{s}=4.0\,$GeV. The systematic uncertainty due to the π^0 reconstruction efficiency is assigned to be 1.0% [12]. The systematic uncertainty associated with the BF of $\pi^0\to\gamma\gamma$ is 0.03% [23], which is negligible. In order to estimate the systematic uncertainties arising from the veto of peaking backgrounds involving Λ, Σ^+, Σ^0, K^0_S, and Σ^-, the corresponding resolutions in the MC simulation samples are corrected to agree with those in data, and the BFs are then re-evaluated with the updated MC simulation samples. The deviations from the baseline BF measurements are taken as the associated systematic uncertainties for $\Lambda_c^+\to n\pi^+\pi^0$ and $\Lambda_c^+\to n\pi^+\pi^-\pi^+$, which are 1.6% and 1.0%, respectively. The systematic uncertainties from the fitted DT yields are 1.3%, 2.2%, and 0.9% for $\Lambda_c^+\to n\pi^+\pi^0$, $n\pi^+\pi^-\pi^+$, and $nK^-\pi^+\pi^+$, respectively, which are estimated from varying the alternative polynomial descriptions for the $q\bar{q}$ and $\Lambda_c^+\bar{\Lambda}_c^-$ backgrounds, respectively. The systematic uncertainty in the total ST $\bar{\Lambda}_c^-$ yield is 0.1% for $\Lambda_c^+\to n\pi^+\pi^0$, and 0.2% for $\Lambda_c^+\to n\pi^+\pi^-\pi^+$, and $nK^-\pi^+\pi^+$. These uncertainties arise from the fluctuation of background together with a component coming from the fit to the $M_{\rm BC}$ distribution. The systematic uncertainties arising from the MC modeling are investigated by reweighting the MC distribution to data, and they are assigned as the efficiency differences between the original and reweighted samples. The systematic uncertainties due to limited sample sizes of the MC samples are estimated to be 0.2%. Assuming that all the sources are uncorrelated, the total uncertainties are then taken to be the quadratic sums of the individual values, which are 3.1%, 5.6%, and 4.5% for $\Lambda_c^+\to n\pi^+\pi^0$, $\Lambda_c^+\to n\pi^+\pi^-\pi^+$, and $\Lambda_c^+\to nK^-\pi^+\pi^+$, respectively. All the above systematic uncertainties are summarized in Table 4.

Table 3. The obtained BFs, where the first uncertainties are statistical and the second are systematic.

Signal decay	$\mathcal{B}$ (%)
$\Lambda_c^+\to n\pi^+\pi^0$	$0.64\pm0.09\pm0.02$
$\Lambda_c^+\to n\pi^+\pi^-\pi^+$	$0.45\pm0.07\pm0.03$
$\Lambda_c^+\to nK^-\pi^+\pi^+$	$1.90\pm0.08\pm0.09$

Table 4. Relative systematic uncertainties in BF measurements. "…" means the uncertainty is negligible. "–" indicates cases where there is no uncertainty.

Source	$n\pi^+\pi^0$ (%)	$n\pi^+\pi^-\pi^+$ (%)	$nK^-\pi^+\pi^+$ (%)
No extra charged track	1.1	1.1	1.1
Tracking	1.0	3.0	3.0
PID	1.0	3.0	3.0
π^0 reconstruction	1.0	–	–
Background veto	1.6	1.0	–
$M_{\rm miss}$ fit	1.3	2.2	0.9
ST $\bar{\Lambda}_c^-$ yield	0.1	0.2	0.2
MC model	1.1	2.6	…
MC sample size	0.2	0.2	0.2
Total	3.1	5.6	4.5

IV. SUMMARY

In summary, by analyzing 4.5 fb^{-1} of data collected at c.m. energies between 4599.53 and 4698.82 MeV, we report for the first time the observation of $\Lambda_c^+ \to n\pi^+\pi^0$, $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ and $\Lambda_c^+ \to nK^-\pi^+\pi^+$. The BFs of these decays are determined to be $\mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^0) = (0.64\pm 0.09 \pm 0.02)\%$, $\mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^-\pi^+) = (0.45 \pm 0.07 \pm 0.03)\%$, and $\mathcal{B}(\Lambda_c^+ \to nK^-\pi^+\pi^+) = (1.90 \pm 0.08 \pm 0.09)\%$, where the first uncertainties are statistical and the second systematic. These observations are important additions to our knowledge of Λ_c^+ decays. Comparison of these results to those of decays involving protons provides crucial inputs for understanding the mechanisms in the charmed baryon decays under the $SU(3)$ flavor symmetry. Taking $\mathcal{B}(\Lambda_c^+ \to p\pi^-\pi^+) = (0.461 \pm 0.028)\%$ from the PDG [23], we can calculate $\mathcal{B}(\Lambda_c^+ \to p\pi^-\pi^+)/\mathcal{B}(\Lambda_c^+ \to n\pi^0\pi^+) = 0.72\pm 0.11$. This result provides useful input to test of isospin symmetry in the charm baryon sector. Taking $\mathcal{B}(\Lambda_c^+ \to n\pi^+) = (6.6 \pm 1.3)\times 10^{-4}$ [18], the ratio $\mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^0)/\mathcal{B}(\Lambda_c^+ \to n\pi^+)$ is calculated to be 9.7 ± 2.4, indicating an order-of-magnitude difference in the rates of the two decays. This ratio is greater than $\mathcal{B}(\Lambda_c^+ \to pK_S^0\pi^0)/\mathcal{B}(\Lambda_c^+ \to pK_S^0) = (1.24 \pm 0.10)$. To further understand this behavior, amplitude analysis will be needed to decouple the intermediate resonances contributions for $\Lambda_c^+ \to n\pi^+\pi^0$ and $\Lambda_c^+ \to pK_S^0\pi^0$. The ratio of $\mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^-\pi^+)/\mathcal{B}(\Lambda_c^+ \to nK^-\pi^+\pi^+) = (0.24 \pm 0.04)$ which is consistent with the ratio of Cabibbo-Kobayashi-Maskawa matrix elements $|V_{cd}|/|V_{cs}| = (0.224 \pm 0.005)$ offers a new constraint on the CS and CF decay dynamics. The results from this analysis provide an essential input for the phenomenological studies on the underlying dynamics of charmed bayond decays.

ACKNOWLEDGEMENTS

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APPENDIX A

Figures A1 and A2 show the topological diagrams of $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ and $\Lambda_c^+ \to nK^-\pi^+\pi^+$, respectively.

APPENDIX B

Figures B1–B6 show the fits for the $M_{\rm BC}$ distributions of the ST $\bar{\Lambda}_c^-$ candidates for various tag modes from the data samples at $\sqrt{s}$ = 4599.53, 4611.86, 4628.00, 4640.91, 4661.24, and 4698.82 MeV, respectively.

Tables B1–B6 show the ST yields, ST and DT efficiencies for various tag modes from the data samples at $\sqrt{s}$ = 4599.53, 4611.86, 4628.00, 4640.91, 4661.24, and 4698.82 MeV, respectively.

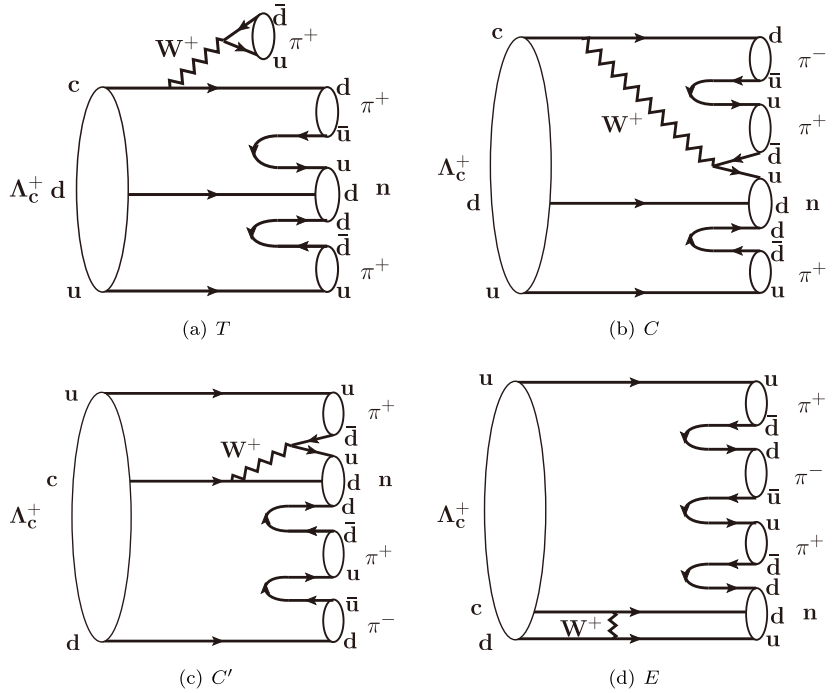

Fig. A1. Topological diagrams of $\Lambda_c^+ \to n\pi^+\pi^-\pi^+$ via (a) external W-emission T, (b) internal W-emission C, (c) inner W-emission C', and (d) W-exchange diagram E.

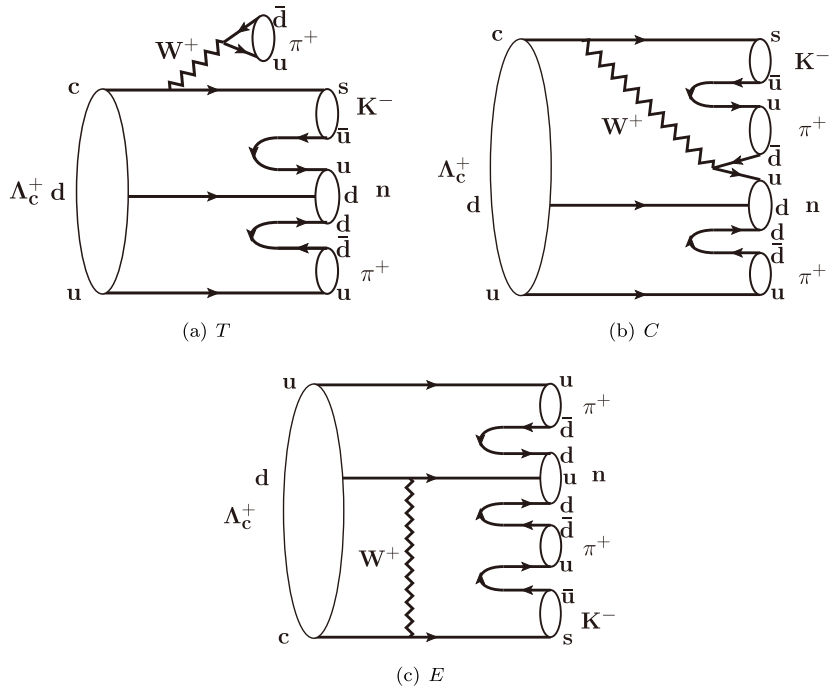

Fig. A2. Topological diagrams of $\Lambda_c^+ \to nK^-\pi^+\pi^+$ via (a) external W-emission T, (b) internal W-emission C, (c) W-exchange diagram E.

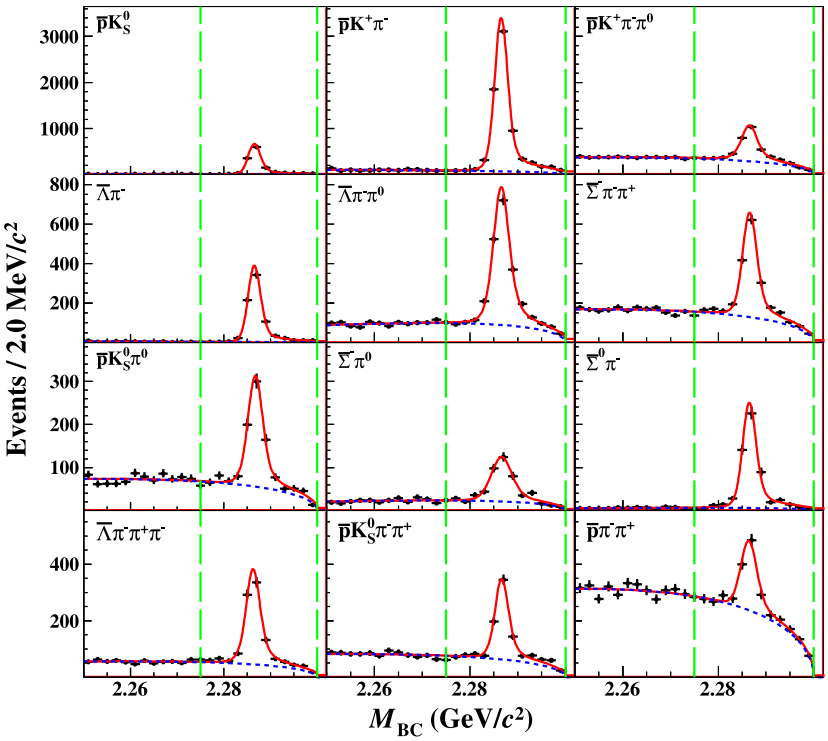

Fig. B1. (color online) The $M_{\rm BC}$ distributions of the ST $\bar{\Lambda}_c^-$ candidates of various tag modes for the data sample at $\sqrt{s} = 4599.53$ MeV. The points with error bars represent data. The red solid curves indicate the fit results and the blue dashed curves describe the background shapes. The ranges between green dashed lines are the signal regions.

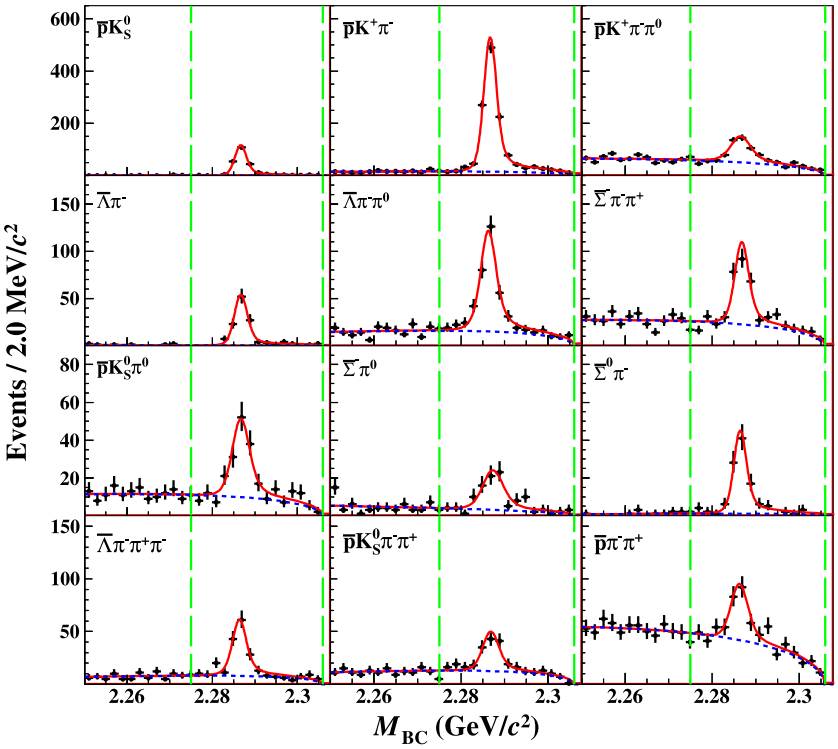

Fig. B2. (color online) The $M_{\rm BC}$ distributions of the ST $\bar{\Lambda}_c^-$ candidates of various tag modes for the data sample at $\sqrt{s} = 4611.86$ MeV. The points with error bars represent data. The red solid curves indicate the fit results and the blue dashed curves describe the background shapes. The ranges between green dashed lines are the signal regions.

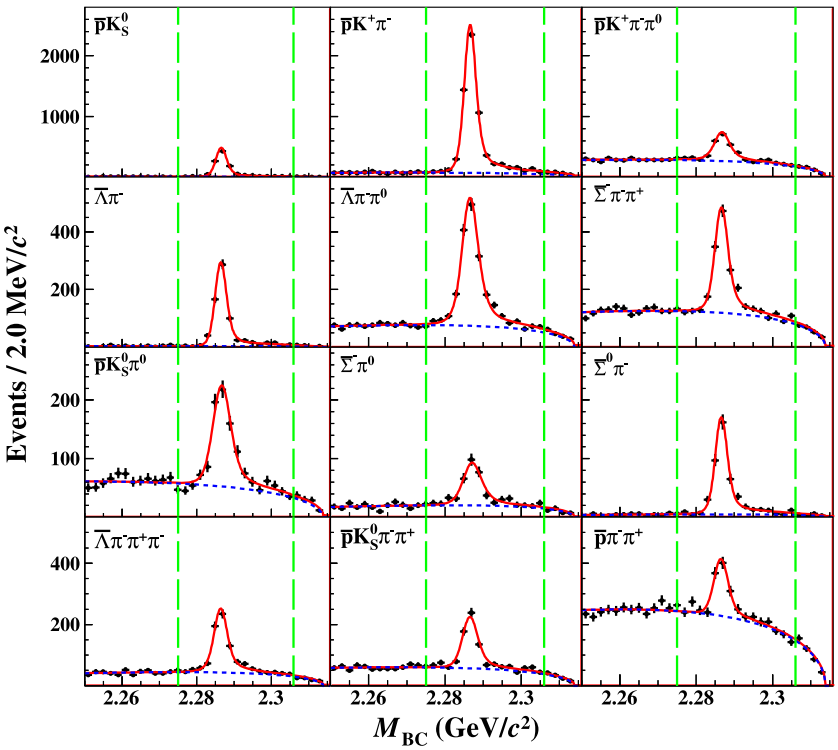

Fig. B3. (color online) The $M_{\rm BC}$ distributions of the ST $\bar{\Lambda}_c^-$ candidates of various tag modes for the data sample at $\sqrt{s} = 4628.00$ MeV. The points with error bars represent data. The red solid curves indicate the fit results and the blue dashed curves describe the background shapes. The ranges between green dashed lines are the signal regions.

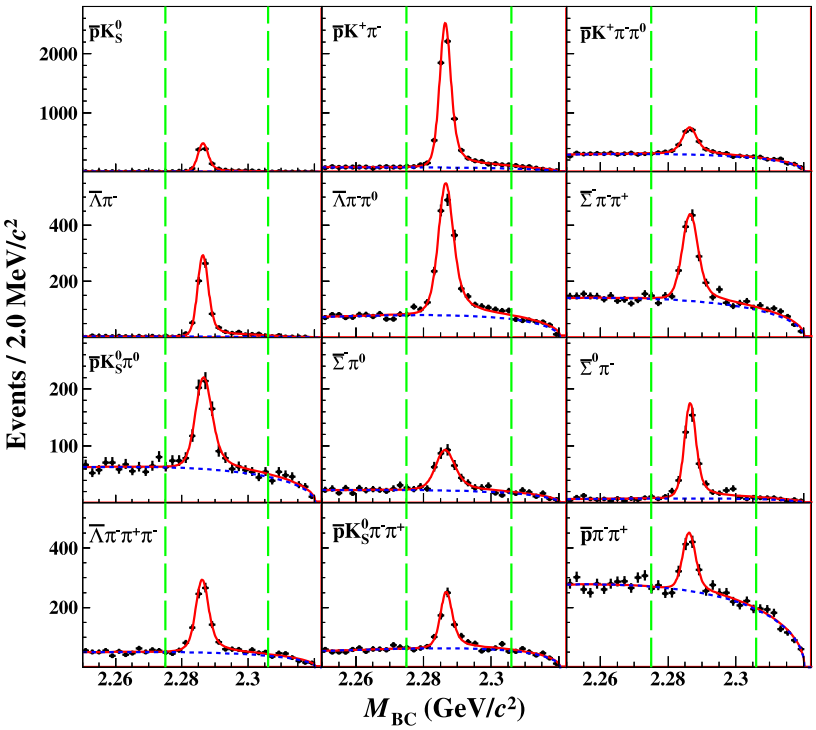

Fig. B4. (color online) The $M_{\rm BC}$ distributions of the ST $\bar{\Lambda}_c^-$ candidates of various tag modes for the data sample at $\sqrt{s} = 4640.91$ MeV. The points with error bars represent data. The red solid curves indicate the fit results and the blue dashed curves describe the background shapes. The ranges between green dashed lines are the signal regions.

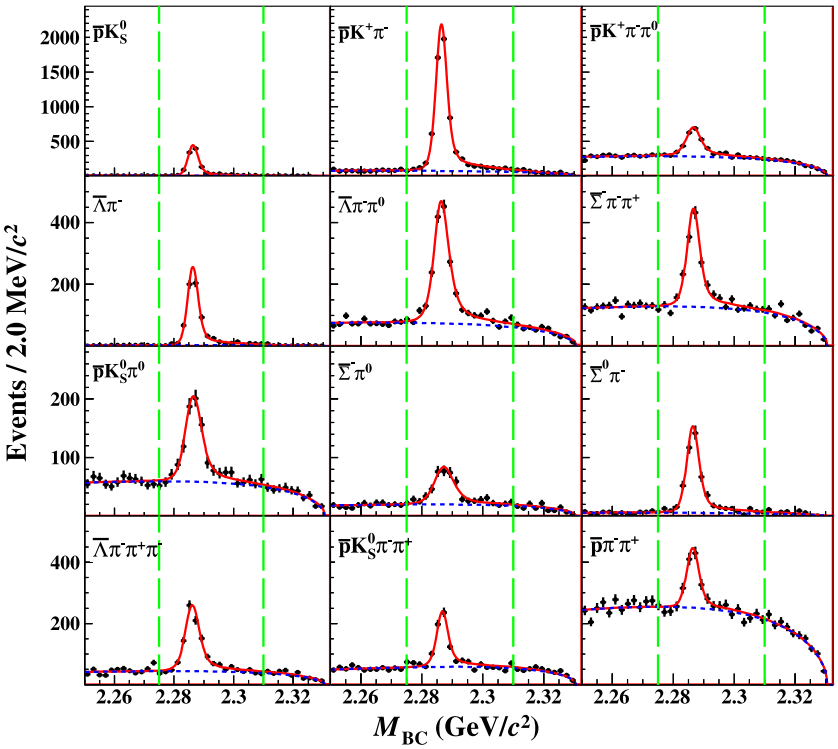

Fig. B5. (color online) The $M_{\rm BC}$ distributions of the ST $\bar{\Lambda}_c^-$ candidates of various tag modes for the data sample at $\sqrt{s} = 4661.24$ MeV. The points with error bars represent data. The red solid curves indicate the fit results and the blue dashed curves describe the background shapes. The ranges between green dashed lines are the signal regions.

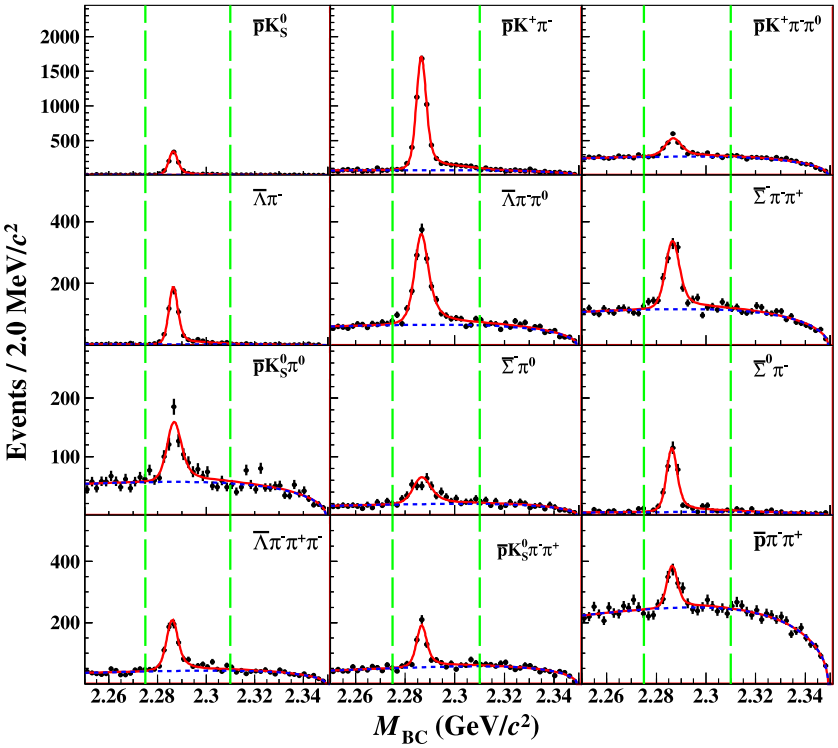

Fig. B6. (color online) The $M_{\rm BC}$ distributions of the ST $\bar{\Lambda}_c^-$ candidates of various tag modes for the data sample at $\sqrt{s} = 4698.82$ MeV. The points with error bars represent data. The red solid curves indicate the fit results and the blue dashed curves describe the background shapes. The ranges between green dashed lines are the signal regions.

Table B1. ST yields, ST and DT efficiencies of various tag modes for the data sample at $\sqrt{s} = 4599.53$ MeV. The uncertainties are statistical only. The quoted efficiencies do not include any subdecay BFs. Entries of "…" are for the cases where knowledge of the DT efficiencies are not required in the analysis.

Tag mode	$N_i^{\rm ST}$	$\varepsilon_i^{\rm ST}(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^0)(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^-\pi^+)(\%)$	$\varepsilon_i^{\rm DT}(nK^-\pi^+\pi^-)(\%)$
$\bar{p}K_S^0$	1277 ± 36	56.1	14.27	13.49	20.87
$\bar{p}K^+\pi^-$	6806 ± 91	51.5	11.25	11.53	17.64
$\bar{p}K_S^0\pi^0$	606 ± 34	23.0	5.17	5.25	7.34
$\bar{p}K_S^0\pi^-\pi^+$	613 ± 34	23.5	4.51	4.76	6.89
$\bar{p}K^+\pi^-\pi^0$	2197 ± 78	20.6	…	4.57	6.83
$\bar{\Lambda}\pi^-$	757 ± 28	48.4	11.62	10.9	17.08
$\bar{\Lambda}\pi^-\pi^0$	1742 ± 56	21.6	4.52	4.43	6.42
$\bar{\Lambda}\pi^-\pi^+\pi^-$	769 ± 36	15.6	2.86	2.90	4.17
$\bar{\Sigma}^0\pi^-$	520 ± 26	29.4	7.30	6.96	11.20
$\bar{\Sigma}^-\pi^0$	320 ± 25	23.7	6.47	6.05	8.38
$\bar{\Sigma}^-\pi^-\pi^+$	1186 ± 49	25.4	5.84	5.82	8.81
$\bar{p}\pi^-\pi^+$	598 ± 47	64.3	…	…	22.04

Table B2. ST yields, ST and DT efficiencies of various tag modes for the data sample at $\sqrt{s} = 4611.86\,\text{MeV}$. The uncertainties are statistical only. The quoted efficiencies do not include any subdecay BFs. Entries of "⋯" are for the cases where knowledge of the DT efficiencies are not required in the analysis.

Tag mode	$N_i^{\rm ST}$	$\varepsilon_i^{\rm ST}(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^0)(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^-\pi^+)(\%)$	$\varepsilon_i^{\rm DT}(nK^-\pi^+\pi^-)(\%)$
$\bar{p}K_S^0$	239 ± 16	53.7	13.89	12.71	20.09
$\bar{p}K^+\pi^-$	1166 ± 39	51.0	11.34	11.27	16.98
$\bar{p}K_S^0\pi^0$	127 ± 17	22.2	5.04	4.91	7.20
$\bar{p}K_S^0\pi^-\pi^+$	106 ± 16	21.9	4.35	4.33	6.04
$\bar{p}K^+\pi^-\pi^0$	364 ± 34	19.9	⋯	4.40	6.25
$\bar{\Lambda}\pi^-$	123 ± 11	46.9	11.29	10.22	15.27
$\bar{\Lambda}\pi^-\pi^0$	302 ± 23	19.8	4.34	4.17	6.08
$\bar{\Lambda}\pi^-\pi^+\pi^-$	139 ± 15	13.6	2.68	2.75	3.69
$\bar{\Sigma}^0\pi^-$	102 ± 13	26.6	6.89	6.53	8.65
$\bar{\Sigma}^-\pi^0$	73 ± 10	22.6	6.47	5.9	8.46
$\bar{\Sigma}^-\pi^-\pi^+$	218 ± 22	25.5	5.74	5.63	8.47
$\bar{p}\pi^-\pi^+$	155 ± 22	71.4	⋯	⋯	21.82

Table B3. ST yields, ST and DT efficiencies of various tag modes for the data sample at $\sqrt{s} = 4628.00\,\text{MeV}$. The uncertainties are statistical only. The quoted efficiencies do not include any subdecay BFs.Entries of "⋯" are for the cases where knowledge of the DT efficiencies are not required in the analysis.

Tag mode	$N_i^{\rm ST}$	$\varepsilon_i^{\rm ST}(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^0)(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^-\pi^+)(\%)$	$\varepsilon_i^{\rm DT}(nK^-\pi^+\pi^-)(\%)$
$\bar{p}K_S^0$	1054 ± 35	51.8	13.27	12.17	18.98
$\bar{p}K^+\pi^-$	5886 ± 39	49.2	10.98	10.97	16.77
$\bar{p}K_S^0\pi^0$	616 ± 36	20.7	4.94	4.79	6.73
$\bar{p}K_S^0\pi^-\pi^+$	510 ± 32	20.6	4.21	4.17	6.01
$\bar{p}K^+\pi^-\pi^0$	1589 ± 69	18.7	⋯	4.32	6.95
$\bar{\Lambda}\pi^-$	675 ± 28	43.2	10.69	9.84	14.72
$\bar{\Lambda}\pi^-\pi^0$	1454 ± 54	19.1	4.18	3.98	5.90
$\bar{\Lambda}\pi^-\pi^+\pi^-$	587 ± 33	13.6	2.70	2.65	3.67
$\bar{\Sigma}^0\pi^-$	413 ± 23	27.2	6.62	6.24	8.41
$\bar{\Sigma}^-\pi^0$	263 ± 23	23.4	6.20	5.65	9.10
$\bar{\Sigma}^-\pi^-\pi^+$	994 ± 20	23.6	5.54	5.47	8.24
$\bar{p}\pi^-\pi^+$	517 ± 45	61.6	⋯	⋯	19.75

Table B4. ST yields, ST and DT efficiencies of various tag modes for the data sample at $\sqrt{s} = 4640.91\,\text{MeV}$. The uncertainties are statistical only. The quoted efficiencies do not include any subdecay BFs. Entries of "⋯" are for the cases where knowledge of the DT efficiencies are not required in the analysis.

Tag mode	$N_i^{\rm ST}$	$\varepsilon_i^{\rm ST}(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^0)(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^-\pi^+)(\%)$	$\varepsilon_i^{\rm DT}(nK^-\pi^+\pi^-)(\%)$
$\bar{p}K_S^0$	1107 ± 36	50.7	13.14	12.08	19.12
$\bar{p}K^+\pi^-$	6250 ± 89	48.5	10.91	10.86	16.83
$\bar{p}K_S^0\pi^0$	599 ± 36	20.7	4.82	4.77	7.02
$\bar{p}K_S^0\pi^-\pi^+$	522 ± 33	20.8	4.21	4.15	6.50
$\bar{p}K^+\pi^-\pi^0$	1632 ± 70	18.1	⋯	4.23	6.62
$\bar{\Lambda}\pi^-$	705 ± 29	42.7	10.51	9.55	14.24
$\bar{\Lambda}\pi^-\pi^0$	1613 ± 54	19.1	4.14	3.95	5.46
$\bar{\Lambda}\pi^-\pi^+\pi^-$	745 ± 36	14.2	2.70	2.67	4.20
$\bar{\Sigma}^0\pi^-$	445 ± 25	26.2	6.43	6.09	8.97
$\bar{\Sigma}^-\pi^0$	298 ± 24	24.6	6.01	5.52	7.70
$\bar{\Sigma}^-\pi^-\pi^+$	1077 ± 49	23.4	5.43	5.38	8.01
$\bar{p}\pi^-\pi^+$	552 ± 47	59.7	⋯	⋯	22.97

Table B5. ST yields, ST and DT efficiencies of various tag modes for the data sample at $\sqrt{s} = 4661.24$ MeV. The uncertainties are statistical only. The quoted efficiencies do not include any subdecay BFs. Entries of "⋯" are for the cases where knowledge of the DT efficiencies are not required in the analysis.

Tag mode	$N_i^{\rm ST}$	$\varepsilon_i^{\rm ST}(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^0)(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^-\pi^+)(\%)$	$\varepsilon_i^{\rm DT}(nK^-\pi^+\pi^-)(\%)$
$\bar{p}K_S^0$	1119 ± 35	49.6	12.82	11.78	18.58
$\bar{p}K^+\pi^-$	5938 ± 86	48.2	10.93	10.85	16.47
$\bar{p}K_S^0\pi^0$	594 ± 36	20.1	4.86	4.75	7.12
$\bar{p}K_S^0\pi^-\pi^+$	537 ± 33	20.2	4.29	4.17	6.27
$\bar{p}K^+\pi^-\pi^0$	1700 ± 73	18.1	⋯	4.15	6.36
$\bar{\Lambda}\pi^-$	668 ± 27	41.7	10.15	9.41	14.98
$\bar{\Lambda}\pi^-\pi^0$	1491 ± 51	18.9	4.15	3.91	5.91
$\bar{\Lambda}\pi^-\pi^+\pi^-$	780 ± 36	14.1	2.71	2.63	4.04
$\bar{\Sigma}^0\pi^-$	454 ± 25	26.3	6.56	5.94	9.20
$\bar{\Sigma}^-\pi^0$	298 ± 25	23.2	5.96	5.35	8.22
$\bar{\Sigma}^-\pi^-\pi^+$	1066 ± 49	23.2	5.48	5.34	8.04
$\bar{p}\pi^-\pi^+$	590 ± 48	60.2	⋯	⋯	20.06

Table B6. ST yields, ST and DT efficiencies of various tag modes for the data sample at $\sqrt{s} = 4698.82$ MeV. The uncertainties are statistical only. The quoted efficiencies do not include any subdecay BFs. Entries of "⋯" are for the cases where knowledge of the DT efficiencies are not required in the analysis.

Tag mode	$N_i^{\rm ST}$	$\varepsilon_i^{\rm ST}(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^0)(\%)$	$\varepsilon_i^{\rm DT}(n\pi^+\pi^-\pi^+)(\%)$	$\varepsilon_i^{\rm DT}(nK^-\pi^+\pi^-)(\%)$
$\bar{p}K_S^0$	958 ± 33	47.5	12.29	11.25	17.86
$\bar{p}K^+\pi^-$	5167 ± 80	46.3	10.62	10.49	16.16
$\bar{p}K_S^0\pi^0$	471 ± 34	18.9	4.66	4.52	6.62
$\bar{p}K_S^0\pi^-\pi^+$	462 ± 31	19.5	4.21	4.11	6.33
$\bar{p}K^+\pi^-\pi^0$	1389 ± 74	17.5	⋯	4.02	6.63
$\bar{\Lambda}\pi^-$	538 ± 25	39.4	9.80	8.81	13.64
$\bar{\Lambda}\pi^-\pi^0$	1301 ± 49	17.9	3.93	3.73	5.82
$\bar{\Lambda}\pi^-\pi^+\pi^-$	639 ± 34	14.6	2.73	2.63	3.82
$\bar{\Sigma}^0\pi^-$	371 ± 22	24.3	6.12	5.75	8.4
$\bar{\Sigma}^-\pi^0$	251 ± 24	22.2	5.62	5.18	7.90
$\bar{\Sigma}^-\pi^-\pi^+$	956 ± 48	22.2	5.25	5.07	8.16
$\bar{p}\pi^-\pi^+$	459 ± 47	55.9	⋯	⋯	19.66

References

[1] L. L. Chau, H. Y. Cheng, and B. Tseng, Phys. Rev. D **54**, 2132 (1996)
[2] H. Y. Cheng, Chin. J. Phys. **78**, 324 (2022)
[3] C. D. Lü, W. Wang, and F. S. Yu, Phys. Rev. D **93**, 056008 (2016)
[4] C. Q. Geng, Y. K. Hsiao, C. W. Liu *et al.*, J. High Energy Phys. **11**, 147 (2017)
[5] C. Q. Geng, Y. K. Hsiao, C. W. Liu *et al.*, Eur. Phys. J. C **78**, 593 (2018)
[6] C. Q. Geng, Y. K. Hsiao, C. W. Liu *et al.*, Phys. Rev. D **97**, 073006 (2018)
[7] C. Q. Geng, Y. K. Hsiao, C. W. Liu *et al.*, Phys. Rev. D **99**, 073003 (2019)
[8] C. Q. Geng, C. W. Liu, T. H. Tsai *et al.*, Phys. Rev. D **99**, 114022 (2019)
[9] X. G. He, Y. J. Shi, and W. Wang, Eur. Phys. J. C **80**, 359 (2020)
[10] H. J. Zhao, Y. L. Wang, Y. K. Hsiao *et al.*, J. High Energy Phys. **02**, 165 (2020)
[11] H. B. Li and X. R. Lyu, Natl. Sci. Rev. **8**, nwab181 (2021)
[12] M. Ablikim *et al.* (BESIII Collaboration), Phys. Rev. Lett. **116**, 052001 (2016)

[13] M. Ablikim *et al*. (BESIII Collaboration), Phys. Rev. D **95**, 111102 (2017)
[14] M. Ablikim *et al*. (BESIII Collaboration), Phys. Rev. Lett. **118**, 112001 (2017)
[15] M. Ablikim *et al*. (BESIII Collaboration), Phys. Lett. B **783**, 200 (2018)
[16] M. Ablikim *et al*. (BESIII Collaboration), Chin. Phys. C **43**, 083002 (2019)
[17] M. Ablikim *et al*. (BESIII Collaboration), Phys. Rev. D **99**, 032010 (2019)
[18] M. Ablikim *et al*. (BESIII Collaboration), Phys. Rev. Lett. **128**, 142001 (2022)
[19] R. Aaij *et al*. (LHCb Collaboration), Phys. Rev. D **97**, 091101 (2018)
[20] S. X. Li *et al*. (Belle Collaboration), Phys. Rev. D **103**, 072004 (2021)
[21] S. X. Li *et al*. (Belle Collaboration), Phys. Rev. D **104**, 072008 (2021)
[22] M. Gronau and J. L. Rosner, Phys. Rev. D **97**, 116015 (2018)
[23] P. A. Zyla *et al*. (Particle Data Group), Prog. Theor. Exp. Phys. **2022**, 083C01 (2022)
[24] C. P. Jia, D. Wang, and F. S. Yu, Nucl. Phys. B **956**, 115048 (2020)
[25] M. Saur and F. S. Yu, Sci. Bull. **65**, 1428 (2020)
[26] K. K. Sharma and R. C. Verma, Phys. Rev. D **55**, 7067 (1997)
[27] J. Adler *et al*. (MARKIII Collaboration), Phys. Rev. Lett. **62**, 1821 (1989)
[28] M. Ablikim *et al*. (BESIII Collaboration), Chin. Phys. C **40**, 063001 (2016)
[29] M. Ablikim *et al*.(BESIII Collaboration), Chin. Phys. C **39**, 093001 (2015)
[30] M. Ablikim *et al*. (BESIII Collaboration), Chin. Phys. C **46**, 113003 (2022)
[31] M. Ablikim *et al*. (BESIII Collaboration), Chin. Phys. C **46**, 113002 (2022)
[32] M. Ablikim *et al*. (BESIII Collaboration), Nucl. Instrum. Meth. A **614**, 345 (2010)
[33] C. H. Yu *et al.*, Proceedings of IPAC2016, Busan, Korea, 2016
[34] M. Ablikim *et al*. (BESIII Collaboration), Chin. Phys. C **44**, 040001 (2020)
[35] X. Li *et al*., Radiat. Detect. Technol. Methods **1**, 13 (2017)
[36] Y. X. Guo *et al*., Radiat. Detect. Technol. Methods **1**, 15 (2017)
[37] P. Cao *et al*., Nucl. Instrum. Meth. A **953**, 163053 (2020)
[38] S. Jadach, B. F. L. Ward and Z. Was, Phys. Rev. D **63**, 113009 (2001)
[39] D. J. Lange, Nucl. Instrum. Meth. A **462**, 152 (2001)
[40] R. G. Ping, Chin. Phys. C **32**, 599 (2008)
[41] J. C. Chen, G. S. Huang, X. R. Qi, D. H. Zhang and Y. S. Zhu, Phys. Rev. D **62**, 034003 (2000)
[42] S. Agostinelli *et al*. (GEANT4 Collaboration), Nucl. Instrum. Meth. A **506**, 250 (2003)
[43] J. Allison *et al*., IEEE Trans. Nucl. Sci. **53**, 270 (2006)
[44] H. Albrecht *et al*. (ARGUS Collaboration), Phys. Lett. B **241**, 278 (1990)