

Competing nonlinearities in NLS equations as source of threshold phenomena on star graphs

Original

Competing nonlinearities in NLS equations as source of threshold phenomena on star graphs / Adami, Riccardo; Boni, Filippo; Dovetta, Simone. - In: JOURNAL OF FUNCTIONAL ANALYSIS. - ISSN 0022-1236. - ELETTRONICO. - 283:1(2022), pp. 1-34. [10.1016/j.jfa.2022.109483]

Availability:

This version is available at: 11583/2959508 since: 2022-03-25T15:09:53Z

Publisher:

Elsevier

Published

DOI:10.1016/j.jfa.2022.109483

Terms of use:

This article is made available under terms and conditions as specified in the corresponding bibliographic description in the repository

Publisher copyright

Elsevier postprint/Author's Accepted Manuscript

© 2022. This manuscript version is made available under the CC-BY-NC-ND 4.0 license
<http://creativecommons.org/licenses/by-nc-nd/4.0/>. The final authenticated version is available online at:
<http://dx.doi.org/10.1016/j.jfa.2022.109483>

(Article begins on next page)

COMPETING NONLINEARITIES IN NLS EQUATIONS AS SOURCE OF THRESHOLD PHENOMENA ON STAR GRAPHS

RICCARDO ADAMI¹, FILIPPO BONI^{2,3}, AND SIMONE DOVETTA⁴

ABSTRACT. We investigate the existence of ground states for the nonlinear Schrödinger Equation on star graphs with two subcritical focusing nonlinear terms: a standard power nonlinearity, and a delta-type nonlinearity located at the vertex. We find that if the standard nonlinearity is stronger than the pointwise one, then ground states exist for small mass only. On the contrary, if the pointwise nonlinearity prevails, then ground states exist for large mass only. All ground states are radial, in the sense that their restriction to each half-line is always the same function, and coincides with a soliton tail. Finally, if the two nonlinearities are of the same size, then the existence of ground states is insensitive to the value of the mass, and holds only on graphs with a small number of half-lines.

Furthermore, we establish the orbital stability of the branch of radial stationary states to which the ground states belong, also in the mass regimes in which there is no ground state.

AMS Subject Classification: 35R02, 35Q55, 81Q35, 35Q40, 49J40.

Keywords: nonlinear Schrödinger, star graphs, nonlinear delta, standing waves, threshold phenomena

1. INTRODUCTION

In this paper we investigate existence and uniqueness of ground states for the NLS energy functional

$$F_{p,q}(u) = \frac{1}{2} \int_{S_N} |u'|^2 dx - \frac{1}{p} \int_{S_N} |u|^p dx - \frac{1}{q} |u(0)|^q, \quad (1)$$

with the mass constraint

$$\int_{S_N} |u|^2 dx = \mu, \quad (2)$$

where S_N is the star graph made of N half-lines glued together at their common origin, which in the following will be denoted by v or, alternatively, by 0 .

Each half-line \mathcal{H}_i of S_N is provided with a coordinate $x_i : \mathcal{H}_i \rightarrow [0, +\infty)$, so that $x_i = 0$ corresponds to the origin of \mathcal{H}_i , for every $i = 1, \dots, N$. A function u on S_N is given by the collection of its restrictions u_i to each half-line \mathcal{H}_i , i.e. $u = (u_1, \dots, u_N)$. We then define $L^p(S_N)$ as the space of all functions u on S_N such that, for every i , $u_i \in L^p(\mathcal{H}_i)$ with respect to the standard Lebesgue measure on \mathcal{H}_i . We define

$$\|u\|_{L^p(S_N)}^p := \sum_{i=1}^N \|u_i\|_{L^p(\mathcal{H}_i)}^p.$$

Similarly, we set $H^1(S_N)$ to be the space of all functions u on S_N that are continuous on the graph, in particular at v , and such that $u_i \in H^1(\mathcal{H}_i)$ for every i , endowed with the

norm

$$\|u\|_{H^1(S_N)}^2 := \sum_{i=1}^N \|u_i\|_{H^1(\mathcal{H}_i)}^2.$$

Introducing the notation

$$H_\mu^1(S_N) := \{u \in H^1(S_N) : \|u\|_{L^2(S_N)}^2 = \mu\}$$

for the mass constrained space, and

$$\mathcal{F}_{p,q}(\mu) := \inf_{v \in H_\mu^1(S_N)} F_{p,q}(v) \quad (3)$$

for the ground state energy level of (1) in $H_\mu^1(S_N)$, we define a ground state of (1) at mass μ as a minimizer of the energy among functions with mass μ , i.e. $u \in H_\mu^1(S_N)$ such that $F_{p,q}(u) = \mathcal{F}_{p,q}(\mu)$.

The aim of the paper is thus to understand whether ground states exist and are unique in the L^2 subcritical regime for both nonlinearities, namely $p \in (2, 6)$ and $q \in (2, 4)$.

Since their first appearance as a model for valence electrons in naphthalene molecules [66], the study of dynamics on metric graphs (or networks) has grown through the decades to become a prominent line of research. To date, models involving Schrödinger operators have been gathering a considerable interest, both in the linear setting (see for instance [19, 24, 42, 50, 51] as well as the monograph [25] and references therein) and in the nonlinear case [1, 10–13, 26, 30, 31, 36–40, 44, 48, 49, 59–61, 63–65]. Recent investigations are now available also for the KdV equation [57] and the Dirac equation [27].

Within this framework, star graphs provide a prototypical model. Particularly, the issue of the existence of NLS ground states on star graphs has been widely investigated for the last ten years. As a first step, nonexistence on star graphs made of at least three half-lines with free (or Kirchhoff's) conditions at the origin was established in [2], so that non-trivial conditions are required in order to have existence. The effect of an attractive linear delta interaction at the vertex, that corresponds to $q = 2$ in (1), was then studied in [3, 4], finding that ground states exist for small values of the mass only, and they bifurcate from the bound state of the corresponding linear Schrödinger Equation. Such ground states are always radial, in the sense that their restriction to every half-line of the graph always coincide with the same function. Moreover, the stability of the family of radial stationary states was proved even for the values of the mass for which such functions are no longer ground states. Further stability analysis on star graphs with linear pointwise interaction at the origin have then been developed in [17, 18, 43]. In [32] the result of existence of a nonlinear ground state bifurcating from the linear one was extended to the presence of a further linear potential.

Here we introduce nonlinear vertex conditions, more specifically conditions mimicking the nonlinear delta potential introduced in [16], and recently studied for the same problem on the line [28].

The concentrated nonlinearity is nowadays a widely accepted model of the net effect of the confinement of charges in small regions [47, 54], as well as in the study of resonant tunneling [58]. Related models have been originally discussed in dimension one [15, 16] and three [8, 9], and more recently the analysis has been broadened to the two-dimensional case [6, 7, 33] and to non-compact metric graphs [41, 67, 68, 72]. On the other hand, starting with the seminal paper [71], the study of the interplay between different nonlinearities for NLS equations has been recently carried out for instance in [35, 46, 52, 53, 55, 56, 69, 70],

focusing on the case of two different power nonlinearities. Particularly, [46, 69, 70] are devoted to the problem of the existence, the shape and the stability of prescribed mass ground states, defined as minimizers of the energy among all stationary states at given mass (see also the series of works [20–23] and [62] for the case of NLS systems).

In the present paper, we describe the effect of the combined action of two focusing nonlinearities again of the power type, with a substantial difference with respect to the cited ones, namely the fact that the second nonlinearity is concentrated at a point: the vertex of the graph.

The interaction between this two kinds of nonlinearity has already been explored in the case of the line, keeping the two powers at the subcritical or L^2 -critical level [28]. A non trivial interplay emerges in the critical case, with a modification of the critical mass. For the graph model treated here, the emerging scenario is considerably richer and surprising, as the first of our main theorems starts unravelling.

Theorem 1.1. *Let $N \geq 3$, $p \in (2, 6)$, $q \in (2, 4)$ and $q \neq \frac{p}{2} + 1$. Then there exists a critical mass $\mu_{p,q} > 0$ such that*

- (i) *if $q < \frac{p}{2} + 1$, then ground states of (1) at mass μ exist if and only if $\mu \leq \mu_{p,q}$;*
- (ii) *if $q > \frac{p}{2} + 1$, then ground states of (1) at mass μ exist if and only if $\mu \geq \mu_{p,q}$.*

Furthermore, whenever they exist, ground states at prescribed mass are unique and they are radial and decreasing on S_N , in the sense that their restriction to each half-line of the graph corresponds to the same decreasing function on \mathbb{R}^+ .

This result highlights the emergence of a natural comparison between the strength of the two nonlinearities. If $q < p/2 + 1$, then the point interaction is weak, so that it does not change the qualitative information gained in [4], that ground states exist only for masses below a critical value. Conversely, if $q > p/2 + 1$, then the point interaction is strong enough to reverse the result, so that ground states exist only for masses above a critical value.

The threshold phenomena in Theorem 1.1 reveals a natural scaling for the doubly non-linear problem on graphs. Heuristically, one may interpret such a feature as the result of the competition between the two terms defining $F_{p,q}$: the standard NLS energy on the one side

$$\frac{1}{2} \int_{\mathcal{G}} |u'|^2 dx - \frac{1}{p} \int_{\mathcal{G}} |u|^p dx, \quad (4)$$

and the delta nonlinearity on the other side

$$-\frac{1}{q} |u(0)|^q.$$

It is well-known (see for instance [11]) that the unique minimizer of the standard NLS energy at mass μ on the real line, the so-called soliton ϕ as in (10) below, verifies

$$\frac{1}{2} \int_{\mathbb{R}} |\phi'|^2 dx - \frac{1}{p} \int_{\mathbb{R}} |\phi|^p dx \sim \mu^{2\beta+1}, \quad \beta = \frac{p-2}{6-p},$$

whereas

$$\|\phi\|_{L^\infty(\mathbb{R})} \sim \mu^\alpha, \quad \alpha = \frac{2}{6-p}.$$

An elementary calculation then shows that

$$\alpha q = 2\beta + 1 \iff q = \frac{p}{2} + 1,$$

so that the comparison between the standard and the pointwise nonlinearity reported in Theorem 1.1 corresponds to the balance between the standard NLS energy and the maximal delta nonlinearity for the soliton on the line. However, at a first sight, the detailed result in Theorem 1.1, with that unprecedented dependence of the existence of ground states on the mass, sounds puzzling and requires a qualitative explanation. As well understood since [2], fixed a value for the mass, the existence of ground states for the NLSE on metric graphs is ruled by the competition between the soliton on the line and the standing wave with the lowest energy among all stationary states. Indeed, even though the presence of the vertex together with the prescription of continuity exclude the solitons from the family of possible competitors, one can always approximate its standard energy (4) arbitrarily well, for instance by a sequence of solitons supported on one half-line and truncated near the vertex. The approximation is made better and better by moving faraway from the vertex. This reasoning shows that, if the lowest-energy standing wave wins the competition, namely, if its energy is lower than the standard energy of the soliton on the line (4), then the ground state exists. Viceversa, if the energy of the soliton is lower than the energy of every stationary state, then the ground state does not exist. This is nowadays well understood.

Of course, the outcome of the competition between solitons and standing waves can depend on the mass, and this is the new phenomenon put in evidence, in a surprising way, by Theorem 1.1. We recall that such a phenomenon is not present in the analogous model on the line, since in that case the soliton is an admissible competitor and its doubly nonlinear energy is strictly smaller than the standard one (4). As a consequence, for the problem on the line a ground state exists for every value of the mass and is centred at the origin [28]. On the contrary, in Theorem 1.1 a dramatic dependence on the mass emerges as the graph structure plays a crucial role. Indeed, in order to lower the energy, from one side it is convenient to exploit the presence of the interaction at the vertex, but from the other it is convenient to escape it, as the proximity to the vertex increases the kinetic energy (quantitative estimates of this effect can be found by using rearrangement theory, see [11, Proposition 3.1]). By this observation, it becomes possible to interpret the role of the mass in Theorem 1.1: if $q > p/2 + 1$, then for large masses, and consequently large values at the vertex of the radial standing waves, the point interaction prevails and a ground state exists. On the other hand, for small mass, the high power on a small value at the vertex makes the contribution of the point interaction tiny, so that the standard nonlinearity prevails, escaping the vertex becomes convenient, and then there is no ground state. Viceversa, in the case of weaker delta interactions, namely $q < p/2 + 1$, for small mass, and then small value at the vertex, the lower power in the pointwise term gives to the interaction a larger size than that of the standard nonlinearity, so it is convenient to stay on the vertex and a ground state exists. On the other hand, if the mass is large, then the effect of the standard nonlinearity is magnified more, so that it is convenient to escape the vertex and there is no ground state. Of course, this description is qualitative, but is made rigorous by the computations in Section 4.

It remains to investigate what happens at the threshold $q = p/2 + 1$, and the answer is given by the following theorem, bearing another unexpected feature: ground states exist only for a small number of half-lines, and, notably, such a number does not depend on the mass.

Theorem 1.2. *Let $p \in (2, 6)$ and $q = \frac{p}{2} + 1$. Then there exists a critical number of half-lines $N_p \geq 2$ such that*

- (i) if $N \leq N_p$, then ground states of (1) at mass μ exist for every μ ;
- (ii) if $N > N_p$, then ground state of (1) at mass μ never exist.

Furthermore, whenever they exist, ground states at prescribed mass are unique and they are radial and decreasing on S_N , in the sense that their restriction to each half-line of the graph corresponds to the same decreasing function on \mathbb{R}^+ .

The number N_p of half-lines above which there is never a ground state is not presently known, but on the basis of numerical simulations (see Section 4) we have evidence that it is at least three for every value of p . However, at the time being we can prove analytically only the following result.

Proposition 1.3. *For every $p \in (2, 6)$, let N_p be as in Theorem 1.2. Then:*

- (i) $\lim_{p \rightarrow 2^+} N_p = +\infty$;
- (ii) if $p \geq 4$, then $N_p \in \{2, 3, 4\}$. Furthermore, there exists $\delta > 0$ so that $N_p = 3$ for every $p \in (4 - \delta, 4 + \delta) \cup (6 - \delta, 6)$.

To conclude, recall that to every critical point $u \in H_\mu^1(S_N)$ of (1) subject to the mass constraint (2) it corresponds a standing wave solution $\psi(x, t) := e^{i\omega t}u(x)$ to the time-dependent NLS equation on S_N

$$i\partial_t \psi(t, x) = -\partial_{xx}^2 \psi(t, x) - |\psi(t, x)|^{p-2} \psi(t, x) - |\psi(t, x)|^{q-2} \delta_0 \psi(t, x). \quad (5)$$

The standard stability theory in [45] guarantees that ψ is orbitally stable if and only if u is a local minimizer of $F_{p,q}$ in $H_\mu^1(S_N)$, so that Theorems 1.1–1.2 imply the existence of an orbitally stable standing wave of (5) whenever ground states exist. Adapting the argument originally developed in [3], the last proposition of this paper improves this result, by proving the orbital stability for all elements of the branch of radial stationary states, parametrized by the mass. Theorems 1.1–1.2 establish that for some values of the mass there is no ground state. However, even for those values the branch of radial standing waves exists, and we prove that they are always orbitally stable.

Proposition 1.4. *For every $p \in (2, 6)$, $q \in (2, 4)$ and every $\mu > 0$, the unique radially-symmetric stationary state of (5) is a local minimizer of $F_{p,q}$ in $H_\mu^1(S_N)$, and the associated standing wave is always orbitally stable.*

The paper is organized as follows. Section 2 reviews some preliminary results on nonlinear Schrödinger equations with the standard nonlinearity only, and it develops the analysis of the stationary states of the doubly nonlinear model on S_N . Within Section 3 we derive a existence criterion and we show that, whenever they exist, ground states coincide with the unique radial and decreasing stationary state. Finally, Section 4 completes the proofs of Theorems 1.1–1.2 and of Proposition 1.3, while Section 5 addresses that of Proposition 1.4.

Notation. In what follows, when denoting a norm, we omit the domain of integration whenever it is understood, writing for instance $\|u\|_p$ instead of $\|u\|_{L^p(S_N)}$. The complete notation will be used if needed to avoid ambiguity.

2. PRELIMINARIES

In this section we briefly recall some well-known facts about nonlinear Schrödinger equations on the real line and we discuss some preliminary properties of stationary solutions on star graphs that will be important in the forthcoming analysis.

2.1. NLSE with standard nonlinearity on the real line. The minimization problem on the real line

$$\mathcal{E}(\mu) := \inf_{v \in H_\mu^1(\mathbb{R})} E(v, \mathbb{R}),$$

where $E : H^1(\mathbb{R}) \rightarrow \mathbb{R}$ is the NLS energy functional involving the standard nonlinearity only

$$E(v) := \frac{1}{2} \int_{\mathbb{R}} |v'(x)|^2 dx - \frac{1}{p} \int_{\mathbb{R}} |v(x)|^p dx, \quad (6)$$

is nowadays classical (see for instance [34]).

Standard variational arguments show that ground states are solutions to the stationary nonlinear Schrödinger equation

$$u'' + |u|^{p-2}u = \omega u \quad \text{on } \mathbb{R} \quad (7)$$

for some $\omega > 0$. In fact, for every $\omega > 0$ the unique (up to translations) positive solution $\phi_\omega \in H^1(\mathbb{R})$ of (7) is

$$\phi_\omega(x) = \left[\frac{p}{2} \omega \left(1 - \tanh^2 \left(\left(\frac{p}{2} - 1 \right) \sqrt{\omega} |x| \right) \right) \right]^{\frac{1}{p-2}}. \quad (8)$$

The mass of ϕ_ω is given explicitly by

$$\|\phi_\omega\|_2^2 = \frac{4 \left(\frac{p}{2} \right)^{\frac{2}{p-2}} \omega^{\frac{6-p}{2(p-2)}}}{p-2} \int_0^1 (1-s^2)^{\frac{4-p}{p-2}} ds, \quad (9)$$

which is a continuous, strictly increasing and unbounded function of ω . Therefore, for every $\mu > 0$ there exists a unique $\omega(\mu)$ such that $\phi_{\omega(\mu)}$ is the unique (up to translations) positive ground state of E in $H_\mu^1(\mathbb{R})$. Such ground states are called solitons, and their dependence on μ is given by

$$\phi_{\omega(\mu)}(x) = C_p \mu^\alpha \operatorname{sech}^{\frac{\alpha}{\beta}} \left(c_p \mu^\beta x \right), \quad \alpha = \frac{2}{6-p}, \beta = \frac{p-2}{6-p}, \quad (10)$$

where $C_p, c_p > 0$ depends on p only, and one can easily compute

$$\mathcal{E}(\mu) = E(\phi_{\omega(\mu)}) = -\theta_p \mu^{2\beta+1} \quad (11)$$

where $\theta_p > 0$ depends on p only.

In the following, for $u \in H^1(S_N)$, we will denote by $E(u, S_N)$ the analogous standard NLS energy as in (6)

$$E(u, S_N) := \frac{1}{2} \int_{S_N} |u'|^2 dx - \frac{1}{p} \int_{S_N} |u|^p dx = \frac{1}{2} \sum_{i=1}^N \int_{\mathcal{H}_i} |u'_i|^2 dx - \frac{1}{p} \sum_{i=1}^N \int_{\mathcal{H}_i} |u_i|^p dx.$$

Note that, for every positive $u \in H_\mu^1(S_N)$, it holds

$$E(u, S_N) > \mathcal{E}(\mu). \quad (12)$$

Indeed, if u is compactly supported on a unique half-line of S_N , then one can regard it as a compactly supported function in $H_\mu^1(\mathbb{R})$ and (12) is immediate since only solitons attain $\mathcal{E}(\mu)$. On the contrary, if $u \not\equiv 0$ on at least two half-lines of S_N , then there exist infinitely many values t in the image of u such that the number of pre-images

$$N(t) := \#\{x \in S_N : u(x) = t\} \geq 2.$$

If $u(0) \neq 0$ then this is trivial, as u tends to 0 along each half-line. Similarly, if $u(0) = 0$, then all values realized by u in a suitably small neighbourhood of the origin are attained

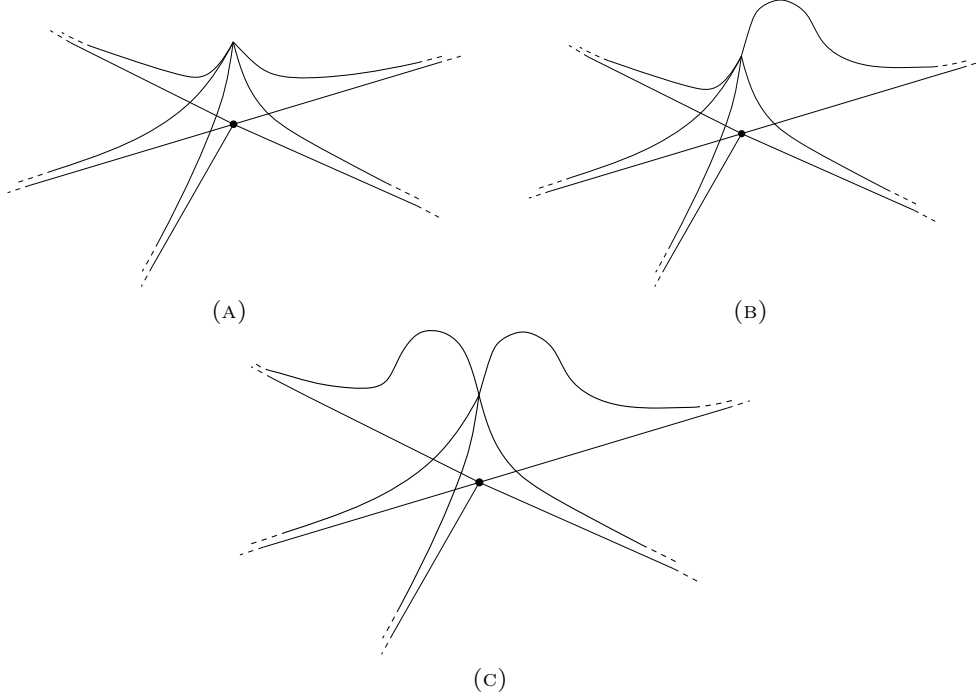


FIGURE 1. The stationary states of Proposition 2.1 on S_5 : η_0^ω (A), η_1^ω (B) and η_2^ω (C).

at least twice the number of half-lines that belong to the support of u . Hence, letting $\hat{u} \in H_\mu^1(\mathbb{R})$ denote the symmetric rearrangement of u , it follows (see [11, Proposition 3.1])

$$\|u'\|_{L^2(S_N)} > \|\hat{u}'\|_{L^2(\mathbb{R})} \quad \text{and} \quad \|u\|_{L^p(S_N)} = \|\hat{u}\|_{L^p(\mathbb{R})}, \quad p \geq 1,$$

yielding again (12).

2.2. Some remarks on stationary states. Computing the Euler–Lagrange equations associated to the energy (1) and to the mass constraint (2), it turns out that ground states of (1) at given mass are solutions to the system

$$\begin{cases} u_i'' + u_i|u_i|^{p-2} = \omega u_i & \text{on } \mathcal{H}_i \text{ for all } i = 1, \dots, N, \\ \sum_{i=1}^N \frac{du_i}{dx_i}(0^+) = -u(0)|u(0)|^{q-2}, \end{cases} \quad (13)$$

for some Lagrange multiplier $\omega > 0$. Note that here the expression $\frac{du_i}{dx_i}(0^+)$ naturally corresponds to the right derivative of u_i at $x = 0$, since each edge of the graph is identified with the interval $[0, +\infty)$ so that the vertex always corresponds to the point $x = 0$.

The next proposition provides a complete characterization of the set of positive solutions of (13) in $H^1(S_N)$.

Proposition 2.1. *Let $\omega > 0$. Then the set \mathcal{S}_ω of positive solutions of (13) in $H^1(S_N)$ is given by*

$$\mathcal{S}_\omega = \left\{ \eta_{J,J}^\omega : J \in \mathbb{N}, 0 \leq J \leq \frac{N-1}{2} \right\}.$$

Here, $\eta_J^\omega \in H^1(S_N)$ is such that

$$\eta_J^\omega(\cdot) = \begin{cases} \phi_\omega(\cdot - a) & \text{on } J \text{ half-lines of } S_N \\ \phi_\omega(\cdot + a) & \text{on the other } N - J \text{ half-lines of } S_N, \end{cases} \quad (14)$$

where $a > 0$ is the unique positive solution to

$$\frac{\tanh\left(\left(\frac{p}{2} - 1\right)\sqrt{\omega}a\right)}{\left(1 - \tanh^2\left(\left(\frac{p}{2} - 1\right)\sqrt{\omega}a\right)\right)^{\frac{q-2}{p-2}}} = \frac{\left(\frac{p}{2}\right)^{\frac{q-2}{p-2}} \omega^{\frac{2q-2-p}{2(p-2)}}}{N - 2J}. \quad (15)$$

Furthermore, for every J , the function $\omega \mapsto \|\eta_J^\omega\|_2^2$ is continuous and strictly increasing.

Proof. Relying on the discussion of the previous subsection, it is immediate to see that, on each half-line \mathcal{H}_i , $i = 1, \dots, N$, solutions of (13) coincide with the restriction $\phi_\omega(x - a_i)$ of the soliton ϕ_ω , for suitable $a_i \in \mathbb{R}$. Since the continuity condition at the origin is prescribed, then $\phi_\omega(-a_i) = \phi_\omega(-a_j)$ for $1 \leq i, j \leq N$, therefore $a_i = \epsilon_i a$, where $a > 0$ and $\epsilon_i = \text{sgn}(a_i)$ for all $1 \leq i \leq N$.

Moreover, by (8), the second line of (13) becomes

$$\sqrt{\omega} \tanh\left(\left(\frac{p}{2} - 1\right)\sqrt{\omega}a\right) \sum_{i=1}^N \epsilon_i = -\phi_\omega(a)^{q-2} < 0, \quad (16)$$

thus implying $\sum_{i=1}^N \epsilon_i < 0$, namely the number of positive ϵ_i cannot exceed $\frac{N-1}{2}$.

Let then J be a given integer such that $0 \leq J \leq \frac{N-1}{2}$ and let $\epsilon_i = 1$ if and only if $i \leq J$. Since then $\sum_{i=1}^N \epsilon_i = 2J - N$, (16) reads

$$\tanh\left(\left(\frac{p}{2} - 1\right)\sqrt{\omega}a\right) = \frac{\phi_\omega(a)^{q-2}}{\sqrt{\omega}(N - 2J)}, \quad (17)$$

so that (15) is proved.

Setting $t := \tanh\left(\left(\frac{p}{2} - 1\right)\sqrt{\omega}a\right)$, then $t \in [0, 1)$ and we can rewrite (15) as

$$f(t) := \frac{t}{(1 - t^2)^{\frac{q-2}{p-2}}} = \frac{\left(\frac{p}{2}\right)^{\frac{q-2}{p-2}} \omega^{\frac{2q-2-p}{2(p-2)}}}{N - 2J}. \quad (18)$$

As $f(0) = 0$, $\lim_{t \rightarrow 1^-} f(t) = +\infty$ and $f'(t) > 0$ for every $t \in (0, 1)$, it follows that there exists a unique $t \in (0, 1)$ for which (18) is satisfied, in turn showing that for every $\omega > 0$ and every integer $0 \leq J \leq \frac{N-1}{2}$, there exists a unique positive solution η_J^ω of (13), up to exchange of edges.

Relying again on (8), we get

$$\|\eta_J^\omega\|_2^2 = \frac{2\left(\frac{p}{2}\right)^{\frac{2}{p-2}} \omega^{\frac{6-p}{2(p-2)}}}{p-2} \left[2J \int_0^1 (1 - s^2)^{\frac{4-p}{p-2}} ds + (N - 2J) \int_t^1 (1 - s^2)^{\frac{4-p}{p-2}} ds \right]. \quad (19)$$

Differentiating with respect to ω yields

$$\begin{aligned} \frac{d}{d\omega} \|\eta_J^\omega\|_2^2 &= \frac{\left(\frac{p}{2}\right)^{\frac{2}{p-2}} (6-p)}{(p-2)^2} 2J \omega^{\frac{6-p}{2(p-2)}-1} \int_0^1 (1 - s^2)^{\frac{4-p}{p-2}} ds \\ &\quad + \frac{2\left(\frac{p}{2}\right)^{\frac{2}{p-2}}}{p-2} (N - 2J) \omega^{\frac{6-p}{2(p-2)}-1} \left[\frac{6-p}{2(p-2)} \int_t^1 (1 - s^2)^{\frac{4-p}{p-2}} ds - \omega(1 - t^2)^{\frac{4-p}{p-2}} t'(\omega) \right], \end{aligned} \quad (20)$$

where, by (18),

$$t'(\omega) = \frac{\left(\frac{p}{2}\right)^{\frac{q-2}{p-2}} (2q-2-p)}{2(p-2)(N-2J)} \omega^{\frac{2q-2-p}{2(p-2)}-1} \frac{(1-t^2)^{\frac{q-2}{p-2}+1}}{t^2 \left(2\frac{q-2}{p-2} - 1\right) + 1}. \quad (21)$$

Note that the first term in the sum on the right hand side of (20) is strictly positive for every $\omega > 0$. Furthermore, plugging (21) into (20) and making use of (18) allows to rewrite the term between square brackets as

$$\frac{6-p}{2(p-2)} \int_t^1 (1-s^2)^{\frac{4-p}{p-2}} ds - \frac{2q-2-p}{2(p-2)} \frac{t(1-t^2)^{\frac{2}{p-2}}}{t^2 \left(2\frac{q-2}{p-2} - 1\right) + 1},$$

which can be proved to be strictly positive for every $t \in (0, 1)$ by the same calculations in [28, pp.11–12]. Hence, $\frac{d}{d\omega} \|\eta_J^\omega\|_2^2 > 0$ for every $\omega > 0$ and $\|\eta_J^\omega\|_2^2$ is a strictly increasing function of ω . \square

The next straightforward corollary describes the set of stationary states at prescribed mass.

Corollary 2.2. *Let $\mu > 0$ be fixed. Then for every integer J such that $0 \leq J \leq \frac{N-1}{2}$ there exists a unique $\omega > 0$ such that $\eta_J^\omega \in H_\mu^1(S_N)$. Moreover, it holds*

$$F_{p,q}(\eta_J^\omega) = -\frac{6-p}{2(p+2)} \omega \mu + \left(\frac{2}{p+2} - \frac{1}{q}\right) |\eta_J^\omega(0)|^q. \quad (22)$$

Proof. Fix J integer such that $0 \leq J \leq \frac{N-1}{2}$. By Proposition 2.1, $\|\eta_J^\omega\|_2^2$ is an unbounded, strictly increasing continuous function of ω , thus implying that for every $\mu > 0$ there exists a unique value of ω for which $\|\eta_J^\omega\|_2^2 = \mu$, i.e. $\eta_J^\omega \in H_\mu^1(S_N)$.

Since η_J^ω solves (13), multiplying the equation by $(\eta_J^\omega)'$ and, for every $x \in S_N$, integrating on $[x, +\infty)$ one obtains

$$\frac{1}{2} ((\eta_J^\omega)')^2(x) = \frac{\omega}{2} (\eta_J^\omega)^2(x) - \frac{1}{p} (\eta_J^\omega)^p(x),$$

so that integrating on S_N gives

$$\frac{1}{2} \|(\eta_J^\omega)'\|_2^2 = \frac{\omega}{2} \|\eta_J^\omega\|_2^2 - \frac{1}{p} \|\eta_J^\omega\|_p^p. \quad (23)$$

Furthermore, multiplying the first line of (13) by η_J^ω , integrating over S_N and making use of the second line of (13) yields

$$|\eta_J^\omega(0)|^q - \|(\eta_J^\omega)'\|_2^2 + \|\eta_J^\omega\|_p^p - \omega \|\eta_J^\omega\|_2^2 = 0. \quad (24)$$

Combining (23) and (24) leads to (22) and we conclude. \square

3. EXISTENCE CRITERION AND CHARACTERIZATION OF GROUND STATES

In this section we provide a sufficient condition granting existence of ground states of $F_{p,q}$ in $H_\mu^1(S_N)$ and prove that, whenever they exist, such ground states must be monotonically decreasing radial functions on S_N .

Let us begin with a compactness result. To this aim we recall the Gagliardo–Nirenberg inequalities

$$\|u\|_p^p \leq K_p \|u\|_2^{\frac{p}{2}+1} \|u'\|_2^{\frac{p}{2}-1}, \quad p \geq 2, \quad (25)$$

where $K_p > 0$ depends only on p , and

$$\|u\|_\infty^2 \leq \|u\|_2 \|u'\|_2, \quad (26)$$

holding for every $u \in H^1(S_N)$ (we refer to [12] for a proof of these inequalities on general metric graphs).

Proposition 3.1. *For every $\mu > 0$ it holds*

$$\mathcal{F}_{p,q}(\mu) \leq \mathcal{E}(\mu). \quad (27)$$

Furthermore, if $\mathcal{F}_{p,q}(\mu) < \mathcal{E}(\mu)$, then ground states of (1) at mass μ exist.

Proof. Let us first prove (27). For every $\varepsilon > 0$, let $v_\varepsilon := \kappa_\varepsilon(\phi_\mu - \varepsilon)_+$, where ϕ_μ is as in (10) and $\kappa_\varepsilon > 0$ is chosen to guarantee $\|v_\varepsilon\|_{L^2(\mathbb{R})}^2 = \mu$, so that $v_\varepsilon \in H_\mu^1(\mathbb{R})$. Since $\|v_\varepsilon\|_{L^q(\mathbb{R})} \rightarrow \|\phi_\omega\|_{L^q(\mathbb{R})}$ as $\varepsilon \rightarrow 0$, for every $q \geq 1$, then $\kappa_\varepsilon \rightarrow 1$ for $\varepsilon \rightarrow 0$, and we get

$$\begin{aligned} \mathcal{E}(\mu) \leq E(v_\varepsilon, \mathbb{R}) &= \frac{1}{2} \kappa_\varepsilon^2 \int_{\mathbb{R}} |[(\phi_\mu - \varepsilon)_+]'|^2 dx - \frac{1}{p} \kappa_\varepsilon^p \int_{\mathbb{R}} |(\phi_\mu - \varepsilon)_+|^p dx \\ &\leq E(\phi_\omega, \mathbb{R}) + o(1) = \mathcal{E}(\mu) + o(1) \end{aligned}$$

for ε small enough, making use also of $\|v'_\varepsilon\|_{L^2(\mathbb{R})} \leq \|\phi_\omega\|_{L^2(\mathbb{R})}$. Hence, $E(v_\varepsilon, \mathbb{R}) \rightarrow \mathcal{E}(\mu)$ as $\varepsilon \rightarrow 0$. Moreover, v_ε has compact support, so that one can think of it as supported on any given half-line of S_N . We thus have

$$\mathcal{E}(\mu) = \lim_{\varepsilon \rightarrow 0^+} E(v_\varepsilon, \mathbb{R}) = \lim_{\varepsilon \rightarrow 0^+} F_{p,q}(v_\varepsilon) \geq \mathcal{F}_{p,q}(\mu),$$

so (27) is proved.

Assume now that $\mathcal{F}_{p,q}(\mu) < \mathcal{E}(\mu)$ and let $(u_n) \subset H_\mu^1(S_N)$ be a minimizing sequence for $F_{p,q}$. Plugging (25) and (26) into the definition of $F_{p,q}$ gives

$$F_{p,q}(u_n) \geq \frac{1}{2} \|u'_n\|_2^2 - \frac{K_p}{p} \mu^{\frac{p+2}{4}} \|u'_n\|_2^{\frac{p}{2}-1} - \frac{1}{q} \mu^{\frac{q}{4}} \|u'_n\|_2^{\frac{q}{2}}$$

which, since $p \in (2, 6)$, $q \in (2, 4)$, ensures that (u_n) is bounded in $H^1(S_N)$. Therefore there exists $u \in H^1(S_N)$ such that, up to subsequences, $u_n \rightharpoonup u$ weakly in $H^1(S_N)$, $u_n \rightarrow u$ in $L_{loc}^\infty(S_N)$ and consequently $u_n \rightarrow u$ a.e. in S_N .

Set $m := \|u\|_2^2$. By weak lower semicontinuity, we have $m \leq \mu$.

Assume $m = 0$, that is $u \equiv 0$. Then $u_n(0) \rightarrow 0$ as $n \rightarrow +\infty$, so that recalling (12) leads to

$$\mathcal{E}(\mu) > \mathcal{F}_{p,q}(\mu) = \lim_n F_{p,q}(u_n) = \lim_n E(u_n, S_N) \geq \mathcal{E}(\mu),$$

i.e., a contradiction. Hence, $u \not\equiv 0$ on S_N .

Suppose then that $0 < m < \mu$. By weak convergence in $H^1(S_N)$ of u_n to u , we get $\|u_n - u\|_2^2 = \mu - m + o(1)$ for $n \rightarrow +\infty$. On the one hand, since $p, q > 2$ and $\frac{\mu}{\|u_n - u\|_2^2} > 1$ for n sufficiently large,

$$\begin{aligned} \mathcal{F}_{p,q}(\mu) &\leq F_{p,q} \left(\sqrt{\frac{\mu}{\|u_n - u\|_2^2}} (u_n - u) \right) \\ &= \frac{1}{2} \frac{\mu}{\|u_n - u\|_2^2} \|u'_n - u'_n\|_2^2 - \frac{1}{p} \left(\frac{\mu}{\|u_n - u\|_2^2} \right)^{\frac{p}{2}} \|u_n - u\|_p^p \\ &\quad - \frac{1}{q} \left(\frac{\mu}{\|u_n - u\|_2^2} \right)^{\frac{q}{2}} |u_n(0) - u(0)|^q < \frac{\mu}{\|u_n - u\|_2^2} F_{p,q}(u_n - u), \end{aligned}$$

so that

$$\liminf_n F_{p,q}(u_n - u) \geq \frac{\mu - m}{\mu} \mathcal{F}_{p,q}(\mu). \quad (28)$$

On the other hand, an analogous reasoning leads to

$$\mathcal{F}_{p,q}(\mu) \leq F_{p,q} \left(\sqrt{\frac{\mu}{\|u\|_2^2}} u \right) < \frac{\mu}{\|u\|_2^2} F_{p,q}(u),$$

so

$$F_{p,q}(u) > \frac{m}{\mu} \mathcal{F}_{p,q}(\mu). \quad (29)$$

Moreover, it holds

$$F_{p,q}(u_n) = F_{p,q}(u_n - u) + F_{p,q}(u) + o(1). \quad (30)$$

Indeed, by $u'_n \rightharpoonup u'$ weakly in $L^2(S_N)$ and $u_n \rightarrow u$ in $L^\infty_{\text{loc}}(S_N)$, we have $\|u'_n - u'\|_2^2 = \|u'_n\|_2^2 - \|u'\|_2^2 + o(1)$ and $|(u_n - u)(0)|^q = o(1)$ as n is large enough. Furthermore, owing to the Brezis-Lieb lemma [29],

$$\|u_n\|_p^p = \|u_n - u\|_p^p + \|u\|_p^p + o(1).$$

Using now (28), (29) and (30), we get

$$\begin{aligned} \mathcal{F}_{p,q}(\mu) &= \lim_n F_{p,q}(u_n) = \lim_n F_{p,q}(u_n - u) + F_{p,q}(u) \\ &> \frac{\mu - m}{\mu} \mathcal{F}_{p,q}(\mu) + \frac{m}{\mu} \mathcal{F}_{p,q}(\mu) = \mathcal{F}_{p,q}(\mu), \end{aligned}$$

which is again a contradiction.

Henceforth, $m = \mu$ and $u \in H_\mu^1(S_N)$. In particular, $u_n \rightarrow u$ in $L^2(S_N)$ so that, (u_n) being bounded in $L^\infty(S_N)$, $u_n \rightarrow u$ in $L^p(S_N)$ as $n \rightarrow +\infty$. Thus, by weak lower semicontinuity

$$F_{p,q}(u) \leq \lim_n F_{p,q}(u_n) = \mathcal{F}_{p,q}(\mu),$$

that is u is a ground state of $F_{p,q}$ at mass μ . \square

Corollary 3.2. *Let $\mu > 0$ be fixed. If there exists $u \in H_\mu^1(S_N)$ such that $F_{p,q}(u) \leq \mathcal{E}(\mu)$, then ground states of (1) at mass μ exist.*

Proof. If $\mathcal{F}_{p,q}(\mu) = F_{p,q}(u)$ then u is a ground state at mass μ . Otherwise, $\mathcal{F}_{p,q}(\mu) < F_{p,q}(u) \leq \mathcal{E}(\mu)$ and a ground state of (1) at mass μ exists by Proposition 3.1. \square

Once existence of ground states is granted, the following proposition ensures uniqueness and provides a complete characterization of their symmetry properties.

Proposition 3.3. *Let $\mu > 0$ be such that $\mathcal{F}_{p,q}(\mu)$ is attained. Then the unique positive ground state of $F_{p,q}$ at mass μ is the stationary state $\eta_0^{\omega(\mu)}$ such that $\eta_0^{\omega(\mu)} \in H_\mu^1(S_N)$.*

To prove this proposition we need two auxiliary lemmas.

Lemma 3.4. *Let $u \in H_\mu^1(S_N)$, $u \geq 0$. Then there exists $u^* \in H_\mu^1(S_N)$, $u^* \geq 0$, such that $F_{p,q}(u^*) \leq F_{p,q}(u)$ and either*

- (i) u^* is symmetric with respect to the vertex and monotonically decreasing on each half-line, or
- (ii) u^* is symmetric with respect to the vertex and monotonically decreasing on $N - 1$ half-lines of S_N , whereas on the remaining half-line it is non-decreasing from the origin to a unique maximum point and then non-increasing from this point to infinity.

Proof. Given a non-negative function $u \in H_\mu^1(S_N)$, suppose first that $\|u\|_\infty$ is attained at least once on every half-line of S_N (this hypothesis includes the particular case in which the maximum is attained at the origin). Then all the values in the image of u are attained at least N times on the graph. Thus, letting $u^* \in H_\mu^1(S_N)$ be the symmetric rearrangement of u on S_N as defined in [3, Appendix A], standard properties of rearrangements give

$$\|u'\|_2 \geq \|(u^*)'\|_2, \quad \|u\|_p = \|u^*\|_p, \quad |u^*(0)| = \|u\|_\infty,$$

the inequality being strict unless u is symmetric with respect to the vertex and monotonically decreasing on each half-lines. Hence, $F_{p,q}(u^*) \leq F_{p,q}(u)$ and u^* is as in (i).

Assume now that $\|u\|_\infty$ is not attained on every half-line of the graph. Let us discuss separately the cases $u(0) = 0$ and $u(0) > 0$.

If $u(0) = 0$, then let $\|u\|_\infty$ be attained at $x_1 \in \mathcal{H}_1$ and denote by I the interval connecting the origin to x_1 along \mathcal{H}_1 . Consider the following construction.

First, let $\tilde{u} \in H^1(0, x_1)$ be the monotone rearrangement of the restriction $u|_I$ of u to I . By definition of monotone rearrangement and the Pólya–Szegő inequality, we have $\tilde{u}(0) = \|u\|_\infty$, $\tilde{u}(x_1) = 0$ and

$$\|u'\|_{L^2(I)} \geq \|\tilde{u}'\|_{L^2(0, x_1)}, \quad \|u\|_{L^p(I)} = \|\tilde{u}\|_{L^p(0, x_1)} \quad \forall p \geq 1.$$

Secondly, let $\bar{u} \in H^1(\mathbb{R}^+)$ be the monotone rearrangement of the restriction $u|_{S_N \setminus I}$ of u to $S_N \setminus I$, so that $\bar{u}(0) = \|u\|_\infty$ and, by usual estimates that hold for monotone rearrangements,

$$\|u'\|_{L^2(S_N \setminus I)} \geq \|\bar{u}'\|_{L^2(\mathbb{R}^+)}, \quad \|u\|_{L^p(S_N \setminus I)} = \|\bar{u}\|_{L^p(\mathbb{R}^+)} \quad \forall p \geq 1.$$

Define then $u^* : S_N \rightarrow \mathbb{R}$ to be

$$u^*(x) := \begin{cases} \tilde{u}(x_1 - x) & x \in I, \\ \bar{u}(x - x_1) & x \in \mathcal{H}_1 \setminus I, \\ 0 & \text{otherwise.} \end{cases}$$

By construction, it follows that $u^* \in H_\mu^1(S_N)$, u^* is as in (ii) and $F_{p,q}(u^*) \leq F_{p,q}(u)$.

To conclude, it remains to deal with the case $u(0) > 0$. Given this, let $J := \{x \in S_N : u(x) > u(0)\}$ and let $u|_J$ be the restriction of u to J . Note that $u(J) = (u(0), \|u\|_\infty]$ is connected and every $t \in u(J)$ is attained at least twice on S_N , except possibly $\|u\|_\infty$. Hence, denoting by $\hat{u} \in H^1(-L, L)$ the symmetric rearrangement of $u|_J$ on the interval $(-L, L)$, with $L := \frac{|J|}{2}$, we have (see [12, Proposition 3.1])

$$\|u'\|_{L^2(J)} \geq \|\hat{u}'\|_{L^2(-L, L)}, \quad \|u\|_{L^p(J)} = \|\hat{u}\|_{L^p(-L, L)} \quad \forall p \geq 1, \quad \hat{u}(0) = \|u\|_\infty.$$

Similarly, $u(S_N \setminus J)$ is connected and every value $t \in u(S_N \setminus J)$ is attained at least N times (once on each half-line). Therefore, letting $u^\dagger \in H^1(S_N)$ be the symmetric rearrangement on S_N of $u|_{S_N \setminus J}$ as in [3, Appendix A], we get

$$\|u'\|_{L^2(S_N \setminus J)} \geq \|(u^\dagger)'\|_{L^2(S_N)}, \quad \|u\|_{L^p(S_N \setminus J)} = \|u^\dagger\|_{L^p(S_N)} \quad p \geq 1, \quad u^\dagger(0) = u(0).$$

Let then I be the interval $[0, 2L]$ along \mathcal{H}_1 , and set $u^* : S_N \rightarrow \mathbb{R}$ to be

$$u^*(x) := \begin{cases} \hat{u}(x - L) & x \in I \\ u^\dagger(x - 2L) & x \in \mathcal{H}_1 \setminus I \\ u^\dagger(x) & \text{otherwise.} \end{cases}$$

By construction, $u^* \in H_\mu^1(S_N)$, u^* is as in (ii) and $F_{p,q}(u^*) \leq F_{p,q}(u)$, so that the proof is complete. \square

Lemma 3.5. *Consider $u \in H_\mu^1(S_N)$ that does not have a local maximum point at the origin. Then there exists $v \in H_\mu^1(S_N)$ such that $F_{p,q}(v) < F_{p,q}(u)$.*

Proof. Given $u \in H_\mu^1(S_N)$, let \bar{u} be the restriction of u to $\mathcal{H}_1 \cup \mathcal{H}_2$ and \tilde{u} be the restriction of u to $S_N \setminus (\mathcal{H}_1 \cup \mathcal{H}_2)$. Set also

$$\bar{\mu} := \int_{\mathcal{H}_1 \cup \mathcal{H}_2} |\bar{u}|^2 dx, \quad \tilde{\mu} := \int_{S_N \setminus (\mathcal{H}_1 \cup \mathcal{H}_2)} |\tilde{u}|^2 dx.$$

For $\varepsilon > 0$ small enough, let $\nu \in (-\varepsilon, \varepsilon)$. Since u has no maximum point at the origin, define

$$u_\nu(x) := \begin{cases} \sqrt{\frac{\bar{\mu}+\nu}{\bar{\mu}}} \bar{u}(x + T(\nu)) & x \in \mathcal{H}_1 \cup \mathcal{H}_2 \\ \sqrt{\frac{\tilde{\mu}-\nu}{\tilde{\mu}}} \tilde{u}(x) & x \in S_N \setminus (\mathcal{H}_1 \cup \mathcal{H}_2), \end{cases}$$

where the shift $T(\nu)$ is such that $T(0) = 0$ and u_ν is continuous at the origin. One has $\|u_\nu\|_2^2 = \mu$ for every ν and

$$\begin{aligned} \frac{d^2}{d\nu^2} F_{p,q}(u_\nu) \Big|_{\nu=0} &= -\frac{p-2}{4} \left[\frac{1}{\bar{\mu}^2} \int_{\mathcal{H}_1 \cup \mathcal{H}_2} |\bar{u}|^p dx + \frac{1}{\tilde{\mu}^2} \int_{S_N \setminus (\mathcal{H}_1 \cup \mathcal{H}_2)} |\tilde{u}|^p dx \right] - \frac{q-2}{4\tilde{\mu}^2} |u(0)|^q \\ &\leq -\frac{1}{4\mu^2} \left[(p-2) \int_{S_N} |u|^p dx + (q-2) |u(0)|^q \right] < 0, \end{aligned}$$

so that, choosing ε small enough and setting $v := u_\nu$ for any $\nu \in (-\varepsilon, \varepsilon)$, we conclude. \square

Proof of Proposition 3.3. Note that, by Lemma 3.4–3.5, any ground state $u \in H_\mu^1(S_N)$ of $F_{p,q}$ at mass μ has to be symmetric with respect to the origin and non-increasing on each half-line. Since the unique solution to (13) fulfilling these properties is η_0^ω and given that, by Corollary 2.2, there exists a unique $\omega > 0$ in (13) for which η_0^ω belongs to $H_\mu^1(S_N)$, the proof is complete. \square

Note that, whenever $\mathcal{F}_{p,q}(\mu)$ is attained, Proposition 3.3 entails

$$\mathcal{F}_{p,q}(\mu) = \inf_{v \in H_{\mu,\text{rad}}^1(S_N)} F_{p,q}(v)$$

where $H_{\mu,\text{rad}}^1 = \{v \in H_\mu^1(S_N) : v \text{ is symmetric with respect to the origin of } S_N\}$. We conclude this section with the next proposition concerning radial ground states of $F_{p,q}$, establishing some properties that will be useful in what follows.

Proposition 3.6. *For every $\mu > 0$, the minimization problem*

$$\mathcal{F}_{p,q}^{\text{rad}}(\mu) := \inf_{v \in H_{\mu,\text{rad}}^1(S_N)} F_{p,q}(v) \tag{31}$$

is always attained by the unique stationary state $\eta_0^{\omega(\mu)}$.

Proof. First notice that

$$\mathcal{F}_{p,q}^{\text{rad}}(\mu) < 0 \tag{32}$$

for every $\mu > 0$. Indeed, let us denote by $\varphi \in H^1(\mathbb{R}^+)$ the restriction to \mathbb{R}^+ of the soliton $\phi_{\omega(\frac{2\mu}{N})}$ of mass $\frac{2\mu}{N}$, so that $\|\varphi\|_{L^2(\mathbb{R}^+)}^2 = \frac{\mu}{N}$. Setting $v := \varphi$ on each half-line of S_N , we get $v \in H_{\mu,\text{rad}}^1(S_N)$ and

$$\mathcal{F}_{p,q}^{\text{rad}}(\mu) \leq F_{p,q}(v) < NE(\varphi, \mathbb{R}^+) = -\left(\frac{2}{N}\right)^{2\beta} \theta_p \mu^{2\beta+1} < 0.$$

Moreover, the minimization problem (31) is equivalent to minimize $F_{p,q}$ among all functions $u \in H_{\mu,\text{rad}}^1(S_N)$ non-increasing on each half-line. Indeed, arguing as in the first part of the proof of Lemma 3.3, it is possible to construct $u^* \in H_{\mu,\text{rad}}^1(S_N)$, monotonically decreasing on each half-line and such that $F_{p,q}(u^*) < F_{p,q}(u)$.

Therefore, let $(u_n) \subset H_{\mu,\text{rad}}^1(S_N)$ be a minimizing sequence for (31) and, due to Lemma 3.5, assume without loss of generality that $\|u_n\|_\infty = u_n(0)$ and u_n is non-increasing on each half-line. By Gagliardo–Nirenberg inequalities (25)–(26) it follows that (u_n) is bounded in $H^1(S_N)$, so that $u_n \rightharpoonup u$ in $H^1(S_N)$ and $u_n \rightarrow u$ in $L_{\text{loc}}^\infty(S_N)$, for some $u \in H^1(S_N)$.

Assume by contradiction that $u \equiv 0$ on S_N . Then $u_n \rightarrow 0$ in $L_{\text{loc}}^\infty(S_N)$, that is $u_n \rightarrow 0$ in $L^\infty(S_N)$ since u_n attains its L^∞ norm at the origin. Thus $u_n \rightarrow 0$ strongly in $L^p(S_N)$ and by weak lower semicontinuity

$$\mathcal{F}_{p,q}^{\text{rad}}(\mu) = \lim_n F_{p,q}(u_n) \geq 0$$

which is impossible by (32). Thus $u \not\equiv 0$ on S_N .

Let then $m := \|u\|_2^2$, so that $\|u_n - u\|_2^2 \rightarrow \mu - m$ as $n \rightarrow +\infty$ by weak convergence of u_n to u in $L^2(S_N)$, and assume by contradiction that $0 < m < \mu$.

Since $u_n \in H_{\mu,\text{rad}}^1(S_N)$ is non-increasing on each half-line, then for every n we have that $u_n - u$ is symmetric with respect to the origin and $(u_n - u)(0) \rightarrow 0$ as $n \rightarrow +\infty$ since $u_n \rightarrow u$ in $L_{\text{loc}}^\infty(S_N)$. Hence, one can argue again as in the first part of the proof of Lemma 3.3 to construct $(u_n - u)^* \in H^1(S_N)$ symmetric with respect to the origin, non-increasing on each half-line and such that $\|(u_n - u)^*\|_2 = \|u_n - u\|_2$ and $F_{p,q}((u_n - u)^*) \leq F_{p,q}(u_n - u)$.

Thus, following the steps in the proof of Proposition 3.1, we have

$$\begin{aligned} \mathcal{F}_{p,q}^{\text{rad}}(\mu) &\leq F_{p,q}\left(\sqrt{\frac{\mu}{\|(u_n - u)^*\|_2^2}}(u_n - u)^*\right) \\ &< \frac{\mu}{\|(u_n - u)^*\|_2^2} F_{p,q}((u_n - u)^*) \leq \frac{\mu}{\|u_n - u\|_2^2} F_{p,q}(u_n - u), \end{aligned}$$

that is

$$\liminf_n F_{p,q}(u_n - u) \geq \frac{\mu - m}{\mu} \mathcal{F}_{p,q}^{\text{rad}}(\mu).$$

Similarly, since $u \in H_{m,\text{rad}}^1(S_N)$,

$$F_{p,q}(u) > \frac{m}{\mu} \mathcal{F}_{p,q}^{\text{rad}}(\mu).$$

Relying again on $u'_n \rightharpoonup u'$ in $L^2(S_N)$, $u_n \rightarrow u$ in $L_{\text{loc}}^\infty(S_N)$ and on the Brezis–Lieb lemma [29] as in the proof of Proposition 3.1, we obtain

$$\begin{aligned} \mathcal{F}_{p,q}^{\text{rad}}(\mu) &= \lim_n F_{p,q}(u_n) = \lim_n F_{p,q}(u_n - u) + F_{p,q}(u) \\ &> \frac{\mu - m}{\mu} \mathcal{F}_{p,q}^{\text{rad}}(\mu) + \frac{m}{\mu} \mathcal{F}_{p,q}^{\text{rad}}(\mu) = \mathcal{F}_{p,q}^{\text{rad}}(\mu), \end{aligned}$$

i.e. a contradiction. Henceforth, $u \in H_{\mu,\text{rad}}^1(S_N)$ and by weak lower semicontinuity $F_{p,q}(u) = \mathcal{F}_{p,q}^{\text{rad}}(\mu)$. In particular, u is solution of (13) for some $\omega > 0$, so that it must coincide with the unique solution $\eta_0^{\omega(\mu)}$ to (13) that verifies $\|\eta_0^{\omega(\mu)}\|_2^2 = \mu$. \square

4. PROOF OF THEOREM 1.1–1.2 AND OF PROPOSITION 1.3

This section is devoted to the details of the proof of the main results of the paper. We begin with the following preliminary lemma.

Lemma 4.1. *Let $\bar{\mu} > 0$ be given.*

- (i) *If $q < \frac{p}{2} + 1$ and ground states of $F_{p,q}$ exist at mass $\bar{\mu}$, then, for every $\mu \leq \bar{\mu}$, ground states at mass μ exist too.*
- (ii) *if $q > \frac{p}{2} + 1$ and ground states of $F_{p,q}$ exist at mass $\bar{\mu}$, then, for every $\mu \geq \bar{\mu}$, ground states at mass μ exist too.*

Proof. Recall that, by Proposition 3.3, whenever ground states exist they coincide with the unique stationary state $\eta_0^{\omega(\mu)}$ such that $\|\eta_0^{\omega(\mu)}\|_2^2$ coincides with the prescribed mass, so that in this case $\mathcal{F}_{p,q}(\mu) = F_{p,q}(\eta_0^{\omega(\mu)})$. Therefore, relying also on Proposition 3.1 and Corollary 3.2, we have that ground states of $F_{p,q}$ at mass μ exist if and only if

$$F_{p,q}(\eta_0^{\omega(\mu)}) \leq \mathcal{E}(\mu)$$

which by (11) may be rewritten as

$$\frac{\mathcal{F}_{p,q}^{\text{rad}}(\mu)}{\mu^{2\beta+1}} \leq -\theta_p,$$

where, as defined in (10), $\beta = \frac{p-2}{6-p}$ and $\mathcal{F}_{p,q}^{\text{rad}}(\mu) = F_{p,q}(\eta_0^{\omega(\mu)})$ by Proposition 3.6.

Set $K(\mu) := \frac{\mathcal{F}_{p,q}^{\text{rad}}(\mu)}{\mu^{2\beta+1}}$ for every $\mu > 0$. Since $\mathcal{F}_{p,q}^{\text{rad}}(\mu)$ is a differentiable function of μ by Proposition 3.6 and formula (22), it follows that K' exists for every $\mu > 0$ and it verifies

$$K'(\mu) = \frac{1}{\mu^{2\beta+1}} \left((\mathcal{F}_{p,q}^{\text{rad}})'(\mu) - (2\beta+1) \frac{\mathcal{F}_{p,q}^{\text{rad}}(\mu)}{\mu} \right), \quad (33)$$

where $(\mathcal{F}_{p,q}^{\text{rad}})'$ denotes the derivative of $\mathcal{F}_{p,q}^{\text{rad}}$.

Let us now prove statement (i). Assume $q < \frac{p}{2} + 1$. Letting $\eta_0^{\omega(\mu)}$ be the radial, monotonically decreasing ground state at mass μ , we have

$$\begin{aligned} \mathcal{F}_{p,q}^{\text{rad}}(\mu - \varepsilon) - \mathcal{F}_{p,q}^{\text{rad}}(\mu) &\leq F_{p,q} \left(\sqrt{\frac{\mu - \varepsilon}{\mu}} \eta_0^{\omega(\mu)} \right) - F_{p,q}(\eta_0^{\omega(\mu)}) \\ &= \frac{1}{2} \frac{\varepsilon}{\mu} \left(-\|(\eta_0^{\omega(\mu)})'\|_2^2 + \|\eta_0^{\omega(\mu)}\|_p^p + |\eta_0^{\omega(\mu)}(0)|^q \right) + o(\varepsilon) = \frac{1}{2} \varepsilon \omega(\mu) + o(\varepsilon), \end{aligned}$$

where $\varepsilon > 0$ is sufficiently small. Therefore

$$(\mathcal{F}_{p,q}^{\text{rad}})'(\mu^-) \geq -\frac{1}{2} \omega(\mu) \quad \forall \mu > 0,$$

where $(\mathcal{F}_{p,q}^{\text{rad}})'(\mu^-)$ denotes the left derivative of $\mathcal{F}_{p,q}^{\text{rad}}$ at μ .

Conversely, denoting by $(\mathcal{F}_{p,q}^{\text{rad}})'(\mu^+)$ the right derivative of $\mathcal{F}_{p,q}^{\text{rad}}$ at μ , the same argument leads to

$$(\mathcal{F}_{p,q}^{\text{rad}})'(\mu^+) \leq -\frac{1}{2} \omega(\mu),$$

so that, since $\mathcal{F}_{p,q}^{\text{rad}}$ is differentiable at every $\mu > 0$,

$$(\mathcal{F}_{p,q}^{\text{rad}})'(\mu) = -\frac{1}{2} \omega(\mu).$$

Coupling with (33) and making use of the explicit expression of $F_{p,q}(\eta_0^{\omega(\mu)})$ as in (22) yields

$$K'(\mu) = \frac{p+2-2q}{q(6-p)} \frac{|\eta_0(0)|^q}{\mu^{2\beta+2}} > 0$$

for every $\mu > 0$, and (i) is proved.

The proof of statement (ii) is analogous. \square

Proof of Theorem 1.1. Set

$$\mu_{p,q} := \begin{cases} \sup\{\mu > 0 : \mathcal{F}_{p,q}(\mu) \leq \mathcal{E}(\mu)\} & \text{if } q < \frac{p}{2} + 1 \\ \inf\{\mu > 0 : \mathcal{F}_{p,q}(\mu) \leq \mathcal{E}(\mu)\} & \text{if } q > \frac{p}{2} + 1. \end{cases}$$

By Proposition 3.1, Corollary 3.2 and Lemma 4.1, it follows that if $q < \frac{p}{2} + 1$, then ground states of $F_{p,q}$ at mass μ exist if and only if $\mu \leq \mu_{p,q}$, whereas if $q > \frac{p}{2} + 1$, then ground states $F_{p,q}$ at mass μ exist if and only if $\mu \geq \mu_{p,q}$. Furthermore, Proposition 3.6 ensures that, whenever they exist, ground states at prescribed mass are also unique. Thus, to complete the proof of Theorem 1.1 it is enough to show that

$$0 < \mu_{p,q} < +\infty.$$

The proof is divided in two steps.

Step 1. Existence. Let $u = (u_i)_{i=1}^N \subset H_\mu^1(S_N)$ be given by

$$u_i(x) = Ae^{-Bx} \quad \text{on } \mathcal{H}_i, i = 1, \dots, N.$$

Imposing the boundary condition in (13), we get that

$$NB = A^{q-2}. \quad (34)$$

Furthermore,

$$\mu = \frac{NA^2}{2B} \quad (35)$$

and

$$F_{p,q}(u) = \frac{N}{4}A^2B - \frac{N}{p^2}\frac{A^p}{B} - \frac{1}{q}A^q. \quad (36)$$

Combining (34), (35) and (36), we get

$$F_{p,q}(u) = -\left(\frac{1}{q} - \frac{1}{4}\right) \left(\frac{2}{N^2}\right)^{\frac{q}{4-q}} \mu^{\frac{q}{4-q}} - \frac{N^2}{p^2} \left(\frac{2}{N^2}\right)^{\frac{p-q+2}{4-q}} \mu^{\frac{p-q+2}{4-q}}.$$

Now, if $q < \frac{p}{2} + 1$, then

$$\frac{q}{4-q} < \frac{2q}{6-p} < \frac{p+2}{6-p} = 2\beta + 1$$

so that, recalling (11),

$$\mathcal{F}_{p,q}(\mu) \leq F_{p,q}(u) < \mathcal{E}(\mu) \quad \text{for } \mu \text{ small enough,}$$

which shows that $\mu_{p,q} > 0$.

On the contrary, if $q > \frac{p}{2} + 1$, then

$$\frac{q}{4-q} > \frac{2q}{6-p} > \frac{p+2}{6-p} = 2\beta + 1$$

and consequently

$$\mathcal{F}_{p,q}(\mu) \leq F_{p,q}(u) < \mathcal{E}(\mu) \quad \text{for } \mu \text{ large enough,}$$

i.e. $\mu_{p,q} < +\infty$.

Step 2. Non-existence. As Proposition 3.3 ensures that if a ground state of (1) at mass μ exists, then it coincides with $\eta_0^{\omega(\mu)}$, relations (18) and (19) become respectively

$$\frac{t}{(1-t^2)^{\frac{q-2}{p-2}}} = \frac{\left(\frac{p}{2}\right)^{\frac{q-2}{p-2}} \omega(\mu)^{\frac{2q-2-p}{2(p-2)}}}{N} \quad (37)$$

and

$$\mu = 2N \frac{\left(\frac{p}{2}\right)^{\frac{2}{p-2}} \omega(\mu)^{\frac{6-p}{2(p-2)}}}{p-2} \int_t^1 (1-s^2)^{\frac{4-p}{p-2}} ds. \quad (38)$$

Assume now $q < \frac{p}{2} + 1$. Since

$$0 \leq \int_t^1 (1-s^2)^{\frac{4-p}{p-2}} \leq \int_0^1 (1-s^2)^{\frac{4-p}{p-2}} < +\infty \quad \text{for } 2 < p < 6,$$

by (38) we get $\omega(\mu) \rightarrow \infty$ as $\mu \rightarrow +\infty$, and consequently $t \rightarrow 0^+$ by (37). Hence it follows that

$$\omega(\mu) \sim \mu^{\frac{2(p-2)}{6-p}} \quad \text{as } \mu \rightarrow +\infty.$$

Let then $m := \|\phi_{\omega(\mu)}\|_{L^2(\mathbb{R})}^2$, where $\phi_{\omega(\mu)}$ is the soliton (8) associated to the Lagrange multiplier $\omega(\mu)$. Since $t \rightarrow 0$ as $\mu \rightarrow +\infty$, recalling (9) shows that

$$\lim_{\mu \rightarrow +\infty} \frac{\mu}{m} = \frac{N}{2}$$

and combining with (10) gives

$$\left| \eta_0^{\omega(\mu)}(0) \right|^q \leq \|\phi_{\omega(\mu)}\|_{L^\infty(\mathbb{R})}^q = C_p m^{\frac{2q}{6-p}} \sim \mu^{\frac{2q}{6-p}} \quad \text{as } \mu \rightarrow +\infty. \quad (39)$$

Conversely, since almost every value in the range of $\eta_0^{\omega(\mu)}$ is attained N times on S_N , by [14, Lemma 2.1]

$$E\left(\eta_0^{\omega(\mu)}, S_N\right) \geq -\theta_p \left(\frac{2}{N}\right)^{\frac{2(p-2)}{6-p}} \mu^{\frac{p+2}{6-p}} \quad (40)$$

for every $\mu > 0$.

Combining with (39) and the fact that $q < \frac{p}{2} + 1$ entails

$$F_{p,q}\left(\eta_0^{\omega(\mu)}\right) \geq -\theta_p \left(\frac{2}{N}\right)^{\frac{2(p-2)}{6-p}} \mu^{\frac{p+2}{6-p}} - C \mu^{\frac{2q}{6-p}} \sim -\theta_p \left(\frac{2}{N}\right)^{\frac{2(p-2)}{6-p}} \mu^{\frac{p+2}{6-p}} > \mathcal{E}(\mu)$$

for μ sufficiently large and $N \geq 3$, i.e. $\mu_{p,q} < +\infty$.

If on the contrary $q > \frac{p}{2} + 1$, then a similar argument shows that

$$\omega(\mu) \sim \mu^{\frac{2(p-2)}{6-p}} \quad \text{as } \mu \rightarrow 0^+$$

and

$$\left| \eta_0^{\omega(\mu)}(0) \right|^q \leq C \mu^{\frac{2q}{6-p}}$$

for μ small enough.

Coupling again with (40) then leads to

$$F_{p,q}\left(\eta_0^{\omega(\mu)}\right) > \mathcal{E}(\mu)$$

provided μ is sufficiently small, yielding $\mu_{p,q} > 0$ and we conclude. \square

In the final part of the section we prove the main result in the case $q = \frac{p}{2} + 1$.

Proof of Theorem 1.2. If $q = \frac{p}{2} + 1$, then (37) reduces to

$$t = \frac{\sqrt{p}}{\sqrt{p + 2N^2}} \quad (41)$$

and (38) can be rewritten as

$$\mu = 2N \frac{\left(\frac{p}{2}\right)^{\frac{2}{p-2}} \omega^{\frac{6-p}{2(p-2)}}}{p-2} \mathcal{I}(t),$$

where

$$\mathcal{I}(x) := \int_x^1 (1-s^2)^{\frac{4-p}{p-2}} ds \quad \forall x \in [0, 1].$$

Then

$$\omega(\mu) = \left(\frac{(p-2)\mu}{2N\mathcal{I}(t) \left(\frac{p}{2}\right)^{\frac{2}{p-2}}} \right)^{\frac{2(p-2)}{6-p}}, \quad (42)$$

and plugging (41) and (42) into (22) for $\eta_0^{\omega(\mu)}$ gives

$$F_{p,q}(\eta_0^{\omega(\mu)}) = -\frac{2^{-\frac{p-2}{6-p}}(6-p)}{(p+2)p^{\frac{4}{6-p}}} \left(\frac{p-2}{N\mathcal{I}(t)} \right)^{\frac{2(p-2)}{6-p}} \mu^{\frac{p+2}{6-p}},$$

whereas making use of (10) one can rewrite $\mathcal{E}(\mu)$ as

$$\mathcal{E}(\mu) = -\frac{2^{-\frac{p-2}{6-p}}(6-p)}{(p+2)p^{\frac{4}{6-p}}} \left(\frac{p-2}{\mathcal{I}(0)} \right)^{\frac{2(p-2)}{6-p}} \mu^{\frac{p+2}{6-p}}.$$

By Proposition 3.1, Corollary 3.2 and Proposition 3.6, ground states of (1) at mass μ exist if and only if

$$F_{p,q}(\eta_0^{\omega(\mu)}) \leq \mathcal{E}(\mu)$$

that, thanks to the previous expressions, can be reduced to

$$N \frac{\mathcal{I}\left(\sqrt{\frac{p}{p+2N^2}}\right)}{\mathcal{I}(0)} \leq 2. \quad (43)$$

Note that the function $h : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ defined as

$$h(x) := x \int_{\sqrt{\frac{p}{p+2x^2}}}^1 (1-s^2)^{\frac{4-p}{p-2}} ds \quad (44)$$

is non-decreasing on \mathbb{R}^+ . Indeed, differentiating (44), we have

$$h'(x) := \int_{\sqrt{\frac{p}{p+2x^2}}}^1 (1-s^2)^{\frac{4-p}{p-2}} ds + \frac{2^{-\frac{2}{p-2}} \sqrt{p} x^{\frac{4}{p-2}}}{(p+2x^2)^{\frac{p+2}{2(p-2)}}} > 0 \quad \forall x > 0.$$

Therefore, observing that the left-hand side of (43) is strictly less than 2 for $N = 2$, it is non-decreasing in N and diverges to $+\infty$ as $N \rightarrow +\infty$, since \mathcal{I} is continuous and bounded and $\frac{p}{p+2N^2} \rightarrow 0$, then there exists $N_p \geq 2$ such that existence of ground states is guaranteed for $N \leq N_p$, while non-existence holds for $N > N_p$. \square

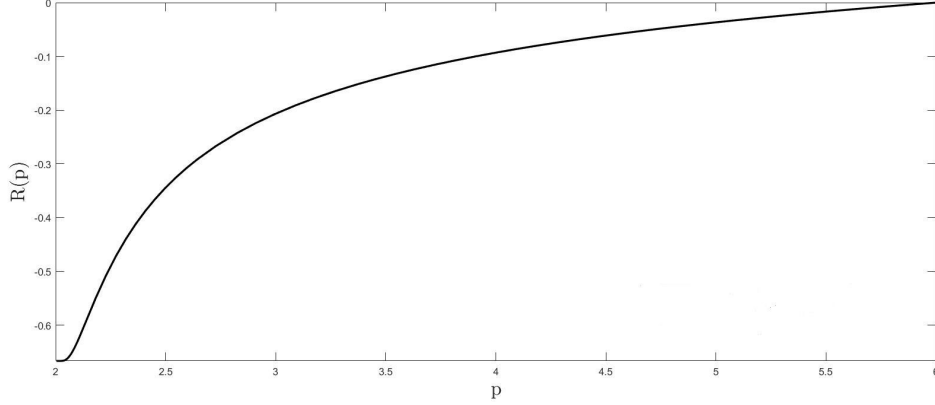


FIGURE 2. The graph of $R(p) := \frac{\mathcal{I}(\sqrt{\frac{p}{p+18}})}{\mathcal{I}(0)}$ as a function of $p \in (2, 6)$. The validity of (43) at $N = 3$ is equivalent to $R(p) \leq 0$. These numerical simulations suggest that condition (43) always holds at $N = 3$, for every $p \in (2, 6)$.

As already pointed out in the Introduction and displayed clearly in the previous proof, without further assumption on $p \in (2, 6)$, at the moment we can prove only that N_p in Theorem 1.2 satisfies $N_p \geq 2$ for every $p \in (2, 6)$. However, numerical simulations (see Figure 2) strongly suggest that actually $N_p \geq 3$ for every $p \in (2, 6)$. To conclude, we thus provide the proof of Proposition 1.3.

Proof of Proposition 1.3. We split the proof in two parts.

Statement (i). Fix $N > 0$. We need to show that $N_p \geq N$ as soon as p is sufficiently close to 2. To this end, we verify that, given N , condition (43) is always satisfied when p approaches 2.

Note that if $p \in (2, 4)$, then $(1 - s^2)^{\frac{4-p}{p-2}}$ is a decreasing function of $s \in [0, 1]$. Hence,

$$\mathcal{I}\left(\sqrt{\frac{p}{p+2N^2}}\right) = \int_{\sqrt{\frac{p}{p+2N^2}}}^1 (1 - s^2)^{\frac{4-p}{p-2}} ds \leq \left(\frac{2N^2}{p+2N^2}\right)^{\frac{4-p}{p-2}} \left(1 - \sqrt{\frac{p}{p+2N^2}}\right).$$

Conversely,

$$\mathcal{I}(0) = \int_0^1 (1 - s^2)^{\frac{4-p}{p-2}} ds = \int_0^1 (1 + s)^{\frac{4-p}{p-2}} (1 - s)^{\frac{4-p}{p-2}} ds \geq \int_0^1 (1 - s)^{\frac{4-p}{p-2}} ds = \frac{p-2}{2}.$$

Therefore we get

$$N \frac{\mathcal{I}\left(\sqrt{\frac{p}{p+2N^2}}\right)}{\mathcal{I}(0)} \leq \frac{2N}{p-2} \left(\frac{2N^2}{p+2N^2}\right)^{\frac{4-p}{p-2}} \left(1 - \sqrt{\frac{p}{p+2N^2}}\right) \rightarrow 0 \quad \text{as } p \rightarrow 2^+,$$

so that (43) holds true as soon as p is close enough to 2. This shows that there exists $\delta = \delta(N)$ so that for every $p \in (2, 2 + \delta)$ we have $N_p \geq N$, i.e. Proposition 1.3(i) is proved.

Statement (ii). Let $p \in [4, 6)$. We begin by showing that $N_p < 5$, proving that (43) fails whenever $p \in [4, 6)$ and $N \geq 5$.

Note that assuming $p \geq 4$ implies that $(1 - s^2)^{\frac{4-p}{p-2}}$ is increasing as a function of s and that $(1 - s^2)^{\frac{4-p}{p-2}} \geq 1$ on $[0, 1]$. Henceforth, we get

$$\mathcal{I}\left(\sqrt{\frac{p}{p+2N^2}}\right) \geq 1 - \sqrt{\frac{p}{p+2N^2}}.$$

Moreover, fixing $s \in [0, 1)$ and regarding $(1 - s^2)^{\frac{4-p}{p-2}}$ as a function of $p \in [4, 6)$, we have

$$\frac{d}{dp}(1 - s^2)^{\frac{4-p}{p-2}} = -(1 - s^2)^{\frac{4-p}{p-2}} \ln(1 - s^2) \frac{2}{(p-2)^2} > 0,$$

i.e. for every given $s \in (0, 1)$, $(1 - s^2)^{\frac{4-p}{p-2}}$ is an increasing function of p . Therefore, for every $s \in [0, 1)$ and $p \in [4, 6)$

$$(1 - s^2)^{\frac{4-p}{p-2}} \leq (1 - s^2)^{-\frac{1}{2}},$$

so that integrating over $[0, 1]$ gives

$$\mathcal{I}(0) \leq \int_0^1 (1 - s^2)^{-\frac{1}{2}} ds = \frac{\pi}{2}.$$

We thus obtain

$$N \frac{\mathcal{I}\left(\sqrt{\frac{p}{p+2N^2}}\right)}{\mathcal{I}(0)} \geq \frac{2N}{\pi} \left(1 - \sqrt{\frac{p}{p+2N^2}}\right) \geq \frac{2N}{\pi} \left(1 - \sqrt{\frac{3}{3+N^2}}\right),$$

where the last inequality follows from the fact that, for every given N , $\sqrt{\frac{p}{p+2N^2}}$ is an increasing function of p .

In view of (43), to prove that $N_p \in \{2, 3, 4\}$ for every $p \in [4, 6)$, it is then enough to show that, for every $N \geq 5$

$$\frac{2N}{\pi} \left(1 - \sqrt{\frac{3}{3+N^2}}\right) > 2.$$

that can be equivalently rewritten as

$$G(N) := N^4 - 2\pi N^3 + \pi^2 N^2 - 6\pi N + 3\pi^2 > 0. \quad (45)$$

To prove that (45) holds for every $N \geq 5$, we will show that it is true when $N = 5$ and that $G(N)$ is increasing function of N on $[5, +\infty)$.

On the one hand, direct calculations immediately show that $G(5) > 0$. On the other hand, differentiating (45) with respect to N gives

$$G'(N) = 4N^3 - 6\pi N^2 + 2\pi^2 N - 6\pi,$$

and again we directly see that $G'(5) > 0$. A further differentiation leads to

$$G''(N) = 12N^2 - 12\pi N + 2\pi^2,$$

yielding

$$G''(N) > 0 \quad \text{for every } N \geq \frac{3 + \sqrt{3}}{6}\pi.$$

Since $5 > \frac{3 + \sqrt{3}}{6}\pi$ and $G'(5) > 0$, this implies that $G'(N) \geq 0$ for every $N \geq 5$, and coupling with $G(5) > 0$ this ensures that $G(N) \geq 0$ for every $N \geq 5$. This concludes the proof of the first part of Proposition 1.3(ii).

We are then left to show that there exists $\delta > 0$ such that, if $p \in (4 - \delta, 4 + \delta) \cup (6 - \delta, 6)$, then $N_p = 3$.

We start by proving the results for p in a suitable neighbourhood of 4. Note that, for every given N , we have

$$N \frac{\mathcal{I}\left(\sqrt{\frac{4}{4+2N^2}}\right)}{\mathcal{I}(0)} = N \frac{\int_0^1 \sqrt{\frac{4}{4+2N^2}} ds}{\int_0^1 ds} = N \left(1 - \sqrt{\frac{2}{2+N^2}}\right).$$

Hence, evaluating the previous expression at $N = 3$ and $N = 4$ respectively, we obtain

$$3 \left(1 - \sqrt{\frac{2}{11}}\right) < 2, \quad 4 \left(1 - \sqrt{\frac{2}{18}}\right) = \frac{8}{3} > 2,$$

so that condition (43) is satisfied at $p = 4$, $N = 3$, whereas it fails at $p = 4$, $N = 4$, in turn implying $N_4 = 3$. Since both previous inequalities are strict, by continuity with respect to p , we conclude that $N_p = 3$ for every $p \in (4 - \delta, 4 + \delta)$, for some $\delta > 0$.

Let us now concentrate on the case $p \in (6 - \delta, 6)$. When $p = 6$, evaluating the left hand side of (43) at $N = 3$ and $N = 4$ gives

$$\begin{aligned} 3 \frac{\int_{\frac{1}{2}}^1 (1-s^2)^{-\frac{1}{2}} ds}{\int_0^1 (1-s^2)^{-\frac{1}{2}} ds} &= 3 \frac{\frac{\pi}{2} - \frac{\pi}{6}}{\frac{\pi}{2}} = 2, \\ 4 \frac{\int_{\sqrt{\frac{3}{19}}}^1 (1-s^2)^{-\frac{1}{2}} ds}{\int_0^1 (1-s^2)^{-\frac{1}{2}} ds} &= 4 \frac{\frac{\pi}{2} - \arcsin\left(\sqrt{\frac{3}{19}}\right)}{\frac{\pi}{2}} > 2. \end{aligned}$$

On the one hand, the second inequality being strict shows that (43) is violated when $N = 4$ and $p \in (6 - \delta, 6)$, for suitable $\delta > 0$, so that for all these exponents it must be $N_p \leq 3$.

On the other hand, the first line of the previous equation shows that (43) becomes an equality when $p = 6$ and $N = 3$. Therefore, to show that $N_p = 3$, we need to further analyse the behaviour of the left hand side of (43) when $N = 3$ and p approaches 6. To do this, fix $N = 3$ and set

$$R(p) := \frac{\mathcal{I}\left(\sqrt{\frac{p}{p+18}}\right)}{\mathcal{I}(0)}. \quad (46)$$

We will conclude the proof by showing that $R'(6) > 0$. By continuity, this eventually guarantees the existence of $\delta > 0$ so that $N_p = 3$ for every $p \in (6 - \delta, 6)$.

Differentiating (46) with respect to p we obtain

$$\begin{aligned} R'(p) &= \frac{\left(-\frac{9\sqrt{p+18}}{\sqrt{p}(p+18)^2} \left(\frac{18}{p+18}\right)^{\frac{4-p}{p-2}} - \frac{2}{(p-2)^2} \int_0^1 \sqrt{\frac{p}{p+18}} (1-s^2)^{\frac{4-p}{p-2}} \ln(1-s^2) ds\right) \int_0^1 (1-s^2)^{\frac{4-p}{p-2}}}{\left(\int_0^1 (1-s^2)^{\frac{4-p}{p-2}}\right)^2} \\ &\quad + \frac{\frac{2}{(p-2)^2} \int_0^1 \sqrt{\frac{p}{p+18}} (1-s^2)^{\frac{4-p}{p-2}} ds \int_0^1 (1-s^2)^{\frac{4-p}{p-2}} \ln(1-s^2) ds}{\left(\int_0^1 (1-s^2)^{\frac{4-p}{p-2}}\right)^2}. \end{aligned} \quad (47)$$

According to the previous expression, the numerator of $R'(6)$ reads

$$\left(-\frac{1}{16\sqrt{3}} - \frac{1}{8} \int_{\frac{1}{2}}^1 (1-s^2)^{-\frac{1}{2}} \ln(1-s^2) ds \right) \frac{\pi}{2} + \frac{1}{8} \left(\frac{\pi}{2} - \frac{\pi}{6} \right) \int_0^1 (1-s^2)^{-\frac{1}{2}} \ln(1-s^2) ds,$$

that can be rewritten as

$$\frac{\pi}{16} \left(-\frac{1}{3} \int_0^1 (1-s^2)^{-\frac{1}{2}} \ln(1-s^2) ds - \frac{1}{2\sqrt{3}} + \int_0^{\frac{1}{2}} (1-s^2)^{-\frac{1}{2}} \ln(1-s^2) ds \right). \quad (48)$$

The first integral in the above bracket can be computed explicitly making use of polylogarithmic functions

$$\int_0^1 (1-s^2)^{-\frac{1}{2}} \ln(1-s^2) ds = -\pi \ln 2. \quad (49)$$

Furthermore, since $(1-s^2)^{-\frac{1}{2}} \ln(1-s^2)$ is a decreasing function of s on $[0, \frac{1}{2}]$, we have

$$\int_0^{\frac{1}{2}} (1-s^2)^{-\frac{1}{2}} \ln(1-s^2) ds \geq \int_0^{\frac{1}{2}} \frac{2}{\sqrt{3}} \ln\left(\frac{3}{4}\right) ds = \frac{1}{\sqrt{3}} \ln\left(\frac{3}{4}\right).$$

As it holds

$$\frac{\pi \ln 2}{3} > \frac{1}{2\sqrt{3}} - \frac{1}{\sqrt{3}} \ln\left(\frac{3}{4}\right),$$

combining with (47), (48) and (49) we have $R'(6) > 0$, and the proof is complete. \square

5. PROOF OF PROPOSITION 1.4

In this section we prove the orbital stability of the radial stationary state $\eta_0^{\omega(\mu)}$, even for values of the mass for which there is no ground state. The proof of Proposition 1.4 strongly relies on a method introduced in [5], which is essentially based on the reduction of an infinite-dimensional problem to a finite-dimensional one. The argument of [5] extends almost straightforwardly to our setting. The idea of the method is the following. We map continuously every function in the mass constrained space into another function made of pieces of solitons, whose energy is lower than the one of the original function. Then, we prove that the radial stationary states are local minimum points for the energy among functions with the same mass and made of pieces of solitons. Thus, a fortiori, they are local minima in the whole mass constrained energy space, and therefore, due to the general stability theory, they are orbitally stable. Here we limit ourselves to sketch the main steps of the proof, explicitly pointing out the minor differences with respect to [5] whenever occurring.

Proof of Proposition 1.4. Owing to Theorem 3 in [45], the orbital stability of $\eta_0^{\omega(\mu)}$ is equivalent to its local minimality for $F_{p,q}$ in $H_\mu^1(S_N)$, hence we prove that $\eta_0^{\omega(\mu)}$ is a local minimum for the energy in $H_\mu^1(S_N)$.

Following [5, Definition 2.1], we fix $\mu > 0$ and define the multi-soliton manifold \mathcal{M} as the subspace of $H_\mu^1(S_N)$ made of all the functions whose restriction to each half-line \mathcal{H}_j of S_N gives a piece of soliton, i.e.

$$\mathcal{M} := \{u \in H_\mu^1(S_N) : u_j = \phi_{\omega_j}(\cdot + a_j), \text{ for some } \omega_j, a_j \in \mathbb{R}, j = 1, \dots, N\}.$$

Given a function $\eta \in H_\mu^1(S_N)$ such that $\eta(0) \neq 0$, we define the multi-soliton transformation of η as the unique function $\Sigma\eta \in \mathcal{M}$ so that the restriction $(\Sigma\eta)_j$ of $\Sigma\eta$ to the half-line \mathcal{H}_j satisfies

$$(\Sigma\eta)_j := \phi_{\omega(m_j, h)}(\cdot + a(m_j, h)),$$

where

$$m_j := \int_{\mathcal{H}_j} \eta_j^2 dx_j, \quad j = 1, \dots, N, \quad \text{and} \quad h = |\eta(0)|,$$

and $\phi_{\omega(m_j, h)}(\cdot + a(m_j, h))$ is the unique piece of soliton with mass m_j and $\phi_{\omega(m_j, h)}(a(m_j, h)) = h$. For every given values of m_j, h , the uniqueness of such $\phi_{\omega(m_j, h)}(\cdot + a(m_j, h))$ has been proved in [11, Theorem 4.1]. Furthermore, by [5, Remark 3.4] we have for every η so that $\eta(0) \neq 0$

$$F_{p,q}(\Sigma\eta) \leq F_{p,q}(\eta),$$

and equality holds if and only if $\eta \in \mathcal{M}$, that is $\Sigma\eta \equiv \eta$. In light of this and of the continuity of Σ [5, Proposition 3.2], to show that $\eta_0^{\omega(\mu)}$ is a local minimizer of $F_{p,q}$ in $H_\mu^1(S_N)$, it is enough to prove that it locally minimizes the energy in \mathcal{M} . Note that $\eta_0^{\omega(\mu)} \in \mathcal{M}$ for every $\mu > 0$.

We observe that any given function $\varphi \in \mathcal{M}$ corresponds to a point $P = (m_1, \dots, m_{N-1}, h) \in (0, +\infty)^N$, where m_j is the mass of the restriction φ_j of φ to \mathcal{H}_j , $j = 1, \dots, N-1$, and $h = |\varphi(0)|$. Therefore, it is natural to define the reduced energy function $r : (0, +\infty)^N \rightarrow \mathbb{R}$ as

$$r(P) := F_{p,q}(\varphi),$$

which can be conveniently decomposed as follows

$$r(P) = \sum_{i=1}^{N-1} e(m_i, h) + e\left(\mu - \sum_{i=1}^{N-1} m_i, h\right),$$

where $e : (0, +\infty) \times \mathbb{R}^+ \rightarrow \mathbb{R}$ is given by

$$e(m, h) := \frac{1}{2} \|\phi'_{\omega(m, h)}(\cdot + a(m, h))\|_{L^2(\mathbb{R}^+)}^2 - \frac{1}{p} \|\phi_{\omega(m, h)}(\cdot + a(m, h))\|_{L^p(\mathbb{R}^+)}^p - \frac{1}{qN} h^q.$$

Thus, setting $\bar{h} = |\eta_0^{\omega(\mu)}(0)|$, the local minimality of $\eta_0^{\omega(\mu)}$ in \mathcal{M} is equivalent to the local minimality for r of the point $\bar{P} = (\frac{\mu}{N}, \dots, \frac{\mu}{N}, \bar{h})$.

Since \bar{P} is an internal point of $(0, +\infty)^N$ and it is a stationary point for r as $\eta_0^{\omega(\mu)}$ is a critical point for $F_{p,q}$, to conclude it is then sufficient to prove that the Hessian matrix of r evaluated at \bar{P} is positive definite. By straightforward computations,

$$\begin{aligned} \frac{\partial^2 r}{\partial m_i \partial m_j}(\bar{P}) &= (1 + \delta_{ij}) \frac{\partial^2 e}{\partial m^2} \left(\frac{\mu}{N}, \bar{h} \right), \\ \frac{\partial^2 r}{\partial h^2}(\bar{P}) &= N \frac{\partial^2 e}{\partial h^2} \left(\frac{\mu}{N}, \bar{h} \right), \\ \frac{\partial^2 r}{\partial m_i \partial h}(\bar{P}) &= 0, \end{aligned}$$

where δ_{ij} denotes as usual the Kronecker's symbol of i, j . By elementary linear algebra, one easily sees that the Hessian matrix has three eigenvalues: $N \frac{\partial^2 e}{\partial m^2} \left(\frac{\mu}{N}, \bar{h} \right)$ with multiplicity

1, $\frac{\partial^2 e}{\partial m^2} \left(\frac{\mu}{N}, \bar{h} \right)$ with multiplicity $N - 2$ and $\frac{\partial^2 e}{\partial h^2} \left(\frac{\mu}{N}, \bar{h} \right)$ with multiplicity 1. Therefore, to show that the Hessian matrix is positive definite, we need to prove that

$$\frac{\partial^2 e}{\partial m^2} \left(\frac{\mu}{N}, \bar{h} \right) > 0 \quad (50)$$

and

$$\frac{\partial^2 e}{\partial h^2} \left(\frac{\mu}{N}, \bar{h} \right) > 0. \quad (51)$$

The proof of (50)–(51) is analogous to the one of inequalities (4.2)–(4.3) of [5]. The main idea is to consider the variations

$$f_1(t) := r(\underline{m}(t), \bar{h}) \quad \text{and} \quad f_2(t) := r \left(\frac{\mu}{N}, \dots, \frac{\mu}{N}, \bar{h} + t \right), \quad t \in (-\varepsilon, \varepsilon)$$

where $\underline{m}(t) = \left(\frac{\mu}{N} + t, \frac{\mu}{N}, \dots, \frac{\mu}{N} \right)$.

It is plainly seen that

$$f_1''(0) = 2 \frac{\partial^2 e}{\partial m^2} \left(\frac{\mu}{N}, \bar{h} \right) \quad \text{and} \quad f_2''(0) = N \frac{\partial^2 e}{\partial h^2} \left(\frac{\mu}{N}, \bar{h} \right).$$

Note that $f_1(t)$ corresponds to an exchange of mass t between the first and the N -th half-lines, without involving the remaining $N - 2$ ones. Hence,

$$f_1(t) - f_1(0) = f_1(t) - r(\bar{P}) = \tilde{F}_{p,q}(\varphi_t, \mathbb{R}) - \tilde{F}_{p,q}(\varphi_0, \mathbb{R})$$

where $\varphi_t, \varphi_0 \in H_{\frac{2\mu}{N}}^1(\mathbb{R})$ denote respectively the restriction to the line $\mathcal{H}_1 \cup \mathcal{H}_N$ of the function $\eta_t \in \mathcal{M}$ corresponding to the point $(\underline{m}(t), \bar{h})$ and of the stationary state $\eta_0^{\omega(\mu)}$, and $\tilde{F}_{p,q} : H_{\frac{2\mu}{N}}^1(\mathbb{R}) \rightarrow \mathbb{R}$ is given by

$$\tilde{F}(u, \mathbb{R}) = \frac{1}{2} \|u'\|_{L^2(\mathbb{R})}^2 - \frac{1}{p} \|u\|_{L^p(\mathbb{R})}^p - \frac{2}{qN} |u(0)|^q.$$

Here is the point where we need to argue slightly differently with respect to [5]. Indeed, we now rely on the results of [28], which guarantees that φ_0 as above is a global minimizer of $\tilde{F}_{p,q}$ in $H_{\frac{2\mu}{N}}^1(\mathbb{R})$. Coupling with the stability result in [45, Theorem 3.4], this immediately yields

$$f_1(t) - f_1(0) \geq C \|\varphi_t - \varphi_0\|_{L^2(\mathbb{R})}^2$$

for some constant $C > 0$ and t small enough. Moving from the previous inequality, and repeating the same calculations as in [5, pp. 7411], we eventually obtain

$$f_1(t) - f_1(0) \geq Ct^2,$$

that coupled with $f_1'(0) = 0$ ensures $f_1''(0) \geq C > 0$. This proves (50). Since the same argument developed for f_2 leads to (51), we conclude. \square

ACKNOWLEDGEMENTS

The first and the second authors acknowledge that the present research has been partially supported by MIUR grant Dipartimenti di Eccellenza 2018-2022 (E11G18000350001). All the authors wish to thank Enrico Serra and Paolo Tilli for fruitful discussions and suggestions.

REFERENCES

- [1] Adami R., Boni F., Ruighi A., *Non-Kirchhoff vertices and nonlinear Schrödinger ground states on graphs*, Mathematics **8**(4) (2020), 617.
- [2] Adami R., Cacciapuoti C., Finco D., Noja D., *Stationary states of NLS on star graphs*, EPL (Europhysics Letters) **100**(1) (2012), 10003.
- [3] Adami R., Cacciapuoti C., Finco D., Noja D., *Variational properties and orbital stability of standing waves for NLS equation on a star graph*, J. Differential Equations **257**(10) (2014), 3738–3777.
- [4] Adami R., Cacciapuoti C., Finco D., Noja D., *Constrained energy minimization and orbital stability for the NLS equation on a star graph*, Ann. Inst. Poincaré (C) An. Non. Lin. **31**(6) (2014), 1289–1310.
- [5] Adami R., Cacciapuoti C., Finco D., Noja D., *Stable standing waves for a NLS on star graphs as local minimizers of the constrained energy*, J. Differential Equations **260**(10) (2016), 7397–7415.
- [6] Adami R., Carlone R., Correggi M., Tentarelli L., *Blow up for the pointwise NLS in dimension two: absence of critical power*, J. Differential Equations **269**(1) (2020), 1–37.
- [7] Adami R., Carlone R., Correggi M., Tentarelli L., *Stability of the standing waves of the concentrated NLSE in dimension two*, Math. Eng. **3**(2) (2021), 1–15.
- [8] Adami R., Dell’Antonio G., Figari R., Teta A., *The Cauchy problem for the Schrödinger equation in dimension three with concentrated nonlinearity*, Ann. Inst. Poincaré (C) An. Non. Lin. **20**(3) (2003), 477–500.
- [9] Adami R., Dell’Antonio G., Figari R., Teta A., *Blow-up solutions for the Schrödinger equation in dimension three with a concentrated nonlinearity*, Ann. Inst. Poincaré (C) An. Non. Lin. **21**(1) (2004), 121–137.
- [10] Adami R., Dovetta S., Serra E., Tilli P., *Dimensional crossover with a continuum of critical exponents for NLS on doubly periodic metric graphs*, Anal. PDE **12**(6) (2019), 1597–1612.
- [11] Adami R., Serra E., Tilli P., *NLS ground states on graphs*, Calc. Var. PDE **54** (2015), 743–761.
- [12] Adami R., Serra E., Tilli P., *Threshold phenomena and existence results for NLS ground states on graphs*, J. Funct. An. **271**(1) (2016), 201–223.
- [13] Adami R., Serra E., Tilli P., *Negative energy ground states for the L^2 -critical NLSE on metric graphs*, Comm. Math. Phys. **352**(1) (2017), 387–406.
- [14] Adami R., Serra E., Tilli P., *Multiple positive bound states for the subcritical NLS equation on metric graphs*, Calc. Var. PDE **58**(5) (2019), 16pp.
- [15] Adami R., Teta A., *A simple model of concentrated nonlinearity*, Oper. Theory: Adv. Appl. **108** (1999), 183–189.
- [16] Adami R., Teta A., *A Class of nonlinear Schrödinger equations with concentrated nonlinearity*, J. Funct. An. **180**(1) (2001), 148–175.
- [17] Angulo J., Goloshchapova N., *Extension theory approach in the stability of the standing waves for the NLS equation with point interactions on a star graph*, Adv. Diff. Equations **23** (2018), 793–846.
- [18] Angulo J., Goloshchapova N., *On the orbital instability of excited states for the NLS equation with the δ -interaction on a star graph*, Disc. Cont. Dyn. Systems **38**(10) (2018), 5039–5066.
- [19] Band R., Lévy G., *Quantum Graphs which Optimize the Spectral Gap*, Ann. Henri Poincaré **18** (2017), 3269–3323.
- [20] Bartsch T., Jeanjean L., *Normalized solutions for nonlinear Schrödinger systems*, Proc. Roy. Soc. Edinburgh Sect. A: Math. **148**(2) (2018), 225–242.
- [21] Bartsch T., Jeanjean L., Soave N., *Normalized solutions for a system of coupled cubic Schrödinger equations on \mathbb{R}^3* , J. Math. Pures Appl. **106**(4) (2016), 583–614.
- [22] Bartsch T., Soave N., *A natural constraint approach to normalized solutions of nonlinear Schrödinger equations and systems*, J. Funct. An. **272**(12) (2017), 4998–5037.
- [23] Bartsch T., Soave N., *Multiple normalized solutions for a competing system of Schrödinger equations*, Calc.Var. PDE **58**(1) (2019), art. n. 22.
- [24] Berkolaiko G., Kennedy J.B., Kurasov P., Mugnolo D., *Surgery principles for the spectral analysis of quantum graphs* Trans. Amer. Math. Soc. **372** (2019), 5153–5197.
- [25] Berkolaiko G., Kuchment P., *Introduction to quantum graphs*, Mathematical Surveys and Monographs 186, American Mathematical Society, Providence, RI, 2013.
- [26] Berkolaiko G., Marzuola J.L., Pelinovsky D.E., *Edge-localized states on quantum graphs in the limit of large mass*, Ann. Inst. Poincaré (C) An. Non. Lin. **38**(5) (2021), 1295–1335.

- [27] Borrelli W., Carlone R., Tentarelli L., *Nonlinear Dirac equation on graphs with localized nonlinearities: bound states and nonrelativistic limit*, SIAM J. Math. An. **51**(2) (2019), 1046–1081.
- [28] Boni F., Dovetta S., *Ground states for a doubly nonlinear Schrödinger equation in dimension one*, J. Math. Anal. Appl. **496**(1) (2021), 124797.
- [29] Brezis H., Lieb E., *A relation between pointwise convergence of functions and convergence of functionals*, Proc. Amer. Math. Soc. **88**(3) (1983), 486–490.
- [30] Cacciapuoti C., Dovetta S., Serra E., *Variational and stability properties of constant solutions to the NLS equation on compact metric graphs*, Milan J. Math. **86**(2) (2018), 305–327.
- [31] Cacciapuoti C., Finco, D., *Graph-like models for thin waveguides with Robin boundary conditions*, Asymptot. Anal. **70** (2010), 199–230.
- [32] Cacciapuoti C., Finco D., Noja D., *Ground state and orbital stability for the NLS equation on a general starlike graph with potentials*, Nonlinearity **30** (2017), 3271–3303.
- [33] Carlone R., Correggi M., Tentarelli L., *Well-posedness of the two-dimensional nonlinear Schrödinger equation with concentrated nonlinearity*, Ann. Inst. Poincaré (C) An. Non. Lin. **36**(1) (2019), 257–294.
- [34] Cazenave T., *Semilinear Schrödinger Equations*, Courant Lecture Notes 10. American Mathematical Society, Providence, RI, 2003.
- [35] Cheng X., Miao C., Zhao L., *Global well-posedness and scattering for nonlinear Schrödinger equations with combined nonlinearities in the radial case*, J. Differential Equations **261**(6) (2016), 2881–2934.
- [36] Dovetta S., *Existence of infinitely many stationary solutions of the L^2 -subcritical and critical NLSE on compact metric graphs*, J. Differential Equations **264**(7) (2018), 4806–4821.
- [37] Dovetta S., *Mass-constrained ground states of the stationary NLSE on periodic metric graphs*, Nonlinear Diff. Eq. Appl. (2019) 26:30.
- [38] Dovetta S., Ghimenti M., Micheletti A.M., Pistoia A., *Peaked and low action solutions of NLS equations on graphs with terminal edges*, SIAM J. Math. Anal. **52**(3) (2020), 2874–2894.
- [39] Dovetta S., Serra E., Tilli P., *Uniqueness and non-uniqueness of prescribed mass NLS ground states on metric graphs*, Adv. Math. **374** (2020), 107352.
- [40] Dovetta S., Serra E., Tilli P., *NLS ground states on metric trees: existence results and open questions*, J. London Math. Soc. **102**(3) (2020), 1223–1240.
- [41] Dovetta S., Tentarelli L., *L^2 -critical NLS on noncompact metric graphs with localized nonlinearity: topological and metric features*, Calc. Var. PDE **58**(3) (2019), 58:108.
- [42] Ekholm T., Frank R.L., Kovarik H., *Eigenvalue estimates for Schrödinger operators on metric trees*, Adv. Math. **226**(6) (2011), 5165–5197.
- [43] Goloshchapova N., Ohta M., *Blow-up and strong instability of standing waves for the NLS- δ equation on a star graph*, Nonlinear Anal. **196** (2020), 111753.
- [44] Goodman R.H., Kairzhan A., Pelinovsky D.E., *Drift of spectrally stable shifted states on star graphs*, SIAM J. Appl. Dyn. Syst. **18**(4) (2019), 1723–1755.
- [45] Grillakis M., Shatah J., Strauss W., *Stability theory of solitary waves in the presence of symmetry, I*, J. Funct. An. **74** (1987), 160–197.
- [46] Jeanjean L., Jendrej J., Le T.T., Visciglia N., *Orbital stability of ground states for a Sobolev critical Schrödinger equation*, arXiv:2008.12084 [math.AP] (2020).
- [47] Jona-Lasinio G., Presilla C., Sjöstrand J., *On Schrödinger equations with concentrated nonlinearities*, Ann. Phys. **240** (1995), 1–21.
- [48] Kairzhan A., Pelinovsky D.E., *Nonlinear instability of half-solitons on star graphs*, J. Differential Equations **264**(12) (2018), 7357–7383.
- [49] Kairzhan A., Pelinovsky D.E., *Spectral stability of shifted states on star graphs*, J. Phys. A: Math. Theor. **51** (2018) 095203.
- [50] Kennedy J.B., *A family of diameter-based eigenvalue bounds for quantum graphs*, in: Atay F., Kurasov P., Mugnolo D. (eds) Discrete and Continuous Models in the Theory of Networks, Operator Theory: Advances and Applications **281** (2020), Birkhäuser, Cham.
- [51] Kennedy J.B., Kurasov P., Malenová G., Mugnolo D., *On the spectral gap of a quantum graph*, Ann. Henri Poincaré **17** (2016), 2439–2473.
- [52] Killip R., Oh T., Pocovnicu O., Visan M., *Solitons and scattering for the cubic-quintic nonlinear Schrödinger equation on \mathbb{R}^3* , Arch. Ration. Mech. Anal. **225**(1) (2017), 469–548.
- [53] Le Coz S., Martel Y., Raphaël P., *Minimal mass blow up solutions for a double power nonlinear Schrödinger equation*, Rev. Mat. Iberoam. **32**(3) (2016), 795–833.

- [54] Malomed B., Azbel B., *Modulational instability of a wave scattered by a nonlinear center*, Phys. Rev. B **47** (1993), 16.
- [55] Miao C., Xu G., Zhao L., *The dynamics of the 3D radial NLS with the combined terms*, Comm. Math. Phys. **318**(3) (2013), 767–808.
- [56] Miao C., Zhao T., Zheng J., *On the 4D nonlinear Schrödinger equation with combined terms under the energy threshold*, Calc. Var. PDE **56**(6) (2017), 179:39.
- [57] Mugnolo D., Noja D., Seifert C., *Airy-type evolution equations on star graphs*, Anal. PDE **11**(7) (2018), 1625–1652.
- [58] Nier F., *The Dynamics of some Quantum Open System with Short-Range Nonlinearities*, Nonlinearity **11** (1998), 1127.
- [59] Noja D., *Nonlinear Schrödinger equation on graphs: recent results and open problems*, Phil. Trans. R. Soc. A **372** (2014), 20130002 (20 pages).
- [60] Noja D., Pelinovsky D., *Standing waves of the quintic NLS equation on the tadpole graph*, Calc. Var. PDE **59**(5) (2020), art.n. 173.
- [61] Noja D., Pelinovsky D.E., Shaikhova G., *Bifurcations and stability of standing waves in the nonlinear Schrödinger equation on the tadpole graph*, Nonlinearity **28** (2015), 2343–2378.
- [62] Noris B., Tavares H., Verzini G., *Normalized solutions for nonlinear Schrödinger systems on bounded domains*, Nonlinearity **32**(3) (2019), 1044–1072.
- [63] Pankov A., *Nonlinear Schrödinger equations on periodic metric graphs*, Discrete Contin. Dyn. Syst. **38**(2) (2018), 697–714.
- [64] Pelinovsky D.E., Schneider G., *Bifurcations of Standing Localized Waves on Periodic Graphs*, Ann. Henri Poincaré **18**(4) (2017), 1185–1211.
- [65] Pierotti D., Soave N., Verzini G., *Local minimizers in absence of ground states for the critical NLS energy on metric graphs*, Proc. Royal Soc. Edinb. Section A: Math. **151**(2) (2021), 705–733.
- [66] Ruedenberg K., Scherr C. W., *Free-Electron Network Model for Conjugated Systems. I. Theory*, J. Chem. Phys. **21**(9) (1953), 1565–1581.
- [67] Serra E., Tentarelli L., *Bound states of the NLS equation on metric graphs with localized nonlinearities*, J. Differential Equations **260**(7) (2016), 5627–5644.
- [68] Serra E., Tentarelli L., *On the lack of bound states for certain NLS equations on metric graphs*, Nonlinear Anal. **145** (2016), 68–82.
- [69] Soave N., *Normalized ground states for the NLS equation with combined nonlinearities*, J. Differential Equations **269**(9) (2020), 6941–6987.
- [70] Soave N., *Normalized ground states for the NLS equation with combined nonlinearities: the Sobolev critical case*, J. Funct. An. **279**(6) (2020), 108610.
- [71] Tao T., Visan M., Zhang X., *The nonlinear Schrödinger equation with combined power-type nonlinearities*, Comm. PDE **32**(7-9) (2007), 1281–1343.
- [72] Tentarelli L., *NLS ground states on metric graphs with localized nonlinearities*, J. Math. Anal. Appl. **433**(1) (2016), 291–304.

¹POLITECNICO DI TORINO, DIPARTIMENTO DI SCIENZE MATEMATICHE “G.L. LAGRANGE”, CORSO DUCA DEGLI ABRUZZI, 24, 10129, TORINO, ITALY
Email address: `riccardo.adami@polito.it`

²POLITECNICO DI TORINO, DIPARTIMENTO DI SCIENZE MATEMATICHE “G.L. LAGRANGE”, CORSO DUCA DEGLI ABRUZZI, 24, 10129, TORINO, ITALY
Email address: `filippo.boni@polito.it`

³UNIVERSITÀ DEGLI STUDI DI TORINO, DIPARTIMENTO DI MATEMATICA “G. PEANO”, VIA CARLO ALBERTO, 10, 10123, TORINO, ITALY
Email address: `filippo.boni@unito.it`

⁴POLITECNICO DI TORINO, DIPARTIMENTO DI SCIENZE MATEMATICHE “G.L. LAGRANGE”, CORSO DUCA DEGLI ABRUZZI, 24, 10129, TORINO, ITALY
Email address: `simone.dovetta@polito.it`