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Production of ω mesons in pp collisions at $\sqrt{s} = 7$ TeV

ALICE Collaboration[⋆]

CERN, 1211 Geneva 23, Switzerland

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Abstract The invariant differential cross section of inclusive $\omega(782)$ meson production at midrapidity ($|y| < 0.5$) in pp collisions at $\sqrt{s} = 7$ TeV was measured with the ALICE detector at the LHC over a transverse momentum range of $2 < p_{\mathrm{T}} < 17\,\mathrm{GeV}/c$. The ω meson was reconstructed via its $\omega \to \pi^+\pi^-\pi^0$ decay channel. The measured ω production cross section is compared to various calculations: PYTHIA 8.2 Monash 2013 describes the data, while PYTHIA 8.2 Tune 4C overestimates the data by about 50%. A recent NLO calculation, which includes a model describing the fragmentation of the whole vector-meson nonet, describes the data within uncertainties below $6\,\mathrm{GeV}/c$, while it overestimates the data by up to 50% for higher p_{T}. The ω/π^0 ratio is in agreement with previous measurements at lower collision energies and the PYTHIA calculations. In addition, the measurement is compatible with transverse mass scaling within the measured p_{T} range and the ratio is constant with $C^{\omega/\pi^0} = 0.67 \pm 0.03\,(\mathrm{stat}) \pm 0.04\,(\mathrm{sys})$ above a transverse momentum of $2.5\,\mathrm{GeV}/c$.

1 Introduction

Measurements of hadron production cross sections in proton–proton (pp) collisions at high energies are important to test our understanding of strong interaction and its underlying theory of quantum chromodynamics (QCD) [1]. Its perturbative treatment (pQCD) becomes feasible for predictions of particle production in hard scattering processes that have a sufficiently high momentum transfer Q^2. This is possible by factorizing [2] the scattering process into three contributions: a QCD matrix element describing the scattering of partons, a parton distribution function (PDF) [3] describing the probability to find a scattering parton within each colliding hadron, and a fragmentation function (FF) [4] that relates the final-state parton momentum to the momentum of an observed hadron. While the QCD matrix element can be calculated in pQCD for sufficiently hard scales, the FFs and PDFs are obtained by global fits of experimental data at various collision energies [5]. However, most particles are produced in soft scattering processes that involve small momentum transfers and therefore can not be calculated within pQCD. In this regime, calculations rely on phenomenological models that also require experimental verification.

Comparison of measured particle spectra with calculations is essential to test their underlying assumptions and provide constraints for the FFs and the PDFs. For example, recent measurements of π^0 and η mesons [6–8] at several LHC collision energies constrained gluon fragmentation [9] in a regime not accessible by measurements at lower collision energies. Like the π^0 and η mesons, the ω meson is comprised mainly of light valence quarks and hence has similar flavor content. However, it has spin 1 and is heavier than the π^0 and η with a mass of $782\,\mathrm{MeV}/c^2$ [10]. These differences make the ω meson an interesting complementary probe to improve our understanding of hadron production in high-energy collisions. Even though there have been several theoretical efforts to describe the fragmentation into pseudoscalar mesons and baryons such as π, K, η and protons [11,12], only a few theoretical models exist to describe the fragmentation into vector mesons, due to a lack of experimental data. Nonetheless, recent efforts [13,14] have been made to describe the fragmentation into the entire vector meson nonet using a model with broken SU(3) symmetry by analysing RHIC (pp) and LEP (e^+e^-) data.

This article presents the invariant differential cross section of inclusive ω meson production at mid-rapidity ($|y| < 0.5$) in pp collisions at $\sqrt{s} = 7$ TeV. The cross section of ω production in hadronic interactions has been measured at collision energies of $\sqrt{s} = 62$ GeV [15] and $\sqrt{s} = 200$ GeV [16–18] at ISR and RHIC respectively. At LHC energies, ω production has only been measured by ALICE at forward rapidities ($2.5 < y < 4.0$) in pp collision at 7 TeV [19] in a transverse momentum (p_{T}) range of $1 < p_{\mathrm{T}} < 5\,\mathrm{GeV}/c$. The results reported here provide the first measurement of ω production at mid-rapidity at LHC energies, and in a wide p_{T} range of $2 < p_{\mathrm{T}} < 17\,\mathrm{GeV}/c$, which tests existing cal-

⋆ e-mail: alice-publications@cern.ch

culations in this regime and provides input for future theoretical studies of vector meson fragmentation functions. In addition, the ω/π^0 production ratio as a function of p_{T} is compared to results of measurements at lower collision energies. This ratio also tests the validity of transverse mass (m_{T}) scaling [20] for ω mesons at LHC energies, which is typically applied to estimate hadronic backgrounds in direct photon or di-electron measurements in situations where no measured hadron spectra are available. The empirical scaling rule, which was established in measurements of identified particle spectra at lower collision energies at ISR and RHIC [21], states that the p_{T}-differential yields of most particles can be described as $E\mathrm{d}^3\sigma/\mathrm{d}p^3 = C^h f(m_{\mathrm{T}})$, where $f(m_{\mathrm{T}})$ is a universal function for all hadron species and C^h is a constant normalisation factor.

The article is structured as follows: Sect. 2 briefly describes the ALICE sub-detectors, with a focus on those relevant for the measurement. Details on the event selection and signal extraction are given in Sects. 3–5. Sources of systematic uncertainties are discussed in Sect. 6. The data and comparisons to model predictions are presented in Sect. 7. Finally, conclusions are provided in Sect. 8.

2 ALICE detector

The ω meson was reconstructed via its decay to $\pi^+\pi^-\pi^0$, where in turn the π^0 decays to two photons. This strategy required the reconstruction of charged tracks in the ALICE central tracking system, composed of the inner tracking system (ITS) [22] and the time projection Chamber (TPC) [23], and the reconstruction of photons using the electromagnetic calorimeter (EMCal) [24,25] and the photon spectrometer (PHOS) [26]. In addition, photons were reconstructed using the photon conversion method (PCM) [27], which exploits the capability of the central tracking system to reconstruct photons from electron–positron track pairs. A detailed description of the ALICE detector system and its performance can be found in Refs. [27,28], respectively. Below, a brief overview of the previously mentioned detectors and the V0 detector [29], used for the minimum bias trigger, is given.

The ITS is positioned closest to the nominal interaction point and consists of two layers of Silicon Pixel Detectors (SPD), two layers of silicon drift detectors (SDD) and two outermost layers of silicon strip detectors (SSD). The layers are positioned between 3.9 cm and 43.0 cm radial distance from the beamline, where the two SPD layers cover a pseudorapidity range of $|\eta| < 2$ and $|\eta| < 1.4$, respectively. The SDD and SSD have a pseudorapidity coverage of $|\eta| < 0.9$ and $|\eta| < 1.0$, respectively. The ITS is used for the tracking of charged particles and the reconstruction of the primary vertex.

The TPC is a large (90 m^3) cylindrical drift detector, which allows for the measurement of charged particles and their identification via specific energy loss ($\mathrm{d}E/\mathrm{d}x$) measurements. The TPC covers a pseudorapidity range of $|\eta| < 0.9$ over the full azimuth and enables the measurement of up to 159 space points per track. A large solenoidal magnet surrounding the central barrel detectors provides a magnetic field of $B = 0.5\,\mathrm{T}$, which allows one to reconstruct tracks down to $p_{\mathrm{T}} \approx 100\,\mathrm{MeV}/c$. For the reconstruction of charged particles in the ITS and TPC, a transverse momentum resolution of about 1% at $1\,\mathrm{GeV}/c$ is achieved, which decreases to about 3% at $10\,\mathrm{GeV}/c$ [23].

The EMCal is a Pb-scintillator sampling calorimeter, which covered an azimuthal range of $\Delta\varphi = 40°$ and $|\eta| < 0.67$ in pseudorapidity during 2010 data taking. In that period, it was comprised of 4 super modules, each consisting of 288 modules. The module consists of four towers with a size of $\approx 6 \times 6\,\mathrm{cm}^2$, corresponding to approximately twice the Molière radius. Each tower is made up of 140 alternating lead and scintillator layers, where the latter are connected to avalanche photo diodes (APDs) that measure the scintillation light of the electromagnetic showers produced by particles traversing the lead absorber. The energy resolution is given by $\sigma_E/E = 4.8\%/E \oplus 11.3\%/\sqrt{E} \oplus 1.7\%$ with energy E in units of GeV [25].

The PHOS is an electromagnetic calorimeter with high granularity based on lead-tungstate ($PbWO_4$) scintillation crystals. At the time these data were collected, it had an acceptance of $\Delta\varphi = 60°$ and $|\eta| < 0.12$, divided into three modules, each consisting of 3584 crystals that are connected to APDs. A high granularity is achieved by small crystal size of $\approx 2.2 \times 2.2\,\mathrm{cm}^2$, where the lateral dimensions of the cells are only slightly larger than the $PbWO_4$ Molière radius of 2 cm. The high light yield of the $PbWO_4$ crystals operated at $-25\,°\mathrm{C}$ results in an energy resolution of $\sigma_E/E = 1.3\%/E \oplus 3.6\%/\sqrt{E} \oplus 1.1\%$ with energy E in units of GeV [30].

The V0 detector provides the minimum bias triggers and is employed to reduce background events, such as beam-gas interactions and out-of-bunch pileup. It consists of two scintillator arrays located in the forward and backward rapidity regions of the ALICE apparatus, covering a pseudorapidity of $2.8 < \eta < 5.1$ and $-3.7 < \eta < -1.7$, respectively.

3 Event and track selection

The pp collision data used for the ω meson measurement were recorded by the ALICE experiment in 2010 at a centre-of-mass energy of $\sqrt{s} = 7\,\mathrm{TeV}$. In 2010, a minimum bias trigger $\mathrm{MB_{OR}}$, which required a signal either in the SPD or in one of the V0 scintillator arrays, was used. The total inelastic pp collision cross section was determined on

the basis of the van der Meer scan and was found to be $\sigma_{\text{inel}} = 73.2^{+2.0}_{-4.6}\,(\text{model}) \pm 2.6\,(\text{lumi})\,\text{mb}$ [31]. The corresponding cross section of the MB_{OR} trigger was $\sigma_{\text{MB}_{\text{OR}}} = (62.4 \pm 2.2)\,\text{mb}$. Beam-induced background events, such as beam-gas interactions or out-of-bunch pileup, are rejected offline by using the timing information from the V0 detectors and the number of reconstructed hit points and track segments in the SPD, which are expected to be uncorrelated for background events. The rejection of in-bunch pileup events, where multiple interactions occur per bunch crossing, was achieved by requiring that only a single primary vertex is reconstructed with the SPD per event. Moreover, collision events with a reconstructed vertex more than 10 cm away from the nominal interaction point along the beam axis were rejected. The integrated luminosities $\mathscr{L}_{\text{int}} = N_{\text{evt}}/\sigma_{\text{MB}_{\text{OR}}}$ were determined to be $\mathscr{L}_{\text{int}}^{\text{EMCal}} = (6.4 \pm 0.2)\,\text{nb}^{-1}$ and $\mathscr{L}_{\text{int}}^{\text{PHOS}} = (6.0 \pm 0.2)\,\text{nb}^{-1}$ for the measurement involving the EMCal and PHOS, respectively. The integrated luminosity of the sample using only the PCM for photon reconstruction amounts to $\mathscr{L}_{\text{int}}^{\text{PCM}} = (7.4 \pm 0.3)\,\text{nb}^{-1}$.

Charged pion trajectories (tracks) with $|\eta| < 0.9$ were reconstructed in the ITS and TPC, requiring at least 70 crossed cathode pad rows in the TPC and at least one hit in any of the layers of the ITS. Furthermore, the χ^2 of the track refit procedure per TPC space point was required to be below 4 and tracks with a momentum below 100 MeV/c were rejected. The tracks were loosely constrained to the collision vertex by requiring a maximum distance of closest approach of a few centimeters to the collision vertex in beam direction and transverse plane. The resolution of the transverse distance to the primary vertex for ITS and TPC charged particle tracks is below 150 μm for $p_{\text{T}} \gtrsim 0.5\,\text{GeV}/c$ [27]. Furthermore, charged pions can be identified using the specific energy loss $\text{d}E/\text{d}x$ along the track in the TPC [32].

4 Photon measurement

To enhance the probability of the reconstruction of π^0 mesons, all methods to measure photons and π^0s at midrapidity with ALICE were exploited. The EMCal and the PHOS allow for the measurement of photons via their electromagnetic shower deposits above $\sim$ 0.5 GeV, while the PCM enables the measurement of photons down to lower p_{T} by exploiting the e^+e^- pair creation by a photon within the inner detector material. Looser photon selection criteria as in previous publications, see e.g. Ref. [33], were applied to increase the ω reconstruction efficiency.

The electromagnetic shower produced in the EMCal or PHOS by an incoming particle usually spreads over multiple adjacent towers, requiring the combination of the individual energy depositions to so-called clusters, which is achieved by clusterisation algorithms [27]. Each reconstructed cluster in the EMCal and PHOS was required to have a total energy of $E_{\text{clus}} > 0.7$ GeV and 0.3 GeV respectively to suppress contributions from minimum-ionising particles and noise. Additionally, in case of the EMCal, it ensures a good timing resolution. Detector noise in a single tower was removed by only selecting clusters with at least 2 (EMCal) or 3 (PHOS) towers for analysis. In order to remove clusters from pileup events originating from subsequent bunch crossings, which occur in $\approx$ 150 ns intervals, a cut on the timing of the leading tower for EMCal clusters of $-100\,\text{ns} < t_{\text{cluster}} < 100\,\text{ns}$ with respect to the collision time was applied. Photon clusters were selected according to their cluster shape and, additionally, a track-matching procedure was applied to suppress clusters originating from charged particles reconstructed in the tracking system. The EMCal cluster shape is parametrised by the larger eigenvalue σ^2_{long} of the dispersion matrix of the shower shape ellipse [33,34]. A requirement of $0.1 \leq \sigma^2_{\text{long}} \leq 0.5$ was imposed, where the lower threshold removes contamination from non-physical background. The upper threshold suppresses elongated clusters originating from low-p_{T} electron and hadron tracks that hit the calorimeter surface not perpendicularly and merged clusters. The latter mostly originate from high-p_{T} neutral pions that decay with a small opening angle, resulting in both decay photons to be reconstructed as a single cluster.

Photons traversing the detector material of ALICE convert to an electron–positron pair with a probability of about 8.5% [27] within a radial distance of 180 cm from the beam axis. Such photons can be reconstructed using the PCM, which allows for the measurement of photons converting in the ITS and TPC within the fiducial acceptance of $|\eta| < 0.9$. First, secondary vertices (V^0s) were reconstructed by an algorithm [35] exploiting the distinct topology of two tracks with opposite curvature that originate from a common point within the tracking detectors. Good reconstruction quality of the tracks associated with a secondary vertex was assured by requiring $p_{\text{T}} > 50\,\text{MeV}/c$ and for the track to be comprised of at least 60% of the findable TPC clusters. Tracks originating from electrons were identified via their specific energy loss $\text{d}E/\text{d}x$ in the TPC, which was required to be within -3 to 5 σ_e of the expected energy loss of electrons, where σ_e is the standard deviation of the measured $\text{d}E/\text{d}x$ distribution of electrons. Contamination of charged pion tracks was suppressed by rejecting tracks whose $\text{d}E/\text{d}x$ was within $\pm 1\sigma_{\pi^\pm}$ of the expected energy loss for pions. Several additional selection criteria were applied to identify V^0 candidates originating from photon conversions, exploiting the kinematics and topology of the conversion, as discussed in more detail in Ref. [8]. These include, e.g. selections to assure that the momentum vector of a conversion pair is pointing towards the primary vertex and a selection based on the min-

imal distance between the conversion point and the primary vertex, in order to remove contributions from Dalitz decays. Furthermore, the quality of the obtained V^0 candidates was improved by constraining the reduced χ^2 of the Kalman-filter hypothesis for the track pair. Remaining contamination from K^0_S, Λ and $\bar{\Lambda}$ decays was reduced by a selection based on the decay kinematics in an Armenteros–Podolanski plot [36], where photon conversions contribute as symmetric decays of a particle with vanishing rest mass. Compared to previous PCM measurements [8,33], a p_T dependence of the selection criteria was introduced to further reduce the contamination from K^0_S and Λ decays.

5 Meson reconstruction

In order to reconstruct the ω mesons via their $\pi^+\pi^-\pi^0$ decay channel, where the π^0 decays to two photons with a branching ratio of $\approx 99\%$, a prior selection of π^0 candidates from pairs of photon candidates was applied. For the photons that passed the selection criteria, the two-photon invariant mass ($M_{\gamma\gamma}$) of all possible photon pairs in a given event was calculated. Four different methods were used for the π^0 candidate reconstruction, differing in how the photons entering the $M_{\gamma\gamma}$ calculation were selected. These are referred to as PCM, PHOS and EMC, when *both* photons used for the π^0 reconstruction were measured with the respective method. In addition, a hybrid method (PCM-EMC) was used, where one PCM photon was combined with a photon measured with the EMCal. The resulting invariant mass distributions exhibit a peak of photon pairs originating from π^0 decays on top of combinatorial background. The peak was parametrised in p_T slices with a Gaussian to characterize the mean and width (σ_{π^0}) of the π^0 mass distribution. Photon pairs lying within about $\pm 2\sigma_{\pi^0}$ of the expected π^0 mass were selected as neutral pion candidates for the ω meson reconstruction. For the PHOS measurement [37], π^0 candidates were furthermore required to have both photons in the same PHOS module and to have a minimum transverse momentum of $1.5\,\mathrm{GeV}/c$. Finally, the nominal neutral pion mass was assigned to the mass of selected π^0 candidates in order to improve the ω mass resolution. This was achieved by subtracting the difference between the reconstructed π^0 mass and its nominal mass from the reconstructed ω mass.

Analogously to the π^0 reconstruction, the invariant mass of all $\pi^+\pi^-\pi^0$ combinations in a given event was determined by summing the four-momentum vectors of the candidate decay products passing the selection criteria. While charged pions were identified by requiring a $\mathrm{d}E/\mathrm{d}x$ within $\pm 3\sigma$ of their expected energy loss, no such selection was applied for the ω analysis with the π^0 reconstructed in PHOS to improve the ω reconstruction efficiency.

Figure 1 shows the invariant mass distribution in the vicinity of the ω nominal mass for indicated p_T intervals for the π^0 reconstructed with PCM, EMC and PHOS, where a peak originating from ω meson decays is clearly visible above the combinatorial background. The latter can be described using a second order polynomial for $p_T < 10\,\mathrm{GeV}/c$. At higher momenta, a first order polynomial was used for the PHOS measurement. The signal obtained after background subtraction was fitted with a Gaussian and the raw yield was obtained by adding counts within $\pm 2\sigma_\omega$ ($\pm 3\sigma_\omega$ for PHOS) of the reconstructed ω mass, where σ_ω denotes the standard deviation of the Gaussian ω signal fit. The ω mass resolution was found to be about $15\,\mathrm{MeV}/c^2$ with a slight dependence on p_T and reconstruction technique. This is achieved by the use of the previously mentioned nominal mass assignment for π^0 candidates, which improved the mass resolution by up to 30%.

The obtained raw yields for each reconstruction method were corrected for geometrical acceptance and reconstruction efficiency, which were evaluated using Monte Carlo simulations. The event generator PYTHIA6.2 [38] was used to simulate the minimum bias pp collisions, where the implemented kinematics of the ω three-body decay are weighted assuming the experimentally observed phase space density distributions [39,40]. All final state particles were propagated through the ALICE detector using GEANT 3 [41], taking into account the operating conditions of the detector at the time of data taking. For each calorimeter, PHOS and EMCal, the relative difference in the energy scale and the non-linearity were tuned in the Monte Carlo to ensure agreement between the p_T-dependent reconstructed π^0 mass and width in data. This agreement propagates to the ω candidates, where mass and width in data and Monte Carlo are found to be consistent within the statistical uncertainties. The full correction factors, ε, that were applied to the raw yields for the four different methods are shown in Fig. 2. These factors include the geometrical acceptance evaluated for each method and the reconstruction efficiency, where the former is normalised to unit rapidity and 2π azimuth angle to allow for a direct comparison between the different methods. The use of the four reconstruction techniques combines the strengths of the individual methods and maximizes the accessible p_T reach. The reconstruction with PCM offers a low p_T-reach, however, the efficiency is limited due to the low conversion probability of about 8.5%, while the reconstruction with the two calorimeters complements the measurement at high p_T.

6 Systematic uncertainties

The systematic and statistical uncertainties on the measured ω yield for the four individual reconstruction techniques in exemplary p_T intervals are summarised in Table 1. The uncer-

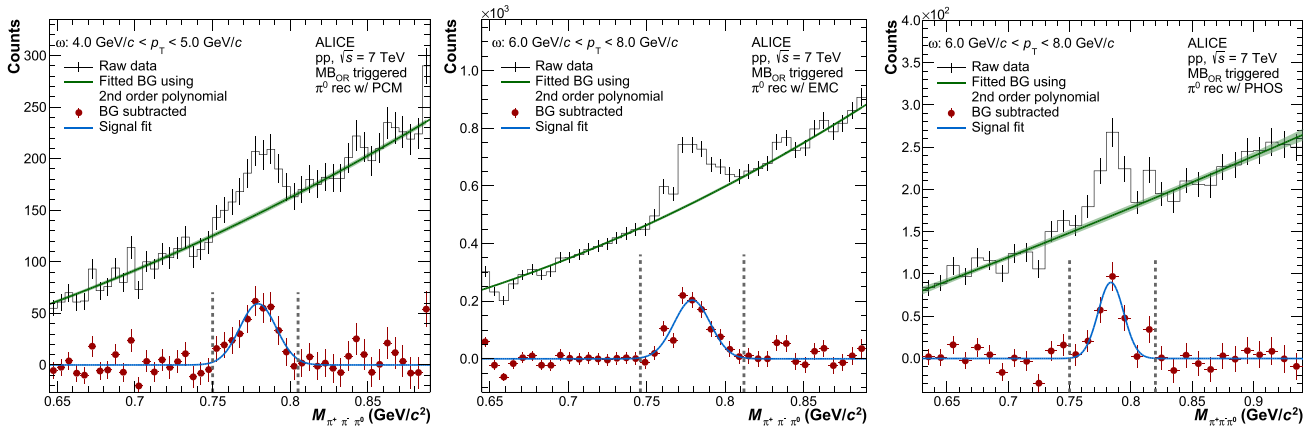


Fig. 1 Invariant mass of $\pi^+\pi^-\pi^0$ candidates shown in the vicinity of the nominal mass of the ω meson for indicated p_T-ranges for π^0 reconstruction with PCM (left), EMC (middle) and PHOS (right). The second order polynomial used for the background description is shown with a band denoting the statistical uncertainties of the fit. The points show the signal obtained after subtraction of the background fit. The signal is fitted with a Gaussian, where the vertical lines indicate the integration range used to obtain the raw yield by bin-by-bin counting, as outlined in Sect. 5

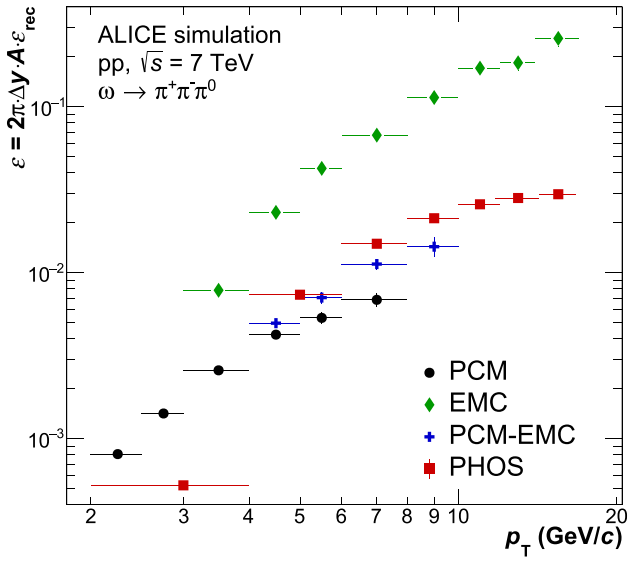


Fig. 2 Correction factors applied to the raw ω yields for each indicated π^0 reconstruction method. The factors include the geometrical acceptance A and the reconstruction efficiency $\epsilon_{\text{rec.}}$. In addition, a normalisation to unit rapidity and 2π azimuth angle is applied to allow for a direct comparison between the different methods

tainties are given as relative uncertainties of the measured values in percent.

The signal extraction dominates the systematic uncertainties of the measurement and includes uncertainties due to the yield extraction. For the PCM, PCM-EMC and EMC techniques the yield extraction uncertainty was estimated by varying simultaneously the bin-counting window used to obtain the raw yield in data and Monte Carlo and the fit range used for the polynomial fit of the combinatorial background. Additionally, for the PHOS analysis, the signal region was excluded from the background fit and the signal was obtained by Gaussian integral instead of bin-by-bin counting. The material budget uncertainty accounts for a possible mismatch between the amount of material present in the ALICE detector and its implementation in GEANT 3. The material budget uncertainty for a conversion photon was studied in Ref. [6], and found to be 4.5% per photon. For the measurements involving the EMCal or the PHOS uncertainties of 3 and 3.5%, respectively, were assigned for the material budget, which is dominated by the material of outer detectors positioned in front of calorimeter modules during data taking in 2010, as outlined in Ref. [33]. The material uncertainty of the inner detectors is negligible for calorimeter photons due to the low conversion probability.

The conversion and calorimeter photon reconstruction uncertainties were evaluated by independently varying the respective selection criteria given in Sect. 4. The conversion photon reconstruction uncertainty was found to be dominated by the reduced χ^2 selection of the electron tracks and the requirement on the number of space points in the TPC for each track. For EMCal related measurements, the cluster energy selection and the cluster shape have most influence on the uncertainty. For PHOS, the photon reconstruction uncertainty was evaluated by variation of the track matching condition and cluster shape selection. Uncertainties arising from the non-linearity and cluster energy scale of the respective calorimeters was taken into account by varying the scheme used to obtain the energy scale calibration and are included in the overall calorimeter photon reconstruction uncertainty. Like the photon reconstruction uncertainties, the systematic uncertainties arising from the charged pion reconstruction were estimated by independent variation of the selection criteria given in Sect. 3. To study the influence of in-bunch

Table 1 Overview of the relative uncertainties given in percent in exemplary p_{T}-intervals for the four individual reconstruction techniques of the ω meson. The given categories summarise systematic uncertainties arising from each analysis step. For each method the statistical and total uncertainties are reported in addition, as well as the uncertainties of the combined measurement. The uncertainty from the $\sigma_{\mathrm{MB_{OR}}}$ determination of 3.5% is independent from the individual measurements and indicated separately in Fig. 3

p_{T} interval	4–5 GeV/c			6–8 GeV/c				12 − 14 GeV/c	
Method	PCM	PCM-EMC	EMC	PCM	PCM-EMC	EMC	PHOS	EMC	PHOS
Signal extraction	12.3	12.6	12.2	13.5	13.5	12.3	6.0	18.9	11.0
Material	9.0	4.7	3.0	9.0	4.7	3.0	3.5	3.0	3.5
Charged pion rec.	6.8	6.8	6.8	6.8	6.8	6.8	6.0	6.8	6.0
Conv photon rec.	4.1	4.1	–	4.1	4.1	–	–	–	–
Calo photon rec.	–	5.0	6.9	–	5.0	6.9	5.2	6.9	9.3
Neutral pion rec.	6.0	6.0	6.0	6.0	6.0	6.0	4.0	6.0	4.0
Pileup	0.5	0.5	0.5	0.5	0.5	0.5	0.5	0.5	0.5
Total syst. uncertainty	18.1	17.7	17.0	19.1	18.3	17.1	11.0	22.4	16.0
Statistical uncertainty	14.5	14.7	9.8	18.9	22.0	9.2	13.0	21.7	32.0
Combined stat. unc.		7.4			7.2			18.0	
Combined syst. unc.		13.7			10.3			16.6	

pileup on the measurement, the rejection criterium was loosened, resulting in a 0.5% systematic uncertainty. The systematic uncertainty due to the selection of neutral pion candidates was estimated by varying the invariant mass selection window. For the PHOS measurement, the selection was additionally varied according to the π^0 candidates transverse momentum. A detailed description of these sources of uncertainty is provided in Refs. [33,37].

Table 1 also shows, for each method, the statistical uncertainty together with the total systematic uncertainty, which is obtained by adding the individual sources in quadrature. In addition, the statistical and systematic uncertainties of the combined measurement are given, which were obtained taken into account correlations across the measurements as elaborated in the following section.

7 Results

The fully corrected invariant cross sections of ω production were obtained for each reconstruction technique using

$$E\frac{\mathrm{d}^3\sigma^{pp\to\omega+X}}{\mathrm{d}p^3} = \frac{1}{2\pi}\frac{1}{p_{\mathrm{T}}}\cdot\frac{1}{\mathscr{L}_{\mathrm{int}}} \cdot\frac{1}{A\cdot\varepsilon_{\mathrm{rec.}}}\cdot\frac{1}{\mathrm{BR}_{\omega\to\pi^+\pi^-\pi^0}}\cdot\frac{N^{\omega}}{\Delta y\Delta p_{\mathrm{T}}}. \quad (1)$$

Here, $\mathscr{L}_{\mathrm{Int}}$ is the integrated luminosity given in Sect. 3, $\varepsilon_{\mathrm{rec.}}$ and A are the reconstruction efficiency and acceptance of the corresponding method and BR $= (89.3 \pm 0.6)\%$ is the branching ratio of the $\omega \to \pi^+\pi^-\pi^0$ decay [10]. Moreover, N^{ω} denotes the number of reconstructed ω mesons in the transverse momentum range Δp_{T} and the given rapidity range Δy.

The production cross sections were measured individually for each reconstruction method and then combined using p_{T}-dependent weights that are calculated according to the best linear unbiased estimate (BLUE) algorithm [42], which uses concepts that are routinely applied in statistical fields. The combination took into account statistical and systematic uncertainties. For the systematic uncertainties, the individual measurements are found to be correlated by about 30%, dominantly originating from the charged-pion selection and the material budget uncertainties. These correlations were taken into account in the combination procedure. The statistical and systematic uncertainties of the combined measurement are given in Table 1.

The cross section of ω meson production for $2 < p_{\mathrm{T}} < 17\,\mathrm{GeV}/c$ at midrapidity in pp collisions at $\sqrt{s} = 7\,\mathrm{TeV}$ is shown in Fig. 3. It was fitted using a Levy–Tsallis function [46] given by

$$E\frac{\mathrm{d}^3\sigma}{\mathrm{d}p^3} = \frac{C}{2\pi}\frac{(n-1)(n-2)}{nT[nT+m(n-2)]}\left(1+\frac{m_T-m}{nT}\right)^{-n}, \quad (2)$$

which describes the cross section over the whole measured transverse momentum range, as demonstrated in the lower panel of the figure. The parameters m and $m_{\mathrm{T}} = \sqrt{m^2+p_{\mathrm{T}}^2}$ correspond to the particle mass and the transverse mass, respectively, while C, T and n are the free parameters of the Levy–Tsallis function.

The values of the fit parameters and the reduced χ^2 of the fit are given in Table 2, where the fit was obtained using only statistical uncertainties, and using the systematic and statistical uncertainties of the measurement added in quadrature. To account for finite p_{T}-interval width, the combined cross section points were assigned to p_{T} values shifted from the bin centre of the p_{T} intervals according to the underlying spectrum [47] described by a Levy–Tsallis function. This correction resulted in a shift below 2% in each p_{T} interval.

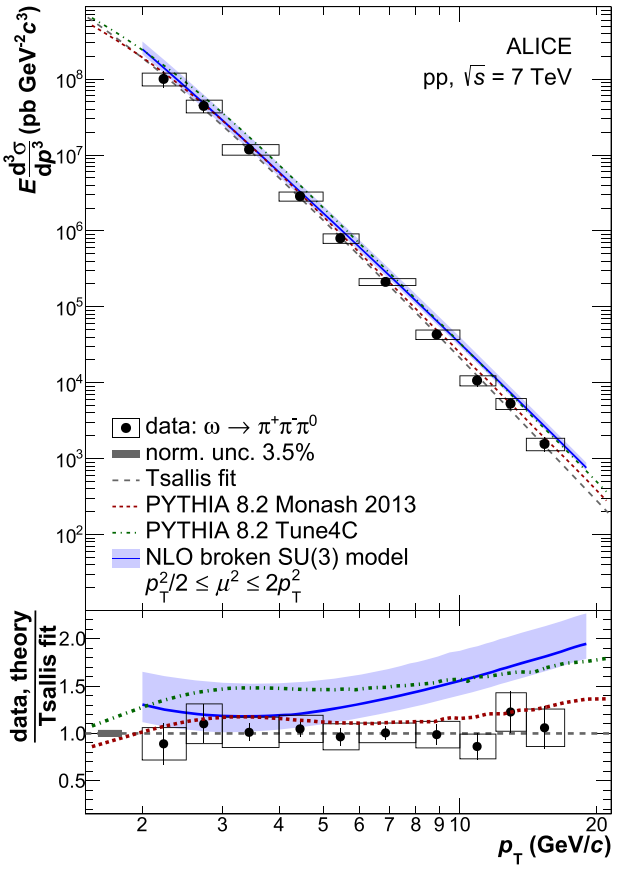


Fig. 3 Invariant cross section of ω meson production in pp collisions at $\sqrt{s} = 7$ TeV compared to theoretical predictions. The statistical and systematic uncertainties are represented by vertical bars and boxes, respectively. A Levy–Tsallis function was used to parametrise the spectrum, where the obtained fit parameters are given in Table 2. The predictions are obtained using PYTHIA 8.2 [43] with the Monash 2013 [44] and 4C [45] tunes. Furthermore, a NLO calculation [14] incorporating a model dedicated to describe vector-meson fragmentation is shown, where the band denotes the uncertainty of the scale μ, which was used for factorisation, renormalisation and fragmentation. In the bottom panel, the ratios of the theoretical estimates to the Levy–Tsallis fit of the measurement are shown; the ratio of the data to the Levy–Tsallis fit is also presented

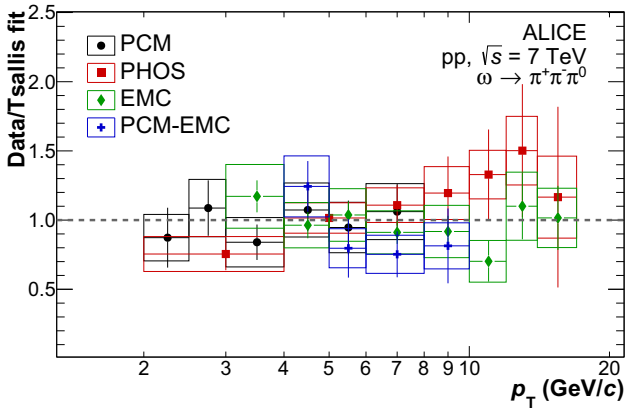


Fig. 4 Ratios of the fully corrected ω spectra obtained with the individual reconstruction methods to the Levy–Tsallis fit of the combined spectrum, where the fit parameters are shown in Table 2. The statistical and systematic uncertainties are represented by vertical bars and boxes, respectively

Figure 4, which shows the ratios of the cross sections for the individual reconstruction methods to the Levy–Tsallis fit of the combined measurement, demonstrates the agreement between all methods within the statistical and systematic uncertainties, justifying the combination of the individual results as discussed earlier.

The measured differential cross section of ω production is compared to several calculations in Fig. 3. The ratio of each prediction to the Levy–Tsallis fit of the measurement is shown in the bottom panel of the figure. Two PYTHIA 8.2 [43] Monte Carlo event generator calculations were considered for comparison, which are based on the Monash 2013 [44] and the 4C [45] tunes, respectively. The Monash 2013 tune describes the measurement over the full reported p_{T} range within the uncertainties, while the Tune 4C overestimates the data by about 50%. The Monash 2013 tune includes more recent experimental results than Tune 4C and thus a more refined set of parameters. In particular, the rate of light flavor vector meson production used in hadronisation process was revised and lowered, improving the description of ω meson yields [44].

The measurement is also compared to a next-to-leading order (NLO) calculation using a model with broken SU(3) symmetry to describe vector meson production [14], where the model parameters have been constrained using ω production data measured by PHENIX in pp collisions at $\sqrt{s} = 200$ GeV [16]. The same scale $\mu = p_{\mathrm{T}}$ was used for factorisation, renormalisation and fragmentation for the calculation and the shaded band reported in Fig. 3 denotes the scale variation of ${p_{\mathrm{T}}}^2/2 \leq \mu^2 \leq 2{p_{\mathrm{T}}}^2$. The calculation describes the measurement within the uncertainties below 6 GeV/c, and overestimates the data by up to 50% for higher p_{T}.

The ratio of ω relative to π^0 meson production is shown as a function of p_{T} in Fig. 5, where data points for the π^0 measurement were taken from Ref. [6]. The ratio is observed to be constant above 2.5 GeV/c with a value of

Table 2 Parameters and χ^2/NDF of the fit to the ω invariant cross section using the Levy–Tsallis function [46] from Eq. 2

Levy–Tsallis	C ($\times 10^{10}$ pb)	T (GeV)	n	χ^2/NDF	NDF
ω	$4.01^{\pm 2.47\ (\mathrm{stat})}_{\pm 3.41\ (\mathrm{tot})}$	$0.182^{\pm 0.042\ (\mathrm{stat})}_{\pm 0.061\ (\mathrm{tot})}$	$6.46^{\pm 0.37\ (\mathrm{stat})}_{\pm 0.55\ (\mathrm{tot})}$	$0.45\ (\mathrm{stat})$ $0.22\ (\mathrm{tot})$	7

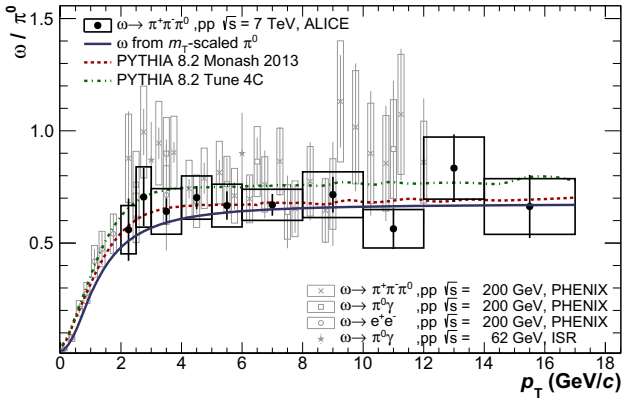


Fig. 5 Ratio of ω/π^0 production as a function of transverse momentum p_{T} for pp collisions at $\sqrt{s} = 7$ TeV (black) compared to various lower collision energies ranging from $\sqrt{s} = 62-200$ GeV [15–18] (gray). In addition, PYTHIA predictions at $\sqrt{s} = 7$ TeV and the ω/π^0 ratio obtained from m_{T}-scaling are shown with lines

$C^{\omega/\pi^0} = 0.69 \pm 0.03$ (stat) ± 0.04 (sys). Within the uncertainties, the ω/π^0 ratio is described by the PYTHIA predictions. Even though the Tune 4C overestimates the ω production, it describes the ω/π^0 ratio due to a similar overestimation of π^0 production, which was reported in Ref. [8].

The measured ω/π^0 ratio at $\sqrt{s} = 7$ TeV is compared to data from lower collision energies at $\sqrt{s} = 62$ [15] and 200 GeV [16–18]. The ω/π^0 ratios measured at the different collision energies agree within the uncertainties. In order to test the validity of m_{T}-scaling, the Levy–Tsallis parametrisation $f_{\pi^0}(p_{\mathrm{T},\pi^0})$ of the π^0 spectrum reported in Ref. [6] was scaled using the ratio $C^{\omega/\pi^0} = 0.67$, following the procedure discussed in detail in Ref. [20]. The scaled parametrisation $f_{\omega}(p_{\mathrm{T},\omega})$ was used to calculate the ω/π^0 ratio via $f_{\omega}(p_{\mathrm{T},\omega})/f_{\pi^0}(p_{\mathrm{T},\omega})$, where the relation $p^2_{\mathrm{T},\omega} + m^2_{0,\omega} = p^2_{\mathrm{T},\pi^0} + m^2_{0,\pi^0}$ was used to ensure the evaluation of both spectra at the same transverse mass. The obtained m_{T}-scaling prediction of the ω/π^0 ratio is shown in Fig. 5 and found to be consistent with the measurement. Unlike in the case of the η/π^0 ratio measured at $\sqrt{s} = 2.76$, 7 and 8 TeV [6,8,33], where a violation of m_{T}-scaling was observed below 3.5 GeV, no such violation is observed within the uncertainties for the ω meson in the entire measured momentum range. However, while the measurement is compatible with the m_{T}-scaling prediction at low-p_{T}, the sensitivity of the measurement to a possible m_{T}-scaling violation is limited by the uncertainties and p_{T} reach. Here, future studies with increased precision could provide further insights and more stringent tests of m_{T}-scaling for low-p_{T} ω mesons. Interestingly, the PYTHIA calculations and the m_{T}-scaled prediction both describe the ω/π^0 ratio at lower collision energies even below $p_{\mathrm{T}} = 2$ GeV/c, suggesting a universal feature of meson production.

8 Conclusion

The invariant differential cross section of ω meson production at midrapidity in pp collisions $\sqrt{s} = 7$ TeV was measured with the ALICE detector, covering a transverse-momentum range of 2 to 17 GeV/c. Within the uncertainties, PYTHIA 8.2 predictions for the Monash 2013 tune describes the measurement over the whole p_{T} range, while Tune 4C overestimates the data by about 50%. A NLO calculation using a model dedicated to describing fragmentation into the entire vector meson nonet describes the data below 6 GeV/c, while it overestimates the data by up to 50% at higher p_{T}. Above 2.5 GeV/c the ω/π^0 ratio is found to be constant with a value of $C^{\omega/\pi^0} = 0.67 \pm 0.03$ (stat) ± 0.04 (sys) and agrees with measurements at lower collision energies and with PYTHIA predictions over the whole reported p_{T} range. Within the uncertainties, the m_{T}-scaling prediction for ω mesons is consistent with the measured spectrum above 2 GeV/c.

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ALICE Collaboration

S. Acharya[141], D. Adamová[95], A. Adler[74], J. Adolfsson[81], M. M. Aggarwal[100], S. Agha[14], G. Aglieri Rinella[34], M. Agnello[30], N. Agrawal[10,54], Z. Ahammed[141], S. Ahmad[16], S. U. Ahn[76], Z. Akbar[51], A. Akindinov[92], M. Al-Turany[107], S. N. Alam[40], D. S. D. Albuquerque[122], D. Aleksandrov[88], B. Alessandro[59], H. M. Alfanda[6], R. Alfaro Molina[71], B. Ali[16], Y. Ali[14], A. Alici[10,26,54], N. Alizadehvandchali[125], A. Alkin[2,34], J. Alme[21], T. Alt[68], L. Altenkamper[21], I. Altsybeev[113], M. N. Anaam[6], C. Andrei[48], D. Andreou[34], A. Andronic[144], M. Angeletti[34], V. Anguelov[104], T. Antičić[108], F. Antinori[57], P. Antonioli[54], N. Apadula[80], L. Aphecetche[115], H. Appelshäuser[68], S. Arcelli[26], R. Arnaldi[59], M. Arratia[80], I. C. Arsene[20], M. Arslandok[104], A. Augustinus[34], R. Averbeck[107], S. Aziz[78], M. D. Azmi[16], A. Badalà[56], Y. W. Baek[41], S. Bagnasco[59], X. Bai[107], R. Bailhache[68], R. Bala[101], A. Balbino[30], A. Baldisseri[137], M. Ball[43], S. Balouza[105], D. Banerjee[3], R. Barbera[27], L. Barioglio[25], G. G. Barnaföldi[145], L. S. Barnby[94], V. Barret[134], P. Bartalini[6], C. Bartels[127], K. Barth[34], E. Bartsch[68], F. Baruffaldi[28], N. Bastid[134], S. Basu[143], G. Batigne[115], B. Batyunya[75], D. Bauri[49], J. L. Bazo Alba[112], I. G. Bearden[89], C. Beattie[146], C. Bedda[63], I. Belikov[136], A. D. C. Bell Hechavarria[144], F. Bellini[34], R. Bellwied[125], V. Belyaev[93], G. Bencedi[145], S. Beole[25], A. Bercuci[48], Y. Berdnikov[98], D. Berenyi[145], R. A. Bertens[130], D. Berzano[59], M. G. Besoiu[67], L. Betev[34], A. Bhasin[101], I. R. Bhat[101], M. A. Bhat[3], H. Bhatt[49], B. Bhattacharjee[42], A. Bianchi[25], L. Bianchi[25], N. Bianchi[52], J. Bielčík[37], J. Bielčíková[95], A. Bilandzic[105], G. Biro[145], R. Biswas[3], S. Biswas[3], J. T. Blair[119], D. Blau[88], C. Blume[68], G. Boca[139], F. Bock[96], A. Bogdanov[93], S. Boi[23], J. Bok[61], L. Boldizsár[145], A. Bolozdynya[93], M. Bombara[38], G. Bonomi[140], H. Borel[137], A. Borissov[93], H. Bossi[146], E. Botta[25], L. Bratrud[68], P. Braun-Munzinger[107], M. Bregant[121], M. Broz[37], E. Bruna[59], G. E. Bruno[33,106], M. D. Buckland[127], D. Budnikov[109], H. Buesching[68], S. Bufalino[30], O. Bugnon[115], P. Buhler[114], P. Buncic[34], Z. Buthelezi[72,131], J. B. Butt[14], S. A. Bysiak[118], D. Caffarri[90], A. Caliva[107], E. Calvo Villar[112], J. M. M. Camacho[120], R. S. Camacho[45], P. Camerini[24], F. D. M. Canedo[121], A. A. Capon[114], F. Carnesecchi[26], R. Caron[137], J. Castillo Castellanos[137], A. J. Castro[130], E. A. R. Casula[55], F. Catalano[30], C. Ceballos Sanchez[75], P. Chakraborty[49], S. Chandra[141], W. Chang[6], S. Chapeland[34], M. Chartier[127], S. Chattopadhyay[141], S. Chattopadhyay[110], A. Chauvin[23], C. Cheshkov[135], B. Cheynis[135], V. Chibante Barroso[34], D. D. Chinellato[122], S. Cho[61], P. Chochula[34], T. Chowdhury[134], P. Christakoglou[90], C. H. Christensen[89], P. Christiansen[81], T. Chujo[133], C. Cicalo[55], L. Cifarelli[10,26], F. Cindolo[54], M. R. Ciupek[107], G. Clai[54,b], J. Cleymans[124], F. Colamaria[53], D. Colella[53], A. Collu[80], M. Colocci[26], M. Concas[59,c], G. Conesa Balbastre[79], Z. Conesa del Valle[78], G. Contin[24,60], J. G. Contreras[37], T. M. Cormier[96], Y. Corrales Morales[25], P. Cortese[31], M. R. Cosentino[123], F. Costa[34], S. Costanza[139], P. Crochet[134], E. Cuautle[69], P. Cui[6], L. Cunqueiro[96], D. Dabrowski[142], T. Dahms[105], A. Dainese[57], F. P. A. Damas[115,137], M. C. Danisch[104], A. Danu[67], D. Das[110], I. Das[110], P. Das[86], P. Das[3], S. Das[3], A. Dash[86], S. Dash[49], S. De[86], A. De Caro[29], G. de Cataldo[53], L. De Cilladi[25], J. de Cuveland[39], A. De Falco[23], D. De Gruttola[10], N. De Marco[59], C. De Martin[24], S. De Pasquale[29], S. Deb[50], H. F. Degenhardt[121], K. R. Deja[142], A. Deloff[85], S. Delsanto[25,131], W. Deng[6], P. Dhankher[49], D. Di Bari[33], A. Di Mauro[34], R. A. Diaz[8], T. Dietel[124], P. Dillenseger[68], Y. Ding[6], R. Divià[34], D. U. Dixit[19], Ø. Djuvsland[21], U. Dmitrieva[62], A. Dobrin[67], B. Dönigus[68], O. Dordic[20], A. K. Dubey[141], A. Dubla[90,107], S. Dudi[100], M. Dukhishyam[86], P. Dupieux[134], R. J. Ehlers[96], V. N. Eikeland[21], D. Elia[53], B. Erazmus[115], F. Erhardt[99], A. Erokhin[113], M. R. Ersdal[21], B. Espagnon[78], G. Eulisse[34], D. Evans[111], S. Evdokimov[91], L. Fabbietti[105], M. Faggin[28], J. Faivre[79], F. Fan[6], A. Fantoni[52], M. Fasel[96], P. Fecchio[30], A. Feliciello[59], G. Feofilov[113], A. Fernández Téllez[45], A. Ferrero[137], A. Ferretti[25], A. Festanti[34], V. J. G. Feuillard[104], J. Figiel[118], S. Filchagin[109], D. Finogeev[62], F. M. Fionda[21], G. Fiorenza[53], F. Flor[125], A. N. Flores[119], S. Foertsch[72], P. Foka[107], S. Fokin[88], E. Fragiacomo[60], U. Frankenfeld[107], U. Fuchs[34], C. Furget[79], A. Furs[62], M. Fusco Girard[29], J. J. Gaardhøje[89], M. Gagliardi[25], A. M. Gago[112], A. Gal[136], C. D. Galvan[120], P. Ganoti[84], C. Garabatos[107], J. R. A. Garcia[45], E. Garcia-Solis[11], K. Garg[115], C. Gargiulo[34], A. Garibli[87], K. Garner[144], P. Gasik[105,107], E. F. Gauger[119], M. B. Gay Ducati[70], M. Germain[115], J. Ghosh[110], P. Ghosh[141], S. K. Ghosh[3], M. Giacalone[26], P. Gianotti[52], P. Giubellino[59,107], P. Giubilato[28], A. M. C. Glaenzer[137], P. Glässel[104], A. Gomez Ramirez[74], V. Gonzalez[107,143], L.H. González-Trueba[71], S. Gorbunov[39], L. Görlich[118], A. Goswami[49], S. Gotovac[35], V. Grabski[71], L. K. Graczykowski[142], K. L. Graham[111], L. Greiner[80], A. Grelli[63], C. Grigoras[34], V. Grigoriev[93], A. Grigoryan[1], S. Grigoryan[75], O. S. Groettvik[21], F. Grosa[30,59], J. F. Grosse-Oetringhaus[34], R. Grosso[107], R. Guernane[79], M. Guittiere[115], K. Gulbrandsen[89], T. Gunji[132], A. Gupta[101], R. Gupta[101], I. B. Guzman[45], R. Haake[146], M. K. Habib[107], C. Hadjidakis[78], H. Hamagaki[82], G. Hamar[145], M. Hamid[6], R. Hannigan[119], M. R. Haque[86], A. Harlenderova[107], J. W. Harris[146], A. Harton[11], J. A. Hasenbichler[34], H. Hassan[96], Q. U. Hassan[14], D. Hatzifotiadou[10,54], P. Hauer[43], L. B. Havener[146], S. Hayashi[132], S. T. Heckel[105], E. Hellbär[68], H. Helstrup[36], A. Herghelegiu[48], T. Herman[37], E. G. Hernandez[45], G. Herrera Corral[9], F. Herrmann[144], K. F. Hetland[36], H. Hillemanns[34], C. Hills[127], B. Hippolyte[136], B. Hohlweger[105], J. Honermann[144], D. Horak[37], A. Hornung[68], S. Hornung[107], R. Hosokawa[15,133], P. Hristov[34], C. Huang[78], C. Hughes[130], P. Huhn[68], T. J. Humanic[97], H. Hushnud[110], L. A. Husova[144], N. Hussain[42], S. A. Hussain[14], D. Hutter[39], J. P. Iddon[34,127],

R. Ilkaev[109], H. Ilyas[14], M. Inaba[133], G. M. Innocenti[34], M. Ippolitov[88], A. Isakov[95], M. S. Islam[110], M. Ivanov[107], V. Ivanov[98], V. Izucheev[91], B. Jacak[80], N. Jacazio[34,54], P. M. Jacobs[80], S. Jadlovska[117], J. Jadlovsky[117], S. Jaelani[63], C. Jahnke[121], M. J. Jakubowska[142], M. A. Janik[142], T. Janson[74], M. Jercic[99], O. Jevons[111], M. Jin[125], F. Jonas[96,144], P. G. Jones[111], J. Jung[68], M. Jung[68], A. Jusko[111], P. Kalinak[64], A. Kalweit[34], V. Kaplin[93], S. Kar[6], A. Karasu Uysal[77], D. Karatovic[99], O. Karavichev[62], T. Karavicheva[62], P. Karczmarczyk[142], E. Karpechev[62], A. Kazantsev[88], U. Kebschull[74], R. Keidel[47], M. Keil[34], B. Ketzer[43], Z. Khabanova[90], A. M. Khan[6], S. Khan[16], A. Khanzadeev[98], Y. Kharlov[91], A. Khatun[16], A. Khuntia[118], B. Kileng[36], B. Kim[61], B. Kim[133], D. Kim[147], D. J. Kim[126], E. J. Kim[73], H. Kim[17], J. Kim[147], J. S. Kim[41], J. Kim[104], J. Kim[147], J. Kim[73], M. Kim[104], S. Kim[18], T. Kim[147], T. Kim[147], S. Kirsch[68], I. Kisel[39], S. Kiselev[92], A. Kisiel[142], J. L. Klay[5], C. Klein[68], J. Klein[34,59], S. Klein[80], C. Klein-Bösing[144], M. Kleiner[68], T. Klemenz[105], A. Kluge[34], M. L. Knichel[34], A. G. Knospe[125], C. Kobdaj[116], M. K. Köhler[104], T. Kollegger[107], A. Kondratyev[75], N. Kondratyeva[93], E. Kondratyuk[91], J. Konig[68], S. A. Konigstorfer[105], P. J. Konopka[34], G. Kornakov[142], L. Koska[117], O. Kovalenko[85], V. Kovalenko[113], M. Kowalski[118], I. Králik[64], A. Kravčáková[38], L. Kreis[107], M. Krivda[64,111], F. Krizek[95], K. Krizkova Gajdosova[37], M. Krüger[68], E. Kryshen[98], M. Krzewicki[39], A. M. Kubera[97], V. Kučera[34,61], C. Kuhn[136], P. G. Kuijer[90], L. Kumar[100], S. Kundu[86], P. Kurashvili[85], A. Kurepin[62], A. B. Kurepin[62], A. Kuryakin[109], S. Kushpil[95], J. Kvapil[111], M. J. Kweon[61], J. Y. Kwon[61], Y. Kwon[147], S. L. La Pointe[39], P. La Rocca[27], Y. S. Lai[80], A. Lakrathok[116], M. Lamanna[34], R. Langoy[129], K. Lapidus[34], A. Lardeux[20], P. Larionov[52], E. Laudi[34], R. Lavicka[37], T. Lazareva[113], R. Lea[24], L. Leardini[104], J. Lee[133], S. Lee[147], S. Lehner[114], J. Lehrbach[39], R. C. Lemmon[94], I. León Monzón[120], E. D. Lesser[19], M. Lettrich[34], P. Lévai[145], X. Li[12], X. L. Li[6], J. Lien[129], R. Lietava[111], B. Lim[17], V. Lindenstruth[39], A. Lindner[48], C. Lippmann[107], M. A. Lisa[97], A. Liu[19], J. Liu[127], S. Liu[97], W. J. Llope[143], I. M. Lofnes[21], V. Loginov[93], C. Loizides[96], P. Loncar[35], J. A. Lopez[104], X. Lopez[134], E. López Torres[8], J. R. Luhder[144], M. Lunardon[28], G. Luparello[60], Y. G. Ma[40], A. Maevskaya[62], M. Mager[34], S. M. Mahmood[20], T. Mahmoud[43], A. Maire[136], R. D. Majka[146,a], M. Malaev[98], Q. W. Malik[20], L. Malinina[75,d], D. Mal'Kevich[92], P. Malzacher[107], G. Mandaglio[32,56], V. Manko[88], F. Manso[134], V. Manzari[53], Y. Mao[6], M. Marchisone[135], J. Mareš[66], G. V. Margagliotti[24], A. Margotti[54], A. Marín[107], C. Markert[119], M. Marquard[68], N. A. Martin[104], P. Martinengo[34], J. L. Martinez[125], M. I. Martínez[45], G. Martínez García[115], S. Masciocchi[107], M. Masera[25], A. Masoni[55], L. Massacrier[78], E. Masson[115], A. Mastroserio[53,138], A. M. Mathis[105], O. Matonoha[81], P. F. T. Matuoka[121], A. Matyja[118], C. Mayer[118], F. Mazzaschi[25], M. Mazzilli[53], M. A. Mazzoni[58], A. F. Mechler[68], F. Meddi[22], Y. Melikyan[62,93], A. Menchaca-Rocha[71], C. Mengke[6], E. Meninno[29,114], A. S. Menon[125], M. Meres[13], S. Mhlanga[124], Y. Miake[133], L. Micheletti[25], L. C. Migliorin[135], D. L. Mihaylov[105], K. Mikhaylov[75,92], A. N. Mishra[69], D. Miśkowiec[107], A. Modak[3], N. Mohammadi[34], A. P. Mohanty[63], B. Mohanty[86], M. Mohisin Khan[16,e], Z. Moravcova[89], C. Mordasini[105], D. A. Moreira De Godoy[144], L. A. P. Moreno[45], I. Morozov[62], A. Morsch[34], T. Mrnjavac[34], V. Muccifora[52], E. Mudnic[35], D. Mühlheim[144], S. Muhuri[141], J. D. Mulligan[80], A. Mulliri[23,55], M. G. Munhoz[121], R. H. Munzer[68], H. Murakami[132], S. Murray[124], L. Musa[34], J. Musinsky[64], C. J. Myers[125], J. W. Myrcha[142], B. Naik[49], R. Nair[85], B. K. Nandi[49], R. Nania[10,54], E. Nappi[53], M. U. Naru[14], A. F. Nassirpour[81], C. Nattrass[130], R. Nayak[49], T. K. Nayak[86], S. Nazarenko[109], A. Neagu[20], R. A. Negrao De Oliveira[68], L. Nellen[69], S. V. Nesbo[36], G. Neskovic[39], D. Nesterov[113], L. T. Neumann[142], B. S. Nielsen[89], S. Nikolaev[88], S. Nikulin[88], V. Nikulin[98], F. Noferini[10,54], P. Nomokonov[75], J. Norman[79,127], N. Novitzky[133], P. Nowakowski[142], A. Nyanin[88], J. Nystrand[21], M. Ogino[82], A. Ohlson[81], J. Oleniacz[142], A. C. Oliveira Da Silva[130], M. H. Oliver[146], C. Oppedisano[59], A. Ortiz Velasquez[69], T. Osako[46], A. Oskarsson[81], J. Otwinowski[118], K. Oyama[82], Y. Pachmayer[104], V. Pacik[89], S. Padhan[49], D. Pagano[140], G. Paić[69], J. Pan[143], S. Panebianco[137], P. Pareek[50,141], J. Park[61], J. E. Parkkila[126], S. Parmar[100], S. P. Pathak[125], B. Paul[23], J. Pazzini[140], H. Pei[6], T. Peitzmann[63], X. Peng[6], L. G. Pereira[70], H. Pereira Da Costa[137], D. Peresunko[88], G. M. Perez[8], S. Perrin[137], Y. Pestov[4], V. Petráček[37], M. Petrovici[48], R. P. Pezzi[70], S. Piano[60], M. Pikna[13], P. Pillot[115], O. Pinazza[34,54], L. Pinsky[125], C. Pinto[27], S. Pisano[10,52], D. Pistone[56], M. Płoskoń[80], M. Planinic[99], F. Pliquett[68], M. G. Poghosyan[96], B. Polichtchouk[91], N. Poljak[99], A. Pop[48], S. Porteboeuf-Houssais[134], V. Pozdniakov[75], S. K. Prasad[3], R. Preghenella[54], F. Prino[59], C. A. Pruneau[143], I. Pshenichnov[62], M. Puccio[34], J. Putschke[143], S. Qiu[90], L. Quaglia[25], R. E. Quishpe[125], S. Ragoni[111], S. Raha[3], S. Rajput[101], J. Rak[126], A. Rakotozafindrabe[137], L. Ramello[31], F. Rami[136], S. A. R. Ramirez[45], R. Raniwala[102], S. Raniwala[102], S. S. Räsänen[44], R. Rath[50], V. Ratza[43], I. Ravasenga[90], K. F. Read[96,130], A. R. Redelbach[39], K. Redlich[85,f], A. Rehman[21], P. Reichelt[68], F. Reidt[34], X. Ren[6], R. Renfordt[68], Z. Rescakova[38], K. Reygers[104], A. Riabov[98], V. Riabov[98], T. Richert[81,89], M. Richter[20], P. Riedler[34], W. Riegler[34], F. Riggi[27], C. Ristea[67], S. P. Rode[50], M. Rodríguez Cahuantzi[45], K. Røed[20], R. Rogalev[91], E. Rogochaya[75], D. Rohr[34], D. Röhrich[21], P. F. Rojas[45], P. S. Rokita[142], F. Ronchetti[52], A. Rosano[56], E. D. Rosas[69], K. Roslon[142], A. Rossi[57], A. Rotondi[139], A. Roy[50], P. Roy[110], O. V. Rueda[81], R. Rui[24], B. Rumyantsev[75], A. Rustamov[87], E. Ryabinkin[88], Y. Ryabov[98], A. Rybicki[118], H. Rytkonen[126], O. A. M. Saarimaki[44], R. Sadek[115], S. Sadhu[141], S. Sadovsky[91], K. Šafařík[37], S. K. Saha[141], B. Sahoo[49], P. Sahoo[49], R. Sahoo[50], S. Sahoo[65], P. K. Sahu[65], J. Saini[141], S. Sakai[133],

S. Sambyal[101], V. Samsonov[93,98], D. Sarkar[143], N. Sarkar[141], P. Sarma[42], V. M. Sarti[105], M. H. P. Sas[63], E. Scapparone[54], J. Schambach[119], H. S. Scheid[68], C. Schiaua[48], R. Schicker[104], A. Schmah[104], C. Schmidt[107], H. R. Schmidt[103], M. O. Schmidt[104], M. Schmidt[103], N. V. Schmidt[68,96], A. R. Schmier[130], J. Schukraft[89], Y. Schutz[136], K. Schwarz[107], K. Schweda[107], G. Scioli[26], E. Scomparin[59], J. E. Seger[15], Y. Sekiguchi[132], D. Sekihata[132], I. Selyuzhenkov[93,107], S. Senyukov[136], D. Serebryakov[62], A. Sevcenco[67], A. Shabanov[62], A. Shabetai[115], R. Shahoyan[34], W. Shaikh[110], A. Shangaraev[91], A. Sharma[100], A. Sharma[101], H. Sharma[118], M. Sharma[101], N. Sharma[100], S. Sharma[101], O. Sheibani[125], K. Shigaki[46], M. Shimomura[83], S. Shirinkin[92], Q. Shou[40], Y. Sibiriak[88], S. Siddhanta[55], T. Siemiarczuk[85], D. Silvermyr[81], G. Simatovic[90], G. Simonetti[34], B. Singh[105], R. Singh[86], R. Singh[101], R. Singh[50], V. K. Singh[141], V. Singhal[141], T. Sinha[110], B. Sitar[13], M. Sitta[31], T. B. Skaali[20], M. Slupecki[44], N. Smirnov[146], R. J. M. Snellings[63], C. Soncco[112], J. Song[125], A. Songmoolnak[116], F. Soramel[28], S. Sorensen[130], I. Sputowska[118], J. Stachel[104], I. Stan[67], P. J. Steffanic[130], E. Stenlund[81], S. F. Stiefelmaier[104], D. Stocco[115], M. M. Storetvedt[36], L. D. Stritto[29], A. A. P. Suaide[121], T. Sugitate[46], C. Suire[78], M. Suleymanov[14], M. Suljic[34], R. Sultanov[92], M. Šumbera[95], V. Sumberia[101], S. Sumowidagdo[51], S. Swain[65], A. Szabo[13], I. Szarka[13], U. Tabassam[14], S. F. Taghavi[105], G. Taillepied[134], J. Takahashi[122], G. J. Tambave[21], S. Tang[6,134], M. Tarhini[115], M. G. Tarzila[48], A. Tauro[34], G. Tejeda Muñoz[45], A. Telesca[34], L. Terlizzi[25], C. Terrevoli[125], D. Thakur[50], S. Thakur[141], D. Thomas[119], F. Thoresen[89], R. Tieulent[135], A. Tikhonov[62], A. R. Timmins[125], A. Toia[68], N. Topilskaya[62], M. Toppi[52], F. Torales-Acosta[19], S. R. Torres[37], A. Trifiró[32,56], S. Tripathy[50,69], T. Tripathy[49], S. Trogolo[28], G. Trombetta[33], L. Tropp[38], V. Trubnikov[2], W. H. Trzaska[126], T. P. Trzcinski[142], B. A. Trzeciak[37,63], A. Tumkin[109], R. Turrisi[57], T. S. Tveter[20], K. Ullaland[21], E. N. Umaka[125], A. Uras[135], G. L. Usai[23], M. Vala[38], N. Valle[139], S. Vallero[59], N. van der Kolk[63], L. V. R. van Doremalen[63], M. van Leeuwen[63], P. Vande Vyvre[34], D. Varga[145], Z. Varga[145], M. Varga-Kofarago[145], A. Vargas[45], M. Vasileiou[84], A. Vasiliev[88], O. Vázquez Doce[105], V. Vechernin[113], E. Vercellin[25], S. Vergara Limón[45], L. Vermunt[63], R. Vernet[7], R. Vértesi[145], M. Verweij[63], L. Vickovic[35], Z. Vilakazi[131], O. Villalobos Baillie[111], G. Vino[53], A. Vinogradov[88], T. Virgili[29], V. Vislavicius[89], A. Vodopyanov[75], B. Volkel[34], M. A. Völkl[103], K. Voloshin[92], S. A. Voloshin[143], G. Volpe[33], B. von Haller[34], I. Vorobyev[105], D. Voscek[117], J. Vrláková[38], B. Wagner[21], M. Weber[114], S. G. Weber[144], A. Wegrzynek[34], S. C. Wenzel[34], J. P. Wessels[144], J. Wiechula[68], J. Wikne[20], G. Wilk[85], J. Wilkinson[10], G. A. Willems[144], E. Willsher[111], B. Windelband[104], M. Winn[137], W. E. Witt[130], J. R. Wright[119], Y. Wu[128], R. Xu[6], S. Yalcin[77], Y. Yamaguchi[46], K. Yamakawa[46], S. Yang[21], S. Yano[137], Z. Yin[6], H. Yokoyama[63], I.-K. Yoo[17], J. H. Yoon[61], S. Yuan[21], A. Yuncu[104], V. Yurchenko[2], V. Zaccolo[24], A. Zaman[14], C. Zampolli[34], H. J. C. Zanoli[63], N. Zardoshti[34], A. Zarochentsev[113], P. Závada[66], N. Zaviyalov[109], H. Zbroszczyk[142], M. Zhalov[98], S. Zhang[40], X. Zhang[6], Z. Zhang[6], V. Zherebchevskii[113], Y. Zhi[12], D. Zhou[6], Y. Zhou[89], Z. Zhou[21], J. Zhu[6,107], Y. Zhu[6], A. Zichichi[10,26], G. Zinovjev[2], N. Zurlo[140]

[1] A.I. Alikhanyan National Science Laboratory (Yerevan Physics Institute) Foundation, Yerevan, Armenia
[2] Bogolyubov Institute for Theoretical Physics, National Academy of Sciences of Ukraine, Kiev, Ukraine
[3] Department of Physics and Centre for Astroparticle Physics and Space Science (CAPSS), Bose Institute, Kolkata, India
[4] Budker Institute for Nuclear Physics, Novosibirsk, Russia
[5] California Polytechnic State University, San Luis Obispo, CA, USA
[6] Central China Normal University, Wuhan, China
[7] Centre de Calcul de l'IN2P3, Villeurbanne, Lyon, France
[8] Centro de Aplicaciones Tecnológicas y Desarrollo Nuclear (CEADEN), Havana, Cuba
[9] Centro de Investigación y de Estudios Avanzados (CINVESTAV), Mexico City and Mérida, Mexico
[10] Centro Fermi, Museo Storico della Fisica e Centro Studi e Ricerche "Enrico Fermi', Rome, Italy
[11] Chicago State University, Chicago, IL, USA
[12] China Institute of Atomic Energy, Beijing, China
[13] Faculty of Mathematics, Physics and Informatics, Comenius University Bratislava, Bratislava, Slovakia
[14] COMSATS University Islamabad, Islamabad, Pakistan
[15] Creighton University, Omaha, NE, USA
[16] Department of Physics, Aligarh Muslim University, Aligarh, India
[17] Department of Physics, Pusan National University, Pusan, Republic of Korea
[18] Department of Physics, Sejong University, Seoul, Republic of Korea
[19] Department of Physics, University of California, Berkeley, CA, USA
[20] Department of Physics, University of Oslo, Oslo, Norway
[21] Department of Physics and Technology, University of Bergen, Bergen, Norway
[22] Dipartimento di Fisica dell'Università 'La Sapienza' and Sezione INFN, Rome, Italy

[23] Dipartimento di Fisica dell'Università and Sezione INFN, Cagliari, Italy
[24] Dipartimento di Fisica dell'Università and Sezione INFN, Trieste, Italy
[25] Dipartimento di Fisica dell'Università and Sezione INFN, Turin, Italy
[26] Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Bologna, Italy
[27] Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Catania, Italy
[28] Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Padua, Italy
[29] Dipartimento di Fisica 'E.R. Caianiello' dell'Università and Gruppo Collegato INFN, Salerno, Italy
[30] Dipartimento DISAT del Politecnico and Sezione INFN, Turin, Italy
[31] Dipartimento di Scienze e Innovazione Tecnologica dell'Università del Piemonte Orientale and INFN Sezione di Torino, Alessandria, Italy
[32] Dipartimento di Scienze MIFT, Università di Messina, Messina, Italy
[33] Dipartimento Interateneo di Fisica 'M. Merlin' and Sezione INFN, Bari, Italy
[34] European Organization for Nuclear Research (CERN), Geneva, Switzerland
[35] Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, University of Split, Split, Croatia
[36] Faculty of Engineering and Science, Western Norway University of Applied Sciences, Bergen, Norway
[37] Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague, Prague, Czech Republic
[38] Faculty of Science, P.J. Šafárik University, Košice, Slovakia
[39] Frankfurt Institute for Advanced Studies, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
[40] Fudan University, Shanghai, China
[41] Gangneung-Wonju National University, Gangneung, Republic of Korea
[42] Gauhati University, Department of Physics, Guwahati, India
[43] Helmholtz-Institut für Strahlen- und Kernphysik, Rheinische Friedrich-Wilhelms-Universität Bonn, Bonn, Germany
[44] Helsinki Institute of Physics (HIP), Helsinki, Finland
[45] High Energy Physics Group, Universidad Autónoma de Puebla, Puebla, Mexico
[46] Hiroshima University, Hiroshima, Japan
[47] Hochschule Worms, Zentrum für Technologietransfer und Telekommunikation (ZTT), Worms, Germany
[48] Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania
[49] Indian Institute of Technology Bombay (IIT), Mumbai, India
[50] Indian Institute of Technology Indore, Indore, India
[51] Indonesian Institute of Sciences, Jakarta, Indonesia
[52] INFN, Laboratori Nazionali di Frascati, Frascati, Italy
[53] INFN, Sezione di Bari, Bari, Italy
[54] INFN, Sezione di Bologna, Bologna, Italy
[55] INFN, Sezione di Cagliari, Cagliari, Italy
[56] INFN, Sezione di Catania, Catania, Italy
[57] INFN, Sezione di Padova, Padua, Italy
[58] INFN, Sezione di Roma, Rome, Italy
[59] INFN, Sezione di Torino, Turin, Italy
[60] INFN, Sezione di Trieste, Trieste, Italy
[61] Inha University, Incheon, Republic of Korea
[62] Institute for Nuclear Research, Academy of Sciences, Moscow, Russia
[63] Institute for Subatomic Physics, Utrecht University/Nikhef, Utrecht, The Netherlands
[64] Institute of Experimental Physics, Slovak Academy of Sciences, Košice, Slovakia
[65] Institute of Physics, Homi Bhabha National Institute, Bhubaneswar, India
[66] Institute of Physics of the Czech Academy of Sciences, Prague, Czech Republic
[67] Institute of Space Science (ISS), Bucharest, Romania
[68] Institut für Kernphysik, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
[69] Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Mexico City, Mexico
[70] Instituto de Física, Universidade Federal do Rio Grande do Sul (UFRGS), Porto Alegre, Brazil
[71] Instituto de Física, Universidad Nacional Autónoma de México, Mexico City, Mexico
[72] iThemba LABS, National Research Foundation, Somerset West, South Africa
[73] Jeonbuk National University, Jeonju, Republic of Korea

[74] Johann-Wolfgang-Goethe Universität Frankfurt Institut für Informatik, Fachbereich Informatik und Mathematik, Frankfurt, Germany
[75] Joint Institute for Nuclear Research (JINR), Dubna, Russia
[76] Korea Institute of Science and Technology Information, Daejeon, Republic of Korea
[77] KTO Karatay University, Konya, Turkey
[78] Laboratoire de Physique des 2 Infinis, Irène Joliot-Curie, Orsay, France
[79] Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS-IN2P3, Grenoble, France
[80] Lawrence Berkeley National Laboratory, Berkeley, CA, USA
[81] Lund University Department of Physics, Division of Particle Physics, Lund, Sweden
[82] Nagasaki Institute of Applied Science, Nagasaki, Japan
[83] Nara Women's University (NWU), Nara, Japan
[84] National and Kapodistrian University of Athens, School of Science, Department of Physics , Athens, Greece
[85] National Centre for Nuclear Research, Warsaw, Poland
[86] National Institute of Science Education and Research, Homi Bhabha National Institute, Jatni, India
[87] National Nuclear Research Center, Baku, Azerbaijan
[88] National Research Centre Kurchatov Institute, Moscow, Russia
[89] Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
[90] Nikhef, National institute for subatomic physics, Amsterdam, The Netherlands
[91] NRC Kurchatov Institute IHEP, Protvino, Russia
[92] NRC «Kurchatov»Institute - ITEP, Moscow, Russia
[93] NRNU Moscow Engineering Physics Institute, Moscow, Russia
[94] Nuclear Physics Group, STFC Daresbury Laboratory, Daresbury, UK
[95] Nuclear Physics Institute of the Czech Academy of Sciences, Řež u Prahy, Czech Republic
[96] Oak Ridge National Laboratory, Oak Ridge, TN, USA
[97] Ohio State University, Columbus, OH, USA
[98] Petersburg Nuclear Physics Institute, Gatchina, Russia
[99] Physics department, Faculty of science, University of Zagreb, Zagreb, Croatia
[100] Physics Department, Panjab University, Chandigarh, India
[101] Physics Department, University of Jammu, Jammu, India
[102] Physics Department, University of Rajasthan, Jaipur, India
[103] Physikalisches Institut, Eberhard-Karls-Universität Tübingen, Tübingen, Germany
[104] Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
[105] Physik Department, Technische Universität München, Munich, Germany
[106] Politecnico di Bari, Bari, Italy
[107] Research Division and ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung GmbH, Darmstadt, Germany
[108] Rudjer Bošković Institute, Zagreb, Croatia
[109] Russian Federal Nuclear Center (VNIIEF), Sarov, Russia
[110] Saha Institute of Nuclear Physics, Homi Bhabha National Institute, Kolkata, India
[111] School of Physics and Astronomy, University of Birmingham, Birmingham, UK
[112] Sección Física, Departamento de Ciencias, Pontificia Universidad Católica del Perú, Lima, Peru
[113] St. Petersburg State University, St. Petersburg, Russia
[114] Stefan Meyer Institut für Subatomare Physik (SMI), Vienna, Austria
[115] SUBATECH, IMT Atlantique, Université de Nantes, CNRS-IN2P3, Nantes, France
[116] Suranaree University of Technology, Nakhon Ratchasima, Thailand
[117] Technical University of Košice, Košice, Slovakia
[118] The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Kraków, Poland
[119] The University of Texas at Austin, Austin, TX, USA
[120] Universidad Autónoma de Sinaloa, Culiacán, Mexico
[121] Universidade de São Paulo (USP), São Paulo, Brazil
[122] Universidade Estadual de Campinas (UNICAMP), Campinas, Brazil
[123] Universidade Federal do ABC, Santo Andre, Brazil
[124] University of Cape Town, Cape Town, South Africa

[125] University of Houston, Houston, TX, USA
[126] University of Jyväskylä, Jyväskylä, Finland
[127] University of Liverpool, Liverpool, UK
[128] University of Science and Technology of China, Hefei, China
[129] University of South-Eastern Norway, Tonsberg, Norway
[130] University of Tennessee, Knoxville, TN, USA
[131] University of the Witwatersrand, Johannesburg, South Africa
[132] University of Tokyo, Tokyo, Japan
[133] University of Tsukuba, Tsukuba, Japan
[134] Université Clermont Auvergne, CNRS/IN2P3, LPC, Clermont-Ferrand, France
[135] Université de Lyon, Université Lyon 1, CNRS/IN2P3, IPN-Lyon, Villeurbanne, Lyon, France
[136] Université de Strasbourg, CNRS, IPHC UMR 7178, F-67000 Strasbourg; France, Strasbourg, France
[137] Université Paris-Saclay Centre d'Etudes de Saclay (CEA), IRFU, Départment de Physique Nucléaire (DPhN), Saclay, France
[138] Università degli Studi di Foggia, Foggia, Italy
[139] Università degli Studi di Pavia, Pavia, Italy
[140] Università di Brescia, Brescia, Italy
[141] Variable Energy Cyclotron Centre, Homi Bhabha National Institute, Kolkata, India
[142] Warsaw University of Technology, Warsaw, Poland
[143] Wayne State University, Detroit, MI, USA
[144] Westfälische Wilhelms-Universität Münster, Institut für Kernphysik, Muenster, Germany
[145] Wigner Research Centre for Physics, Budapest, Hungary
[146] Yale University, New Haven, CT, USA
[147] Yonsei University, Seoul, Republic of Korea

[a] Deceased
[b] Italian National Agency for New Technologies, Energy and Sustainable Economic Development (ENEA), Bologna, Italy
[c] Dipartimento DET del Politecnico di Torino, Turin, Italy
[d] M.V. Lomonosov Moscow State University, D.V. Skobeltsyn Institute of Nuclear, Physics, Moscow, Russia
[e] Department of Applied Physics, Aligarh Muslim University, Aligarh, India
[f] Institute of Theoretical Physics, University of Wroclaw, Wrocław, Poland