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# Polynomial ring representations of endomorphisms of exterior powers

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## Abstract

An explicit description of the ring of the rational polynomials in  $r$  indeterminates as a representation of the Lie algebra of the endomorphisms of the  $k$ -th exterior power of a countably infinite-dimensional vector space is given. Our description is based on results by Laksov and Throup concerning the symmetric structure of the exterior power of a polynomial ring. Our results are based on approximate versions of the vertex operators occurring in the celebrated bosonic vertex representation, due to Date, Jimbo, Kashiwara and Miwa, of the Lie algebra of all matrices of infinite size, whose entries are all zero but finitely many.

**Keywords** Hasse–Schmidt derivations and vertex operators on exterior algebras · Bosonic and fermionic representations by Date–Jimbo–Kashiwara–Miwa · Symmetric functions

**Mathematics Subject Classification** 14M15 · 15A75 · 05E05 · 17B69

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# 1 Introduction

## 1.1 Statement of the main result

The purpose of this paper is to supply an explicit description of the polynomial ring  $B_r := \mathbb{Q}[e_1, \dots, e_r]$  as a module over the Lie algebra of endomorphisms of  $k$ -th exterior powers of a vector space  $V := \bigoplus_{i \geq 0} \mathbb{Q} \cdot b_i$  of infinite countable dimension.

Let  $\beta_j : V \rightarrow \mathbb{Q}$  be the unique linear form such that  $\beta_j(b_i) = \delta_{ji}$  so that  $V^* := \bigoplus_{j \geq 0} \mathbb{Q} \cdot \beta_j$  is the *restricted dual* of  $V$ . Write the  $r$ -th exterior powers of  $V$  and  $V^*$ , respectively, as  $\bigwedge^r V := \bigoplus_{\lambda \in \mathcal{P}_r} \mathbb{Q} \cdot [\mathbf{b}]_\lambda^r$  and  $\bigwedge^r V^* := \bigoplus_{\lambda \in \mathcal{P}_r} \mathbb{Q} \cdot [\beta]_\lambda^r$ , with  $[\mathbf{b}]_\lambda^r := b_{r-1+\lambda_1} \wedge \dots \wedge b_{\lambda_r}$  and  $[\beta]_\lambda^r([\mathbf{b}]_\mu^r) = \delta_{\lambda,\mu}$ , where  $\lambda$  and  $\mu$  range over the set  $\mathcal{P}_r$  of all the partitions of length at most  $r$ .

A strict relative of  $\bigwedge^r V$  is the vector space  $B_r$ , which can be identified with the ring of symmetric polynomials in  $r$  indeterminates. It is well known that it possesses a  $\mathbb{Q}$ -basis formed by certain Schur determinants  $S_\lambda := S_\lambda(e_1, \dots, e_r)$  (Cf. Sect. 2.4, adopting the notation of [12, p. 41]). The  $\mathbb{Q}$ -linear extension of the set map  $S_\lambda \mapsto [\mathbf{b}]_\lambda^r$  yields a  $\mathbb{Q}$ -vector space isomorphism  $B_r \rightarrow \bigwedge^r V$ , sending  $1 \mapsto [\mathbf{b}]_0^r := b_{r-1} \wedge \dots \wedge b_0$ . It is convenient to phrase it by saying that  $\bigwedge^r V$  is a free  $B_r$ -module of rank 1 generated by  $[\mathbf{b}]_0^r$ , such that  $[\mathbf{b}]_\lambda^r = \Delta_\lambda(H_r)[\mathbf{b}]_0^r$  (Sect. 2.3).

Let  $gl(\bigwedge^k V)$  be the Lie algebra of the endomorphisms of  $\bigwedge^k V$  vanishing at  $[\mathbf{b}]_\lambda^k$  for all but finitely many partitions  $\lambda \in \mathcal{P}_k$  and denote by  $\mathcal{E}_{\mu,\nu}^k$  the *elementary endomorphism*  $[\mathbf{b}]_\mu^k \otimes [\beta]_\nu^k$ , so that:

$$gl\left(\bigwedge^k V\right) = \bigwedge^k V \otimes \bigwedge^k V^* = \bigoplus_{\mu,\nu \in \mathcal{P}_k} \mathbb{Q} \cdot \mathcal{E}_{\mu,\nu}^k.$$

For all  $k, r \geq 0$ , we consider the  $B_r$ -representation of  $gl(\bigwedge^k V)$ , which we understand as the action:

$$(\mathcal{E}_{\mu,\nu}^k \Delta_\lambda(H_r))[\mathbf{b}]_0^r = [\mathbf{b}]_\mu^k \wedge ([\beta]_{\nu \downarrow}^k [\mathbf{b}]_\lambda^r), \tag{1}$$

where  $[\beta]_{\nu \downarrow}^k : \bigwedge^r V \rightarrow \bigwedge^{r-k} V$  is the standard *contraction* operator (Sect. 2.2).

To express the  $gl(\bigwedge^k V)$ -action (1) on  $B_r$  through a compact formula, a standard philosophy suggests to use generating functions.

To this purpose, let us introduce some notation. Let  $\mathbf{z}_k := (z_1, \dots, z_k)$  and  $\mathbf{w}_k := (w_1, \dots, w_k)$  be two sets of formal variables. The  $k$ -tuples of the formal inverses  $(z_1^{-1}, \dots, z_k^{-1})$  and  $(w_1^{-1}, \dots, w_k^{-1})$  will be denoted by  $\mathbf{z}_k^{-1}$  and  $\mathbf{w}_k^{-1}$  respectively. If  $\mathbf{u} := (u_1, \dots, u_k)$  are arbitrary formal variables, denote by  $p_i(\mathbf{u}_k)$  the power sum  $u_1^i + \dots + u_k^i$  of degree  $i$ . The standard notation  $s_\mu(\mathbf{z}_k)$  and  $s_\nu(\mathbf{w}_k^{-1})$  stands for the symmetric Schur polynomials in the  $\mathbf{z}_k$  and  $\mathbf{w}_k^{-1}$  (See [12, p. 40]). Let  $E_r(z) := 1 - e_1 z + \dots + (-1)^r e_r z^r \in B_r[z]$ , set by convention  $b_j = 0$  if  $j < 0$  and denote by  $\sigma_{-1}$  the locally nilpotent endomorphism of  $V$  mapping  $b_j \mapsto b_{j-1}$  for all  $j \geq 0$ . Let  $\delta : \text{End}(V) \mapsto \text{End}(\bigwedge V)$  be the natural representation of  $\text{End}(V)$  as a Lie algebra of (even) derivations of  $\bigwedge V$ .

The main ingredients to state our main result are certain *vertex operators*  $\Gamma(\mathbf{z}_k), \Gamma^*(\mathbf{z}_k) : \bigwedge V \rightarrow \bigwedge V[\mathbf{z}_k, \mathbf{z}_k^{-1}]$  acting on the exterior algebra  $\bigwedge V$ . They are introduced in Definition 4.3 as products of *Schubert derivations*, and studied in more detail only in Sects. 6 and 7. However, if  $r$  is big with respect to the length of the partition  $\lambda$  they can be explicitly written as

$$\Gamma(\mathbf{z}_k)[\mathbf{b}]_\lambda^r := \prod_{j=1}^k \frac{1}{E_r(z_j)} \exp\left(-\sum_{i \geq 1} \frac{1}{i} \delta(\sigma_{-1}^i) p_i(\mathbf{z}_k^{-1})\right) [\mathbf{b}]_\lambda^{r+k}$$

and

$$\Gamma^*(\mathbf{w}_k)[\mathbf{b}]_\lambda^r := \prod_{j=1}^k E_r(w_j) \exp\left(\sum_{i \geq 1} \frac{1}{i} \delta(\sigma_{-1}^i) p_i(\mathbf{w}_k^{-1})\right) [\mathbf{b}]_\lambda^{r-k}.$$

Consider now the generating formal power series

$$\mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1}) = \sum_{\mu, \nu \in \mathcal{P}_k} \mathcal{E}_{\mu\nu}^k \cdot s_\mu(\mathbf{z}_k) s_\nu(\mathbf{w}_k^{-1}) : B_r \rightarrow B_r[\mathbf{z}_k, \mathbf{w}_k^{-1}],$$

defined by the equality:

$$\left(\mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1}) S_\lambda\right) [\mathbf{b}]_0^r = \sum_{\mu, \nu \in \mathcal{P}_k} s_\mu(\mathbf{z}_k) s_\nu(\mathbf{w}_k^{-1}) [\mathbf{b}]_\lambda^k \wedge ([\beta]_{\nu, \lambda}^k [\mathbf{b}]_\lambda^r).$$

**Main Theorem.** For all  $k, r \geq 0$  and all  $\lambda \in \mathcal{P}_r$ , the action of  $\mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1})$  on the basis element  $S_\lambda$  of  $B_r$  is given by:

$$\left(\mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1}) S_\lambda\right) [\mathbf{b}]_0^r = \prod_{j=1}^k \left(\frac{z_j}{w_j}\right)^{r-k} \cdot \Gamma(\mathbf{z}_k) \Gamma^*(\mathbf{w}_k) [\mathbf{b}]_\lambda^r. \tag{2}$$

The above result corresponds to Theorem 8.5 within the text and supplies the explicit description of the ring  $B_r$  as a module over the Lie algebra  $gl(\bigwedge^k V)$ , for all  $k \geq 0$ . In fact, the  $\mathcal{E}_{\mu, \nu}^k$ -image of  $S_\lambda$  is determined by the coefficient of  $s_\mu(\mathbf{z}_k) s_\nu(\mathbf{w}_k^{-1})$  obtained by the expansion of the right hand side of (2). This may sounds tricky to evaluate, but is nothing else than the coefficient of

$$z_1^{k-1+\mu_1} \dots z_k^{\mu_k} \cdot w_1^{-k+1-\nu_1} \dots w_k^{-\nu_k}$$

of the right hand side of (2), multiplied by the Vandermonde determinants of  $\mathbf{z}_k$  and  $\mathbf{w}_k^{-1}$ .

That  $B_r$  is a representation of  $gl(\bigwedge^k V)$  is easy to see in very special cases. For  $k = 0$ , it is the multiplication by rational numbers, as  $\bigwedge^0 V = \mathbb{Q}$ , while for  $k > r$  is the trivial null representation. The case  $r = k = 1$  recovers the well known general fact that every vector space is a module over the Lie algebra of its own endomorphisms. In fact the linear extension of the set map  $e_i^j \mapsto b_i$  is a vector space isomorphism  $B_1 \rightarrow V$ , making  $B_1$  into a  $gl(V)$ -module, by pulling back that structure from  $V$ . Our Main Theorem then takes into account the general case.

### 1.2 The boson–fermion correspondence and the DJKM representation

The  $gl(\bigwedge^k V)$ -module structure of  $B_r$ , described in Main Theorem, will be referred to as *bosonic representation* of  $gl(\bigwedge^k V)$ , by a possibly strong, but suggestive, abuse of terminology, due to the evident relationship with the pioneering work by Date, Jimbo, Kashiwara and Miwa (DJKM) [8] (see also [24, 25]) which also fits into the more general framework considered in the reference [7].

As a matter of fact, one main motivation of this paper was to better understand a fundamental, although elementary, representation theoretical fact. Let  $\mathcal{V} := \bigoplus_{j \in \mathbb{Z}} \mathbb{Q} \cdot b_j$  be a vector space with basis  $\mathbf{b} := (b_j)_{j \in \mathbb{Z}}$ , parameterized by the integers (one may think of  $\mathcal{V}$  as being the vector space  $\mathbb{Q}[X^{-1}, X]$  of the Laurent polynomials) and  $\mathcal{V}^*$  its restricted dual with basis  $(\beta_j)_{j \in \mathbb{Z}}$ . It is well known that  $\mathcal{V} \oplus \mathcal{V}^*$  supports a canonical structure of Clifford algebra  $\mathcal{C} := \mathcal{C}(\mathcal{V} \oplus \mathcal{V}^*)$  ([9, p. 85] or [18]) and that the *Fermionic Fock space*  $F$  (also called the semi-infinite wedge power and denoted by  $\bigwedge^{\infty/2} \mathcal{V}$ ) is an irreducible representation of  $\mathcal{C}$ . More precisely,  $F$  is an invertible module over the Lie super-algebra  $\mathcal{C}$  generated by a distinguished vector  $|0\rangle$ , the *vacuum*, that in the formalism of the infinite wedge power can be suggestively written as  $b_0 \wedge b_{-1} \wedge b_{-2} \wedge \dots$

The huge Clifford algebra  $\mathcal{C}$ , whose elements are finite linear combinations of words of the form  $b_{i_1} \dots b_{i_h} \beta_{j_1} \dots \beta_{j_k}$ , contains in a natural way all, but not only, the Lie algebras  $gl(\bigwedge^k V)$ , for all  $k \geq 0$ . In particular, it turns out that  $F$  is a  $gl(\bigwedge^k \mathcal{V})$ -module for all  $k \geq 0$ . Then, the bosonic Fock space  $B := B_\infty := \mathbb{Q}[e_1, e_2, \dots]$  gets a  $gl(\bigwedge^k V)$ -module structure, for all  $k \geq 0$ , pulled back from  $F$  via the *boson-fermion correspondence*, a natural module isomorphism  $B \rightarrow F$  over the infinite dimensional Lie Heisenberg algebra. The latter may well be interpreted as a sort of Poincaré duality for infinite dimensional Grassmannians. The case  $k = 1$  recovers precisely the celebrated bosonic vertex operator representation of the Lie algebra  $gl_\infty(\mathbb{Q}) := gl(\mathcal{V}) = gl(\bigwedge^1 V)$  due to Date, Jimbo, Kashiwara and Miwa ([8, 24]) as in e.g. [25, Theorem 5.1] we have alluded above.

Our paper, however, aims to look at more traditional, but relevant, contexts. Exactly as in the case of the Fermionic Fock space  $\mathcal{F}$ , the exterior algebra of  $V \cong \mathbb{Q}[X]$  is an irreducible representation of the canonical Clifford algebra  $\mathcal{C}$  supported on  $V \oplus V^*$ . This occurrence convinced ourselves to give a closer look to the  $gl(\bigwedge V)$ -structure of  $\bigwedge V$ , certainly not treated in any literature we have consulted up to now.

We have so gotten a description of the  $gl(\bigwedge^k V)$ -module structure of  $B_r$ , which generalises the case  $r < \infty$  and  $k = 1$  studied in [20]. The output is that the direct sum  $\bigoplus_{k \geq 0} gl(\bigwedge^k V)$  is a Lie subalgebra of  $gl(\bigwedge V)$ , represented by  $B_r$  for all  $r \geq 0$ . In the case of the fermionic Fock space, the  $gl(\bigwedge^1 V)$ -structure of  $B_\infty$  is the DJKM one [8, 24]. The general case, which amounts to the description of the  $gl(\bigwedge \mathcal{V})$ -module structure of  $\mathcal{F}$ , is faced in the contribution [2] as a best example of the extension of the techniques used in [21].

### 1.3 Methods and their applications

The vertex operators occurring in our description of  $B_r$  as a representation of  $gl(\bigwedge^k V)$  are defined by means of *Schubert derivations*, which are distinguished Hasse–Schmidt (HS) derivations on exterior algebras.

HS derivations were first introduced in [13] and extensively treated in [18]; see also the survey [1] or [6, p. 116], for more discussions. In a finite dimensional context Schubert derivations are related to Chern and Segre polynomials of the tautological bundle over a Grassmannian. The point is that the Segre and Chern polynomials act as a HS-derivation on the exterior algebra of the homology of the projective space, which is the same as saying that to do Schubert calculus on Grassmannians, Bézout theorem suffices.

Hasse-Schmidt derivations on exterior algebras have shown their versatility in applications to improve effectiveness in Schubert Calculus computations (see [4, 5]), to equivariant cohomology of Grassmannians (Cf. [22], but also [26]), to generalise the Cayley–Hamilton theorem [16, 23], with perspective applications to globalise the local Wronskian as in [15, Section 4.2], or, inspired by [14, 17], like in [2, 20, 21] and in the present paper, to

revisit the bosonic vertex representation of Lie algebras of endomorphisms as in [8] (see also [24] and [25, Propositions 5.2–5.3]), providing new methods and new insight.

The Schubert derivations we introduce here, denoted by  $\sigma_+(z)$ ,  $\bar{\sigma}_+(z)$ ,  $\sigma_-(w)$  and  $\bar{\sigma}_-(w)$  enjoy some nice commutation rules. Those with the same sign as subscripts commute in the algebra of endomorphisms of the exterior algebra. However, due to the fact that  $\sigma_-(w)$  and  $\bar{\sigma}_-(w)$  are locally nilpotent, they commute with  $\sigma_+(z)$  and  $\bar{\sigma}_+(z)$  only up to the multiplication by a rational function.

## 1.4 Organization of the paper

To be as much self-contained as possible, we collect most of preliminaries and basic notation in Sect. 2. The first part recalls basics of the theory of symmetric polynomials as, e.g., in [12]. The second part accounts for the invertible Hasse–Schmidt derivation on an exterior algebra, essential in the subsequent sections.

Section 3 contains the explicit expression of the Schubert derivations that makes evident their strong connection with vertex operators.

Section 4 also contains an effective definition of what we have proposed to name vertex operators on a Grassmann algebra, because an obvious relationship with those occurring in the classical boson–fermion correspondence. Also, to check the Main Theorem without neglecting any minimum detail, we state and prove in Sect. 4 relevant commutation rules, some of which can be recognized within the phrasing of the categorical framework for the boson–fermion correspondence, depicted in [10] (see also [28] for a recent update).

In Sect. 5 we instead study commutation rules involving the contraction operator: this is a typical issue in the situation involving a finite wedge power. In fact, the finiteness makes the subject trickier than when working with the infinite wedge power.

Vertex operators in the sense of Definition 4.3 are homogeneous operators on the exterior algebra, one of positive and the other of negative degree. We devote one section to each one of them (Sects. 7 and 8) to dig up their relationship with basic computations in multilinear algebra, such as wedging and contracting.

Certainly this idea is already present in the infinite wedge power context (e.g. [25, Chapter 5]), but the present article, together with [19–21], is the first instance of applications of the techniques and ideas in finite dimensional landscapes.

Finally, last Sect. 8 is concerned with the proof of the Main Theorem together with some of its straightforward declinations in terms of certain familiar objects, like suitable deformations of the same Giambelli’s determinants occurring in classical Schubert Calculus, see Theorem 8.10. To achieve the proof of the Main Theorem, some preliminary lemmas (such as 8.2 and 8.3) are proved. We believe that these lemmas along with Theorems 6.5 and 7.3, are interesting in their own, as pieces of multilinear algebra properties addressed to wider general mathematical audiences.

## 2 Preliminaries and notation

### 2.1 Partitions

A *partition* is a monotonic non increasing sequence  $\lambda := (\lambda_1 \geq \lambda_2 \geq \dots)$  of non negative integers, said to be its *parts*. The *length*  $\ell(\lambda)$  is the number of its non zero parts, and  $|\lambda| = \sum_{i \geq 0} \lambda_i$  is its *weight*. We denote by  $\mathcal{P}_r$  be the set of all partitions of length at most  $r$ .

### 2.2 Exterior powers, exterior algebras and duality pairing

Let  $V := \bigoplus_{i \geq 0} \mathbb{Q} \cdot b_i$  be the vector space with basis  $\mathbf{b} := (b_i)_{i \geq 0}$ . The *restricted dual* of  $V$  is  $V^* := \bigoplus_{j \geq 0} \mathbb{Q} \cdot \beta_j$ , where  $\beta_j(b_i) = \delta_{ij}$ . Denote by  $\mathbf{b}(z)$  and  $\beta(w^{-1})$  the generating series of the basis elements of  $V$  and of  $V^*$  respectively, i.e.:

$$\mathbf{b}(z) := \sum_{i \geq 0} b_i z^i \quad \text{and} \quad \beta(w^{-1}) := \sum_{j \geq 0} \beta_j w^{-j}. \tag{3}$$

The *exterior algebra* of  $V$  is  $\bigwedge V := \bigoplus_{j \geq 0} \bigwedge^j V$ , the direct sum of the *exterior powers*  $\bigwedge^j V$ , where  $\bigwedge^0 V = \mathbb{Q}$  and  $\bigwedge^1 V = V$ . The algebra structure is given by the  $\mathbb{Q}$ -linear extension of the juxtaposition. To each  $\lambda \in \mathcal{P}_r$  we associate

$$[\mathbf{b}]_\lambda^r := b_{r-1+\lambda_1} \wedge b_{r-2+\lambda_2} \wedge \dots \wedge b_{\lambda_r} \in \bigwedge^r V, \tag{4}$$

so that  $([\mathbf{b}]_\lambda^r)_{\lambda \in \mathcal{P}_r}$  is a  $\mathbb{Q}$ -basis of  $\bigwedge^r V$ . The pairing

$$(\beta_{i_1} \wedge \dots \wedge \beta_{i_r})(v_1 \wedge \dots \wedge v_r) = \begin{vmatrix} \beta_{i_1}(v_1) & \dots & \beta_{i_1}(v_r) \\ \vdots & \ddots & \vdots \\ \beta_{i_r}(v_1) & \dots & \beta_{i_r}(v_r) \end{vmatrix} \tag{5}$$

establishes a natural identification between  $\bigwedge^r V^*$  and  $(\bigwedge^r V)^*$ . If one denotes by  $[\beta]_\mu^r$  the basis element

$$\beta_{r-1+\mu_1} \wedge \dots \wedge \beta_{\mu_r} \tag{6}$$

of  $\bigwedge^r V^*$ , an easy check shows that  $[\beta]_\mu^r([\mathbf{b}]_\lambda^r) = \delta_{\mu,\lambda}$ . The pairing (5) enables to attach to any  $\beta \in V^*$  a map  $\beta_\lrcorner : \bigwedge V \rightarrow \bigwedge V$  of degree  $-1$  (with respect to the graduation of the exterior algebra) via the equality

$$\eta(\beta_\lrcorner u) = (\beta \wedge \eta)(u), \quad \forall (u, \eta) \in \bigwedge^r V \times \bigwedge^{r-1} V^*. \tag{7}$$

### 2.3 The ring $B_r$

Let  $r \geq 1$ . The main character of this paper is the polynomial ring  $B_r := \mathbb{Q}[e_1, \dots, e_r]$  in the  $r \geq 1$  indeterminates  $(e_1, \dots, e_r)$  (by convention  $B_0 = \mathbb{Q}$ ).

Given the generic polynomial  $E_r(z) := 1 - e_1 z + \dots + (-1)^r e_r z^r \in B_r[z]$ , one considers the sequence  $H_r := (h_j)_{j \in \mathbb{Z}}$  defined by the equality:

$$\sum_{n \in \mathbb{Z}} h_n z^n := \frac{1}{E_r(z)}. \tag{8}$$

holding in  $B_r[z]$ . In particular  $h_j = 0$  if  $j < 0$  and  $h_0 = 1$ . Moreover for  $j \geq 0$ , the term  $h_j$  is an explicit polynomial in  $(e_1, \dots, e_r)$ , homogeneous of degree  $j$ , once one gives weight  $i$  to  $e_i$ . The *Schur determinants*

$$\Delta_\lambda(H_r) := \det(h_{\lambda_j-j+i})_{1 \leq i,j \leq r} = \begin{vmatrix} h_{\lambda_1} & h_{\lambda_2-1} & \dots & h_{\lambda_r-r+1} \\ h_{\lambda_1+1} & h_{\lambda_2} & \dots & h_{\lambda_r-r+2} \\ \vdots & \vdots & \ddots & \vdots \\ h_{\lambda_1+r-1} & h_{\lambda_2+r-2} & \dots & h_{\lambda_r} \end{vmatrix}, \tag{9}$$

form a  $\mathbb{Q}$ -basis of  $B_r$  parametrized by the partitions of length at most  $r$ :

$$B_r := \bigoplus_{\lambda \in \mathcal{P}_r} \mathbb{Q} \cdot \Delta_\lambda(H_r). \tag{10}$$

It follows that  $B_r$  is naturally isomorphic to  $\bigwedge^r V$  via the  $\mathbb{Q}$ -linear extension of the sets map

$$\Delta_\lambda(H_r) \mapsto [\mathbf{b}]_\lambda^r. \tag{11}$$

### 2.4 Schur polynomials

Especially in the last section we shall be concerned with Schur polynomials in a set of indeterminates. We recall them here. For each partition of length at most  $k$  and any set of  $k$  formal variables  $\mathbf{x}_k := (x_1, \dots, x_k)$ , one defines

$$\Delta_\lambda(\mathbf{x}_k) = \det \left( x_j^{\lambda_{k-i+1}+i-1} \right) = \begin{vmatrix} x_1^{\lambda_k} & x_2^{\lambda_k} & \dots & x_k^{\lambda_k} \\ x_1^{1+\lambda_{k-1}} & x_2^{1+\lambda_{k-1}} & \dots & x_k^{1+\lambda_{k-1}} \\ \vdots & \vdots & \ddots & \vdots \\ x_1^{k-1+\lambda_1} & x_2^{k-1+\lambda_1} & \dots & x_k^{k-1+\lambda_1} \end{vmatrix}.$$

This is an skew symmetric polynomials in  $(x_1, \dots, x_k)$  and therefore divisible by the *Vandermonde determinant*

$$\Delta_0(\mathbf{x}_k) := \begin{vmatrix} 1 & 1 & \dots & 1 \\ x_1 & x_2 & \dots & x_k \\ \vdots & \vdots & \ddots & \vdots \\ x_1^{k-1} & x_2^{k-1} & \dots & x_k^{k-1} \end{vmatrix} = \prod_{i < j} (x_j - x_i).$$

The Schur polynomial associated to  $\mathbf{x}_k$  and the partition  $\lambda$  is defined by the equality

$$\Delta_\lambda(\mathbf{x}_k) = s_\lambda(\mathbf{x}_k) \cdot \Delta_0(\mathbf{x}_k),$$

often said to be the *Jacobi-Trudy* formula.

### 2.5 Hasse–Schmidt derivations on exterior algebras

Let now  $\bigwedge V[z]$  denote the formal power series in the indeterminate  $z$  with coefficients in the exterior algebra  $\bigwedge V$  of  $V$ . If  $\mathcal{S}$  is any set of indeterminates over  $\mathbb{Q}$ , denote by  $\mathbb{Q}[\mathcal{S}]$  the corresponding algebra of formal power series. The following is an extended reformulation of the main definition of the reference [13] (see also [18]). By a *Hasse–Schmidt* derivation on  $\bigwedge V$  we mean any  $\mathbb{Q}[\mathcal{S}]$ -linear extension of a  $\mathbb{Q}$ -linear map  $\mathcal{D}(z) : \bigwedge V \rightarrow \bigwedge V[[z]]$  such that

$$\mathcal{D}(z)(u \wedge v) = \mathcal{D}(z)u \wedge \mathcal{D}(z)v, \quad \forall u, v \in \bigwedge V, \tag{12}$$



which, by abuse of notation, will be denoted by the same symbol

$$\mathcal{D}(z) : \mathbb{Q}[[\mathcal{S}]] \otimes_{\mathbb{Q}} \bigwedge V \rightarrow \mathbb{Q}[[\mathcal{S}]] \otimes_{\mathbb{Q}} \bigwedge V[[z]],$$

(instead of the more precise, but lengthier,  $1_{\mathbb{Q}[[\mathcal{S}]}} \otimes_{\mathbb{Q}} \mathcal{D}(z)$ ).

If  $\mathcal{D}_i \in \text{End}_{\mathbb{Q}}(\bigwedge V)$  are such that  $\sum_{i \geq 0} \mathcal{D}_i z^i := \mathcal{D}(z)$ , then (12) is equivalent to the system of relations

$$\mathcal{D}_i(u \wedge v) = \sum_{j=0}^i \mathcal{D}_j u \wedge \mathcal{D}_{i-j} v, \quad (i \geq 0).$$

The notation

$$\mathcal{D}(z)[\mathbf{b}]_{\lambda}^r = [\mathcal{D}(z)\mathbf{b}]_{\lambda}^r \tag{13}$$

will be used as a shorthand for the equality

$$\mathcal{D}(z)[\mathbf{b}]_{\lambda}^r = \mathcal{D}(z)(b_{r-1+\lambda_1} \wedge \cdots \wedge b_{\lambda_r}) = \mathcal{D}(z)b_{r-1+\lambda_1} \wedge \cdots \wedge \mathcal{D}(z)b_{\lambda_r}$$

meaning that  $\mathcal{D}(z)$  is a HS-derivation. By [23, Proposition 3.3], if  $\mathcal{D}_0$  is invertible in  $\text{End}_{\mathbb{Q}}(\bigwedge V)$ , then  $\mathcal{D}(z)$  is invertible as a  $\text{End}_{\mathbb{Q}}(\bigwedge V)$ -valued formal power series and its inverse,  $\overline{\mathcal{D}}(z)$ , is an HS-derivation as well.

**2.6 Proposition** Let  $\mathcal{D}_0$  be invertible in  $\text{End}_{\mathbb{Q}}(\bigwedge V)$ . Then the *integration by parts formulas* follow for all  $u, v \in \bigwedge V$ :

$$\mathcal{D}(z)(\overline{\mathcal{D}}(z)u \wedge v) = u \wedge \mathcal{D}(z)v, \tag{14}$$

$$\overline{\mathcal{D}}(z)(\mathcal{D}(z)u \wedge v) = u \wedge \overline{\mathcal{D}}(z)v. \tag{15}$$

□

Formulas (14) and (15) are implicitly assuming the  $\mathbb{Q}[[z]]$ -linearity of  $\mathcal{D}(z)$  we alluded to in Definition 2.5. The extension of the linearity of HS-derivations over polynomial algebras will be assumed in the following without any further mention.

**2.6 Transposition**

The transpose  $\mathcal{D}(z)^T : \bigwedge V^* \rightarrow \bigwedge V^*[[z]]$  of the HS derivation  $\mathcal{D}(z)$  is defined via its action on homogeneous elements. For  $\eta \in \bigwedge^r V^*$ , one stipulates that  $\mathcal{D}(z)^T \eta(u) = \eta(\mathcal{D}(z)u)$ , for all  $u \in \bigwedge^r V$ . By [19, Proposition 2.8]  $\mathcal{D}(z)^T$  is a HS-derivation of  $\bigwedge V^*$ .

### 3 Recap on Schubert derivations

#### 3.1 Exterior algebras representations of endomorphisms

Consider the natural representation  $\delta : \text{End}(V) \rightarrow \text{End}(\bigwedge V)$  making any  $\phi \in \text{End}(V)$  into an (even) derivation  $\delta(\phi)$  of  $\bigwedge V$ . In other words,  $\delta(\phi)$  is the unique  $\mathbb{Q}$ -vector space endomorphism of  $\bigwedge V$  such that

$$\delta(\phi)(v \wedge w) = \delta(\phi)v \wedge w + v \wedge \delta(\phi)w$$

for all  $v, w \in \bigwedge V$ , together with the initial condition  $\delta(\phi)u = \phi(u)$  holding for all  $u \in V = \bigwedge^1 V$ . An easy check shows that

$$\mathcal{D}^\phi(z) = \exp\left(\sum_{i \geq 1} \frac{1}{i} \delta(\phi^i)z^i\right),$$

is the unique HS derivation on  $\bigwedge V$  such that  $\mathcal{D}^\phi(z)|_V = \sum_{i \geq 0} \phi^i z^i$ .

Let now  $\sigma_1 : V \rightarrow V$  be such that  $\sigma_1 b_j = b_{j+1}$  and  $\sigma_{-1} : V \rightarrow V$  such that  $\sigma_{-1} b_j = b_{j-1}$ , where by convention we put  $b_k = 0$  if  $k < 0$ .

**3.2 Definition** The *Schubert derivations* on  $\bigwedge V$  are the HS-derivations  $\sigma_+(z) : \bigwedge V \rightarrow \bigwedge V[[z]]$  and  $\sigma_-(z) : \bigwedge V \rightarrow \bigwedge V[z^{-1}]$  defined by

$$\sigma_+(z) = \sum_{i \geq 0} \sigma_i z^i := \exp\left(\sum_{i \geq 1} \frac{1}{i} \delta(\sigma_1^i)z^i\right), \tag{16}$$

$$\sigma_-(z) = \sum_{i \geq 0} (-1)^i \sigma_{-i} z^{-i} := \exp\left(\sum_{i \geq 1} \frac{1}{i} \delta(\sigma_{-1}^i)z^{-i}\right), \tag{17}$$

and their inverses in  $\text{End}_{\mathbb{Q}}(\bigwedge V)[[z]]$  and  $\text{End}_{\mathbb{Q}}(\bigwedge V)[z^{-1}]$  respectively:

$$\bar{\sigma}_+(z) = \sum_{i \geq 0} (-1)^i \bar{\sigma}_i z^i := \exp\left(-\sum_{i \geq 1} \frac{1}{i} \delta(\sigma_1^i)z^i\right), \tag{18}$$

$$\bar{\sigma}_-(z) = \sum_{i \geq 0} (-1)^i \bar{\sigma}_{-i} z^{-i} := \exp\left(-\sum_{i \geq 1} \frac{1}{i} \delta(\sigma_{-1}^i)z^{-i}\right). \tag{19}$$

In particular:

$$\bar{\sigma}_\pm(z)u = u - \sigma_{\pm 1}u \cdot z^{\pm 1}, \quad \forall u \in V = \bigwedge^1 V.$$

**3.3 Remark** It is easily seen that  $\sigma_\pm(z)$  and  $\bar{\sigma}_\pm(z)$  are the unique HS-derivations on  $\bigwedge V$  such that

$$\sigma_+(z)b_j = \sum_{i \geq 0} b_{j+i}z^i, \quad \bar{\sigma}_+(z)b_j = b_j - b_{j+1}z, \tag{20}$$

and

$$\sigma_-(z)b_j = \sum_{i \geq 0} \frac{b_{j-i}}{z^i}, \quad \bar{\sigma}_-(z)b_j = b_j - \frac{b_{j-1}}{z}, \tag{21}$$

putting  $b_i = 0$  for  $i < 0$ .

### 3.2 $B_r$ -module structure of $\bigwedge^r V$

We exploit the Schubert derivation  $\bar{\sigma}_+(z)$  or, equivalently, its inverse  $\sigma_+(z)$ , to endow  $\bigwedge^r V$  with a  $B_r$ -module structure, by declaring that  $e_i u = \bar{\sigma}_i u$  or, equivalently,  $h_i u = \sigma_i u$ , for all  $u \in \bigwedge^r V$ . In particular:

$$\bar{\sigma}_+(z)u = E_r(z) \cdot u \quad \text{and} \quad \sigma_+(z)u := \frac{1}{E_r(z)}u, \quad \forall u \in \bigwedge^r V.$$

The fact that such a product structure is compatible with the natural vector space isomorphism  $B_r \rightarrow \bigwedge^r V$  given by (11) is a consequence of

**3.5 Proposition** Giambelli’s formula for the Schubert derivation  $\sigma_+(z)$  holds:

$$[\mathbf{b}]_\lambda^r = \Delta_\lambda(\sigma_+(z))[\mathbf{b}]_0^r := \det(\sigma_{\lambda_j - j + i})_{1 \leq i, j \leq r} [\mathbf{b}]_0^r. \tag{22}$$

Hence  $\bigwedge^r V$  is a free  $B_r$ -module of rank 1 generated by  $[\mathbf{b}]_0^r$ .

**Proof** Formula (22) may be inferred as a particular case of the general determinantal formula for the exterior power of a polynomial ring due to Laksov and Thorup as in [27, Main Theorem 0.1]. It follows that  $[\mathbf{b}]_\lambda^r = \Delta_\lambda(\sigma_+(z))[\mathbf{b}]_0^r = \Delta_\lambda(H_r)[\mathbf{b}]_0^r$ , proving the second part of the claim.  $\square$

By virtue of 3.5, the set map  $\Delta_\lambda(H_r) \mapsto \Delta_\lambda(H_r)[\mathbf{b}]_0^r$  extends to a well defined vector space isomorphisms  $B_r \rightarrow \bigwedge^r V$ , as it maps the basis  $(\Delta_\lambda(H_r))_{\lambda \in \mathcal{P}_r}$  of  $B_r$  to the basis  $([\mathbf{b}]_\lambda^r)_{\lambda \in \mathcal{P}_r}$  of  $\bigwedge^r V$ . The fact that  $\bigwedge^r V$  is a free  $B_r$ -module of rank 1 generated by  $[\mathbf{b}]_0^r$ , as prescribed by equality (22), shows that the Schubert derivations  $\sigma_-(z), \bar{\sigma}_-(z)$  induce maps  $B_r \rightarrow B_r[z^{-1}]$  which, abusing notation, will be denoted in the same way. Their action on a basis element  $\Delta_\lambda(H_r)$  of  $B_r$  is defined through its action on  $\bigwedge^r V$ :

$$(\bar{\sigma}_-(z)\Delta_\lambda(H_r))[\mathbf{b}]_0^r = \bar{\sigma}_-(z)[\mathbf{b}]_\lambda^r, \tag{23}$$

$$(\sigma_-(z)\Delta_\lambda(H_r))[\mathbf{b}]_0^r = \sigma_-(z)[\mathbf{b}]_\lambda^r. \tag{24}$$

Denote by  $\bar{\sigma}_-(z)H_r$  (respectively  $\sigma_-(z)H_r$ ) the sequence  $(\bar{\sigma}_-(z)h_j)_{j \in \mathbb{Z}}$  (respectively  $(\sigma_-(z)h_j)_{j \in \mathbb{Z}}$ ). By using [19, Theorem 5.7], and exploiting the Laksov & Thorup determinantal formula as in [27, Main Theorem 0.1], one obtains the following statement, which gives a practical way to evaluate the image of  $\Delta_\lambda(H_r)$  through the maps  $\bar{\sigma}_-(z)$  and  $\sigma_-(z)$  defined by (23) and (24).

**3.6 Proposition** ([19, Proposition 5.3]). For all  $r \geq 0$  and all  $\lambda \in \mathcal{P}_r$ ,

$$\sigma_{-}(z)h_j = \sum_{i \geq 0} \frac{h_{j-i}}{z^i} \quad \text{and} \quad \bar{\sigma}_{-}(z)h_j = h_j - \frac{h_{j-1}}{z}. \tag{25}$$

Moreover:

$$\sigma_{-}(z)\Delta_{\lambda}(H_r) = \Delta_{\lambda}(\sigma_{-}(z)H_r) \quad \text{and} \quad \bar{\sigma}_{-}(z)\Delta_{\lambda}(H_r) = \Delta_{\lambda}(\bar{\sigma}_{-}(z)H_r). \tag{26}$$

□

**3.7 Remark** It is important to notice that (26) only holds if  $\ell(\lambda) \leq r$ . For example

$$\Delta_{(1,1)}(\bar{\sigma}_{-}(z)H_1) = \begin{vmatrix} h_1 - \frac{1}{z} & 1 \\ h_2 - \frac{h_1}{z} & h_1 - \frac{1}{z} \end{vmatrix} = -\frac{h_1}{z} + \frac{1}{z^2} \neq 0 = \bar{\sigma}_{-}(z)\Delta_{(1,1)}(H_1).$$

## 4 Commutation rules for Schubert derivations

### 4.1 Product of Schubert derivations

For  $k \geq 1$ , let  $\mathbf{z}_k$  denote the ordered  $k$ -tuple  $(z_1, \dots, z_k)$  of formal variables. By  $\mathbf{z}_k^{-1}$  we shall mean the  $k$ -tuple of the formal inverses  $(z_1^{-1}, \dots, z_k^{-1})$ . Define maps  $\sigma_{\pm}(\mathbf{z}_k), \bar{\sigma}_{\pm}(\mathbf{z}_k) : \bigwedge V \rightarrow \bigwedge V[[\mathbf{z}_k, \mathbf{z}_k^{-1}]$  respectively by

$$\sigma_{\pm}(\mathbf{z}_k) := \sigma_{\pm}(z_1) \cdots \sigma_{\pm}(z_k) \quad \text{and} \quad \bar{\sigma}_{\pm}(\mathbf{z}_k) := \bar{\sigma}_{\pm}(z_1) \cdots \bar{\sigma}_{\pm}(z_k). \tag{27}$$

The maps occurring in formulas (27) are multivariate HS derivations on  $\bigwedge V$ , in the sense that, for instance,  $\sigma_{+}(\mathbf{z}_k)(u \wedge v) = \sigma_{+}(\mathbf{z}_k)u \wedge \sigma_{+}(\mathbf{z}_k)v$ , as it is easy to check and adopting the linear extension of the Schubert derivation to polynomial coefficients according to Definition 2.5. The same holds verbatim for  $\sigma_{-}(\mathbf{z}_k)$  and  $\bar{\sigma}_{\pm}(\mathbf{z}_k)$ . It is an important point that the multivariate HS derivations in (27) are symmetric in the formal variables  $z_i$  and  $w_j$ . This is a consequence of the first of the commutation rules obeyed by the product of Schubert derivations listed in this section and to be used in the sequel.

**4.2 Proposition** Let  $z, w$  be arbitrary formal variables. The equalities

$$\bar{\sigma}_{\pm}(z)\bar{\sigma}_{\pm}(w) = \bar{\sigma}_{\pm}(w)\bar{\sigma}_{\pm}(z), \tag{28}$$

$$\sigma_{\pm}(z)\sigma_{\pm}(w) = \sigma_{\pm}(w)\sigma_{\pm}(z), \tag{29}$$

hold in  $\text{End}_{\mathbb{Q}}(\bigwedge V)[[z^{\pm 1}, w^{\pm 1}]]$ .

**Proof** Formulas (28) and (29) are obvious consequences of the fact that if  $i, j \geq 0$ , then  $\sigma_{\pm i}$  and  $\sigma_{\pm j}$  are pairwise commuting. It is sufficient, then, to show that they commute when restricted to  $V$ , because if they do, then

$$\sigma_{\pm}(z)\sigma_{\pm}(w)[\mathbf{b}]_{\lambda}^r = \left[ \sigma_{\pm}(z)\sigma_{\pm}(w)\mathbf{b} \right]_{\lambda}^r = \left[ \sigma_{\pm}(w)\sigma_{\pm}(z)\mathbf{b} \right]_{\lambda}^r,$$

having used notation as in (13). But  $\sigma_{\pm i}\sigma_{\pm j}u = \sigma_{\pm 1}^{i+j}u = \sigma_{\pm j}\sigma_{\pm i}u$  for all  $u \in V$ , and then the claim follows.  $\square$

In order to give a compact expression of the  $gl(\bigwedge^k V)$ -module structure of  $B_r$ , we shall need to introduce a generalisation of the classical vertex operators arising in the context of the so-called boson–fermion correspondence, like in e.g. [25]. We look at it as a generalisation of the isomorphism  $B_r \rightarrow \bigwedge^r V$ , recalled in Sect. 2.1, reaffirmed and refined in Proposition 3.5.

**4.3 Definition** By *vertex operators* on  $\bigwedge V$  we mean the  $\mathbb{Q}[\mathbf{z}_k, \mathbf{z}_k^{-1}]$ -linear maps  $\Gamma(\mathbf{z}_k), \Gamma^*(\mathbf{z}_k) : \bigwedge V \rightarrow (\bigwedge V)[\mathbf{z}_k, \mathbf{z}_k^{-1}]$  of degree 1 and  $-1$ , with respect to the exterior algebra graduation, given by:

$$\Gamma(\mathbf{z}_k)[\mathbf{b}]_{\lambda}^r = \sigma_+(\mathbf{z}_k)\bar{\sigma}_-(\mathbf{z}_k)[\mathbf{b}]_{\lambda}^{r+k}, \tag{30}$$

$$\Gamma^*(\mathbf{z}_k)[\mathbf{b}]_{\lambda}^r = \left( \bar{\sigma}_+(\mathbf{z}_k)\Delta_{\lambda}(\sigma_-(\mathbf{z}_k)H_{r-k}) \right) [\mathbf{b}]_0^{r-k}. \tag{31}$$

Proposition 4.2 guarantees that the vertex operators  $\Gamma(\mathbf{z}_k)$  and  $\Gamma^*(\mathbf{z}_k)$  are symmetric in the formal variables  $(z_1, \dots, z_k)$ . They will be studied in a more detailed way in Sects. 6 and 7, exploiting further commutation relations, for which we need the preliminary work exposed below. As a matter of fact, we notice that the commutativity of the product of Schubert derivations is granted only if they are of the same kind (both subscripts “+” or both subscripts “-”). In general, for  $i, j > 0$ ,  $\sigma_i$  and  $\sigma_{-j}$  do not commute, because  $\sigma_{-j}$  is locally nilpotent. The simplest example is:  $\sigma_{-1}\sigma_1 b_0 = b_0 \neq 0 = \sigma_1\sigma_{-1}b_0$ . The general pattern is that commutativity only holds up to the multiplication by a rational function.

**4.5 Proposition**

i) If  $\lambda \in \mathcal{P}_r \setminus \mathcal{P}_{r-1}$  (i.e.  $\ell(\lambda) = r$ ), then  $\bar{\sigma}_-(w)$  commutes with both  $\sigma_+(z)$  and  $\bar{\sigma}_+(z)$ , i.e.

$$\bar{\sigma}_-(w)\sigma_+(z) = \sigma_+(z)\bar{\sigma}_-(w), \tag{32}$$

and

$$\bar{\sigma}_-(w)\bar{\sigma}_+(z) = \bar{\sigma}_+(z)\bar{\sigma}_-(w). \tag{33}$$

ii) if  $\lambda \in \mathcal{P}_{r-1}$  (i.e.  $[\mathbf{b}]_{\lambda}^r = [\mathbf{b}]_{\lambda+(1^{r-1})}^{r-1} \wedge b_0$ ):

$$\bar{\sigma}_-(w)\sigma_+(z)[\mathbf{b}]_{\lambda}^r = \left( 1 - \frac{z}{w} \right) \sigma_+(z)\bar{\sigma}_-(w)[\mathbf{b}]_{\lambda}^r. \tag{34}$$

**Proof** As a matter of i), we observe that  $\bar{\sigma}_-(w)\sigma_+(z)b_{\lambda} = \sigma_+(z)\bar{\sigma}_-(z)b_{\lambda}$  if  $\lambda > 0$ . Indeed

$$\begin{aligned} \bar{\sigma}_-(w)\sigma_+(z)b_\lambda &= \bar{\sigma}_-(w)\left(\sum_{j \geq 0} b_{\lambda+j}z^j\right) && \text{(Definition of } \sigma_+(z)b_\lambda) \\ &= \sum_{j \geq 0} \left(b_{\lambda+j} - \frac{b_{\lambda+j-1}}{w}\right)z^j && \text{(Definition of } \bar{\sigma}_-(w)) \\ &= \sigma_+(z)\bar{\sigma}_-(w)b_\lambda. \end{aligned}$$

Similarly

$$\bar{\sigma}_-(w)\bar{\sigma}_+(z)b_\lambda = \bar{\sigma}_+(z)\bar{\sigma}_-(w)b_\lambda,$$

as a direct straightforward computation shows. Therefore, under the hypothesis  $\ell(\lambda) = r$ :

$$\begin{aligned} \bar{\sigma}_-(w)\sigma_+(z)[\mathbf{b}]_\lambda^r &= \bar{\sigma}_-(w)\sigma_+(z)b_{r-1+\lambda_1} \wedge \cdots \wedge \bar{\sigma}_-(w)\sigma_+(z)b_{\lambda_r} \\ &= \sigma_+(z)\bar{\sigma}_-(w)b_{r-1+\lambda_1} \wedge \cdots \wedge \sigma_+(z)\bar{\sigma}_-(w)b_{\lambda_r} = \sigma_+(z)\bar{\sigma}_-(w)[\mathbf{b}]_\lambda^r, \end{aligned}$$

and the same can be argued for the commutation of  $\bar{\sigma}_+(z)$  and  $\bar{\sigma}_-(w)$ .

To prove equality (34), one observes that

$$\begin{aligned} \bar{\sigma}_-(w)\sigma_+(z)b_0 &= \bar{\sigma}_-(w)\sum_{j \geq 0} b_jz^j && \text{(Definition of } \sigma_+(z)b_0) \\ &= b_0 + \sum_{j \geq 1} \left(b_j - \frac{b_{j-1}}{w}\right)z^j && \text{(Definition of } \sigma_-(w)b_j) \\ &= b_0 + \sum_{j \geq 1} b_jz^j - \frac{z}{w}\sum_{j \geq 0} b_jz^j && (35) \\ &= \left(1 - \frac{z}{w}\right)\sigma_+(z)b_0 \\ &= \left(1 - \frac{z}{w}\right)\sigma_+(z)\bar{\sigma}_-(w)b_0, \end{aligned}$$

because, in general,  $\bar{\sigma}_-(w)$  acts on  $[\mathbf{b}]_0^r$  as the identity. So, if  $\ell(\lambda) < r$  (i.e.  $\lambda_r = 0$ ) one obtains:

$$\begin{aligned} \bar{\sigma}_-(w)\sigma_+(z)[\mathbf{b}]_\lambda^r &= \bar{\sigma}_-(w)\sigma_+(z)\left([\mathbf{b}]_{\lambda+(1^{r-1})}^{r-1} \wedge b_0\right) && \text{(Definition of } [\mathbf{b}]_\lambda^r) \\ &= \bar{\sigma}_-(w)\sigma_+(z)[\mathbf{b}]_{\lambda+(1^{r-1})}^{r-1} \wedge \bar{\sigma}_-(w)\sigma_+(z)b_0 && (\bar{\sigma}_-(w)\sigma_+(z) \text{ is a HS derivation)} \\ &= \sigma_+(z)\bar{\sigma}_-(w)[\mathbf{b}]_{\lambda+(1^{r-1})}^{r-1} \wedge \left(1 - \frac{z}{w}\right)\sigma_+(z)\bar{\sigma}_-(w)b_0 && \text{Commutation (35)} \\ &= \left(1 - \frac{z}{w}\right)\sigma_+(z)\bar{\sigma}_-(w)[\mathbf{b}]_\lambda^r. \end{aligned}$$

□

### 5 Commutation rules for contractions

The goal of this section is to prove the following

**5.1 Theorem** For all  $u \in \bigwedge^r V$ , the following commutation rule holds:

$$\beta_{0\lrcorner}\sigma_{-}(w)\bar{\sigma}_{+}(z)u = \left(1 - \frac{z}{w}\right)\bar{\sigma}_{+}(z)(\beta_{0\lrcorner}\sigma_{-}(w)u). \tag{36}$$

To prove Theorem 5.1 some preparation is needed.

**5.1 Diagrams for contractions**

Let us begin to introduce a piece of useful notation. If  $\beta \in V^*$ , we represent the contraction

$$\beta_{\lrcorner}(u_1 \wedge \dots \wedge u_r)$$

as defined by equality (7), via the diagram:

$$\left| \begin{array}{cccc} \beta_{\lrcorner}u_1 & \beta_{\lrcorner}u_2 & \dots & \beta_{\lrcorner}u_r \\ u_1 & u_2 & \dots & u_r \end{array} \right| = \left| \begin{array}{cccc} \beta(u_1) & \beta(u_2) & \dots & \beta(u_r) \\ u_1 & u_2 & \dots & u_r \end{array} \right|, \tag{37}$$

to be read as follows. The scalar  $(-1)^{j+1}\beta_{\lrcorner}u_j := (-1)^j\beta(u_j)$  is the coefficient of the element of  $\bigwedge^{r-1} V$  obtained by removing the wedge factor  $u_j$  from  $u_1 \wedge u_2 \wedge \dots \wedge u_r$ . For example

$$\left| \begin{array}{ccc} \beta_{\lrcorner}u_1 & \beta_{\lrcorner}u_2 & \beta_{\lrcorner}u_3 \\ u_1 & u_2 & u_3 \end{array} \right| = \beta(u_1) \cdot u_2 \wedge u_3 - \beta(u_2)u_1 \wedge u_3 + \beta(u_3)u_1 \wedge u_2.$$

This is exactly the expanded expression of the contraction  $\beta_{\lrcorner}(u_1 \wedge u_2 \wedge u_3)$ . Recall now the generating function  $\beta(w^{-1}) := \sum_{j \geq 0} \beta_j w^{-j}$  introduced in formula (3).

**5.3 Lemma** We have:

$$\beta(w^{-1})_{\lrcorner}[\mathbf{b}]_{\lambda}^r = \left| \begin{array}{cccc} w^{-r+1-\lambda_1} & w^{-r+2-\lambda_2} & \dots & w^{-\lambda_r} \\ b_{r-1+\lambda_1} & b_{r-2+\lambda_2} & \dots & b_{\lambda_r} \end{array} \right|. \tag{38}$$

**Proof** Since  $\beta_j \lrcorner b_i = \beta_j(b_i) = \delta_{ij}$ , it clearly follows that

$$\beta(w^{-1})_{\lrcorner} b_j = \sum_{i \geq 0} \beta_i(b_j)w^{-i} = w^{-j}.$$

Putting  $\beta(w^{-1})$  instead of  $\beta$  in equality (37), gives (38).

**5.4 Proposition** For all  $u \in \bigwedge^r V$ :

$$\beta(w^{-1})_{\lrcorner}\bar{\sigma}_{+}(z)u = \left(1 - \frac{z}{w}\right)\bar{\sigma}_{+}(z)(\beta(w^{-1})_{\lrcorner}u). \tag{39}$$

**Proof** Since each  $u \in \bigwedge^r V$  is a finite linear combination of  $[\mathbf{b}]_{\lambda}^r$ , it is no harm to assume  $u = [\mathbf{b}]_{\lambda}^r$ . Notice that

$$\begin{aligned} \beta(w^{-1})_{\lrcorner}\bar{\sigma}_{+}(z)b_j &= \beta(w^{-1})_{\lrcorner}(b_j - b_{j+1}z) && \text{(Definition of } \bar{\sigma}_{+}(z)\text{)} \\ &= \beta(w^{-1})_{\lrcorner}b_j - \beta(w^{-1})_{\lrcorner}b_{j+1}z && \text{(Action of } \beta(w^{-1})_{\lrcorner}\text{)} \\ &= \frac{1}{w^j} - \frac{z}{w^{j+1}} = \frac{1}{w^j} \left(1 - \frac{z}{w}\right). \end{aligned} \tag{40}$$

By expressing the contraction via diagram (37), one has:

$$\beta(w^{-1})_{\lrcorner} \bar{\sigma}_+(z) [\mathbf{b}]_{\lambda}^r = \begin{vmatrix} \beta(w^{-1})_{\lrcorner} \bar{\sigma}_+(z) b_{r-1+\lambda_1} & \beta(w^{-1})_{\lrcorner} \bar{\sigma}_+(z) b_{r-2+\lambda_2} & \cdots & \beta(w^{-1})_{\lrcorner} \bar{\sigma}_+(z) b_{\lambda_r} \\ \bar{\sigma}_+(z) b_{r-1+\lambda_1} & \bar{\sigma}_+(z) b_{r-2+\lambda_2} & \cdots & \bar{\sigma}_+(z) b_{\lambda_r} \end{vmatrix},$$

which by (40) is equal to:

$$\begin{aligned} &= \begin{vmatrix} \left(1 - \frac{z}{w}\right) \frac{1}{w^{r-1+\lambda_1}} & \left(1 - \frac{z}{w}\right) \frac{1}{w^{r-2+\lambda_2}} & \cdots & \left(1 - \frac{z}{w}\right) \frac{1}{w^{\lambda_r}} \\ \bar{\sigma}_+(z) b_{r-1+\lambda_1} & \bar{\sigma}_+(z) b_{r-2+\lambda_2} & \cdots & \bar{\sigma}_+(z) b_{\lambda_r} \end{vmatrix} \\ &= \left(1 - \frac{z}{w}\right) \begin{vmatrix} \frac{1}{w^{r-1+\lambda_1}} & \frac{1}{w^{r-2+\lambda_2}} & \cdots & \frac{1}{w^{\lambda_r}} \\ \bar{\sigma}_+(z) b_{r-1+\lambda_1} & \bar{\sigma}_+(z) b_{r-2+\lambda_2} & \cdots & \bar{\sigma}_+(z) b_{\lambda_r} \end{vmatrix}. \end{aligned} \tag{41}$$

Since the determinant occurring in (41) is a linear combination of  $[\bar{\sigma}_+(z)\mathbf{b}]_{\lambda^{(j)}}^{r-1} = \bar{\sigma}_+(z)[\mathbf{b}]_{\lambda^{(j)}}^{r-1}$  (because  $\bar{\sigma}_+(z)$  is a HS derivation), where we denoted by  $\lambda^{(j)}$  the partition of length at most  $r - 1$  obtained by omitting the  $j$ -th part, it follows that the action of  $\bar{\sigma}_+(z)$  can be factorized from the bottom row of (41), giving

$$\left(1 - \frac{z}{w}\right) \bar{\sigma}_+(z) \begin{vmatrix} \frac{1}{w^{r-1+\lambda_1}} & \frac{1}{w^{r-2+\lambda_2}} & \cdots & \frac{1}{w^{\lambda_r}} \\ b_{r-1+\lambda_1} & b_{r-2+\lambda_2} & \cdots & b_{\lambda_r} \end{vmatrix} = \left(1 - \frac{z}{w}\right) \bar{\sigma}_+(z) (\beta(w^{-1})_{\lrcorner} [\mathbf{b}]_{\lambda}^r),$$

which ends the proof of the Proposition. □

**5.5 Lemma** For all  $u \in \bigwedge^r V$ ,

$$\beta(w^{-1})_{\lrcorner} u = \bar{\sigma}_-(w)(\beta_0 \lrcorner \sigma_-(w)u). \tag{42}$$

**Proof** It is basically contained in [19, Proposition 4.3] but, because some mild difformity in the notation, we prefer to repeat it here. Recall the definition of transpose of a HS derivation on  $\bigwedge V$ . We observe that  $\beta(w^{-1}) = \sigma_-(w)^T \beta_0$ . Then, for all  $\eta \in \bigwedge^{r-1} V^*$ ,

$$\begin{aligned} \eta(\beta(w^{-1})_{\lrcorner} u) &= (\beta(w^{-1}) \wedge \eta)(u) && \text{(Definition 2.2 of contraction)} \\ &= (\sigma_-(w)^T \beta_0 \wedge \eta)(u) && \text{(by the above observation)} \\ &= \sigma_-(w)^T (\beta_0 \wedge \bar{\sigma}_-(w)^T \eta)(u) && \text{(Integration by parts)} \\ &= \beta_0 \wedge \bar{\sigma}_-(w)^T \eta(\sigma_-(w)u) && \text{(Definition of transpose of } \bar{\sigma}_-(w)^T \text{)} \\ &= \bar{\sigma}_-(w)^T \eta(\beta_0 \lrcorner \sigma_-(w)u) && \text{(Again 2.2)} \\ &= \eta(\bar{\sigma}_-(w)(\beta_0 \lrcorner \sigma_-(w)u)) && \text{(Definition of transpose)}. \end{aligned}$$



The last equality proves (42), due to the arbitrary choice of  $\eta \in \bigwedge^{r-1} V^* \cong (\bigwedge^{r-1} V)^*$ . □

**5.6 Lemma** The operator  $\bar{\sigma}_+(z)$  commutes with contracting against  $\beta_0$ , i.e. for all  $\lambda \in \mathcal{P}_r$

$$\beta_{0\lrcorner}\bar{\sigma}_+(z)[\mathbf{b}]_\lambda^r = \bar{\sigma}_+(z)(\beta_{0\lrcorner}[\mathbf{b}]_\lambda^r). \tag{43}$$

**Proof** There are two cases. If  $\ell(\lambda) = r$ , both members of (43) vanish. If  $\ell(\lambda) \leq r - 1$ , then  $\beta_{0\lrcorner}\bar{\sigma}_+(z)[\mathbf{b}]_\lambda^r = (-1)^{r-1}\bar{\sigma}_+(z)[\mathbf{b}]_{\lambda+(1^{(r-1)})}^{r-1} = \bar{\sigma}_+(z)(\beta_{0\lrcorner}[\mathbf{b}]_\lambda^r)$ . □

We are now in position to provide the

### 5.2 Proof of Theorem 5.1

We have:

$$\begin{aligned} \beta_{0\lrcorner}\sigma_-(w)\bar{\sigma}_+(z)u &= \sigma_-(w)\bar{\sigma}_-(w)(\beta_{0\lrcorner}\sigma_-(w)\bar{\sigma}_+(z)u)(\sigma_-(w)\bar{\sigma}_-(w) = 1) \\ &= \sigma_-(w)(\beta(w^{-1})\lrcorner\bar{\sigma}_+(z)u) && \text{(Lemma 5.5)} \\ &= \left(1 - \frac{z}{w}\right)\sigma_-(w)\bar{\sigma}_+(z)(\beta(w^{-1})\lrcorner u) && \text{(Proposition 5.4)} \\ &= \left(1 - \frac{z}{w}\right)\sigma_-(w)\bar{\sigma}_+(z)\bar{\sigma}_-(w)(\beta_{0\lrcorner}\sigma_-(w)u) && \text{(again Lemma 5.5)} \\ &= \left(1 - \frac{z}{w}\right)\sigma_-(w)\bar{\sigma}_-(w)\bar{\sigma}_+(z)(\beta_{0\lrcorner}\sigma_-(w)u) && \text{(formula (33) of Proposition 4.4)} \\ &= \left(1 - \frac{z}{w}\right)\bar{\sigma}_+(z)(\beta_{0\lrcorner}\sigma_-(w)u) && (\sigma_-(w)\bar{\sigma}_-(w) = 1). \end{aligned}$$

and Theorem 5.1 is thence proven. □

## 6 The vertex operator $\Gamma(\mathbf{z}_k)$

The main purpose of this section is to interpret the vertex operator  $\Gamma(\mathbf{z}_k)$ , introduced in Definition 4.3, formula (30), in terms of wedging operation on the exterior algebra. It will generalise [19, Proposition 4.2]. This will be achieved in Theorem 6.5 below and will be used in our Main Theorem 8.5.

**6.1 Lemma** For all  $j \geq 0$  and all  $k \geq 1$  one has

$$\bar{\sigma}_+(\mathbf{z}_k)b_j = b_j + \sum_{i=1}^k (-1)^i e_i(\mathbf{z}_k)b_{j+i} \tag{44}$$

and

$$\sigma_+(\mathbf{z}_k)b_j = b_j + \sum_{i \geq 1}^k h_i(\mathbf{z}_k)b_{j+i}, \tag{45}$$

where  $e_i(\mathbf{z}_k)$  and  $h_i(\mathbf{z}_k)$  are, respectively, the elementary and complete symmetric polynomial of degree  $i$  in the indeterminates  $\mathbf{z}_k := (z_1, \dots, z_k)$ .

**Proof** Formula (44) is the content of [21, Lemma 5.7] to which we refer to. Formula (45) is a consequence of (44), keeping into account that  $\sigma_+(\mathbf{z}_k)$  and  $\bar{\sigma}_+(\mathbf{z}_k)$  are mutually inverse in  $\text{End}_{\mathbb{Q}}(\wedge V)[\mathbf{z}_k]$ .  $\square$

**6.2 Lemma** One has:

$$\bar{\sigma}_-(\mathbf{z}_k)b_{j+k} = b_{j+k} + \sum_{i=1}^k (-1)^i e_i(\mathbf{z}_k^{-1})b_{j+k-i}, \tag{46}$$

where  $e_i(\mathbf{z}_k^{-1}) = e_i(z_1^{-1}, \dots, z_k^{-1})$  is the elementary symmetric polynomial of degree  $i$  in  $(z_1^{-1}, \dots, z_k^{-1})$ .

**Proof** The proof works the same as in [21, Lemma 6.7]. The formula is true for  $k = 1$ , because

$$\bar{\sigma}_-(z_1)b_{j+1} = b_{j+1} - \frac{b_j}{z_1}.$$

By induction, suppose that (46) holds for  $k - 1 \geq 0$ . Then it holds for  $k$ . Indeed

$$\begin{aligned} & \bar{\sigma}_-(z_1)\bar{\sigma}_-(z_2) \dots \bar{\sigma}_-(z_k)b_{j+k} \\ &= \bar{\sigma}_-(z_1) \left[ b_{j+k} - e_1\left(\frac{1}{z_2}, \dots, \frac{1}{z_k}\right)b_{j+k-1} + \dots + (-1)^{k-1}e_{k-1}\left(\frac{1}{z_2}, \dots, \frac{1}{z_k}\right)b_{j+1} \right] \\ &= b_{j+k} - \frac{b_{j+k-1}}{z_1} - e_1\left(\frac{1}{z_2}, \dots, \frac{1}{z_k}\right)\left(b_{j+k-1} - \frac{b_{j+k-2}}{z_1}\right) + \dots \\ & \quad + (-1)^{k-1}e_{k-1}\left(\frac{1}{z_2}, \dots, \frac{1}{z_k}\right)\left(b_{j+1} - \frac{b_j}{z_1}\right) \tag{definition of  $\bar{\sigma}_-(z_1)$ } \\ &= b_{j+k} + \sum_{i=1}^k (-1)^i e_i\left(\frac{1}{z_1}, \dots, \frac{1}{z_k}\right)b_{j+k-i}, \end{aligned}$$

as desired.  $\square$

**6.3 Lemma** The following equality holds for all  $1 \leq i \leq k$ :

$$\frac{e_i(\mathbf{z}_k)}{z_1 \dots z_k} = e_{k-i}\left(\frac{1}{z_1}, \dots, \frac{1}{z_k}\right). \tag{47}$$

**Proof** Recall the following definition of the elementary symmetric polynomials in  $k$  indeterminates through generating functions:

$$\sum_{i=0}^k e_i(\mathbf{z}_k)t^i = \prod_{i=1}^k (1 + z_i t). \tag{48}$$

By dividing both sides of (48) by  $e_k(\mathbf{z}_k) = z_1 \dots z_k$  we get

$$\sum_{i=0}^k \frac{e_i(\mathbf{z}_k)}{z_1 \dots z_k} t^i = \prod_{i=1}^k \left(\frac{1}{z_i} + t\right). \tag{49}$$

The claim then follows by comparing the coefficient of  $t^i$  in either side of (49). □

**6.4 Lemma** For all  $k \geq 1, r \geq 0$  and  $\lambda \in \mathcal{P}_r$ :

$$[\mathbf{b}]_0^k \wedge \bar{\sigma}_+(\mathbf{z}_k)[\mathbf{b}]_\lambda^r = e_k(\mathbf{z}_k)^r \bar{\sigma}_-(\mathbf{z}_k)[\mathbf{b}]_\lambda^{r+k}. \tag{50}$$

**Proof** Equality (50) holds for  $r = 1$ :

$$\begin{aligned} [\mathbf{b}]_0^k \wedge \bar{\sigma}_+(\mathbf{z}_k)[\mathbf{b}]_\lambda^1 &= [\mathbf{b}]_0^k \wedge \bar{\sigma}_+(\mathbf{z}_k)b_\lambda \\ &= [\mathbf{b}]_0^k \wedge (b_\lambda - e_1(\mathbf{z}_k)b_{\lambda+1} + \dots + (-1)^k e_k(\mathbf{z}_k)b_{\lambda+k}) && \text{(By 44)} \\ &= [\mathbf{b}]_0^k \wedge (-1)^k e_k(\mathbf{z}_k) \left[ b_{\lambda+k} - \frac{e_{k-1}(\mathbf{z}_k)}{e_k(\mathbf{z}_k)} b_{\lambda+k-1} + \dots + (-1)^k \frac{1}{e_k(\mathbf{z}_k)} b_\lambda \right] && \text{(By Factorization)} \\ &= e_k(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{z}_k) b_{\lambda+k} \wedge [\mathbf{b}]_0^k && \text{(By lemma (6.2))} \\ &= e_k(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{z}_k) (b_{\lambda+k} \wedge \bar{\sigma}_-(\mathbf{z}_k)[\mathbf{b}]_0^k) && \text{(Integration by parts)} \\ &= e_k(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{z}_k) (b_{\lambda+k} \wedge b_{k-1} \wedge \dots \wedge b_0) && (\bar{\sigma}_-(\mathbf{z}_k)[\mathbf{b}]_0^k = [\mathbf{b}]_0^k) \\ &= e_k(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{z}_k)[\mathbf{b}]_\lambda^{1+k}. \end{aligned}$$

Therefore the property is true for  $r = 1$ . Assume now (50) holds true for  $r - 1 \geq 0$ . Then

$$\begin{aligned} [\mathbf{b}]_0^k \wedge \bar{\sigma}_+(\mathbf{z}_k)[\mathbf{b}]_\lambda^r &= [\mathbf{b}]_0^k \wedge \bar{\sigma}_+(\mathbf{z}_k)b_{r-1+\lambda_1} \wedge \dots \wedge \bar{\sigma}_+(\mathbf{z}_k)b_{\lambda_r} \\ &= [\mathbf{b}]_0^k \wedge (-1)^k e_k(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{z}_k) b_{r-1+k+\lambda_1} \wedge \dots \wedge (-1)^k e_k(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{z}_k) b_{\lambda_r+k} \\ &= e_k(\mathbf{z}_k)^r [\bar{\sigma}_-(\mathbf{z}_k) b_{r+k-1+\lambda_1} \wedge \dots \wedge \bar{\sigma}_-(\mathbf{z}_k) b_{k+\lambda_r} \wedge \bar{\sigma}_-(\mathbf{z}_k)[\mathbf{b}]_0^k] \\ &= e_k(\mathbf{z}_k)^r \bar{\sigma}_-(\mathbf{z}_k) (b_{r+k-1+\lambda_1} \wedge b_{k+\lambda_r} \wedge b_k \wedge \dots \wedge b_0) \\ &= e_k(\mathbf{z}_k)^r \bar{\sigma}_-(\mathbf{z}_k)[\mathbf{b}]_\lambda^{r+k}, \end{aligned}$$

as claimed. □

**6.5 Theorem** For all  $u \in \bigwedge^r V[\mathbf{w}_k, \mathbf{w}_k^{-1}]$  we have:

$$\sigma_+(z_1, \dots, z_k)[\mathbf{b}]_0^k \wedge u = \prod_{j=1}^r z_j^r \cdot \Gamma(\mathbf{z}_k)u, \tag{51}$$

the equality holding in  $\bigwedge^{r+k} V[[\mathbf{z}_k, \mathbf{w}_k][\mathbf{w}_k^{-1}]]$ .

**Proof** Recall that we consider all the Schubert derivations extended by linearity over rings of formal power series with rational coefficients. See definition 2.5. Then our arbitrary  $u$  is intended as a linear combination of  $[\mathbf{b}]_\lambda^r$  with coefficients being polynomials. Then we can assume with no harm that  $u = [\mathbf{b}]_\lambda^r$ , a basis element of  $\bigwedge^r V$ . Keeping the same notation as in Sect. 2, we first apply integration by parts. Then

$$\begin{aligned} \sigma_+(\mathbf{z}_k)[\mathbf{b}]_0^k \wedge [\mathbf{b}]_\lambda^r &= \sigma_+(\mathbf{z}_k) ([\mathbf{b}]_0^k \wedge \bar{\sigma}_+(\mathbf{z}_k)[\mathbf{b}]_\lambda^r) && \text{(By integration by parts (14))} \\ &= \sigma_+(\mathbf{z}_k) e_k(\mathbf{z}_k)^r \bar{\sigma}_-(\mathbf{z}_k)[\mathbf{b}]_\lambda^{r+k} && \text{(By Lemma 6.4)} \\ &= e_k(\mathbf{z}_k)^r \sigma_+(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{z}_k)[\mathbf{b}]_\lambda^{r+k} \\ &= \prod_{j=1}^k z_j^r \Gamma(\mathbf{z}_k)[\mathbf{b}]_\lambda^r && \text{(Definition of } \Gamma(\mathbf{z}_k)) \end{aligned}$$

as desired. □

If  $k = 1$ , and  $z = z_1$ , one obtains

$$\sigma_+(z)b_0 \wedge [\mathbf{b}]_\lambda^r = z^r \Gamma(z)[\mathbf{b}]_\lambda^r = z^r \sigma_+(z)\bar{\sigma}_-(z)[\mathbf{b}]_\lambda^{r+1},$$

which is precisely [19, Proposition 5.4] or [20, Proposition 3.2]. Their shape there looks more involved because of a different notation.

### 7 The vertex operator $\Gamma^*(\mathbf{w}_k)$

In the same vein of Sect. 6, the present one will be devoted to interpret the action of the vertex operator  $\Gamma^*(\mathbf{w}_k)$  on  $\bigwedge V$  in terms of contraction operators. The output will be Theorem 7.3, stated at the end of the section, another building block of the main Theorem 8.5. We begin with some preparation.

**7.1 Lemma** The following equality holds for all  $r \geq 1$  and all  $\lambda \in \mathcal{P}_r$ :

$$\bar{\sigma}_{-r+1}(\beta_{0\lrcorner}\sigma_-(w)[\mathbf{b}]_\lambda^r) = \Delta_\lambda(\sigma_-(w)H_{r-1})[\mathbf{b}]_0^{r-1}. \tag{52}$$

**Proof** This is [19, Lemma 5.8]. □

$$\beta(w^{-1})\lrcorner[\mathbf{b}]_\lambda^r = w^{-r+1}\bar{\sigma}_+(w)\Delta_\lambda(\sigma_-(w)H_{r-1})[\mathbf{b}]_0^{r-1}. \tag{53}$$

7.2 Lemma

**Proof** Invoking Lemma 5.5,

$$\beta(w^{-1})\lrcorner[\mathbf{b}]_\lambda^r = \bar{\sigma}_-(w)(\beta_{0\lrcorner}\sigma_-(w)[\mathbf{b}]_\lambda^r). \tag{54}$$

Since  $\beta_{0\lrcorner}\sigma_-(w)[\mathbf{b}]_\lambda^r$  is a linear combination of  $[\mathbf{b}]_\mu^{r-1}$  with  $\ell(\mu) = r - 1$  (i.e. no  $b_0$  occurs in the monomial), then by [19, Proposition 4.3]

$$\bar{\sigma}_-(w)(\beta_{0\lrcorner}\sigma_-(w)[\mathbf{b}]_\lambda^r) = w^{-r+1}\bar{\sigma}_+(w)\bar{\sigma}_{-r+1}(\beta_{0\lrcorner}\sigma_-(w)[\mathbf{b}]_\lambda^r).$$

Using Lemma 7.1 one obtains (53). □

**7.3 Theorem** The following equality holds:

$$(\beta(w_k^{-1}) \wedge \beta(w_{k-1}^{-1}) \wedge \dots \wedge \beta(w_1^{-1}))\lrcorner[\mathbf{b}]_\lambda^r = \frac{\Delta_0(\mathbf{w}_k)}{(w_1 \dots w_k)^{r-1}} \Gamma^*(\mathbf{w}_k)[\mathbf{b}]_\lambda^r. \tag{55}$$

where  $\Delta_0(\mathbf{w}_k)$  denotes the Vandermonde determinant  $\prod_{1 \leq i < j \leq k} (w_j - w_i)$ .

**Proof** For  $k = 1$  the property

$$\beta(w_1^{-1})\lrcorner[\mathbf{b}]_\lambda^r = w_1^{-r+1}\bar{\sigma}_+(w_1)\Delta_\lambda(\sigma_-(w_1)H_{r-1})[\mathbf{b}]_0^{r-1}$$

is just Lemma 7.2. Arguing by induction, let us assume the claim holding true for  $0 \leq k - 1 \leq r - 1$  and let us show it holds for  $k$ . We have

$$\beta(w_k^{-1}) \wedge \beta(w_{k-1}^{-1}) \wedge \dots \wedge \beta(w_1^{-1})\lrcorner[\mathbf{b}]_\lambda^r = \beta(w_k^{-1})\lrcorner(\beta(w_{k-1}^{-1}) \wedge \dots \wedge \beta(w_1^{-1})\lrcorner[\mathbf{b}]_\lambda^r).$$

Using the inductive hypothesis:

$$= \beta(w_k^{-1})_{\lrcorner} \frac{\Delta_0(\mathbf{w}_{k-1})}{(w_1 \cdots w_{k-1})^{r-1}} \bar{\sigma}_+(\mathbf{w}_{k-1}) \Delta_\lambda(\sigma_-(\mathbf{w}_{k-1})H_{r-k+1})[\mathbf{b}]_0^{r-k+1}.$$

By applying Lemma 7.2, one gets:

$$= w_k^{r-k+1} \bar{\sigma}_+(w_k) [\beta_{0\lrcorner} \sigma_-(w_k) \bar{\sigma}_+(\mathbf{w}_{k-1}) \Delta_\lambda(\sigma_-(\mathbf{w}_{k-1})H_{r-k+1})[\mathbf{b}]_0^{r-k+1}] \cdot \frac{\Delta_0(\mathbf{w}_{k-1})}{(w_1 \cdots w_{k-1})^{r-1}}$$

Now we use the commutation rules prescribed by Theorem 5.1 and Lemma 5.6:

$$\begin{aligned} &= w_k^{k-r} \bar{\sigma}_+(w_k) \left[ \prod_{j=0}^{k-1} \left( 1 - \frac{w_j}{w_k} \right) \cdot (\beta_{0\lrcorner} \bar{\sigma}_+(\mathbf{w}_{k-1}) \sigma_-(\mathbf{w}_k) \Delta_\lambda(\sigma_-(\mathbf{w}_{k-1})H_{r-k+1})[\mathbf{b}]_0^{r-k+1}) \right] \\ &\quad \cdot \frac{\Delta_0(\mathbf{w}_{k-1})}{(w_1 \cdots w_{k-1})^{r-1}} \\ &= \frac{w_k^{k-r} \Delta_0(\mathbf{w}_{k-1})}{(w_1 \cdots w_{k-1})^{r-1}} \prod_{j=0}^{k-1} \left( 1 - \frac{w_j}{w_k} \right) \cdot \\ &\quad \cdot \bar{\sigma}_+(w_k) \bar{\sigma}_+(\mathbf{w}_{k-1}) [(\beta_{0\lrcorner} \sigma_-(\mathbf{w}_k) \Delta_\lambda(\sigma_-(\mathbf{w}_{k-1})H_{r-k+1})[\mathbf{b}]_0^{r-k+1})] \\ &= \frac{\Delta_0(\mathbf{w}_{k-1}) \prod_{j=1}^{k-1} (w_k - w_j)}{w_k^{r-k} \cdot w_k^{k-1} (w_1 \cdots w_{k-1})^{r-1}} \bar{\sigma}_+(w_k) \Delta_\lambda(\sigma_-(w_k) \sigma_-(\mathbf{w}_{k-1})H_{r-k})[\mathbf{b}]_0^{r-k} \\ &= \frac{\Delta_0(\mathbf{w}_k)}{(w_1 \cdots w_{k-1} \cdot w_k)^{r-1}} \cdot \bar{\sigma}_+(w_k) \Delta_\lambda(\sigma_-(\mathbf{w}_k)H_{r-k})[\mathbf{b}]_0^{r-k} \\ &= \frac{\Delta_0(\mathbf{w}_k)}{(w_1 \cdots w_{k-1} \cdot w_k)^{r-1}} \Gamma^*(w_k)[\mathbf{b}]_\lambda^r, \end{aligned}$$

as claimed. □

### 8 The main theorem and its declinations

In this section we shall be concerned with the several declinations of the main theorem describing the  $B_r$  representation of  $gl(\bigwedge^k V)$ .

#### 8.1 Preparation

Let  $gl(\bigwedge^k V) := \bigwedge^k V \otimes \bigwedge^k V^*$ . It is a Lie sub-algebra of the endomorphisms of  $\bigwedge^k V$ , with respect to the natural commutator. With the same notation as in 1.1, a basis of  $\bigwedge^k V \otimes \bigwedge^k V^*$  is  $([\mathbf{b}]_\mu^k \otimes [\beta]_\nu^k)_{\mu, \nu \in \mathcal{P}_k}$  i.e.

$$\bigwedge^k V \otimes \bigwedge^k V^* = \bigoplus_{\mu, \nu \in \mathcal{P}_k} \mathbb{Q} \cdot [\mathbf{b}]_\mu^k \otimes [\beta]_\nu^k,$$

where  $[\beta]_v^k([\mathbf{b}]_\mu^k) = \delta_{\mu,v}$ .

Then the  $gl(\wedge^k V)$ -module structure of  $B_r$  is defined through the following equality holding in  $\wedge^r V$ :

$$([\mathbf{b}]_\mu^k \otimes [\beta]_v^k \star \Delta_\lambda(H_r))[\mathbf{b}]_0^r = [\mathbf{b}]_\mu^k \wedge ([\beta]_v^k [\mathbf{b}]_\lambda^r). \tag{56}$$

This action is very easy to describe in the case  $k = r$ , but it becomes trickier when  $r - k > 0$ . To describe it we shall consider the generating function

$$\mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1}) := \sum_{\mu, \nu} [\mathbf{b}]_\mu^k \otimes [\beta]_\nu^k \cdot s_\mu(\mathbf{z}_k) s_\nu(\mathbf{w}_k^{-1}) : B_r \rightarrow B_r[[\mathbf{z}_k, \mathbf{w}_k^{-1}]][\mathbf{z}_k^{-1}, \mathbf{w}_k^{-1}].$$

Our main result will consist in the explicit description of  $\mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1})\Delta_\lambda(H_r)$  in case  $k \leq r$  (because otherwise one would obtain the trivial null action), where  $\mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1})\Delta_\lambda(H_r)$  is such that

$$\left(\mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1})\Delta_\lambda(H_r)\right)[\mathbf{b}]_0^r = \sum_{\mu \in \mathcal{P}_k} s_\mu(\mathbf{z}_k)[\mathbf{b}]_\mu^k \wedge (s_\nu(\mathbf{w}_k^{-1})[\beta]_{\nu \perp}^k [\mathbf{b}]_\lambda^r), \tag{57}$$

and where  $s_\mu(\mathbf{z}_k)$  and  $s_\nu(\mathbf{w}_k^{-1})$  denote the Schur symmetric polynomials labeled by the partitions related to the variables  $\mathbf{z}_k$  and  $\mathbf{w}_k^{-1}$  respectively.

**8.2 Lemma** The generating function of the basis  $([\mathbf{b}]_\mu^k)_{\mu \in \mathcal{P}_k}$  of  $\wedge^k V$  is:

$$\sum_{\mu \in \mathcal{P}_k} s_\mu(\mathbf{z}_k)[\mathbf{b}]_\mu^k = \sigma_+(\mathbf{z}_k)[\mathbf{b}]_0^k.$$

**Proof** By exploiting the definition of the  $B_k$ -module structure of  $\wedge^k V$ , we have

$$\begin{aligned} \sum_{\mu \in \mathcal{P}_k} s_\mu(\mathbf{z}_k)[\mathbf{b}]_\mu^k &= \sum_{\mu \in \mathcal{P}_k} s_\mu(\mathbf{z}_k)\Delta_\mu(H_k)[\mathbf{b}]_0^k && \text{(using the } B_k \text{ - module structure of } \wedge^k V) \\ &= \left(\sum_{\mu \in \mathcal{P}_k} s_\mu(\mathbf{z}_k)\Delta_\mu(H_k)\right)[\mathbf{b}]_0^k \\ &= \prod_{j=1}^k (1 + h_1 z_j + h_2 z_j^2 + \dots)[\mathbf{b}]_0^k \\ &= \frac{1}{E_k(z_1)} \cdot \frac{1}{E_k(z_2)} \cdots \frac{1}{E_k(z_k)} [\mathbf{b}]_0^k && \text{(By Cauchy formula as in [12, Proposition 2, (iii)])} \\ &= \sigma_+(\mathbf{z}_k)[\mathbf{b}]_0^k. \end{aligned}$$

□

**8.3 Lemma** The generating function of the basis elements  $\wedge^k V^*$  is:

$$\sum_{\nu \in \mathcal{P}_k} s_\nu(\mathbf{w}_k^{-1})[\beta]_\nu^k = \frac{\prod_{j=1}^k w_j^{k-1}}{\Delta_0(\mathbf{w}_k)} \cdot \beta(w_k^{-1}) \wedge \dots \wedge \beta(w_1^{-1}). \tag{58}$$

**Proof** The one we propose consists in expanding the wedge product of the generating series of the basis  $(\beta_j)_{j \geq 0}$  of  $V^*$ :

$$\begin{aligned}
 & \beta(w_k^{-1}) \wedge \cdots \wedge \beta(w_1^{-1}) \tag{59} \\
 &= \sum_{\nu \in \mathcal{P}_k} [\beta]_{\nu}^k \sum_{\tau \in S_k} \text{sgn}(\tau) w_k^{-k+\tau(1)-\nu_{\tau(1)}} w_{k-1}^{-k+\tau(2)-\nu_{\tau(2)}} \cdots w_1^{-k+\tau(k)-\nu_{\tau(k)}} = \sum_{\nu \in \mathcal{P}_k} [\beta]_{\nu}^k \Delta_{\nu}(w_k^{-1}) \\
 &= \sum_{\nu \in \mathcal{P}_k} [\beta]_{\nu}^k \cdot s_{\nu}(w_k^{-1}) \Delta_0(w_k^{-1}) = \sum_{\nu \in \mathcal{P}_k} s_{\nu}(w_k^{-1}) [\beta]_{\nu}^k \cdot \frac{\Delta_0(w_k)}{w_k^{k-1} \cdots w_1^{k-1}}, \tag{60}
 \end{aligned}$$

whence the claim, obtained by multiplying both (59) and (60) by  $(\prod_{j=1}^k w_j^k) / \Delta_0(w_k)$ .  $\square$

**8.4 Lemma** For all  $u \in \bigwedge^r V[[w_k, w_k^{-1}]]$

$$\sum_{\mu} ([b]_{\mu}^k \wedge u) s_{\mu}(z_k) = \bar{\sigma}_+(z_k) [b]_0^k \wedge u = \prod_{j=1}^k z_j^r \Gamma(z_k) u. \tag{61}$$

**Proof** We specified that the equality holds in  $\bigwedge^r V[[w_k, w_k^{-1}]]$  to emphasize the supposed  $\mathbb{Q}[[w_k, w_k^{-1}]]$  linearity of the Schubert derivation. This said, it is not restrictive to assume that  $u$  is a basis element  $[b]_{\lambda}^r$  of  $\bigwedge^r V$ . The basic remark is that

$$\begin{aligned}
 \frac{1}{E_k(z_1)} \cdot \frac{1}{E_k(z_2)} \cdots \frac{1}{E_k(z_k)} &= \prod_{j=1}^k (1 + h_1 z_j + h_2 z_j^2 + h_3 z_j^3 + \cdots) \\
 &= \sum_{\mu \in \mathcal{P}_k} s_{\mu}(z_k) \Delta_{\lambda}(H_k),
 \end{aligned}$$

having used one of the declination of the celebrated Cauchy formula, as in [11, Proposition 2, (iii)], already employed in the proof of Lemma 8.2. Therefore:

$$\begin{aligned}
 \sum_{\mu} [b]_{\mu}^k s_{\mu}(z_k) &= \sum_{\mu} (\Delta_{\mu}(H_k) s_{\mu}(z_k)) [b]_0^k \\
 &= \frac{1}{E_k(z_1) \cdots E_k(z_k)} [b]_0^k = \sigma_+(z_k) [b]_0^k,
 \end{aligned}$$

where in the last equality we have repeatedly used the module structure of  $\bigwedge^r V$  over  $B_r$ . Then:

$$\sum_{\mu} [b]_{\mu}^k s_{\mu}(z_k) \wedge u = \sigma_+(z_k) [b]_0^k \wedge u,$$

and the result now follows from Theorem 6.5.  $\square$

We can finally express the action of the generating function  $\mathcal{E}_{\mu, \nu}(z_k, w_k^{-1})$  on a basis element of  $B_r$ .

**8.5 Theorem (First Version).** For all  $\lambda \in \mathcal{P}_r$ :

$$(\mathcal{E}(z_k, w_k^{-1}) \Delta_{\lambda}(H_r)) [b]_0^r = \prod_{j=1}^k \left( \frac{z_j}{w_j} \right)^{r-k} \Gamma(z_k) \Gamma^*(w_k) [b]_{\lambda}^r. \tag{62}$$

**Proof** We have

$$\begin{aligned} \mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1})[\mathbf{b}]_\lambda^r &= \sum_{\mu \in \mathcal{P}_k} s_\mu(\mathbf{z}_k) [\mathbf{b}]_\mu^k \wedge \sum_{\nu \in \mathcal{P}_k} s_\nu(\mathbf{w}_k) \left( [\beta]_{\nu \downarrow}^k [\mathbf{b}]_\lambda^r \right) && \text{(definition of } \mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1}) \text{)} \\ &= \sigma_+(\mathbf{z}_k) [\mathbf{b}]_0^k \wedge \frac{\prod_{j=1}^k w_j^{k-1}}{\Delta_0(\mathbf{w}_k)} \cdot \beta(\mathbf{w}_k^{-1}) \wedge \cdots \wedge \beta(\mathbf{w}_1^{-1}) \downarrow [\mathbf{b}]_\lambda^r && \text{(Lemmas 8.2 and 8.3)} \\ &= \sigma_+(\mathbf{z}_k) [\mathbf{b}]_0^k \wedge \frac{1}{\prod_{j=1}^k w_j^{r-k}} \Gamma^*(\mathbf{w}_k) [\mathbf{b}]_\lambda^r. && \text{(Theorem 7.3)} \end{aligned}$$

Now we use the fact that  $\Gamma^*(\mathbf{w}_k) [\mathbf{b}]_\lambda^r$  is a  $\mathbb{Q}[\mathbf{w}_k, \mathbf{w}_k^{-1}]$ -linear combination of  $[\mathbf{b}]_\mu^{r-k}$  and then, by invoking Theorem 6.5, applied to last equality above, we obtain:

$$\prod_{j=1}^k \left( \frac{z_j}{w_j} \right)^{r-k} \Gamma(\mathbf{z}_k) \Gamma^*(\mathbf{w}_k) [\mathbf{b}]_\lambda^r,$$

as announced. □

**8.6 Corollary** If  $r - k \geq \ell(\lambda)$  then

$$\Gamma(\mathbf{z}_k) \Gamma^*(\mathbf{w}_k) [\mathbf{b}]_\lambda^r = \sigma_+(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{z}_k) \bar{\sigma}_+(\mathbf{w}_k) \sigma_-(\mathbf{w}_k) [\mathbf{b}]_\lambda^r.$$

**Proof** First of all notice that for every  $\lambda \in \mathcal{P}_r$ , it turns out that

$$\bar{\sigma}_+(\mathbf{w}_k) \sigma_-(\mathbf{w}_k) [\mathbf{b}]_\lambda^{r-k} = \sum_{\mu \in \mathcal{P}_r} a_\mu(\mathbf{w}_k, \mathbf{w}_k^{-1}) [\mathbf{b}]_\mu^{r-k},$$

where  $a_\mu(\mathbf{w}_k, \mathbf{w}_k^{-1}) \in \mathbb{Q}[\mathbf{w}_k, \mathbf{w}_k^{-1}]$ . Then we have:

$$\begin{aligned} \Gamma(\mathbf{z}_k) \Gamma^*(\mathbf{w}_k) [\mathbf{b}]_\lambda^r &= \Gamma(\mathbf{z}_k) \bar{\sigma}_+(\mathbf{w}_k) \sigma_-(\mathbf{w}_k) [\mathbf{b}]_\lambda^{r-k} \\ &= \sum_{\mu \in \mathcal{P}_r} a_\mu(\mathbf{w}_k, \mathbf{w}_k^{-1}) \Gamma(\mathbf{z}_k) [\mathbf{b}]_\mu^{r-k} \\ &= \sum_{\mu \in \mathcal{P}_r} a_\mu(\mathbf{w}_k, \mathbf{w}_k^{-1}) [\mathbf{b}]_\mu^r \\ &= \sigma_+(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{w}_k) \sum_{\mu \in \mathcal{P}_r} a_\mu(\mathbf{w}_k, \mathbf{w}_k^{-1}) [\mathbf{b}]_\mu^r \\ &= \sigma_+(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{w}_k) \sigma_+(\mathbf{z}_k) \bar{\sigma}_-(\mathbf{w}_k) [\mathbf{b}]_\lambda^r. \end{aligned}$$

**8.7 Remark** If  $\ell(\lambda) > r - k$ , Corollary 8.6 fails. We have however the following uniform way to compute  $\Gamma^*(\mathbf{w}_k) [\mathbf{b}]_\lambda^r$ .

**8.8 Proposition** For all  $\lambda \in \mathcal{P}_r$  and  $k, r \geq 0$ :

$$\Gamma^*(\mathbf{w}_k) [\mathbf{b}]_\lambda^r = \left| \begin{array}{ccc} \frac{1}{w_1^{r-1+\lambda_1}} & \frac{1}{w_1^{r-2+\lambda_2}} & \cdots & \frac{1}{w_1^{\lambda_r}} \\ \vdots & \vdots & \ddots & \\ \frac{1}{w_k^{r-1+\lambda_1}} & \frac{1}{w_k^{r-2+\lambda_2}} & \cdots & \frac{1}{w_k^{\lambda_r}} \\ b_{r-1+\lambda_1} & b_{r-2+\lambda_2} & \cdots & b_{\lambda_r} \end{array} \right| \in \bigwedge^{r-k} V,$$



using the same notation as in 4.1.

**Proof** By Theorem 7.3

$$\Gamma^*(\mathbf{w}_k)[\mathbf{b}]_\lambda^r = \frac{(z_1 \cdots z_k)^{r-1}}{\Delta_0(\mathbf{w}_k)} (\boldsymbol{\beta}(w_k^{-1}) \wedge \boldsymbol{\beta}(w_{k-1}^{-1}) \wedge \cdots \wedge \boldsymbol{\beta}(w_1^{-1})) \lrcorner [\mathbf{b}]_\lambda^r$$

$$= \frac{(z_1 \cdots z_k)^{r-1}}{\Delta_0(\mathbf{w}_k)} \begin{vmatrix} \frac{1}{w_1^{r-1+\lambda_1}} & \frac{1}{w_1^{r-2+\lambda_2}} & \cdots & \frac{1}{w_1^{\lambda_r}} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{1}{w_k^{r-1+\lambda_1}} & \frac{1}{w_k^{r-2+\lambda_2}} & \cdots & \frac{1}{w_k^{\lambda_r}} \\ b_{r-1+\lambda_1} & b_{r-2+\lambda_2} & \cdots & b_{\lambda_r} \end{vmatrix}. \tag{63}$$

### 8.2 Another formulation

Let us agree that

$$w_k^{r-k} \cdots w_2^{r-2} w_1^{r-1} \Delta_\lambda(\mathbf{w}_k^{-1}, H_{r-k})[\mathbf{b}]_0^r = \boldsymbol{\beta}(w_k^{-1}) \wedge \cdots \wedge \boldsymbol{\beta}(w_1^{-1}) \lrcorner [\mathbf{b}]_\lambda^r,$$

defines  $\Delta_\lambda(\mathbf{w}_k, H_{r-k}) \in B_r[w^{-1}]$ . The expansion of (63) as a linear combinations of basis elements of  $\bigwedge^{r-k} V$ , Giambelli’s formula (22) and the expansion rule of a determinant, easily imply that

$$\Delta_\lambda(\mathbf{w}_k^{-1}, H_{r-k}) = \begin{vmatrix} \frac{1}{w_1^{\lambda_1}} & \frac{1}{w_1^{\lambda_2-1}} & \cdots & \frac{1}{w_1^{\lambda_r-r+1}} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{1}{w_k^{\lambda_1+k-1}} & \frac{1}{w_k^{\lambda_2+k-2}} & \cdots & \frac{1}{w_k^{\lambda_r+k-r}} \\ h_{\lambda_1+k} & h_{\lambda_2+k+1} & \cdots & h_{\lambda_r+k+r-1} \\ \vdots & \vdots & \ddots & \vdots \\ h_{\lambda_1+r-1} & h_{\lambda_2+r-2} & \cdots & h_{\lambda_r} \end{vmatrix}. \tag{64}$$

This enables to state a second version of 8.5, which works well for practical purposes and generalises [20, Main Theorem 4.3].

**8.10 Theorem (second version).** *The following equality holds:*

$$\begin{aligned}
 \mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1})\Delta_\lambda(H_r) &= \prod_{j=1}^k \left(\frac{z_j}{w_j}\right)^{r-k} \frac{1}{E_r(z_j)} \Delta_\lambda(\mathbf{w}_k^{-1}, \bar{\sigma}_-(z)H_r) \\
 &= \prod_{j=1}^k \left(\frac{z_j}{w_j}\right)^{r-k} \frac{1}{E_r(z_j)} \begin{vmatrix} \frac{1}{w_1^{\lambda_1}} & \frac{1}{w_1^{\lambda_2-1}} & \cdots & \frac{1}{w_1^{\lambda_r-r+1}} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{1}{w_k^{\lambda_1+k-1}} & \frac{1}{w_k^{\lambda_2+k-2}} & \cdots & \frac{1}{w_k^{\lambda_r+k-r}} \\ \bar{\sigma}_-(\mathbf{z}_k)h_{\lambda_1+k} & \bar{\sigma}_-(\mathbf{z}_k)h_{\lambda_2+k+1} & \cdots & \bar{\sigma}_-(\mathbf{z}_k)h_{\lambda_r+k+r-1} \\ \vdots & \vdots & \ddots & \vdots \\ \bar{\sigma}_-(\mathbf{z}_k)h_{\lambda_1+r-1} & \bar{\sigma}_-(\mathbf{z}_k)h_{\lambda_2+r-2} & \cdots & \bar{\sigma}_-(\mathbf{z}_k)h_{\lambda_r} \end{vmatrix}, \tag{65}
 \end{aligned}$$

where

$$\bar{\sigma}_-(\mathbf{z}_k)h_j = h_j - e_1(\mathbf{z}_k^{-1})h_{j-1} + \cdots + (-1)^k e_k(\mathbf{z}_k^{-1})h_{j-k}. \tag{66}$$

**Proof** By Theorem 8.5 we have:

$$\begin{aligned}
 \mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1})\Delta_\lambda(H_r) &= \prod_{j=1}^k \left(\frac{z_j}{w_j}\right)^{r-k} \Gamma(\mathbf{z}_k)\Gamma^*(\mathbf{w}_k)\Delta_\lambda(H_r) && \text{(Theorem 8.5)} \\
 &= \prod_{j=1}^k \left(\frac{z_j}{w_j}\right)^{r-k} \Gamma(\mathbf{z}_k)\Delta_\lambda(\mathbf{w}_k, H_{r-k}) && \text{(Proposition 8.8 and equation (64))} \\
 &= \prod_{j=1}^k \left(\frac{z_j}{w_j}\right)^{r-k} \sigma_+(\mathbf{z}_k)\bar{\sigma}_-(\mathbf{z}_k)\Delta_\lambda(\mathbf{w}_k, H_r) && \text{(Definition 4.3 – (30) of } \Gamma(\mathbf{z}_k)\text{)} \\
 &= \prod_{j=1}^k \left(\frac{z_j}{w_j}\right)^{r-k} \frac{1}{E_r(z_j)} \cdot \Delta_\lambda(\mathbf{w}_k, \bar{\sigma}_-(\mathbf{z}_k)H_r) && \text{(Definition 3.4 of the } B_r\text{-module structure of } \bigwedge^r V \text{ and Proposition 3.6)}
 \end{aligned}$$

as desired. Expression (66) for  $\bar{\sigma}_-(\mathbf{z}_k)h_j$  is Lemma 6.2.

**8.11 Remark** It is a good point to remark here that for  $k = 1$ , Theorem 8.10 above is a special case of a more general result due to the first author and Nasrollah Nejad [3], showing that the universal factorisation algebra of a generic polynomial itself carries a structure over the Lie algebra of endomorphism of a suitable module keeping track of the indeterminate coefficients of the polynomial.

### 8.3 Exponential form

Finally, let us define, as it is customary, new formal variable  $(x_j)_{j \geq 1}$  through the equality:

$$\exp\left(\sum_{j \geq 0} x_j z^j\right) = \frac{1}{E_r(z)}.$$

In this case one can write

$$\prod_{j=1}^k \frac{1}{E_r(z_j)} = \exp\left(\sum_{j=0} x_j p_j(\mathbf{z}_k)\right),$$

where  $p_j(\mathbf{z}_k) = z_1^j + \dots + z_k^j$  is the  $i$ -th power sum symmetric polynomial in  $(z_1, \dots, z_k)$  and where  $x_i$  is precisely the  $i$ -th degree power sum in the  $r$  universal roots  $(y_1, \dots, y_r)$  of the polynomial  $E_r(z)$ , i.e.  $E_r(z) = \prod_{i=1}^r (1 - y_i z)$  in the universal splitting  $\mathbb{Q}$ -algebra for the polynomial  $E_r(z) \in \mathbb{Q}[z]$ . This allows to shape our result in the form

**8.13 Corollary** We have:

$$\mathcal{E}(\mathbf{z}_k, \mathbf{w}_k^{-1}) \Delta_\lambda(H_r) = \prod_{j=1}^k \left( \frac{z_j}{w_j} \right)^{r-k} \exp \left( \sum_{j=0} x_j p_j(\mathbf{z}_k) \right) \Delta_\lambda(\mathbf{w}_k, \bar{\sigma}_-(\mathbf{z}_k) H_{r-k}). \quad (67)$$

Formula (67) is easy to use for practical computations of the  $gl(\bigwedge^k V)$ -representation of  $B_r$ , for those special case of  $k$  and  $r$  that everybody may possibly need.

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