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(Article begins on next page)

1 **AFM characterization of free-standing diamond microcantilevers for the investigation of the**  
2 **damage-induced variation of Young's modulus**

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15  
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17 microscopy

18  
19 **Abstract**

20 Due to their outstanding mechanical properties, diamond and diamond-like materials find  
21 significant technological applications ranging from well-established industrial fields (cutting tools,  
22 coatings, etc.) to more advanced mechanical devices as micro- and nano-electromechanical systems.  
23 The use of energetic ions is a powerful and versatile tool to fabricate three-dimensional  
24 micro-mechanical structures. In this context, it is of paramount importance to have an accurate  
25 knowledge of the effects of ion-induced structural damage on the mechanical properties of this

26 material, firstly to predict potential undesired side-effects of the ion implantation process, and  
27 possibly to tailor the desired mechanical properties of the fabricated devices. We present an Atomic  
28 Force Microscopy (AFM) characterization of free-standing cantilevers in single-crystal diamond  
29 obtained by a FIB-assisted lift-off technique, which allows a determination of the Young's modulus  
30 of the diamond crystal after the MeV ion irradiation process concurrent to the fabrication of the  
31 microstructures. The AFM measurements were performed with the beam-bending technique,  
32 showing that the thermal annealing process adopted during the microfabrication process allows for  
33 an effective recovery of the mechanical properties of the pristine crystal.

34

## 35 **1. Introduction**

36 MeV ion implantation has been widely exploited in recent years for the micro-fabrication of  
37 single-crystal diamond, through the implementation of the so-called "lift-off technique" [1-4]. This  
38 technique can be effectively adopted to fabricate micro-mechanical structures in single-crystal  
39 diamond, with applications ranging from high-frequency MEMS devices [5-11] to opto-mechanical  
40 resonators [12], thus taking advantage of the extreme mechanical properties of diamond [13].  
41 Recently, the latter topic attracted significant interest due to the outstanding properties of  
42 nitrogen-vacancy centers in diamond [14], whose spin-dependent optical transition can effectively  
43 couple with local mechanical stresses [15-17]. To this end, various different techniques have been  
44 employed to fabricate opto-mechanical resonators in diamond [18-21].

45 In the case of the lift-off technique, the fabrication process is based on the local conversion of  
46 diamond to a sacrificial graphitic layer through MeV-ion-induced damage [4]. The fabrication  
47 technique is very versatile, because the local induced damage density can be controlled by varying  
48 implantation parameters (namely, ion energy, species and fluence). Nonetheless, a residual damage  
49 density (and related mechanical stress) is induced in the non-sacrificial regions as a side-effect of  
50 the fabrication technique [22]. Similarly, also with other fabrication techniques [18-21], a residual

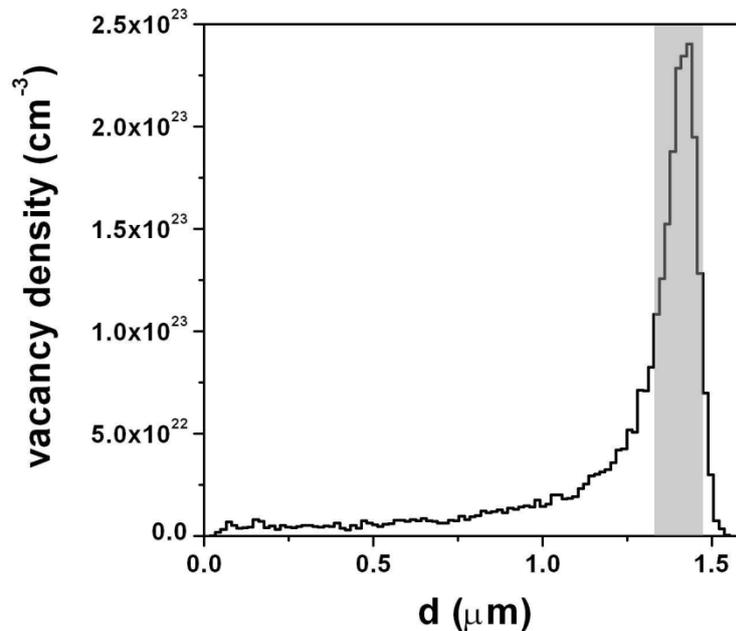
51 damage can be induced in the fabricated opto-mechanical microstructures, particularly if ion  
52 implantation is adopted to induce the formation of nitrogen-vacancy centers [23]  
53 For these reasons, it is necessary to accurately estimate deformation and stress levels to reliably  
54 design and fabricate MEMS structures. Moreover, the variation of elastic properties of damaged  
55 diamond as a function of induced damage density and post-processing (annealing) parameters  
56 remains to be clarified. In particular the Young's modulus of ion-implanted diamond can potentially  
57 vary between that of pristine diamond ( $>1$  TPa, in the presence of no damage) to that of amorphous  
58 carbon ( $\sim 10$  GPa, for full amorphization), i.e. over two orders of magnitude. Clearly, this large  
59 variation in elastic properties is likely to strongly affect modelling results in the fabrication of  
60 mechanical and opto-mechanical resonators. Attempts have been made to experimentally derive the  
61 variation of elastic properties of diamond as a function of induced damage, but only indirect  
62 estimations with limited accuracy have been obtained [24]. This lack of experimental data is partly  
63 due to its high Young's modulus, which makes it difficult to perform indentation experiments.  
64 Here, we perform a study of the elastic properties of ion-implanted diamond by means of Atomic  
65 Force Microscope (AFM) measurements on free-standing cantilever structures microfabricated in  
66 single-crystal diamond with FIB-assisted lift-off technique [3, 4].

67

## 68 **2. Micro-fabrication**

69 An artificial diamond sample grown by High Pressure High Temperature (HPHT) by ElementSix  
70 (UK) was employed in this work. The sample is  $2.6 \times 2.6 \times 0.5$  mm<sup>3</sup> in size and is classified as type  
71 Ib, on the basis of a nominal concentration of substitutional nitrogen  $\sim 500$  ppm. The sample is cut  
72 along the [100] crystal direction and is optically polished on one of the two opposite large faces.  
73 The sample was implanted at room temperature across one of the polished surfaces with 800 keV  
74 He<sup>+</sup> ions at the AN2000 accelerator of the INFN National Laboratories of Legnaro with a focused  
75 ion beam, in order to deliver a fluence of  $1 \times 10^{17}$  cm<sup>-2</sup>. The microbeam spot was  $\sim 10$   $\mu$ m in

76 diameter, and was raster-scanned to implant a rectangular area of  $\sim 500 \times 200 \mu\text{m}^2$ . The high density  
77 of damage induced by ion implantation promotes the conversion of the diamond lattice into an  
78 amorphous phase within a layer which is located at  $\sim 1.4 \mu\text{m}$  below the sample surface, as shown in  
79 Fig. 1.



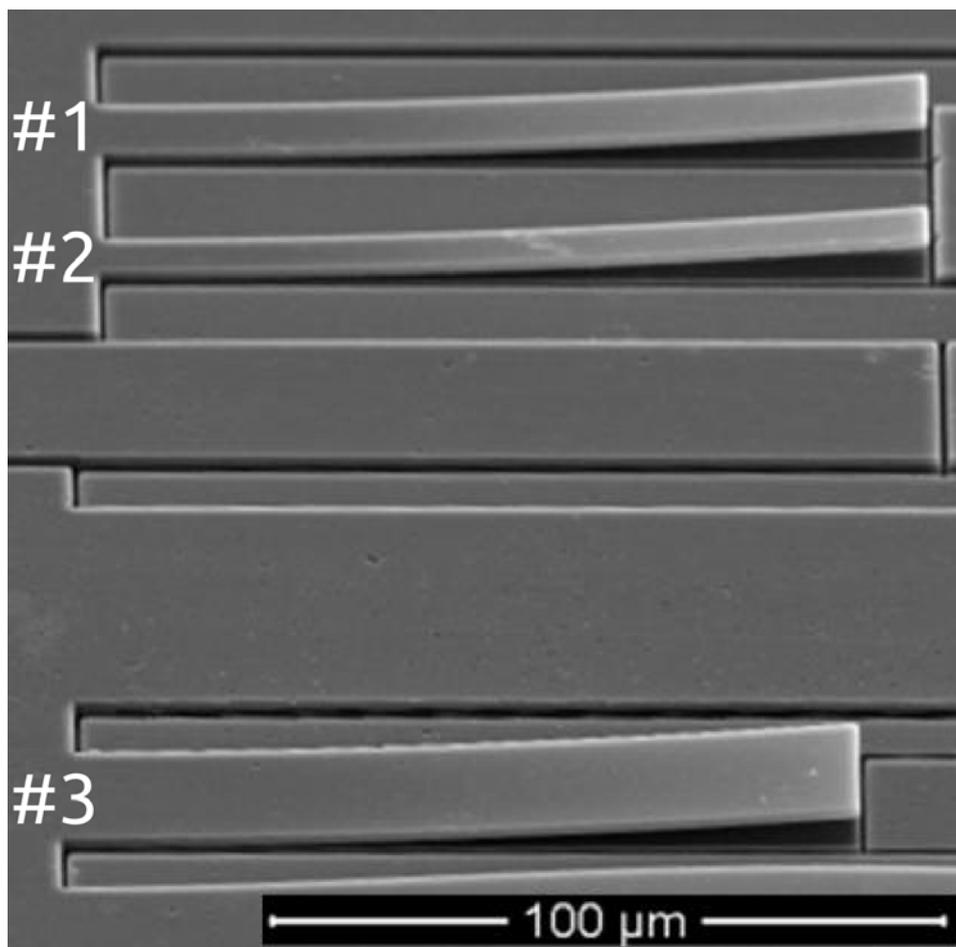
80

81 *Fig. 1: depth profile of the vacancy density induced by 800 keV He<sup>+</sup> ion implantation at fluence*  
82  *$1 \times 10^{17} \text{ cm}^{-2}$ , as evaluated with SRIM2013.00 Monte Carlo code [25] and assuming a linear*  
83 *dependence of the induced vacancy density from the implantation fluence [26].*

84

85 The sample was then annealed in high vacuum ( $\sim 10^{-6}$  mbar) at 1000 °C for 1 h, to convert the  
86 highly-damaged regions located at the ion end of range to a graphitic phase while removing the  
87 structural sub-threshold damage introduced in the layer overlying the damaged region. Following  
88 the fabrication scheme described in [1], FIB milling with 30 keV Ga<sup>+</sup> ions was subsequently  
89 performed on the implanted area, to expose the sub-superficial graphitic layer to the subsequent  
90 etching step, while defining the geometries of three different cantilever structures characterized by  
91 different widths (see Fig. 2). A thin Au film was deposited on the sample surface to avoid charge  
92 effects during FIB micro-machining. The sample was then exposed to contact-less electrochemical

93 etching [27]: the sample was immersed for several hours in de-ionized water with the region of  
94 interest comprised between two close (i.e. few millimeters) Pt electrodes kept at a DC voltage  
95 difference of  $\sim 100$  V. The process resulted in the selective removal of the sacrificial graphitic layer  
96 and in the creation of free-standing cantilever structures with a lateral geometry defined by the  
97 previous FIB micromachining, i.e. a length of  $117\ \mu\text{m}$  for cantilevers #1 and #2 and of  $111\ \mu\text{m}$  for  
98 cantilever #3. The widths of cantilevers #1, #2 and #3 were respectively  $13\ \mu\text{m}$ ,  $9\ \mu\text{m}$  and  $22\ \mu\text{m}$ .  
99 The thickness of all cantilevers corresponded to the penetration depth of the employed  $800\ \text{keV}$   
100 ions, i.e.  $1.3\ \mu\text{m}$ . As shown in Fig. 2, all cantilevers are slightly bent by the inner stresses caused by  
101 residual damage induced during the fabrication process within the "cap layer" comprised between  
102 the sub-superficial graphitic layer and the sample surface.

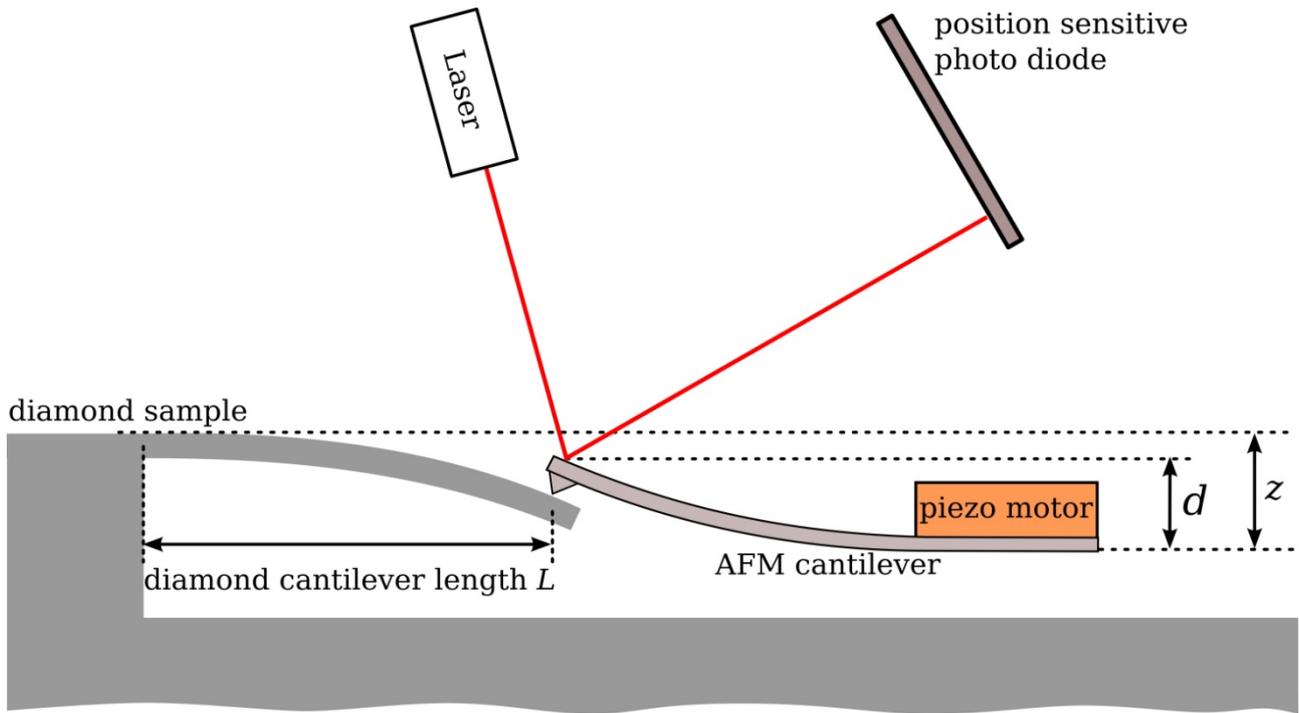


103

104 *Fig. 2: SEM micrograph of the three free-standing cantilever structures fabricated in single-crystal*  
105 *diamond by means of the FIB-assisted lift-off technique.*

### 106 3. AFM Characterisation

107 In order to determine the Young's modulus of the diamond, a beam-bending method was employed  
 108 [28-30]. The method consists in loading the microstructures using an AFM cantilever. As shown in  
 109 Fig. 3, the deflection  $d$  of the probing AFM cantilever for the displacement  $z$  of the piezomotor is  
 110 measured by means of a laser diode and a position-sensitive photodiode (Veeco Dimension 3100).



111

112 *Fig. 3: Schematic representation (not to scale) of the beam-bending technique employed to measure*  
 113 *the Young's modulus of the microfabricated diamond cantilevers. An AFM cantilever loads the*  
 114 *diamond microstructure at a length  $L$ , while the displacement  $d$  of the AFM cantilever is measured*  
 115 *as a function of the piezomotor displacement  $z$ .*

116

117 The effective stiffness  $k_{eff}$  of the system based on the coupling of the two cantilever structures is  
 118 measured by recording approach curves. The effective stiffness  $k_{eq}$  is equal to:

119

$$120 \quad k_{eff} = \frac{1}{\frac{1}{k_{AFM}} + \frac{1}{k_{diam}}} \quad (1)$$

121

122 where  $k_{AFM}$  and  $k_{diam}$  are respectively the stiffness values of the probing AFM cantilever and of the  
 123 diamond cantilever under test. The  $d/z$  value is equal to [28]:

124

$$125 \quad \frac{d}{z} = \frac{k_{eff}}{k_{AFM}} \quad (2)$$

126

127 Thus, the stiffness of the beam under test can be determined, if the stiffness of the AFM cantilever  
 128 is known. The adopted AFM cantilever is a single-crystalline silicon cantilever (NCLR,  
 129 Nanoworld). Its geometry is precisely determined by means of SEM microscopy and its stiffness  
 130 was evaluated as  $k_{AFM} = (57.0 \pm 1.2) \text{ N m}^{-1}$ .

131 The three diamond cantilevers mentioned above were investigated, and for each of them the  
 132 stiffness values were measured in correspondence of several (i.e. >5) different positions along the  
 133 cantilever axis. The stiffness of the diamond cantilever is determined by the following formula:

134

$$135 \quad k_{diam} = \frac{3 \cdot E \cdot I}{(1 - \nu^2) \cdot L^3} \quad (3)$$

136

137 where  $E$  is its Young's modulus,  $I$  is its areal moment of inertia,  $\nu$  is its Poisson's ratio, and  $L$  is its  
 138 length, i.e. the distance from the clamping point where the load is applied [31]. Due to a possible  
 139 systematic error in the determination of the cantilever length and the finite stiffness of the cantilever  
 140 fixture, we used the following correction to fit the  $k_{diam}$  vs  $L$  trend [32, 33]:

141

$$142 \quad k_{diam} = \frac{3 \cdot E \cdot I}{(1 - \nu^2) \cdot (L + L_0)^3} \quad (4)$$

143

144 The Poisson's ratio was estimated with the value corresponding to pristine undamaged diamond, i.e.  
 145 0.105 [34], considering bending along the (100) direction, consistently with the well-defined  
 146 orientation of the structure with respect to the crystal orientation. The geometry of the diamond

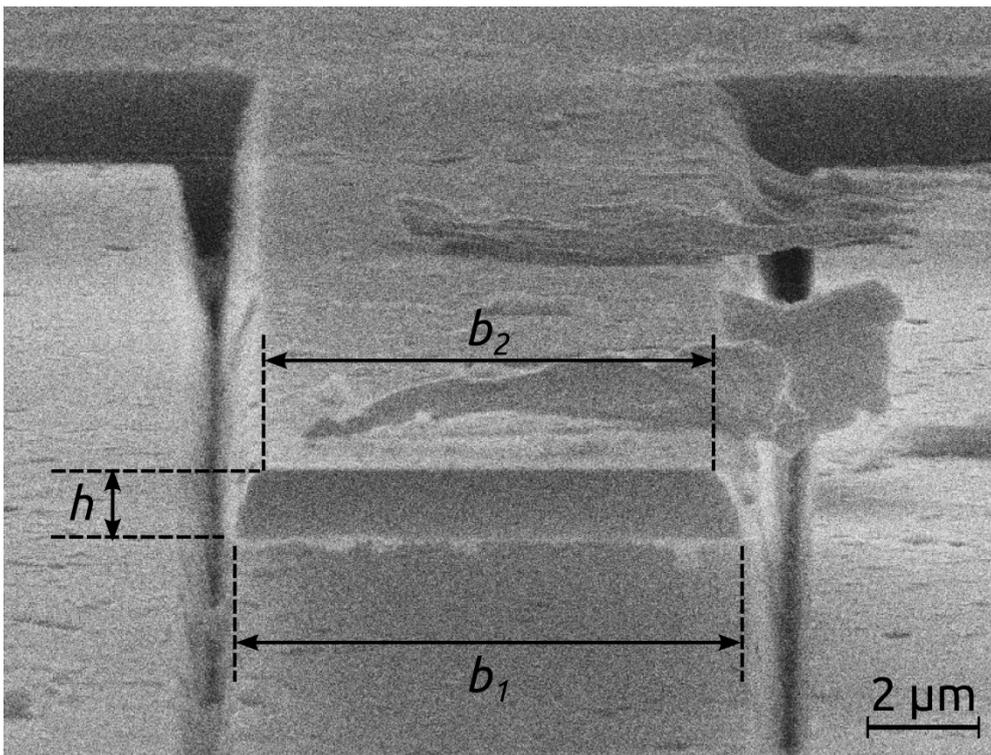
147 beams was measured by means of SEM microscopy and the moment of inertia  $I$  for each cantilever  
148 was calculated by taking into consideration their slightly trapezoidal beam cross-section (see Fig.  
149 4), as follows [35]:

150

$$151 \quad I = \frac{h^3}{36} \cdot \frac{(b_1^2 + 4 \cdot b_1 \cdot b_2 + b_2^2)}{(b_1 + b_2)} \quad (5)$$

152

153 where  $h$  is the thickness of the cantilever and  $b_1$  and  $b_2$  are the two widths of the trapezoidal cross  
154 section (see Fig. 4). The Young's modulus  $E$  of the diamond cantilever is then calculated using by  
155 fitting the  $k_{diam}$  vs  $L$  trend with Eq. (4), in which the moment of inertia  $I$  is estimated as reported in  
156 Eq. (5).



157

158 *Fig. 4: SEM micrograph of the trapezoidal cross-section of cantilever #2.  $h$  is the thickness of the*  
159 *cantilever, while  $b_1$  and  $b_2$  are respectively the lower and upper widths of the cross section.*

160

161 Representative results of the  $k_{diam}$  vs  $L$  measurements relative to cantilever #3 are reported in  
162 Fig. 5a, together with the fitted curve (see using Eq. (4)). The fit of the experimental data is very

163 satisfactory, yielding a Young's modulus value of  $E = (9.6 \pm 1.1) \times 10^2$  GPa. The Young's moduli  
 164 for the three cantilevers are reported in Fig. 5b, together with their weighted average value and its  
 165 relevant uncertainty, as well as with a comparison with the reference value of pristine single-crystal  
 166 diamond. In the case of cantilever #1, the discrepancy between the measured Young's modulus and  
 167 the reference value is statistically significant, and the calculated stiffness value is exceedingly high.  
 168 This can be tentatively attributed to a non-ideal detachment of the beam from the substrate during  
 169 the etching fabrication process, which potentially increases the stiffness of the structure and yields  
 170 incorrect values when using the above formulas. Nevertheless, the weighted average of the three  
 171 Young's modulus values yields an estimation of  $E = (1.11 \pm 0.08)$  TPa, which is statistically  
 172 compatible with the reference value in literature for single-crystal diamond, i.e. 1.05 TPa [36]. This  
 173 compatibility is remarkable, particularly if it is considered that we assumed a value of the Poisson's  
 174 ratio corresponding to the pristine diamond, thus introducing a potential systematic error.  
 175 Since single-crystal diamond is mechanically anisotropic [34], with, the Young's modulus in the  
 176 (100) direction can be calculated as follows:

177

$$178 \quad E_{(100)} = \frac{1}{S_{11}} \quad (6)$$

179

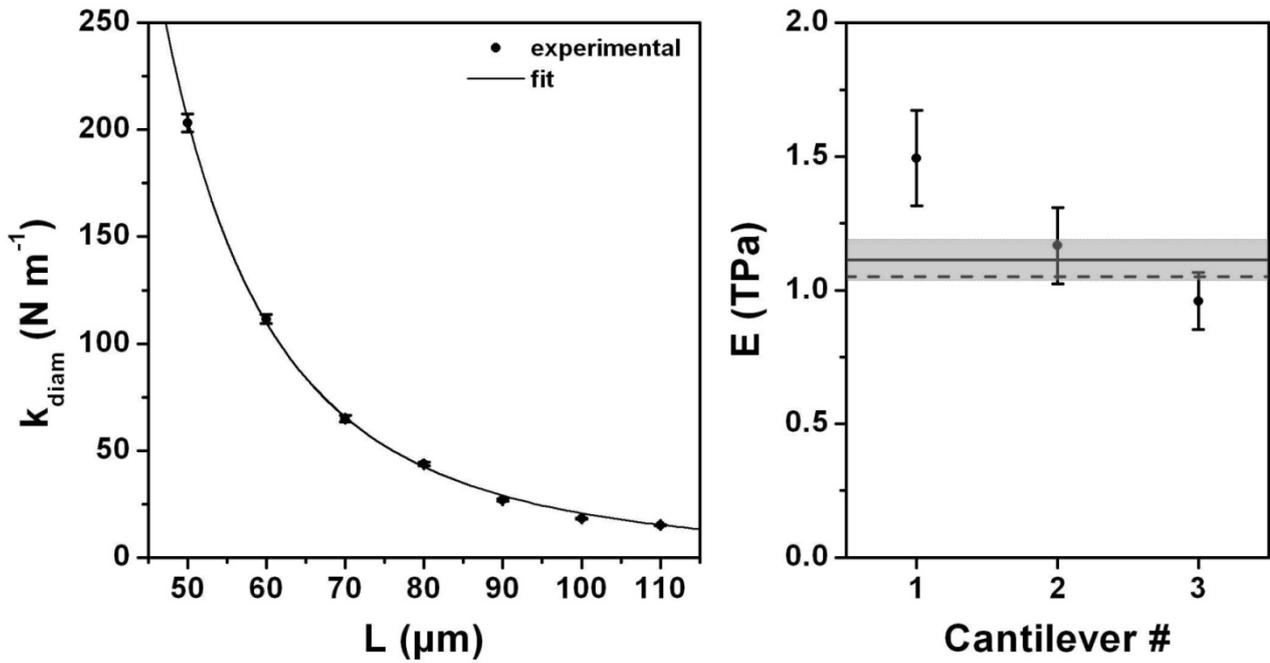
180 where  $S_{11}$  is the 11 component of elastic coefficient tensor  $S_{ij}$  according to the standard notation.  
 181 Thus, the Poisson's ratio describing the contraction in the (001) direction due to tension in the (100)  
 182 direction can be calculated as follows:

183

$$184 \quad \nu_{(100)} = -S_{12} \cdot E_{100} \quad (7)$$

185

186 From the values of the  $S_{ij}$  elastic coefficients reported in [37], we obtain  $E_{(100)} = 1.05$  TPa and  
 187  $\nu_{(100)} = 0.11$ . Again, these values are close to the literature values for single-crystal diamond [36].



188

189 Fig. 5: a) experimental values (dots) of the stiffness of cantilever #3 at different beam lengths  
 190 together with the relevant fitting curve (line). b) Young's moduli for the three cantilevers estimated  
 191 from the fit of the experimental data (dots), together with the weighted average value (solid line)  
 192 and the relevant uncertainty (grey box); the Young's modulus value for pristine single-crystal  
 193 diamond taken from literature [36] is reported for comparison (dashed line).

194

#### 195 4. Conclusions

196 We demonstrated the feasibility and reliability of an AFM-based beam-bending technique to  
 197 determine the mechanical properties single-crystal diamond cantilevers, and investigated the effects  
 198 of MeV ion implantation and subsequent high-temperature annealing on these mechanical  
 199 properties. The obtained results provide direct evidence that both the Young's modulus and  
 200 Poisson's ratio values of diamond after sub-graphitization-threshold irradiation and high-  
 201 temperature annealing are fully recovered to their pristine values. These results provide useful  
 202 information for the reliable design of (opto-) mechanical resonators in single-crystal diamond. We  
 203 envisage to repeat these measurements on micro-structures subjected to controlled ion irradiation,

204 with the purpose of directly investigating the effect of ion-induced damage on their mechanical  
205 properties, thus allowing the fine-tuning of their resonance frequencies.

206

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215

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