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Effect of defect size on P-S-N curves in Very-High-Cycle Fatigue

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Abstract

It is well-known that internal defects play a key role in the Very-High-Cycle Fatigue (VHCF) response of metallic materials. VHCF failures generally nucleate from internal defects, whose size strongly affects the material strength and life. Therefore, S-N curves in the VHCF regime are defect size dependent and the scatter of fatigue data is significantly influenced by the statistical distribution of the defect size within the material.

The present paper proposes an innovative approach for the statistical modeling of Probabilistic-S-N (P-S-N) curves in the VHCF regime. The proposed model considers conditional P-S-N curves that depend on a specific value of the initial defect size. From the statistical distribution of the initial defect size, marginal P-S-N curves are estimated and the effect of the risk-volume on the VHCF response is also modeled. Finally, the paper reports a numerical example that quantitatively illustrates the concepts of conditional and marginal P-S-N curves and that shows the effect of the risk-volume on the VHCF response.

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Keywords: Random fatigue limit; Crack growth; Paris' law

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1. Introduction

In many industrial fields (aerospace, railway, energy, automotive, marine), machinery components experience Very-High-Cycle Fatigue (VHCF) in-service conditions. A reliable design against VHCF failure is of utmost importance in these cases. In the last decades, a great effort has been made in order to collect VHCF data and to provide designers with reliable information on the VHCF response of different metallic materials. A number of models have been proposed in the literature (see, e.g., the review by Li, 2012), in order to quantify either the VHCF strength or the material fatigue limit. However, randomness of VHCF data has been rarely taken into account and the available statistical models have been prevalently devoted to the description of the random transition between High-Cycle Fatigue (HCF) and VHCF (Sakai et al, 2010; Harlow, 2011). A general probabilistic model, which can model both the transition HCF-VHCF and the possible presence of a fatigue limit, has been first proposed in Paolino et al. (2013) and Paolino et al. (2016). The model has been recently exploited in Paolino et al. (2016) for the statistical description of the VHCF response, on the basis of the hydrogen embrittlement theory proposed by Murakami (Murakami, 2002).

The present paper extends the approach in Paolino et al. (2016) to the different weakening mechanisms that have been proposed in the literature for explaining the formation of the so-called Fine Granular Area (FGA) around the initial defect. The FGA (also called Optically Dark Area or ODA by Murakami, Granular Bright Facet or GBF by Shiozawa and Rough Surface Area or RSA by Ochi) plays a major role in the VHCF response, since its formation consumes more than the 95% of the VHCF life. Researchers dispute about the actual mechanism behind the FGA formation (see, e.g., Li, 2016), but they unanimously accept that, within the FGA, crack can grow even if the Stress Intensity Factor (SIF) is below the SIF threshold for crack growth.

In the present paper, the reduction of the SIF threshold within the FGA is mechanistically modeled according to Paolino et al. (2017) and in agreement with the different weakening mechanisms proposed in the literature. Furthermore, since fatigue is scattered by nature, randomness of VHCF data is also taken into account and statistically modeled through originally defined Probabilistic-S-N (P-S-N) curves. The proposed model considers conditional P-S-N curves that depend on a specific value of the initial defect size. From the statistical distribution of the initial defect size, marginal P-S-N curves are estimated and the effect of the risk-volume on the VHCF response is also modeled. Finally, the paper reports a numerical example that quantitatively illustrates the concepts of conditional and marginal P-S-N curves and that shows the effect of the risk-volume on the VHCF response.

Nomenclature

FGA	Fine Granular Area
SIF	Stress Intensity Factor
LEV	Largest Extreme Value distribution
HV	Vickers Hardness
$a_c, a_d, a_0, a_{FGA}, a_{FiE}$	projected areas of defects
$c_{th,g}, \alpha_{th,g}, c_{th,r}, \alpha_{th,r}$	parameters involved in SIF thresholds
$c_Y, m_Y, n_Y, \mu_{\sqrt{A}}, \sigma_{\sqrt{A}}, \sigma_{K_{th,g}}, \sigma_Y$	parameters of statistical distributions
$c_I, m_I, c_{II}, m_{II}, c_{III}, m_{III}, c_S, m_S$	Paris' constants in the three stages of crack growth
$F_{\sqrt{A_0,V}}, f_{\sqrt{A_0,V}}, F_{S_I \sqrt{a_{d,0}}}$	statistical distributions
$k_d, k_{th,g}, k_{th,l}, k_{th,r}$	relevant SIF values
n_f, N_f	number of cycles to failure
n_I, n_{II}, n_{III}	number of cycles in the three stages of crack growth
s	applied stress amplitude
V, V_{exp}	risk-volumes

2. Methods

In Section 2.1, a general expression for modeling the crack growth rate from the initial defect up to the VHCF failure is presented. In Section 2.2, the statistical distribution of the defect size and the related size-effects are analyzed and discussed. Finally, in Section 2.3, the statistical distribution of the fatigue limit is analytically defined and a model for the fatigue limit as a function of the risk-volume is proposed.

In the following, according to Paolino et al. (2017), $k_{th,g}$ denotes the global SIF threshold, $k_{th,r}$ denotes the reduction SIF, k_d denotes the SIF for an internal defect, $k_{th,l}$ denotes the local SIF threshold (i.e., $k_{th,l} = k_{th,g} - k_{th,r}$), a_0 is the projected area of the initial defect and a_{FGA} is the projected area of the FGA.

2.1. Crack growth rate within the FGA

In the VHCF literature (Tanaka and Akiniwa, 2002; Marines-Garcia et al., 2008; Su et al., 2017), the crack growth rate within the FGA is commonly modeled with the Paris’ law. Three stages can be present in the crack growth rate diagram related to a VHCF failure from internal defect (Fig. 1):

- Stage I: the below-threshold region within the FGA, from k_{a_0} (SIF associated to the initial defect) up to $k_{th,g}$;
- Stage II: the steady stage, from the border of the FGA (SIF equal to $k_{th,g}$) up to the border of the fish-eye (with SIF equal to k_{FiE});
- Stage III: the unsteady stage, beyond the fish-eye border (with SIF larger than k_{FiE} , up to the failure).

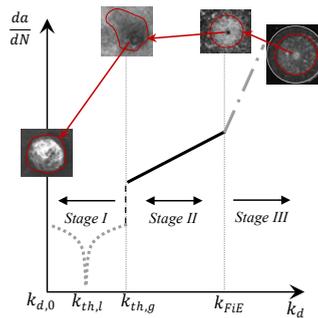


Fig. 1. The three stages of crack propagation in a crack growth rate diagram for VHCF failures from internal defects.

For a stress ratio equal to -1, the modified Paris’ law proposed by Donahue et al. (1972) is considered to model the crack growth within the FGA:

$$\frac{da}{dN} = c_I (k_d - k_{th,l})^{m_I}, \tag{1}$$

where c_I and m_I are the Paris’ constants related to Stage I, from the initial defect size $\sqrt{a_0}$ up to $\sqrt{a_{FGA}}$.

According to (Tanaka and Akiniwa, 2002; Marines-Garcia et al., 2008; Su et al., 2017), in Stage II, from the border of the FGA up to the border of the fish-eye (with size $\sqrt{a_{FiE}}$), the crack growth rate follows the conventional Paris’ law:

$$\frac{da}{dN} = c_{II} k_d^{m_{II}}, \tag{2}$$

where c_{II} and m_{II} are the two Paris’ constants related to Stage II, from $\sqrt{a_{FGA}}$ up to $\sqrt{a_{FiE}}$.

Generally, crack propagates beyond the fish-eye border until it reaches the border of the final fracture, with size $\sqrt{a_c}$. In these cases, a third stage of crack propagation is visible on the fracture surfaces and it can be modeled again with the conventional Paris' law (Su et al., 2017):

$$\frac{da}{dN} = c_{III} k_d^{m_{III}}, \quad (3)$$

where c_{III} and m_{III} are the two Paris' constants related to Stage III, from $\sqrt{a_{FIE}}$ up to $\sqrt{a_c}$. It is worth to note that, if the final fracture occurs when the crack size reaches the border of the fish-eye, Stage III can be neglected.

The number of cycles to failure, N_f , can be expressed as the sum of the number of cycles consumed within the three stages of propagation:

$$n_f = n_I + n_{II} + n_{III}. \quad (4)$$

Following the VHCF literature (see, e.g., Su et al., 2017), n_I can be estimated by subtracting, from the experimental n_f , the numbers of cycles consumed in Stages II and III, which are obtained through integration of Eqs. (2) and (3), respectively.

According to Paolino et al. (2017), the experimental n_I values ($n_{I,exp}$) can be used for the estimation of the Paris' constants c_I , m_I and of the parameters $c_{th,r}$ and $\alpha_{th,r}$ involved in the expression of the $k_{th,l}$. Parameter estimates are obtained through the nonlinear Least Squares Method by minimizing the sum of squared percent errors between the experimental $\log_{10}(n_{I,exp})$ values and the estimated $\log_{10}(n_{I,est})$ values computed through integration of Eq. (1). Finally, an estimate of the fatigue limit can be defined, according to the procedure described in Paolino et al. (2017).

2.2. Statistical distribution of the initial defect size and related size-effect

Size-effects significantly affect the VHCF response of materials (Furuya, 2011; Tridello et al., 2015): the larger the risk-volume (volume of material subjected to a stress larger than the 90% of the maximum stress), the larger the probability of large defects, with a subsequent reduction of the VHCF strength. The dependency between the fatigue limit and the risk-volume is generally modeled by taking into account the statistical distribution of the internal defect size. According to the VHCF literature (Murakami, 2002), the defect originating failure can be considered as the largest defect present within the specimen risk-volume and, therefore, in a statistical framework the internal defect size random variable (rv) follows a Largest Extreme Value (LEV) distribution.

The cumulative distribution function (cdf) of the LEV distribution is reported in Eq. (5): it provides the probability of an internal defect with size smaller than $\sqrt{a_0}$ in a volume V larger than the risk-volume of the tested specimens (V_{exp}).

$$F_{\sqrt{a_0}|V}(\sqrt{a_0}; V) = \left(\exp \left(- \exp \left(- \frac{\sqrt{a_0} - \mu_{\sqrt{A}}}{\sigma_{\sqrt{A}}} \right) \right) \right)^{\frac{V}{V_{exp}}}, \quad (5)$$

where $\mu_{\sqrt{A}}$ and $\sigma_{\sqrt{A}}$ are the two constant parameters of the distribution. The probability density function (pdf) of the LEV distribution, $f_{\sqrt{a_0}|V}$, is expressed by:

$$f_{\sqrt{a_0}|V}(\sqrt{a_0}; V) = \frac{1}{\sigma_{\sqrt{A}}} \exp \left(- \frac{\sqrt{a_0} - (\mu_{\sqrt{A}} + \log(V/V_{exp}))}{\sigma_{\sqrt{A}}} \right) - \exp \left(- \frac{\sqrt{a_0} - (\mu_{\sqrt{A}} + \log(V/V_{exp}))}{\sigma_{\sqrt{A}}} \right). \quad (6)$$

According to Murakami (2002), the parameters of the LEV distribution can be estimated from the defect sizes $\sqrt{a_0}$ measured on the fracture surfaces of the tested specimens (i.e., with $V = V_{exp}$ in Eqs. (5) and (6)).

2.3. Statistical distribution of the fatigue limit and related size-effect

According to Paolino et al. (2016) the cdf of the fatigue limit for a given defect size (conditional VHCF limit) can be expressed as:

$$F_{S_{l|\sqrt{a_0}}}(s_l; \sqrt{a_0}) = \Phi \left(\frac{1}{\sigma_{K_{th,g}}} \left(\log_{10}(s_l) - \log_{10} \left(c_{th,g} \frac{c_{s_l}(HV+120)}{\sqrt{a_0}^{1/2-\alpha_{th,g}}} \right) \right) \right) = \Phi \left(\frac{\log_{10}(s_l) - \mu_l(\sqrt{a_0})}{\sigma_l} \right), \quad (7)$$

where $\Phi(\cdot)$ denotes the standardized Normal cdf, $c_{th,g}$, $\alpha_{th,g}$ and $\sigma_{K_{th,g}}$ are the parameters involved in the statistical distribution of the global SIF threshold (Paolino et al., 2016; Paolino et al., 2017), HV is the Vickers hardness of the

material and $c_{s_l} = \left(\frac{(1/2-\alpha_{th,g})0.5\sqrt{\pi}}{(\alpha_{th,g}-\alpha_{th,r})c_{th,r}} \right)^{\frac{1/2-\alpha_{th,g}}{1/2-\alpha_{th,r}}} \frac{\alpha_{th,g}-\alpha_{th,r}}{0.5\sqrt{\pi}(1/2-\alpha_{th,r})}$.

The cdf of the fatigue limit as a function of the risk-volume (marginal VHCF limit) can be obtained from the definition of marginal cdf and by taking into account the defect size distribution in Eq. (6):

$$F_{S_{l|V}}(s_l; V) = \int_0^\infty F_{S_{l|\sqrt{a_0}}}(s_l; \sqrt{a_0}) f_{\sqrt{A_0}|V}(\sqrt{a_0}; V) d\sqrt{a_0}. \quad (8)$$

2.4. P-S-N curves and related size-effect

The P-S-N curves statistically model the VHCF material response in the fatigue limit region and in the finite fatigue life region.

Eqs. (7) and (8) model the randomness in the fatigue limit region. According to (Paolino et al., 2016), the cdf of the finite fatigue life for given initial defect and applied stress s (conditional finite VHCF life), $N_{f,<\infty}|(s, \sqrt{a_0})$, can be expressed as:

$$F_{N_{f,<\infty}|(s,\sqrt{a_0})}(n_f; s, \sqrt{a_0}) = \Phi \left(\frac{\log_{10}(n_f) - (c_Y + m_Y \log_{10}(s) + n_Y \log_{10}(\sqrt{a_0}))}{\sigma_Y} \right) = \Phi \left(\frac{\log_{10}(n_f) - \mu_Y(s, \sqrt{a_0})}{\sigma_Y} \right), \quad (7)$$

where σ_Y denotes a constant standard deviation and c_Y , m_Y and n_Y are three constant parameters that can be estimated from the experimental failures through the Least Squares Method.

According to the probabilistic model “One failure mode due to one cause with fatigue limit” described in (Paolino et al., 2013), the cdf of the fatigue life (finite as well as infinite) for given initial defect and applied stress s (conditional VHCF life), $N_f|(s, \sqrt{a_0})$, is given by:

$$F_{N_f|(s,\sqrt{a_0})}(n_f; s, \sqrt{a_0}) = F_{S_{l|\sqrt{a_0}}}(s; \sqrt{a_0}) F_{N_{f,<\infty}|(s,\sqrt{a_0})}(n_f; s, \sqrt{a_0}). \quad (8)$$

The α -th quantile of the conditional VHCF life, $n_{f,(s,\sqrt{a_0}),\alpha}$, can be obtained by substituting $F_{N_f|(s,\sqrt{a_0})}(n_f; s, \sqrt{a_0})$ with α and by solving the equation with respect to n_f for different values of s :

$$n_{f,(s,\sqrt{a_0}),\alpha} = 10^{\mu_Y(s,\sqrt{a_0}) + \sigma_Y \Phi^{-1} \left(\alpha / \Phi \left(\frac{\log_{10}(s) - \mu_l(\sqrt{a_0})}{\sigma_l} \right) \right)}, \quad (9)$$

where s must be larger than $10^{\mu_l(\sqrt{a_0}) + \sigma_l \Phi^{-1}(\alpha)}$ in order to have finite values of $n_{f,(s,\sqrt{a_0}),\alpha}$. Eq. (9) thus provides the P-S-N curves given the initial defect size (conditional P-S-N curves).

The cdf of the VHCF life as a function of the risk-volume (marginal VHCF life) can be obtained from the definition of marginal cdf and by taking into account the defect size distribution in Eq. (6):

$$F_{N_f|(s,V)}(n_f; s, V) = \int_0^\infty F_{N_f|(s,\sqrt{a_0})}(n_f; s, \sqrt{a_0}) f_{\sqrt{A_0}|V}(\sqrt{a_0}; V) d\sqrt{a_0}. \quad (10)$$

The α -th quantile of the fatigue life can be obtained by substituting $F_{N_f|(s,V)}(n_f; s, V)$ with α and by solving the equation with respect to s for different values of n_f . Eq. (10) thus provides the P-S-N curves of the material for a given risk-volume (marginal P-S-N curves).

3. Application to an experimental dataset

The models proposed in Section 2 are here applied to an experimental dataset. VHCF tests are carried out on Gaussian specimens (Tridello et al., 2015) made of an AISI H13 steel with Vickers hardness 560 kgf/mm^2 and $V_{exp} = 2300 \text{ mm}^3$. Details on the testing setup and on the tested material are reported in Tridello et al. (2015) and in Tridello et al. (2016) and they will not be recalled here for the sake of brevity. Twelve specimens are loaded at constant stress amplitude up to failure. The number of cycles to failure ranges from $4.2 \cdot 10^7$ to $3.85 \cdot 10^9$ cycles. The initial defect sizes ($\sqrt{a_0}$) and the FGA sizes ($\sqrt{a_{FGA}}$) are measured from pictures taken by a Scanning Electron Microscope (SEM) and by an optical microscope. In order to take into account the stress variation within the V_{exp} , the local stress amplitude in the vicinity of the initial defect is considered as the stress amplitude applied during the test. The local stress amplitudes are in the range 500 - 635 MPa.

The parameters $c_{th,g}$, $\alpha_{th,g}$, $\sigma_{K_{th,g}}$ and c_{s_l} , which are involved in the fatigue limit expressions (Eqs. (7) and (8)), are estimated according to the procedure described in Paolino et al. (2017). Fig. 2a shows the $k_{th,g}$ values with respect to $\sqrt{a_{FGA}}$ together with the estimated model ($k_{th,g,\alpha}$ is the α -quantile of the Global SIF threshold, HV is the material Vickers Hardness and Φ^{-1} is the inverse cumulative distribution function of a standardized Normal distribution). Fig. 2b shows the conditional VHCF limit curves as a function of the initial defect size ($s_{l,\sqrt{a_0},\alpha}$ is the α -quantile of the conditional fatigue limit). The 0.1-th and the 0.9-th quantiles are also depicted in Fig. 2.

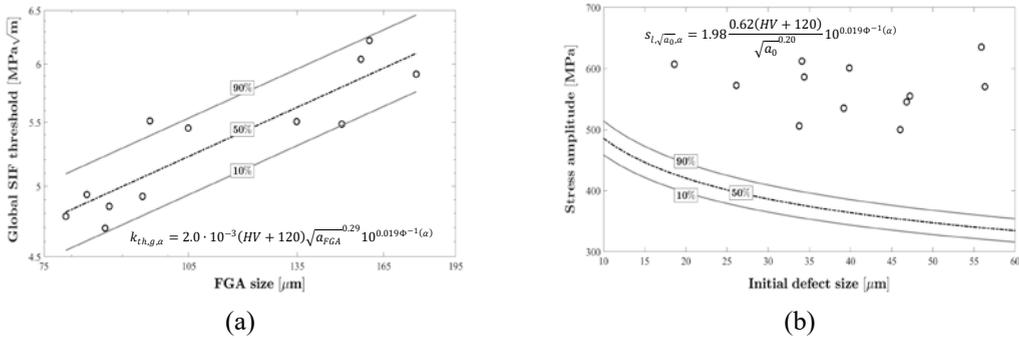


Fig. 2. (a) Global SIF threshold vs. FGA size. (b) Conditional VHCF limit vs. initial defect size.

According to Fig. 2b, the fatigue limit decreases with the initial defect size (Murakami, 2002; Furuya, 2011). The estimated fatigue limit curves are below the experimental failures, as expected from the definition of fatigue limit. The proposed model is therefore effective in the estimation of the fatigue limit variation with respect to the initial defect size and ensures a reliable safety margin with respect to the experimental failures.

The distribution of initial defect size is estimated according to Murakami (2002). Fig. 3a shows the Gumbel plot of the measured $\sqrt{a_0}$ values together with the estimated LEV cdf. Parameter estimation is carried out by considering $V = V_{exp} = 2300 \text{ mm}^3$. From the initial defect size distribution and according to Eq. (8), the 0.1-th, 0.5-th and the 0.9-th quantiles of the fatigue limit are estimated for risk-volumes larger than V_{exp} and then depicted in Fig. 3b.

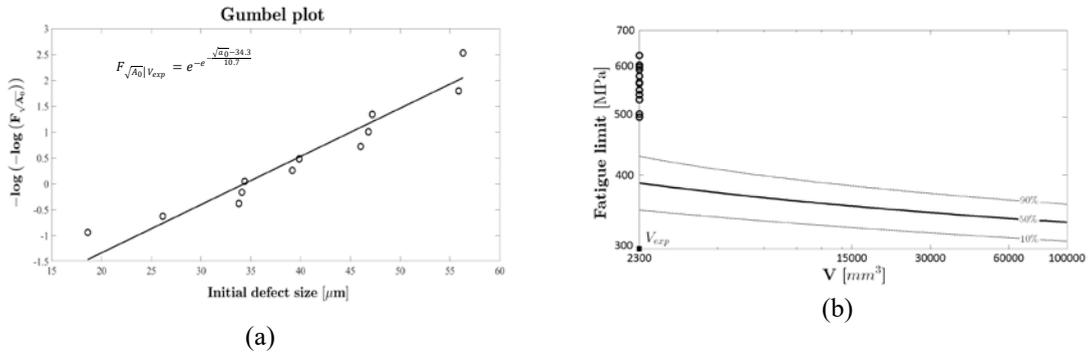


Fig. 3. (a) Gumbel plot of the initial defect size. (b) Variation of the marginal VHCF limit with the risk-volume.

Fig. 4 shows a plot of the conditional P-S-N curves. Fig. 4a plots the 0.1-th, the 0.5-th and the 0.9-th curves corresponding to an initial defect with median size; whereas Fig. 4b depicts the median curves for values of the initial defect size ranging from 15 μm to 60 μm .

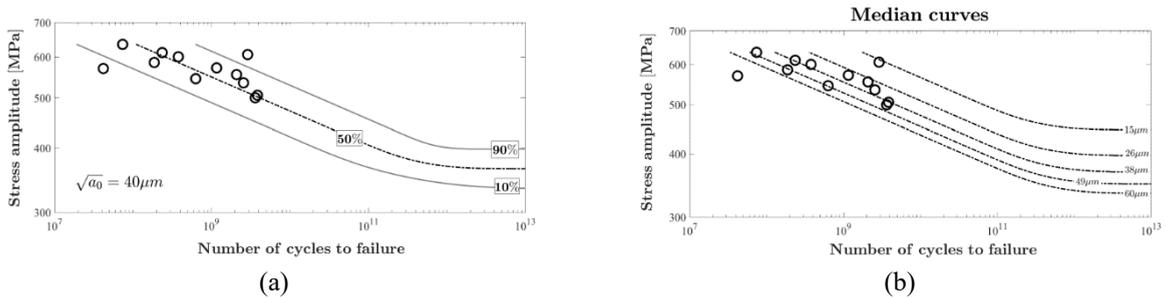


Fig. 4. Conditional P-S-N curves: (a) 80% confidence band for a median defect size; (b) median curves for different defect sizes.

Finally, by taking into account the defect size distribution, the marginal P-S-N curves (no more conditioned to the inclusion size, as in Fig. 4) are estimated. Fig. 5a plots the 0.1-th, the 0.5-th and the 0.9-th curves corresponding to a risk-volume equal to $V_{exp} = 2300 \text{ mm}^3$; whereas Fig. 5b depicts the median curves the risk-volume larger than V_{exp} .

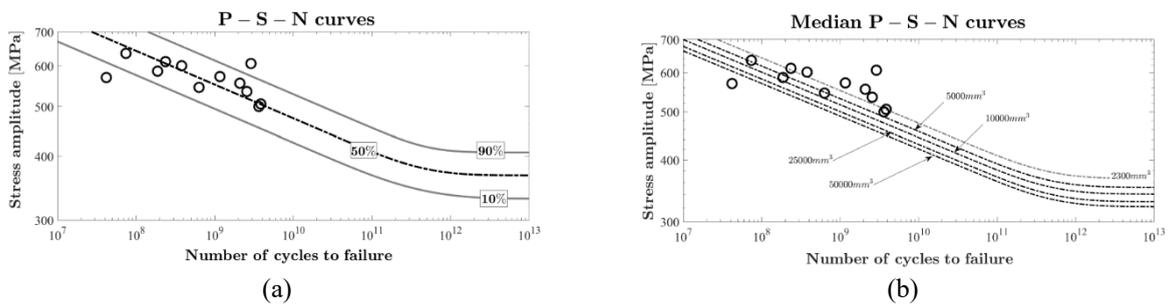


Fig. 5. Marginal P-S-N curves: (a) 80% confidence band for a risk-volume equal to V_{exp} ; (b) median curves for larger risk-volumes.

Marginal P-S-N curves reported in Fig. 5 can be used for component design, since they are not conditioned to a specific value of inclusion size (Fig. 4). In particular, the model in Fig. 5b can be used to predict the VHCF life of components characterized by risk-volumes (up to 100000 mm^3) significantly larger than risk-volumes of specimens commonly tested.

4. Conclusions

In the present paper, a general model for the P-S-N curves in the VHCF regime was proposed. The statistical distribution of the P-S-N curves for a given initial defect size (conditional P-S-N curves) and the P-S-N curves as a function of the risk-volume (marginal P-S-N curves) were defined. The proposed model permitted to take into account defect size-effects on the VHCF response and to predict the P-S-N curves of components characterized by large risk-volumes.

The model was successfully applied to an experimental dataset. The P-S-N curves as a function of the defect size and as a function of the risk-volume were estimated and were in agreement with the experimental data: for the tested risk-volume, about 80% of data were within the estimated 80% confidence band.

The proposed model could be effectively adopted for the estimation of the P-S-N curves when designing large components subjected to VHCF.

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