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Abstract

We introduce and study a mean-field model for a system of spatially distributed players interacting through an evolutionary game driven by a replicator dynamics. Strategies evolve by a replicator dynamics influenced by the position and the interaction between different players and return a feedback on the velocity field guiding their motion.

One of the main novelties of our approach concerns the description of the whole system, which can be represent-dimensional state space (pairs (x, σ) of position and distribution of strategies). We provide a Lagrangian and a Eulerian description of the evolution, and we prove their equivalence, together with existence, uniqueness, and stability of the solution. As a byproduct of the stability result, we also obtain convergence of the finite agents model to our mean-field formulation, when the number N of the players goes to infinity, and the initial discrete distribution of positions and strategies converge.

To this aim we develop some basic functional analytic tools to deal with interaction dynamics and continuity equations in Banach spaces that could be of independent interest. © 2021 The Authors. *Communications on Pure and Applied Mathematics* published by Wiley Periodicals LLC.

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1 Introduction

1.1 Evolutionary games

Physical systems naturally tend to minimize the potential energy. For this funda-
mental reason the study of steady states in physical systems is of utmost relevance,
given the expected frequency for such states to occur. This is also the rationale
according to which game theorists have focused on the characterization of game
equilibria. Very celebrated is in fact the work of John F. Nash [33], where a new no-
tion of noncooperative equilibrium is introduced. The main result of [33], building
on John von Neumann's notion of mixed strategy (necessary to ensure existence
of saddle points in zero sum games with two players), is the existence of mixed
strategy equilibria for noncooperative games with any finite number of players.

However, already at the origin of game theory, Oskar Morgenstern and John von
Neumann pointed out in their classical treatise on game theory [42] the desirability
of a "dynamical" approach to complement their "static" game solution concept. In
fact, while in physical systems evolutions towards minima of the potential energy
are explained according to Newton's law (for which evolutions are, for conserva-
tive forces, the gradient flows of the potential energy), it is not at all clear whether
and how in dynamical games equilibria can emerge. Certainly John Nash had an-
ticipated this issue, when (in an unpublished section of his thesis [33]) he sketched
a "mass action approach" to his equilibrium notion, which was rediscovered in the
mid 1990s by Weibull [44]. Additionally, while Nash equilibria are natural "good"
states for noncooperative games, often in cooperative games, such as the one of
the prisoner's dilemma, Nash equilibria are not necessarily the most interesting or
favorable states. For this reason, it would be very desirable for a proper concept of
dynamical game to be able to select either Nash equilibria or other type of steady
states, according to which is more convenient. Evolutionary games are dynamical
processes describing how the distribution of strategies changes in time according
to their individual success.

1.2 Spatially homogeneous replicator dynamics

One of most advocated mechanisms of dynamical choice of strategies is based
on a selection principle, inspired by Darwinian evolution concepts. The main idea
is to reinterpret the probability of picking a certain strategy with the distribution of
a population of players adopting those strategies. The emergence of steady mixed
strategies would be the result of an evolutionary selection: at discrete times play-
ers meet randomly, interact according to their strategies, and obtain a payoff. This
payoff determines how the frequencies in the strategies will evolve. As it is also
observed in [44, 45], the view that games are played over and over again by indi-
viduals who are randomly drawn from large populations was later independently

taken up after Nash by evolutionary biologists, namely Smith and Price [40] and Taylor and Jonker [41]. The discrete time stochastic evolution described above has been formalized in [9] and yields, as an appropriate limit is considered, a continuous time dynamics as follows: In games where players can adopt pure strategies out of a finite set of N choices, we may describe those as $u_1, \ldots, u_N \in U$, where U is the set of strategies. We may denote with σ_i the frequency with which players pick the strategy u_i. The payoff of playing strategy u_i against u_j will be denoted by $J(u_i, u_j)$, where $J \colon U \times U \to \mathbb{R}$. The relative success of the strategy u_i with respect to the strategies played by the population is measured by

$$(1.1) \quad \Delta_N(u_i) = \sum_{j=1}^{N} J(u_i, u_j)\sigma_j - \sum_{\ell=1}^{N}\sum_{j=1}^{N} \sigma_\ell J(u_\ell, u_j)\sigma_j, \quad i = 1, \ldots, N.$$

The relative rate of change of usage of the strategy u_i is then described by

$$\frac{\dot{\sigma}_i}{\sigma_i} = \Delta_N(u_i) = \sum_{j=1}^{N} J(u_i, u_j)\sigma_j - \sum_{\ell=1}^{N}\sum_{j=1}^{N} \sigma_\ell J(u_\ell, u_j)\sigma_j, \quad i = 1, \ldots, N,$$

or

$$(1.2) \quad \dot{\sigma}_i = \left(\sum_{j=1}^{N} J(u_i, u_j)\sigma_j - \sum_{\ell=1}^{N}\sum_{j=1}^{N} \sigma_\ell J(u_\ell, u_j)\sigma_j \right)\sigma_i, \quad i = 1, \ldots, N.$$

The system of ordinary differential equations (1.2) is known as *replicator dynamics*, as introduced first by Taylor and Jonker [41], which is the best known model in the class of imitative adaptation models (see [44, 45]); we also refer to the recent book [25]. It is one of the most popular dynamical game models, because its ω-limit (the set of accumulation points of the dynamics) and steady states are closely related to the Nash equilibria of the game described by the payoff matrix $A = (J(u_i, u_j))_{ij}$ [25, theorem 7.2.1] (the so-called "folk theorem of evolutionary game theory"). Moreover, as discussed in [26], adopting an equilibrium-based viewpoint is often unable to account for the long-term behaviour of realistic players, who adjust their behaviour to maximise their payoff. The replicator dynamics aims at being a more robust model.

1.3 Mean-field replicator dynamics

There is by now a large body of literature addressing the replicator dynamics for infinite or continuous strategies [8, 18, 20, 23, 24, 34–36, 38], which can be viewed as a natural limit for $N \to \infty$ of system (1.2). The way of deducing this limit is by defining the probability measure

$$\sigma_t^N := \sum_{j=1}^{N} \sigma_{j,t}\delta_{u_j} \in \mathscr{P}(U)$$

and its evolution according to

$$\dot{\sigma}_t^N = \left(\int_U J(\cdot, u') \mathrm{d}\sigma_t^N(u') - \int_{U \times U} J(w, u') \mathrm{d}\sigma_t^N(w) \mathrm{d}\sigma_t^N(u') \right) \sigma_t^N,$$

or, in weak form,

$$\frac{\mathrm{d}}{\mathrm{d}t} \int_U \varphi(u) \mathrm{d}\sigma_t^N(u) = \int_U \varphi(u) \left(\int_U J(u, u') \mathrm{d}\sigma_t^N(u') \right.$$
$$\left. - \int_{U \times U} J(w, u') \mathrm{d}\sigma_t^N(w) \mathrm{d}\sigma_t^N(u') \right) \mathrm{d}\sigma_t^N(u)$$

for any $\varphi \in C(U)$. For any $\sigma \in \mathscr{P}(U)$, we may denote

$$\Delta_\sigma(u) := \left(\int_U J(u, u') \mathrm{d}\sigma(u') - \int_{U \times U} J(w, u') \mathrm{d}\sigma(w) \mathrm{d}\sigma(u') \right),$$

so that $\Delta_\sigma(u_i) = \Delta_N(u_i)$ as in (1.1) for $\sigma = \sum_{j=1}^N \sigma_j \delta_{u_j}$. By assuming that the initial conditions $\sigma_0^N \rightharpoonup \bar{\sigma}$ for a given $\bar{\sigma} \in \mathscr{P}(U)$, one can show that $\sigma_t^N \rightharpoonup \sigma_t$ for $N \to \infty$ for any t, where σ is the solution to

$$\frac{\mathrm{d}}{\mathrm{d}t} \int_U \varphi(u) \mathrm{d}\sigma_t(u) = \int_U \varphi(u) \Delta_\sigma(u) \mathrm{d}\sigma_t(u), \quad \sigma(0) = \bar{\sigma}.$$

A result of well-posedness of such an equation for special choices of U is obtained for instance in [35]; see also the more recent work [6]. The notion of equilibrium is well understood in the case of a finite strategy space [25, 45]. On the contrary, the situation of games with infinite strategies has long lacked a general theory due to the technical and conceptual difficulties stemming from understanding which notion of distance between probability measures was the most suitable. We refer the reader to [35], where stability with respect to the weak topology has been considered, and the discussion in [34]. Since the focus of the present paper is not the analysis of the asymptotic behaviour of the replicator dynamics with continuous strategy space, we refer the interested reader to additional and more recent work along the measure theoretical framework, including [10, 16–19]. Our approach uses the classical transport distances and the general framework of evolution problems in the class of probability measures; see, for instance, [2] for a systematic treatment of this topic; see also [13, 37] for recent contributions. Some results of asymptotic behavior and the stability of solutions are given in [6].

1.4 Spatially inhomogeneous replicator dynamics

In this paper, in contrast to spatially homogenous dynamical games, we assume that the population of players is distributed over a position space and that they are each endowed with probability distributions of strategies, which they draw at random to evolve their positions. The positions of the players are assumed to be in the d-dimensional Euclidean space $\mathbb{R}^d$ and the pure strategies u are in a compact metric space U; a probability measure $\sigma \in \mathscr{P}(U)$ denotes a mixed strategy. Notice that this setting is different from the one of spatial evolutionary games [22, 29]

where pure strategies are distributed over a spatial domain and they are updated in time according to the success of neighbor strategies via local interaction.

With these definitions, the space of pairs of positions and mixed strategies is $C := \mathbb{R}^d \times \mathscr{P}(U)$, whose elements are pairs $y = (x, \sigma)$ describing the state of a player. The system will be described by the evolution of a measure $\Sigma \in \mathscr{P}(C) = \mathscr{P}(\mathbb{R}^d \times \mathscr{P}(U))$ on our state space, which represents a distribution of players with strategies. Notice that such a measure Σ well describes the superposition of players with different strategies that at some time occupy the same position.

Omitting for the time being the temporal variable, we proceed to the description of the dynamics of the pair $y := (x, \sigma) \in C$. The player x moves with a velocity that is obtained by averaging over all the strategies a suitable function $e \colon \mathbb{R}^d \times U \to \mathbb{R}^d$, namely,

$$(1.3) \qquad \dot{x} = \int_U e(x, u) \mathrm{d}\sigma(u).$$

For instance, in the simplest case when $e(x, u) = u$ and the set of strategies U consists of a finite subset of $\mathbb{R}^d$, this "mean velocity" dynamics can be thought of as the result of a faster time scale. For convenience, it is useful to see the right-hand side of (1.3) as the result of a map $a \colon C \to \mathbb{R}^d$ by defining

$$(1.4) \qquad a(y) = a(x, \sigma) := \int_U e(x, u) \mathrm{d}\sigma(u).$$

Notice that a depends linearly on σ. In order to write the evolution law for the strategies we will assume that it is driven by an interaction mechanism, which depends only on the state of the system (position and strategies) and does not distinguish two players occupying the same place with the same instantaneous strategy distribution. Thus we consider a Lipschitz function

$$(1.5) \qquad J \colon (\mathbb{R}^d \times U)^2 \to \mathbb{R},$$

and we define an interaction potential

$$\Delta \colon \mathscr{P}(C) \times C \to \mathrm{C}(U), \quad (\Sigma, y) \mapsto \Delta_{\Sigma, y},$$

by setting (notice that the integrals make sense, under suitable moment assumptions on Σ, since J has at most linear growth)

$$
\begin{aligned}
(1.6) \qquad \Delta_{\Sigma, y} &= \Delta_{\Sigma, (x, \sigma)}(u) \\
&:= \int_C \int_U J(x, u, x', u') \mathrm{d}\sigma'(u') \mathrm{d}\Sigma(x', \sigma') \\
&\quad - \int_U \int_C \int_U J(x, w, x', u') \mathrm{d}\sigma'(u') \mathrm{d}\Sigma(x', \sigma') \mathrm{d}\sigma(w).
\end{aligned}
$$

The evolution law for the mixed strategies of the player at x is again according to a replicator dynamics similar to the ones mentioned above, and it can be written as

$$(1.7) \qquad \dot{\sigma} = \Delta_{\Sigma, (x, \sigma)} \sigma.$$

The interaction potential Δ has here a simple interpretation: $J(x, u, x', u')$ represents the contribution to the payoff that the player x gets from the pure strategy u assuming that the player x' acts with pure strategy u'. When player x' acts with mixed strategy σ', we obtain

$$(1.8) \qquad \mathcal{J}(x, u, x', \sigma') := \int_U J(x, u, x', u') d\sigma'(u').$$

Therefore, the full payoff of x given the pure strategy u is

$$\int_C \mathcal{J}(x, u, x', \sigma') d\Sigma(x', \sigma'),$$

corresponding to the first integral in (1.6). The second term is the integral of this quantity with respect to σ; thus the payoff expected by x from the mixed strategy σ, given the full distribution Σ.

We remark that the dependence of the payoff in (1.8) on the full strategy σ' of x', which may imply a certain level of anticipation, is made for the sake of generality: in practical situations $\mathcal{J}$ may only depend on a marginal of σ', and the strength of the interaction between players may be influenced by their distance $|x - x'|$ through the function J. From another point of view, this full dependence (as the expectations of the players on the future in the mean-field games theory; see [14] and the short discussion in Section 1.6) could emerge from a repetition of the evolutionary game; see [9] for a contribution in this direction in the spatially homogeneous case.

Putting together (1.4) and (1.7), we can write the evolution for y as

$$(1.9) \qquad \dot{y} = (\dot{x}, \dot{\sigma}) = (a(x, \sigma), \Delta_{\Sigma,(x,\sigma)}\sigma) =: b_\Sigma(y),$$

which can be interpreted as an ODE in the convex set C. The theory of ODEs in Banach spaces (see Appendices A and B, or [15, partie II] and [21, chap. X] for classical monographs on this topic) will be a useful tool in the study of the well-posedness of (1.9), the existence, uniqueness, and stability of its solutions, and their properties.

In order to do so, we embed C in the Banach space $Y := \mathbb{R}^d \times F(U)$, where

$$F(U) := \overline{\mathrm{span}(\mathscr{P}(U))}^{\|\cdot\|_{\mathrm{BL}}},$$

and the closure is taken in the dual space $(\mathrm{Lip}(U))'$ with respect to the dual norm (also called the bounded Lipschitz norm)

$$(1.10) \qquad \|\ell\|_{\mathrm{BL}} := \sup\{\langle \ell, \varphi \rangle : \varphi \in \mathrm{Lip}(U), \|\varphi\|_{\mathrm{Lip}} \le 1\},$$

where $\|\cdot\|_{\mathrm{Lip}}$ is the Lipschitz norm in (2.5) below. Notice that $F(U) \subset (\mathrm{Lip}(U))'$, and that $(Y, \|\cdot\|_Y)$, with $\|y\|_Y = \|(x, \sigma)\|_Y := |x| + \|\sigma\|_{\mathrm{BL}}$, is a separable Banach space. The space $F(U)$ defined above, known in the literature as the Arens-Eells space [5], is isometric to the predual of $\mathrm{Lip}(U)$; see [3,43] for more details also on the space $(\mathrm{Lip}(U))'$.

1.5 Formal derivation of a nonlinear master equation

We fix a time interval $[0, T]$, a time step $h = T/N$, and an initial datum $\bar{\Sigma} \in \mathscr{P}(\mathbb{R}^d \times \mathscr{P}(U))$, assuming for simplicity that the first marginal of $\bar{\Sigma}$ is compactly supported. Recalling the notation $C = \mathbb{R}^d \times \mathscr{P}(U)$ and its natural structure of convex set, we build a discrete solution

$$M_h \in \mathscr{P}(\mathrm{C}([0, T]; C))$$

concentrated on paths $(x(t), \sigma(t))\colon [0, T] \to C$, which are piecewise affine (in the N intervals). We denote $\Sigma_{t,h} := (\mathrm{ev}_t)_\# M_h$, where $\mathrm{ev}_t\colon \mathrm{C}([0, T]; C) \to C$ defined by $\mathrm{ev}_t(x, \sigma) := (x(t), \sigma(t))$. In particular, $(\mathrm{ev}_0)_\# M_h = \bar{\Sigma}$ is the given initial condition.

The heuristic idea is the following: If the player at time $t = ih$ for $i \in \{0, \dots, N\}$ is in the position $\bar{x}$, with mixed strategy represented by the probability measure $\bar{\sigma}$, first they upgrade their belief on the probability replacing $\bar{\sigma}$ by

$$\bar{\sigma}' := (1 + h\Delta_{\Sigma_{t,h},(\bar{x},\bar{\sigma})})\,\bar{\sigma}.$$

Then, they move to the next position $\bar{x} + he(\bar{x}, u)$ choosing u with probability $\bar{\sigma}'$ and carrying the same probability $\bar{\sigma}'$ to the new position. The probability measure M_h takes all the future stochastic realizations into account. In more formal terms, the conditional probability relative to $M_h|_{[0,t+h]}$ of $(x(t + h), \sigma(t + h))$, given the information that at time t one has $(x, \sigma)(t) = (\bar{x}, \bar{\sigma})$, is $((\bar{x} + he(\bar{x}, \cdot))_\#\bar{\sigma}') \times \delta_{\bar{\sigma}'}$. By iterating this process N times one can build M_h on the whole time interval $[0, T]$ (alternatively, one can view this as a Markov process with the above-defined transition probabilities and build first a measure N_h in $\mathscr{P}(C)^{N+1}$; then M_h by associating to $N + 1$ points in $\mathscr{P}(C)$ a piecewise affine path in $[0, T]$).

Under boundedness assumptions on the field $e\colon \mathbb{R}^d \times U \to \mathbb{R}^d$ and on the interaction potential Δ, it turns out that M_h is concentrated on paths $(x(t), \sigma(t))$ satisfying the equi-Lipschitz property

$$(1.11) \qquad \begin{cases} |x(s) - x(t)| \le L|s - t| \\ \|\sigma(s) - \sigma(t)\|_{\mathrm{BL}} \le L|s - t| \end{cases} \qquad \text{for all } 0 \le s \le t \le T,$$

with $L = L(e, J, U) \ge 0$.

Given now a bounded test function $\Phi\colon Y \to \mathbb{R}$ of class C^1 (in the Fréchet sense) with respect to the $\|\cdot\|_Y$ norm, let us write a discrete continuity equation associated to $\Sigma_{t,h}$, the marginals of M_h in the sense $\Sigma_{t,h} := (\mathrm{ev}_t)_\# M_h$. For

$t = ih, 0 \leq i \leq N - 1$, one has

$$\int \Phi(x, \sigma) \, d(\Sigma_{t+h,h} - \Sigma_{t,h})(x, \sigma)$$

$$= \int \big[\Phi(x(t+h), \sigma(t+h)) - \Phi(x(t), \sigma(t))\big] dM_h(x(\cdot), \sigma(\cdot)) =$$

$$= \int \big[\Phi(x(t+h), (1 + h\Delta_{\Sigma_{t,h}, x(t), (\sigma(t))})\sigma(t))$$

$$- \Phi(x(t), \sigma(t))\big] dM_h(x(\cdot), \sigma(\cdot))$$

$$\sim h \int D\Phi(x(t), \sigma(t)) \cdot \big(a(x(t), \sigma(t)), \Delta_{\Sigma_{t,h}, (x(t), \sigma(t))}\sigma(t)\big) dM_h(x(\cdot), \sigma(\cdot))$$

$$= h \int D\Phi(x, \sigma) \cdot b_{\Sigma_{t,h}}(x, \sigma) d\Sigma_{t,h}(x, \sigma),$$

where b_Σ is given by (1.9), and we used the chain rule for Fréchet differentiation.

Recalling from (1.11) that M_h is concentrated on equi-Lipschitz paths $t \mapsto (x(t), \sigma(t))$, the family $\{M_h: h = T/N, N \geq 1\}$ is weakly compact in the space $\mathscr{P}(C([0, T]; C))$ and by Prokhorov theorem it has limit points as $h \to 0$. Any limit point M is concentrated on Lipschitz paths satisfying (1.11), and this construction builds a continuous map $\Sigma: [0, T] \to C$ by $\Sigma_t := (ev_t)_\# M$, satisfying the equation

$$(1.12) \qquad\qquad \partial_t \Sigma_t + \mathrm{div}(b_{\Sigma_t} \Sigma_t) = 0$$

in the weak sense of (3.10) below, with $\bar{\Sigma}$ as initial condition. In the following, we explore notions of solutions to (1.12) and conditions for their existence and uniqueness, starting from the corresponding finite agents model.

1.6 Interpretation of the model, comparison with other population dynamics models, and mean-field games

The master equation (1.12) is a novel model of spatially nonhomogenous evolutive games, which fuses mean-field theory, optimal transport, and replicator dynamics from evolutionary game theory. We borrow and translate the mechanism of the replicator dynamics of computing optimal mixed strategies of individual agents involved in a spatial-type dynamics; in this sense our model is about a "dynamical/evolutive game" rather than a "selection/evolutionary game" of a population. The Darwinism/selection is not at the level of the population of agents, but at the level of the best strategy applied by the single agent, and it is certainly more in the spirit of "mean-field games" *à la* Lasry-Lions (see the discussion at the end of this subsection), rather than a genuine contribution to evolutionary game theory, for which population evolution/selection is the main focus. However, we think that our model could also have potentially interesting applications for evolutionary population dynamics, even if this was not our motivating viewpoint. We explain

this latter interpretation in Section 3.1, where an evolutionary game is augmented by external parameters.

In the works [12, 13], a population dynamics is considered in terms of (possibly nonlinear) transport equations modeling agents subjected to predetermined (economic or social) forces induced by the other agents (the first principles of economic or social interaction). Such models are very much inspired by physics where particles are subjected to the action of physical forces governed by first principles. The applicability of this type of models to economics or sociology raises several questions, since in economics one assumes that every external intervention must necessarily change the behavior of agents as they anticipate the effect. So the dynamics with predefined interaction forces between agents does not encounter a full acceptance by economists or sociologists. Instead, the model (1.12) where the driving vector field evolves on the basis of the replicator equation is better adapted to an economical viewpoint. It generates a more "complex" dynamics, due to the self-interacting influence of the replicator equation on the motion, and it is closely connected to the underlying principle of maximizing the individual payoff.

There are other approaches towards modeling spatially nonhomogenous games of a large population of indistinguishable agents. Perhaps the most prominent is the so-called theory of *mean-field games*. This class of problems was considered in the economics literature by Boyan Jovanovic and Robert W. Rosenthal [30], in the engineering literature by Peter E. Caines, Minyi Huang, and Roland P. Malhamé [27, 28], and independently and around the same time by the mathematicians Jean-Michel Lasry and Pierre-Louis Lions [31].

In continuous time, a mean-field game is typically composed of a Hamilton-Jacobi-Bellman equation for the optimal control problem of an individual, and a forward Fokker-Planck-Kolmogorov equation for the dynamics of the aggregate distribution of agents. Under fairly general assumptions, it can be proven that a mean-field game is the limit as $N \to \infty$ of an N-player Nash equilibrium. In particular, one can consider the stochastic evolution of N players dictated by equations

$$(1.13) \qquad \mathrm{d}X_t^i = u_t^i \mathrm{d}t + \sqrt{2\nu}\mathrm{d}B_t^i, \quad i = 1, \dots, N,$$

where X_0^i are independently drawn at random according to a probability distribution $\mu_0 \in \mathscr{P}(\mathbb{R}^d)$, B_t^i are independent d-dimensional Brownian motions, and $\nu > 0$ is a noise parameter, so that the limiting case $\nu = 0$ corresponds to the so-called first-order games. The player i can choose a strategy u^i adapted to the filtration

$$(1.14) \qquad \mathcal{F}_t = \sigma\big(X_0^j, B_s^j : s \le t, \, j = 1, \dots, N\big).$$

The payoff of player i is given by

$$(1.15) \qquad J_i^N(u^1, \dots, u^N) = \mathbb{E}\left[\int_0^T \frac{1}{2}|u_t^i|^2 + F\left(X_t^i, \frac{1}{N-1}\sum_{j \ne i}\delta_{X_t^j}\right)\mathrm{d}t\right].$$

The solution of the N-player game is the suitable minimization of (1.15) under the constraints (1.13) for all $i = 1, \ldots, N$. Hence, the appropriate notion of solution may be precisely the Nash equilibrium, i.e., a configuration of strategies $(u^{*,1}, \ldots, u^{*,N})$ such that

$$J_i^N(u^{*,1}, \ldots, u^{*,N}) \le J_i^N(u^{*,1}, \ldots, u^{*,i-1}, u, u^{*,i+1}, \ldots, u^{*,N}) \quad \text{for all } u,$$

for all $i = 1, \ldots, N$. The computation of such equilibria becomes intractable already for a moderate number N of players. However, for N very large a mean-field approximation for $N \to \infty$ may help to obtain approximate Nash equilibria. In particular, for $N \to \infty$ one can approximate the empirical distribution $\mu_t^N = \frac{1}{N} \sum_{i=1}^{N} \delta_{X_t^i}$ supported on realizations of (1.13) by the solution of the forward Fokker-Planck-Kolmogorov equation

$$(1.16) \qquad \partial_t \mu - \nu \Delta \mu + \operatorname{div}(u\mu) = 0$$

for an appropriate control function $u(t, x)$. For a given time-dependent distribution $\widehat{\mu}_t$, the corresponding payoff functional would be given by

$$(1.17) \qquad J(u) = \int_0^T \int_{\mathbb{R}^d} \left(\frac{1}{2} |u(t, x)|^2 + F(x, \widehat{\mu}_t) \right) \mathrm{d}\mu_t \, \mathrm{d}t.$$

The solution of the mean-field game comes from the minimization of (1.17) under the PDE constraints (1.16) and the fixed point condition $\mu_t = \widehat{\mu}_t$, and these yield the mean-field game system

$$(1.18) \qquad \begin{cases} -\partial_t \psi - \nu \Delta \psi + \frac{1}{2} |\nabla \psi|^2 = F(x, \mu), \\ \partial_t \mu - \nu \Delta \mu - \operatorname{div}(\nabla \psi \mu) = 0, \end{cases}$$

with appropriate initial and terminal conditions. The choice $u(t, x) = -\nabla \psi(t, x)$ yields nearly optimal strategies to be inserted in (1.13) by posing $u_t^i = u(t, X_t^i)$. Several comments about differences between our spatially nonhomogeneous evolutionary game and mean-field games are in order:

- Mean-field games are intimately linked to the theory of (stochastic) mean-field control and as such are global-in-time optimizations: in other words, one has to forecast the future behavior of the agents in order to be able to solve the game, and this is mathematically expressed by the need of simultaneously solving a backward-in-time evolution to compute the optimal strategy. Our model is a simple forward-in-time evolution according to the master equation (1.12) in a continuous local-in-time search for pairwise-game equilibria.
- Mean-field games are built around the concept of Nash equilibrium for noncooperative games, one of the notion of solutions considered in evolutionary game theory, which also aims at reaching different situations.
- Well-posedness of the mean-field game system (1.18) has been shown for special choices of costs F, the so-called potential games [31]. Accordingly, the numerical solution of the mean-field game system (see also,

e.g., [1] for alternative approaches) could be based on the iterative solution of a backward-forward system by means of individual solvers for the two equations: starting from a given initial trajectory $t \mapsto \mu_t^{(0)}$, one iterates the numerical solution for $n = 0, 1, 2, \ldots$

$$\begin{cases} -\partial_t \psi^{(n)} - \nu \Delta \psi^{(n)} + \frac{1}{2} |\nabla \psi^{(n)}|^2 = F(x, \bar{\mu}^{(n)}), \\ \partial_t \mu^{(n+1)} - \nu \Delta \mu^{(n+1)} - \text{div}(\nabla \psi^{(n)} \mu^{(n+1)}) = 0, \end{cases}$$

where $\bar{\mu}^{(n)} = \frac{1}{n} \sum_{i=1}^{n} \mu^{(i)}$, and with appropriate initial and terminal conditions. This iterative procedure, also called *learning* or *fictitious play* in [14], provides a justification to the "knowledge about the future" incorporated in the model, and is shown to converge, at least for potential games. For our model, we will prove that the well-posedness (existence, stability, and uniqueness) of the master equation (1.12) is ensured under Lipschitz assumptions on the function e in (1.4) and on our pairwise-game payoff function J defined in (1.5). In addition, the formal derivation of the master equation (1.12) as in Section 1.5 provides already a rather clear path towards a time marching numerical solution.

2 Measure-Theoretic Preliminaries

Before exploring notions of solutions to (1.12) and conditions for their existence and uniqueness, we make more precise the functional setting where the evolutions governed by (1.12) take place. We use here differential and Bochner calculus in separable Banach spaces, and we refer to Appendix A for some related basic notions and results.

2.1 Notation and distances in the space of measures

If (X, d_X) is a metric space, we denote by $\mathscr{M}(X)$ the space of signed Borel measures in X with finite total variation, by $\mathscr{M}_+(X)$, $\mathscr{P}(X)$ the convex subsets of nonnegative measures and probability measures respectively. For $\sigma \in \mathscr{M}(X)$, $|\sigma| \in \mathscr{M}_+(X)$ denotes the total variation measure of σ (see also (2.4) below). We shall also use the notation $\mathscr{M}_0(X)$ for the subset of measures with 0 mean; we shall use the identity $\mathbb{R}(\mathscr{P}(X) - \mathscr{P}(X)) = \mathscr{M}_0(X)$ (in the sense of Minkowski sums) provided by the Hahn decomposition theorem for signed measures.

For a Lipschitz function $f \colon X \to \mathbb{R}$ we denote by

$$\text{Lip}(f) := \sup_{\substack{x,y \in X \\ x \neq y}} \frac{|f(x) - f(y)|}{\mathsf{d}_X(x, y)}$$

the Lipschitz constant, and denote by $\text{Lip}_b(X)$ the space of bounded Lipschitz functions.

Given $\mu \in \mathscr{M}_+(X)$ and $f \colon X \to Y$, with f μ-measurable, we shall denote by $f_\# \mu \in \mathscr{M}_+(Y)$ the push-forward measure, having the same mass as μ and defined by $f_\# \mu(B) = \mu(f^{-1}(B))$ for any Borel set $B \subset Y$ (when $\mu \in \mathscr{P}(X)$, in

probability theory it is also named law of f under μ); we shall also often use the change-of-variables formula

$$\int_Y g \, \mathrm{d} f_\# \mu = \int_X g \circ f \, \mathrm{d}\mu$$

whenever either one of the integrals makes sense.

In a complete and separable metric space (X, d_X), we shall use the Kantorovich-Rubinstein (possibly infinite) distance $W_1(\mu, \nu)$ in the class $\mathscr{P}(X)$; thanks to Kantorovich duality, the definition

$$W_1(\mu, \nu) := \sup\left\{\int_X \varphi \, \mathrm{d}\mu - \int_X \varphi \, \mathrm{d}\nu \colon \varphi \in \mathrm{Lip}_b(X), \mathrm{Lip}(\varphi) \leq 1\right\}$$

is equivalent to

(2.1) $W_1(\mu, \nu) :=$

$$\inf\left\{\int_{X \times X} \mathsf{d}_X(x, y) \, \mathrm{d}\Pi(x, y) \colon \Pi(A \times X) = \mu(A), \ \Pi(X \times B) = \nu(B)\right\}$$

involving couplings Π of μ and ν, but we shall mostly be working with the first one. Notice that $W_1(\mu, \nu)$ is finite if μ and ν belong to the space

(2.2) $$\mathscr{P}_1(X) := \left\{\mu \in \mathscr{P}(X) \colon \int_X \mathsf{d}_X(x, \bar{x}) \mathrm{d}\mu(x) < +\infty \text{ for some } \bar{x} \in X\right\}$$

and that $(\mathscr{P}_1(X), W_1)$ is complete if (X, d_X) is complete. Recall also that the convergence of $\mu_h \in \mathscr{P}_1(X)$ to $\mu \in \mathscr{P}_1(X)$ with respect to the distance W_1 is equivalent to weak convergence in the duality with bounded Lipschitz functions plus convergence of first moments [2], namely, $\int_X \mathsf{d}_X(\cdot, \bar{x}) \mathrm{d}\mu_h \to \int_X \mathsf{d}_X(\cdot, \bar{x}) \mathrm{d}\mu$ for all $\bar{x} \in X$.

We also need a C_b^1 variant of Kantorovich duality, valid in separable Banach spaces, stated below.

LEMMA 2.1 (C^1 duality). *For any separable Banach space Y, one has that*

$$\mathsf{d}_{\mathrm{C}^1}(\mu, \nu) := \sup_{\substack{\phi \in \mathrm{C}_b^1(Y) \\ \mathrm{Lip}(\phi) \leq 1}} \left(\int_Y \phi \, \mathrm{d}\mu - \int_Y \phi \, \mathrm{d}\nu\right)$$

defines a distance in $\mathscr{M}(Y)$, which coincides with W_1 when restricted to $\mathscr{P}_1(Y)$.

PROOF. Symmetry and the triangle inequality are obvious. To prove the nondegeneracy, notice that any cylindrical function $\phi(y) = \psi(\langle y, z_1'\rangle, \ldots, \langle y, z_d'\rangle)$ with $\psi \colon \mathbb{R}^d \to \mathbb{R}$ bounded, continuously differentiable, and Lipschitz and $z_j' \in Y'$, is a C_b^1 function on Y and, after suitable rescaling, satisfies $\mathrm{Lip}(\phi) \leq 1$. If two measures μ and ν have vanishing $\mathsf{d}_{\mathrm{C}^1}$-distance, then necessarily they must coincide on the σ-algebra generated by such cylindrical functions. From [46, lemma on

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pp. 131–132], since Y is separable there exists a sequence $(z_i')_{i \in \mathbb{N}}$ in the unit ball of Y' with

$$\tag{2.3} \|y\|_Y = \sup_i \langle y, z_i' \rangle = \sup_{i,k} \psi_k(\langle y, z_i' \rangle) \quad \text{for all } y \in Y,$$

where $\psi_k(r) = k \tanh(r/k)$.

Hence, balls and then open sets belong to the σ-algebra generated by cylindrical functions. It follows that this σ-algebra coincides with the Borel σ-algebra and $\mu = \nu$.

The proof of the final statement requires a slight refinement of the previous argument. Let $\mu, \nu \in \mathscr{P}_1(Y)$ and for every $a > 0$ let us set $\mathcal{L}_a := \{\phi \in \mathrm{Lip}_b(Y) : \sup_Y |\phi| \leq a, \mathrm{Lip}(\phi) \leq 1\}$. We consider the subset $\mathcal{G}_a$ of functions $\phi \in \mathcal{L}_a$ such that

$$\int_Y \phi \, \mathrm{d}\mu - \int_Y \phi \, \mathrm{d}\nu \leq \mathsf{d}_{\mathrm{C}^1}(\mu, \nu).$$

By definition of $\mathsf{d}_{\mathrm{C}^1}$, $\mathcal{G}_a$ contains $\mathcal{L}_a \cap \mathrm{C}_b^1(Y)$; we want to show that $\mathcal{G}_a$ coincides with $\mathcal{L}_a$. Since $\mathcal{G}_a$ is closed with respect to pointwise convergence, it is sufficient to prove that $\widetilde{\mathcal{G}}_a := \overline{\mathcal{L}_a \cap \mathrm{C}_b^1(Y)}$ coincides with $\mathcal{L}_a$. Notice that the topology of pointwise convergence in $\mathcal{L}_a$ is metrizable: it is sufficient to select a countable dense subset $(y_i)_{i \in \mathbb{N}}$ of Y and consider the distance

$$\mathsf{d}_{\mathcal{L}_a}(\phi, \psi) := \sum_{i=0}^{\infty} 2^{-i} |\phi(y_i) - \psi(y_i)|.$$

By approximating the convex function

$$S_N(r) := r_1 \vee r_2 \vee \cdots \vee r_N, \quad r = (r_1, \ldots, r_N) \in \mathbb{R}^N,$$

by convolution with a symmetric, nonnegative mollifier $\kappa \in C_c^\infty(\mathbb{R}^N)$, $\int \kappa = 1$,

$$S_{N,h}(r) := h^N \int_{\mathbb{R}^N} S_N(r - s) \kappa(hs) \mathrm{d}s, \quad h \in \mathbb{N},$$

we obtain an increasing sequence of smooth functions converging to S_N as $h \to \infty$ and satisfying

$$I_N(r) \leq S_{N,h}(r) \leq S_N(r), \quad |S_{N,h}(r) - S_{N,h}(r')| \leq \sup_{1 \leq n \leq N} |r_n - r_n'|,$$

where $I_N(r) := -S_N(-r) = r_1 \wedge r_2 \wedge \cdots \wedge r_N$. It follows that if $\psi_1, \ldots, \psi_N \in \widetilde{\mathcal{G}}_a$, then $\psi_h := S_{N,h}(\psi_1, \ldots, \psi_N)$ belongs to $\widetilde{\mathcal{G}}_a$ and therefore also $\psi = \psi_1 \vee \psi_2 \vee \cdots \vee \psi_N$ belongs to $\widetilde{\mathcal{G}}_a$. The same property holds for the infimum and extends to a countable family of functions.

By using the representation (2.3) it is then easy to check that any functions of the form $y \mapsto -a \vee (b + \|y - y_0\|) \wedge a$ belong to $\widetilde{\mathcal{G}}_a$. Since every $f \in \mathcal{L}_a$ can be expressed as

$$f(y) = \inf_{y' \in Y_0} -a \vee (f(y') + \|y - y'\|) \wedge a \quad \forall y \in Y$$

for any dense and countable subset Y_0 of Y, we conclude the proof. $\qquad\square$

From now on we fix a compact metric space U with a distance d. The space $\mathscr{M}(U)$, when endowed with the total variation norm

$$\|\sigma\|_{\mathrm{TV}} := \sup\left\{\int_U \varphi\, \mathrm{d}\sigma : \varphi \in \mathrm{C}(U), |\varphi| \le 1\right\}$$

has the structure of Banach space, isometrically isomorphic to the dual of $\mathrm{C}(U)$. We will also use the representation formulas

$$(2.4) \qquad \|\sigma\|_{\mathrm{TV}} = |\sigma|(U) = \sup\left\{\int_U \varphi\, \mathrm{d}\sigma : \varphi \in \mathrm{B}_b(U), |\varphi| \le 1\right\},$$

where $\mathrm{B}_b(U)$ denotes the class of bounded Borel functions $\varphi\colon U \to \mathbb{R}$.

Mixed strategies can be interpreted as Borel probability measures in U. The set of Borel probability measures $\mathscr{P}(U)$ is a convex and weak-star closed subset of $\mathscr{M}(U)$, corresponding to the class of nondecreasing linear functionals $L\colon \mathrm{C}(U) \to \mathbb{R}$ with $L(1) = 1$ in the dual representation of $\mathscr{M}(U)$. On the other hand, $\mathscr{P}(U)$ can also be realized as a compact convex set of another Banach space, the dual $(\mathrm{Lip}(U))'$ of $\mathrm{Lip}(U)$: the norm in $\mathrm{Lip}(U)$ is given by

$$(2.5) \qquad \|\varphi\|_{\mathrm{Lip}} := \sup_{u \in U} |\varphi(u)| + \mathrm{Lip}(\varphi)$$

so that, for $\ell \in (\mathrm{Lip}(U))'$,

$$\|\ell\|_{\mathrm{BL}} := \sup\left\{\langle \ell, \varphi\rangle : \varphi \in \mathrm{Lip}(U), \|\varphi\|_{\mathrm{Lip}} \le 1\right\}.$$

For later use, let us record the property

$$(2.6) \qquad \|\varphi\psi\|_{\mathrm{Lip}} \le \|\varphi\|_{\mathrm{Lip}}\|\psi\|_{\mathrm{Lip}},$$

for all $\varphi, \psi \in \mathrm{Lip}(U)$.

Actually, while the linear structure we need forces us to go beyond $\mathscr{P}(U)$, for our purposes it will be more convenient to work in a closed subspace of $(\mathrm{Lip}(U))'$, namely,

$$F(U) := \overline{\mathrm{span}(\mathscr{P}(U))}^{\|\cdot\|_{\mathrm{BL}}} \subset (\mathrm{Lip}(U))'.$$

This space, also called Arens-Eells space in the literature, is a separable Banach space containing $\mathscr{M}(U)$.

Notice that, for a measure $\nu \in \mathscr{M}_0(U)$, the BL norm is equivalent to the norm induced by the dual formulation of the 1-Wasserstein distance: in fact, for every 1-Lipschitz function $\varphi\colon U \to \mathbb{R}$ and for every $x_0 \in U$ we have

$$\int_U \varphi\, \mathrm{d}\nu = \int_U (\varphi(x) - \varphi(x_0))\mathrm{d}\nu(x) \le \|\nu\|_{\mathrm{BL}}\big(\sup_{x \in U} \mathsf{d}(x, x_0) + 1\big),$$

so that, with $D_U := \min_{x_0 \in U} \max_{x_1 \in U} \mathsf{d}(x_0, x_1) \le \mathrm{diam}(U)$, one has

$$\|\nu\|_{\mathrm{BL}} \le \sup\left\{\int_U \varphi\, \mathrm{d}\nu : \mathrm{Lip}(\varphi) \le 1\right\} \le (1 + D_U)\|\nu\|_{\mathrm{BL}}.$$

In particular, when $\nu = \mu_1 - \mu_2$ with $\mu_1, \mu_2 \in \mathscr{P}(U)$, Kantorovich duality gives

$$(2.7) \qquad \|\mu_1 - \mu_2\|_{\mathrm{BL}} \leq W_1(\mu_1, \mu_2) \leq (1 + D_U)\|\mu_1 - \mu_2\|_{\mathrm{BL}}.$$

We summarize the previous discussion in the following list of properties:

(i) $(\mathscr{M}(U), \|\cdot\|_{\mathrm{TV}})$ is a Banach space continuously embedded in the space $(F(U), \|\cdot\|_{\mathrm{BL}})$ with

$$(2.8) \qquad \|\nu\|_{\mathrm{BL}} \leq \|\nu\|_{\mathrm{TV}} \quad \text{for every } \nu \in \mathscr{M}(U).$$

(ii) $\mathscr{P}(U)$ is a weak-star and closed convex set of $\mathscr{M}(U)$, endowed with the total variation norm; it is also compact in $F(U)$. Thanks to (2.7) and to the fact that W_1 metrizes the weak-star convergence in $\mathscr{P}(U)$, the BL norm induces the weak-star topology in $\mathscr{P}(U)$.

(iii) For every $\mu_1, \mu_2 \in \mathscr{P}(U)$, besides (2.7) and (2.8) with $\nu = \mu_1 - \mu_2$, we have

$$(2.9) \qquad W_1(\mu_1, \mu_2) \leq D_U \|\mu_1 - \mu_2\|_{\mathrm{TV}}.$$

2.2 Differentiable curves in the space of measures

Let us now consider two curves $t \mapsto \sigma_t \in \mathscr{P}(U)$, and $t \mapsto \nu_t \in F(U)$, $t \in [0, T]$. We will assume that σ and ν are continuous with respect to the BL norm, and we want to give a meaning to the differential equation

$$(2.10) \qquad \frac{\mathrm{d}}{\mathrm{d}t}\sigma_t = \nu_t \quad t \in [0, T].$$

It is easy to check that the classical formulation of (2.10) as an ODE in the Banach space $F(U)$ is equivalent to the weak formulation of (2.10), which reads as

$$(2.11) \quad \int_U \varphi \, \mathrm{d}\sigma_t - \int_U \varphi \, \mathrm{d}\sigma_s = \int_s^t \langle \nu_\tau, \varphi \rangle \, \mathrm{d}\tau \quad \text{for every } \varphi \in \mathrm{Lip}(U), \ s, t \in [0, T].$$

Indeed, since ν is continuous, the map

$$(2.12) \qquad t \mapsto N_t = \int_0^t \nu_\tau \, \mathrm{d}\tau$$

is of class C^1 and its derivative exists in the classical sense

$$(2.13) \qquad \lim_{h \to 0} \left\| \frac{1}{h}(N_{t+h} - N_t) - \nu_t \right\|_{\mathrm{BL}} = 0 \quad \text{for every } t \in [0, T].$$

From (2.11) it follows that

$$\int_U \varphi \, \mathrm{d}\sigma_t = \int_U \varphi \, \mathrm{d}\sigma_0 + \langle N_t, \varphi \rangle \quad \text{for every } \varphi \in \mathrm{Lip}(U),$$

and therefore the density of $\mathrm{Lip}(U)$ in $\mathrm{C}(U)$ implies $\sigma_t = \sigma_0 + N_t$ for all $t \in [0, T]$. In particular, (2.13) gives that $t \mapsto \sigma_t \in F(U)$ is of class C^1 and that (2.10) holds in the classical sense.

Remark 2.2 (Vector integral). Since ν is continuous, the integral in (2.12) can be equivalently defined as a Cauchy-Riemann integral or as a Bochner integral (see Section A.2).

Let us now suppose that ν takes its values in the smaller space $\mathscr{M}(U) \subset F(U)$ and that the stronger condition

$$\int_0^T \|\nu_t\|_{\mathrm{TV}} \, \mathrm{d}t < +\infty$$

is satisfied. Notice that the TV norm is lower-semicontinuous with respect to the BL topology, so that the map $t \mapsto \|\nu_t\|_{\mathrm{TV}}$ is lower-semicontinuous and therefore Borel. From the representation formula of σ_t, it follows that

$$\|\sigma_t - \sigma_s\|_{\mathrm{TV}} = \|N_t - N_s\|_{\mathrm{TV}} \leq \int_s^t \|\nu_\tau\|_{\mathrm{TV}} \, \mathrm{d}\tau \quad \text{for all } 0 \leq s \leq t \leq T,$$

so that $t \mapsto \sigma_t \in \mathscr{M}(U)$ is absolutely continuous. In particular, if $\|\nu_t\|_{\mathrm{TV}} \in L^\infty(0, T)$, then σ_t is a Lipschitz curve. Eventually, if ν is even continuous with respect to the total variation norm, then $t \mapsto \sigma_t \in \mathscr{M}(U)$ is of class C^1, namely,

$$\lim_{h \to 0} \left\| \frac{1}{h}(\sigma_{t+h} - \sigma_t) - \nu_t \right\|_{TV} = 0 \quad \text{for every } t \in [0, T].$$

2.3 ODEs in the space of measures

Let us now consider the case when the right-hand side ν_t in (2.10) is given by a time-dependent family of operators $A(t, \sigma) \colon [0, T] \times \mathscr{P}(U) \to F(U)$, not necessarily linear with respect to σ; we assume that A is a continuous map when both $\mathscr{P}(U)$ and $F(U)$ are endowed with the BL topologies. Then, we are considering the ODE

$$(2.14) \qquad \frac{\mathrm{d}}{\mathrm{d}t}\sigma_t = A(t, \sigma_t) \quad t \in [0, T],$$

in $F(U)$. Initially, one can also look for solutions $\sigma \in C([0, T]; (\mathscr{P}(U), \|\cdot\|_{\mathrm{BL}}))$ in the weak sense of (2.11), namely,

$$\int_U \varphi \, \mathrm{d}\sigma_t - \int_U \varphi \, \mathrm{d}\sigma_s = \int_s^t \langle A(\tau, \sigma_\tau), \varphi \rangle \, \mathrm{d}\tau,$$
$$\text{for every } \varphi \in \mathrm{Lip}(U), \ 0 \leq s \leq t \leq T.$$

But, since A is continuous, we deduce from the previous discussion that $t \mapsto \sigma_t \in F(U)$ is of class C^1; therefore (2.14) holds in the classical pointwise sense at every $t \in [0, T]$:

$$(2.15) \qquad \lim_{h \to 0} \left\| \frac{1}{h}(\sigma_{t+h} - \sigma_t) - A(t, \sigma_t) \right\|_{\mathrm{BL}} = 0 \quad \text{for all } t \in [0, T].$$

If moreover A maps $[0, T] \times \mathscr{P}(U)$ to the smaller space $\mathscr{M}(U)$ and we know that

$$\sup_{t \in [0,T]} \|A(t, \sigma_t)\|_{\mathrm{TV}} < +\infty,$$

then we deduce that the curve $t \mapsto \sigma_t$ is Lipschitz with respect to the total variation norm. Finally, if $A \colon [0, T] \times \mathscr{P}(U) \to \mathscr{M}(U)$ is continuous when the target space is endowed with respect to the total variation norm, (2.15) improves to

$$\lim_{h \to 0} \left\| \frac{1}{h} (\sigma_{t+h} - \sigma_t) - A(t, \sigma_t) \right\|_{\mathrm{TV}} = 0 \quad \text{for all } t \in [0, T].$$

3 Finite Agent Model

3.1 1-average-player evolution

In order to explain how to apply the previous discussion to our model, let us first consider the simplest, and somehow degenerate, case of a single player. Recalling the formal derivation of our model in Section 1.5, this case should be considered as a limit of an evolution process, when all the players (whose total mass is conventionally normalized to 1) are initially concentrated in the same initial place $\bar{x}$ with the same initial distribution $\bar{\sigma}$ of strategies.

The evolution is then given by a moving point $y_t = (x_t, \sigma_t) \in \mathbb{R}^d \times \mathscr{P}(U)$, and Σ_t is just the Dirac mass concentrated at y_t. Then y satisfies the differential equation

$$(3.1) \qquad \begin{cases} \dot{x}_t = a(x_t, \sigma_t) = \displaystyle\int_U e(x_t, u) \mathrm{d}\sigma_t(u), \\[2mm] \dot{\sigma}_t = \left(\displaystyle\int_U J(x_t, \cdot, x_t, u') \mathrm{d}\sigma_t(u') \right. \\[4mm] \qquad\qquad \left. - \displaystyle\int_{U \times U} J(x_t, w, x_t, u') \mathrm{d}\sigma_t(u') \mathrm{d}\sigma_t(w) \right) \sigma_t, \end{cases}$$

where we denoted by $e \colon \mathbb{R}^d \times U \to \mathbb{R}^d$ the velocity field driving the motion of the player according to the choice of a strategy u. This can be interpreted as a differential equation in the phase space $C = \mathbb{R}^d \times \mathscr{P}(U) \subset Y = \mathbb{R}^d \times F(U)$ of the form

$$(3.2) \qquad\qquad\qquad \dot{y}_t = b(y_t)$$

where b is the vector field

$$b(x, \sigma) = \left(\int_U e(x, u) \mathrm{d}\sigma(u), \left(\int_U J(x, \cdot, x, u') \mathrm{d}\sigma(u') \right. \right.$$
$$\left. \left. - \int_{U \times U} J(x, w, x, u') \mathrm{d}\sigma(u') \mathrm{d}\sigma(w) \right) \sigma \right).$$

The 1-average-player evolution thus reproduces the mechanism of a replicator equation influenced by a vector parameter x, whose dynamics is in turn affected by the evolving strategy distribution. As we will see in the next more general cases,

the particular structure of the vector field b satisfies the structural assumptions of the main existence Theorem B.1, and the second component σ_t of the curve y_t will be differentiable even with respect to the total variation norm.

Remark 3.1. We already observed that our paper deals with dynamical/evolutive processes of populations more in the spirit of the mean-field games à la Lasry-Lions (see Section 1.6) rather than selection/evolutionary processes of populations. The Darwinian selection aspect is occurring in the selection of the mixed strategies of each individual. Nevertheless, the dynamics (3.1)–(3.2) admits a different interpretation in terms of a usual replicator dynamics and evolutionary game theory, when the evolution law of the population also depends on a parameter $x \in \mathbb{R}$ ($d = 1$, just to consider the scalar case where x may represents some kind of available resources). Suppose that the distribution of x of the initial population is concentrated on a single given value x_1. It is typical of the replicator dynamics that the evolution will remain concentrated on x_1 for all the time, so that it would be useless to enrich the state space U to $U \times \{x_1\}$, and x_1 enters only as a parameter of the equations.

Instead, the transport effects of the first equation in (3.1) may allow a variation of x with respect to time modeled by a simple ODE, affected by the distribution $\sigma_t \in \mathcal{P}(U)$ of the population itself. This transport effect becomes more relevant when the initial distribution of x is concentrated on N values: this is one of the possible interpretations of the case of the N-player systems of the next section 3.2.

In the present paper, we focus our efforts on the correct way of modeling the interaction between transport and replicator dynamics, by choosing a simple transport dynamic for the $\mathbb{R}^d$-component of the evolution, but more general dynamical systems can also be considered in our setting, see, e.g., [32]. We think that the flexibility offered by the present approach could be relevant for interesting and novel evolutionary game models.

3.2 N-average-player system

In the case of N players, we have to follow the evolution of N points $y_{i,t} = (x_{i,t}, \sigma_{i,t})$, $i = 1, \ldots, N$. It is useful to introduce the interaction field $f \colon C \times C \to Y$ between two players (as usual, a player is identified by the position x and a mixed strategy σ): we can write f as a pair (f_x, f_σ) where $f_x \colon C \times C \to \mathbb{R}^d$ and $f_\sigma \colon C \times C \to F(U)$; the first component is in fact independent of the interaction and can be written as

$$(3.3) \qquad f_x(y, y') = f_x(x, \sigma, x', \sigma') = a(x, \sigma) = \int_U e(x, u) \mathrm{d}\sigma(u),$$

whereas the second component f_σ is given by

$$
\begin{aligned}
f_\sigma(y, y') &= f_\sigma(x, \sigma, x', \sigma') \\
&:= \left(\int_U J(x, \cdot, x', u') \mathrm{d}\sigma'(u') \right. \\
&\quad \left. - \int_{U \times U} J(x, w, x', u') \mathrm{d}\sigma'(u') \mathrm{d}\sigma(w) \right) \sigma.
\end{aligned}
$$

(3.4)

The system

$$
\begin{cases}
\dot{x}_{i,t} & = a(x_{i,t}, \sigma_{i,t}) = \int_U e(x_{i,t}, u)\mathrm{d}\sigma_{i,t}(u), \\
\dot{\sigma}_{i,t} & = \left(\dfrac{1}{N} \displaystyle\sum_{j=1}^N \left(\int_U J(x_{i,t}, \cdot, x_{j,t}, u')\mathrm{d}\sigma_{j,t}(u') \right. \right. \\
& \qquad\qquad \left. \left. - \int_U J(x_{i,t}, w, x_{j,t}, u')\mathrm{d}\sigma_{j,t}(u')\mathrm{d}\sigma_{i,t}(w) \right) \right) \sigma_{i,t},
\end{cases}
$$

can be rewritten in the compact form

$$
(3.5) \qquad\qquad \dot{y}_{i,t} = \frac{1}{N} \sum_{j=1}^N f(y_{i,t}, y_{j,t}).
$$

We will see in the next section that such a problem always admits a unique solution, whenever f satisfies the following two conditions:

(i) f is Lipschitz from $C \times C$ to Y when Y is endowed with the product norm induced by the Euclidean norm in $\mathbb{R}^d$ and the BL norm in $F(U)$ and the convex subset C is endowed with the distance induced by the inclusion in $\mathbb{R}^d \times F(U)$;

(ii) for every $R > 0$ there exists a constant $\theta > 0$ such that for every $x, x' \in B_R(0) \subset \mathbb{R}^d$ and $\sigma, \sigma' \in \mathscr{P}(U)$

$$
\sigma + \theta f_\sigma(x, \sigma, x', \sigma') \in \mathscr{P}(U).
$$

This last condition is needed to have condition (B.2) in Theorem B.1 fulfilled.

The above conditions are surely satisfied if the function J is Lipschitz in $(\mathbb{R}^d \times U)^2$.

Notice that (3.5) admits an equivalent formulation by introducing the time-dependent distribution $\Sigma_t^N := \frac{1}{N} \sum_{i=1}^N \delta_{y_{i,t}} \in \mathscr{P}(X)$ and the associated vector field

$$
b_{\Sigma_t^N}(y) = b_{\Sigma^N}(t, y) := \int_Y f(y, y')\mathrm{d}\Sigma_t^N(y') = \frac{1}{N} \sum_{i=1}^N \sum_{i=1}^N f(y, y_{i,t}).
$$

Such a vector field induces a family $\mathbf{Y}_{\Sigma^N}(t, s, \cdot) \colon C \to C$ for $0 \leq s < t \leq T$ of transition operators (also called a flow map) associated to the initial value problem

$$
\dot{y}_r = b_{\Sigma^N}(r, y_r), \quad y_s = y,
$$

namely, $\mathbf{Y}_{\Sigma^N}(t, s, y) := y_s$.

Therefore (3.5) reads as

$$
y_{i,t} = \mathbf{Y}_{\Sigma^N}(t, s, y_{i,s})
$$

so that

$$
\Sigma_t^N = \mathbf{Y}_{\Sigma^N}(t, s, \cdot)_\# \Sigma_s^N \quad \text{for every } 0 \leq s \leq t \leq T.
$$

We conclude this section by further elaborating on the expressions above. By a suitable exchange of integrals, we rewrite the following equation with the equivalent notation

$$
\begin{aligned}
\dot{\sigma}_{i,t} &= \left(\frac{1}{N} \sum_{j=1}^{N} \left(\int_U J(x_{i,t}, \cdot, x_{j,t}, u') \, d\sigma_{j,t}(u') \right. \right. \\
&\qquad\qquad \left. \left. - \int_{U \times U} J(x_{i,t}, w, x_{j,t}, u') d\sigma_{j,t}(u') d\sigma_{i,t}(w) \right) \right) \sigma_{i,t}, \\
&= \left(\int_C \left(\int_U J(x_{i,t}, \cdot, x', u') \, d\sigma'(u') \right. \right. \\
&\qquad\qquad \left. \left. - \int_{U \times U} J(x_{i,t}, w, x', u') d\sigma'(u') d\sigma_{i,t}(w) \right) d\Sigma_t^N(x', \sigma') \right) \sigma_{i,t} \\
&= \left(J * \Sigma_t^N(x_{i,t}, \cdot) - \int_U J * \Sigma_t^N(x_{i,t}, w) d\sigma_{i,t}(w) \right) \sigma_{i,t},
\end{aligned}
$$

where in the following we shall use the compact notation

$$
(3.6) \qquad J * \Lambda(x, u) := \int_C \int_U J(x, u, x', u') d\sigma'(u') d\Lambda(x', \sigma')
$$

for $\Lambda \in \mathscr{P}(C)$.

3.3 Distributed players system: Eulerian and Lagrangian solutions

The general problem associated to an arbitrary initial distribution of players $\bar{\Sigma} \in \mathscr{P}(C)$ can be described as follows. Recall that we are endowing C with the distance

$$
d_C(y_1, y_2) := |x_1 - x_2| + \|\sigma_1 - \sigma_2\|_{\text{BL}}, \quad y_i = (x_i, \sigma_i).
$$

First of all, in order to take care of the lack of compactness of $\mathbb{R}^d$, the first factor of C, we assume that the first moment of the first marginal of $\bar{\Sigma}$ is finite:

$$
(3.7) \qquad \int_C |x| d\bar{\Sigma}(x, \sigma) = M_0 < +\infty.
$$

Since U is compact, (3.7) holds if and only if $\bar{\Sigma} \in \mathscr{P}_1(C)$, with $\mathscr{P}_1(C)$ defined as in (2.2).

We observe that for every continuous curve $t \in [0, T] \mapsto \Sigma_t \in \mathscr{P}_1(C)$ it is possible to define a time-dependent vector field

$$
(3.8) \qquad b_\Sigma(t, y) = b_{\Sigma_t}(y) = \int_C f(y, y') d\Sigma_t(y')
$$

with $f = (f_x, f_\sigma)$ as in (3.3) and (3.4). Since $f(y, \cdot)$ is continuous with linear growth and $\Sigma_t \in \mathscr{P}_1(C)$, the integral above can be interpreted as a Bochner integral; see Section A.2.

We can then associate to b_Σ the transition maps $\mathbf{Y}_\Sigma(t, s, y)$ induced by ODE in Y

$$(3.9) \qquad \dot{y}_t = b_{\Sigma_t}(y_t), \quad y_s = y.$$

A solution $y = (x, \sigma)$ to (3.9) satisfies, with the notation (3.6),

$$\begin{cases} \dot{x}_t = a(x_t, \sigma_t) = \displaystyle\int_U e(x_t, u)\mathrm{d}\sigma_t(u), \\[2mm] \dot{\sigma}_t = \left(\displaystyle\int_C \left(\displaystyle\int_U J(x_t, \cdot, x', u')\mathrm{d}\sigma'(u') \right. \right. \\[2mm] \qquad\qquad \left. \left. - \displaystyle\int_{U \times U} J(x_t, w, x', u')\mathrm{d}\sigma'(u')\mathrm{d}\sigma_t(w) \right)\mathrm{d}\Sigma_t(x', \sigma') \right)\sigma_t \\[2mm] \qquad = \left(J * \Sigma_t(x_t, \cdot) - \displaystyle\int_U J * \Sigma_t(x_t, w)\mathrm{d}\sigma_t(w) \right)\sigma_t, \end{cases}$$

and the existence of a solution to (3.9) follows again by Theorem B.1; see Theorem 4.1 under the structural properties assumed in Section 3.4 below.

Whenever we have at our disposal the flow map $\mathbf{Y}_\Sigma$, the transported measures $\widehat{\Sigma}_t := \mathbf{Y}(t, 0, \cdot)_\# \bar{\Sigma}$ solve an infinite-dimensional continuity equation driven by the vector field b_{Σ_t} given by (3.8), namely (in integral form)

$$\begin{aligned}(3.10) \quad & \int_C \phi(t, y)\mathrm{d}\widehat{\Sigma}_t(y) - \int_C \phi(0, y)\mathrm{d}\widehat{\Sigma}_0(y) \\ & = \int_0^t \int_C \left(\partial_s\phi(s, y) + \int_C \mathrm{D}\phi(s, y)(f(y, y'))\mathrm{d}\Sigma_s(y') \right)\mathrm{d}\widehat{\Sigma}_s(y)\mathrm{d}s \end{aligned}$$

for every $\phi \in \mathrm{C}_b^1([0, T] \times Y)$. Indeed, using the change-of-variables formula for the push-forward measure, the chain rule, and once more the change-of-variables formula, one has

$$\begin{aligned} & \frac{\mathrm{d}}{\mathrm{d}t} \int_C \phi(t, y)\mathrm{d}\widehat{\Sigma}_t(y) \\ & = \int_C \frac{\mathrm{d}}{\mathrm{d}t}\phi\big(t, \mathbf{Y}_\Sigma(t, 0, z)\big)\mathrm{d}\bar{\Sigma}(z) \\ & = \int_C \big(\partial_t\phi(t, \mathbf{Y}_\Sigma(t, 0, z)) \\ & \qquad\qquad + \mathrm{D}\phi(t, \mathbf{Y}_\Sigma(t, 0, z))\big(b_{\Sigma_t}(\mathbf{Y}_\Sigma(t, 0, z))\big)\big)\mathrm{d}\bar{\Sigma}(z) \\ & = \int_C \big(\partial_t\phi(t, y) + \mathrm{D}\phi(t, y)(b_{\Sigma_t}(y))\big)\mathrm{d}\widehat{\Sigma}_t(y) \\ & = \int_C \left(\partial_t\phi(t, y) + \int_C \mathrm{D}\phi(t, y)(f(y, y'))\mathrm{d}\Sigma_t(y') \right)\mathrm{d}\Sigma_t(y). \end{aligned}$$

Formula (3.10) follows now by integration in time.

We look for an evolving distribution Σ that is self-transported by the generated vector field b_Σ, so that $\widehat{\Sigma} = \Sigma$. These facts motivate the following definition.

DEFINITION 3.2 (Lagrangian and Eulerian solutions). Let $\Sigma \in \mathrm{C}^0([0, T]; (\mathscr{P}_1(C), W_1))$ and $\bar{\Sigma} \in \mathscr{P}_1(C)$. We say that Σ is an Eulerian solution of the initial value problem for the master equation starting from $\bar{\Sigma}$ if $\Sigma_0 = \bar{\Sigma}$ and (3.10) holds with $\hat{\Sigma} = \Sigma$. We say that Σ is a Lagrangian solution starting from $\bar{\Sigma}$ if

$$(3.11) \qquad \Sigma_t = \mathbf{Y}_\Sigma(t, 0, \cdot)_\# \bar{\Sigma} \quad \text{for every } 0 \le t \le T,$$

where $\mathbf{Y}_\Sigma(t, s, y)$ are the transition maps associated to the ODE (3.9).

The Lagrangian notion of solution given by the transport identity (3.11) is rather standard and accepted in the literature of multiagent systems and mean-field equations. One can see, for instance, the notion of solution given in [13, def. 3.3]. On the other hand, as in fluid mechanics, when looking at the evolution of spatially averaged quantities it is also important to derive an alternative Eulerian description in terms of a PDE, in our case (3.10).

We shall first address the problem of existence and uniqueness of Lagrangian solutions. Given that, as we illustrated above, Lagrangian solutions are Eulerian, this settles the existence problem also for Eulerian solutions. The uniqueness of Eulerian solutions is technically harder, and it will be dealt with in Section 5.

THEOREM 3.3. *Suppose that* $J: (\mathbb{R}^d \times U)^2 \to \mathbb{R}$ *and* $e: \mathbb{R}^d \times U \to \mathbb{R}$ *are Lipschitz maps and let* $f: C \times C \to Y$ *be defined as in* (3.3) *and* (3.4). *Then, for every* $\bar{\Sigma} \in \mathscr{P}_1(C)$, *there exists a unique Lagrangian solution* Σ, *and its flow map* $\mathbf{Y}_\Sigma(t, s, y) = (\boldsymbol{x}(t, s, y), \boldsymbol{\sigma}(t, s, y))$ *satisfies the additional regularity property that* $\boldsymbol{\sigma}(\cdot, s, y)$ *is of class* C^1 *with values in* $(\mathscr{M}(U), \|\cdot\|_{TV})$, *with*

$$(3.12) \qquad \mathrm{Lip}\big(\boldsymbol{\sigma}(\cdot, s, y), (\mathscr{M}(U), \|\cdot\|_{TV})\big) \le L_J \, \mathrm{diam}(U).$$

Moreover, there exists $L \ge 0$ *such that for every pair* $\bar{\Sigma}^i$, $i = 1, 2$, *of initial data in* $\mathscr{P}_1(C)$, *the corresponding solutions* Σ_t^i *satisfy*

$$W_1\big(\Sigma_t^1, \Sigma_t^2\big) \le \mathrm{e}^{Lt} W_1\big(\bar{\Sigma}^1, \bar{\Sigma}^2\big) \quad \text{for every } t \in [0, T].$$

The proof of Theorem 3.3 will be given in Section 4: it does not depend on the particular structure of f and J, but relies on their Lipschitz property, the convexity of C, and the Banach framework. Notice that (3.12) comes immediately from (3.4) and the definition of the $F(U)$-component $(b_{\Sigma_t})_\sigma(y) = \int_Y f_\sigma(y, y') \mathrm{d}\Sigma_t(y')$ of b_{Σ_t} by using the estimate

$$\int_U \left| \int_U J(x, u, x', u') \mathrm{d}\sigma'(u') - \int_{U \times U} J(x, w, x', u') \mathrm{d}\sigma'(u') \mathrm{d}\sigma(w) \right| \mathrm{d}\sigma(u)$$
$$\le L_J \, \mathrm{diam}(U)$$

which, thanks to the uniformity with respect to $y' = (x', \sigma')$, gives that

$$\|(b_{\Sigma_t})_\sigma\|_{TV} \le L_J \, \mathrm{diam}(U).$$

Then, thanks to the discussion in Section 2.2, we also obtain C^1 regularity with respect to the total variation norm.

3.4 Structural properties of the interaction term f

Recall that $C = \mathbb{R}^d \times \mathscr{P}(U)$ is a closed and convex subset of $Y = \mathbb{R}^d \times F(U)$. We shall denote in the following by L_e and L_J the Lipschitz constants of e and J, respectively.

Remark 3.4. We endow $C \times C$ with the distance induced by the norm in $Y \times Y$

$$\|(y_1, y_2)\| := \|y_1\|_{\mathrm{BL}} + \|y_2\|_{\mathrm{BL}}.$$

With this choice, a function is L-Lipschitz if (and only if) it is L-Lipschitz separately in the components, that is, if for every y, y_1, $y_2 \in C$ one has

$$\|f(y_1, y) - f(y_2, y)\| \le L \|y_1 - y_2\|_{\mathrm{BL}},$$
$$\|f(y, y_1) - f(y, y_2)\| \le L \|y_1 - y_2\|_{\mathrm{BL}},$$

then

$$\|f(y_1, y_2) - f(y_1', y_2')\| \le L \left(\|y_1 - y_1'\|_{\mathrm{BL}} + \|y_2 - y_2'\|_{\mathrm{BL}} \right)$$
$$= L \|(y_1, y_2) - (y_1', y_2')\|.$$

For $y = (x, \sigma) \in C$ and $\Sigma \in \mathscr{P}_1(C)$, let $b_\Sigma(y) = \int_U f(y, y') \mathrm{d}\Sigma(y')$ as in (3.8), where $f \colon C \times C \to Y$ is defined by

$$
\begin{aligned}
(3.13) \qquad f(y, y') = \Bigg(a(y), \Bigg(& \int_U J(x, \cdot, x', u') \mathrm{d}\sigma'(u') \\
& - \int_{U \times U} J(x, w, x', u') \mathrm{d}\sigma'(u') \mathrm{d}\sigma(w) \Bigg) \sigma \Bigg),
\end{aligned}
$$

with $a(y) = a(x, \sigma) = \int_U e(x, u) \mathrm{d}\sigma(u)$ as in (1.4). We define a map $j \colon \mathbb{R}^d \times \mathbb{R}^d \times \mathscr{P}(U) \to \mathrm{Lip}(U)$ by

$$(3.14) \qquad (x, x', \sigma') \mapsto j(x, x', \sigma')(u) := \int_U J(x, u, x', u') \mathrm{d}\sigma'(u').$$

Notice that the map $j(x, x', \sigma')$ depends linearly on σ'. Moreover, recalling that σ' is a probability measure, it is not difficult to see that

$$
\begin{aligned}
\|j(x, x', \sigma')\|_{\mathrm{Lip}(U)} &= \|j(x, x', \sigma')\|_{L^\infty(U)} + \mathrm{Lip}(j(x, x', \sigma')) \\
&\le \max |J(x, \cdot, x', \cdot)| + L_J \operatorname{diam}(U).
\end{aligned}
$$

LEMMA 3.5. *Let* $\sigma \in \mathscr{P}(U)$ *and let* $z \in \mathrm{Lip}(U)$. *Then the following estimate holds*:

$$(3.15) \qquad \|z\sigma\|_{\mathrm{BL}} \le \|z\|_{\mathrm{Lip}} \|\sigma\|_{\mathrm{BL}}.$$

PROOF. It follows directly by (2.6) and by the definition (1.10) of the BL norm that

$$\|z\sigma\|_{\mathrm{BL}} = \sup_{\eta : \|\eta\|_{\mathrm{Lip}} \le 1} \int_U \eta(u) z(u) \mathrm{d}\sigma(u) \le \|z\|_{\mathrm{Lip}} \|\sigma\|_{\mathrm{BL}}. \qquad \square$$

PROPOSITION 3.6. *The map f defined in* (3.13) *is L-Lipschitz* (*in both variables*), *with L depending only on L_e, L_J, and* $\mathrm{diam}(U)$, *and satisfies the compatibility condition* (*see also* (4.1) *below*)

$$(3.16) \qquad \forall R > 0, \ \exists \theta > 0 \colon y, \ y' \in C \cap B_R(0) \quad \Rightarrow \quad y \pm \theta f(y, y') \in C.$$

PROOF. By Remark 3.4, we can study the Lipschitz dependence of f separately with respect to y and y'. Moreover, we can consider the Lipschitz dependence on x and σ separately, keeping the other variable frozen. Let us start with f_x; since it does not depend on y', we only study the Lipschitz dependence on y. We have

$$
\begin{aligned}
|f_x(y_1, y') - f_x(y_2, y')| &= |a(y_1) - a(y_2)| \\
&= \left| \int_U e(x_1, u) \mathrm{d}\sigma_1(u) - \int_U e(x_2, u) \mathrm{d}\sigma_2(u) \right| \\
&\leq \left| \int_U e(x_1, u) \mathrm{d}\sigma_1(u) - \int_U e(x_1, u) \mathrm{d}\sigma_2(u) \right| \\
&\quad + \left| \int_U e(x_1, u) \mathrm{d}\sigma_2(u) - \int_U e(x_2, u) \mathrm{d}\sigma_2(u) \right| \\
&\leq L_e(1 + \mathrm{diam}\, U) \|\sigma_1 - \sigma_2\|_{\mathrm{BL}} + L_e|x_1 - x_2| \\
&\leq L_e(1 + \mathrm{diam}\, U) \|y_1 - y_2\|,
\end{aligned}
$$
(3.17)

where we have used (2.9). To study the Lipschitz dependence of f_σ on its variables, it is convenient to do it for x, x', σ, and σ' separately. The Lipschitz dependence on x is easy to obtain, and it leads to

$$(3.18) \qquad \|f_\sigma(x_1, \sigma, y') - f_\sigma(x_2, \sigma, y')\|_{\mathrm{BL}} \leq 2L_J(1 + \mathrm{diam}\, U)|x_1 - x_2|.$$

Similarly, one can prove that

$$(3.19) \qquad \|f_\sigma(y, x_1', \sigma') - f_\sigma(y, x_2', \sigma')\|_{\mathrm{BL}} \leq 2L_J(1 + \mathrm{diam}\, U)|x_1' - x_2'|.$$

Let us consider the dependence on σ. Using the map $j(x, x', \sigma')$ defined in (3.14), we have to estimate

$$
\begin{aligned}
&\|f_\sigma(x, \sigma_1, y') - f_\sigma(x, \sigma_2, y')\|_{BL} \\
&\quad = \left\| \sigma_1 \left(j(x, x', \sigma')(\cdot) - \int_U j(x, x', \sigma')(v) \mathrm{d}\sigma_1(v) \right) \right. \\
&\qquad \left. - \sigma_2 \left(j(x, x', \sigma')(\cdot) - \int_U j(x, x', \sigma')(v) \mathrm{d}\sigma_2(v) \right) \right\|_{\mathrm{BL}} \\
&\quad \leq \left\| (\sigma_1 - \sigma_2) \left(j(x, x', \sigma')(\cdot) - \int_U j(x, x', \sigma')(v) \mathrm{d}\sigma_2(v) \right) \right\|_{\mathrm{BL}} \\
&\qquad + \left\| \sigma_1 \left(\int_U j(x, x', \sigma')(v) \mathrm{d}(\sigma_2 - \sigma_1)(v) \right) \right\|_{\mathrm{BL}} =: \mathrm{I} + \mathrm{II}.
\end{aligned}
$$
(3.20)

Term I above can be estimated as follows

$$(3.21) \qquad \mathrm{I} = \left\| (\sigma_1 - \sigma_2) \left(\int_U \big(j(x, x', \sigma')(\cdot) - j(x, x', \sigma')(v) \big) \mathrm{d}\sigma_2(v) \right) \right\|_{\mathrm{BL}}$$

$$\leq L_J (1 + \operatorname{diam} U) \| \sigma_1 - \sigma_2 \|_{\mathrm{BL}}.$$

To estimate II, we use the definition of BL norm and the fact that $\sigma_1 - \sigma_2 \in \mathscr{M}_0(U)$:

$$\mathrm{II} = \sup_{\eta : \| \eta \|_{\mathrm{Lip}} \leq 1} \int_U \eta(u) \left(\int_U j(x, x', \sigma')(v) \mathrm{d}(\sigma_2 - \sigma_1)(v) \right) \mathrm{d}\sigma_1(u)$$

$$(3.22) \qquad \leq \left| \int_U j(x, x', \sigma')(v) \mathrm{d}(\sigma_2 - \sigma_1)(v) \right|$$

$$\leq \int_U | j(x, x', \sigma')(v) - j(x, x', \sigma')(u) | \mathrm{d}(\sigma_2 - \sigma_1)(v)$$

$$\leq L_J \operatorname{diam} U \| \sigma_1 - \sigma_2 \|_{\mathrm{BL}}.$$

Putting (3.21) and (3.22) together, we can complete the estimate for (3.20) and obtain

$$(3.23) \qquad \| f_\sigma(x, \sigma_1, y') - f_\sigma(x, \sigma_2, y') \|_{\mathrm{BL}} \leq 2 L_J (1 + \operatorname{diam} U) \| \sigma_1 - \sigma_2 \|_{\mathrm{BL}}.$$

Using (3.15) and that $\| \sigma \|_{\mathrm{BL}} \leq 1$, let us now estimate

$$\| f_\sigma(y, x', \sigma_1') - f_\sigma(y, x', \sigma_2') \|_{\mathrm{BL}}$$

$$= \left\| \sigma \left(j(x, x', \sigma_1')(\cdot) - \int_U j(x, x', \sigma_1')(v) \mathrm{d}\sigma(v) \right) \right.$$

$$\left. - \sigma \left(j(x, x', \sigma_2')(\cdot) - \int_U j(x, x', \sigma_2')(v) \mathrm{d}\sigma(v) \right) \right\|_{\mathrm{BL}}$$

$$\leq \| j(x, x', \sigma_1') - j(x, x', \sigma_2') \|_{\mathrm{Lip}}$$

$$+ \left\| \int_U \big(j(x, x', \sigma_1')(v) - j(x, x', \sigma_2')(v) \big) \mathrm{d}\sigma(v) \right\|_{\mathrm{Lip}} =: \mathrm{I}' + \mathrm{II}'.$$

To estimate I', we use the definition (3.14) of $j(x, x', \sigma')$ and the fact that $\sigma_1' - \sigma_2' \in \mathscr{M}_0(U)$ to obtain

$$\mathrm{I}' = \left\| \int_U J(x, \cdot, x', u') \mathrm{d}\sigma_1'(u') - \int_U J(x, \cdot, x', u') \mathrm{d}\sigma_2'(u') \right\|_{\mathrm{Lip}}$$

$$= \left\| \int_U J(x, \cdot, x', u') \mathrm{d}(\sigma_1' - \sigma_2')(u') \right\|_{\mathrm{Lip}}$$

$$= \left\| \int_U \big(J(x, \cdot, x', u') - J(x, \cdot, x', v') \big) \mathrm{d}(\sigma_1' - \sigma_2')(u') \right\|_{\mathrm{Lip}}$$

$$\leq L_J (1 + \operatorname{diam} U) \| \sigma_1' - \sigma_2' \|_{\mathrm{BL}}.$$

The estimate of II′ follows in a similar way, so that we obtain

$$(3.24) \quad \|f_\sigma(y, x', \sigma_1') - f_\sigma(y, x', \sigma_2')\|_{BL} \le 2L_J(1 + \operatorname{diam} U)\|\sigma_1' - \sigma_2'\|_{BL}.$$

Putting (3.18), (3.19), (3.23), and (3.24) together, we obtain

$$\|f_\sigma(y_1, y_1') - f_\sigma(y_2, y_2')\|_{BL} \le 2L_J(1 + \operatorname{diam} U)(\|y_1 - y_2\| + \|y_1' - y_2'\|),$$

which, together with (3.17) gives the Lipschitz estimate on f.

Let us now discuss the compatibility conditions (4.1) for the f defined in (3.13). It is clear that the first component of $y + \theta f(y, y')$, namely $x + \theta a(y)$, belongs to $\mathbb{R}^d$ for all $\theta \in \mathbb{R}$, so that we are left with checking that the second component $\sigma + \theta f_\sigma(y, y')$, namely,

$$(3.25) \quad \sigma + \theta\sigma\left(\int_U J(x, \cdot, x', u')\mathrm{d}\sigma'(u') - \int_{U \times U} J(x, u, x', u')\mathrm{d}\sigma'(u')\mathrm{d}\sigma(u)\right),$$

belongs to $F(U)$. As a matter of fact, we will prove that (3.25) is an element of $\mathscr{P}(U)$, which means that its integral over U is 1 and that it is positive. The proof that $\sigma + \theta f_\sigma(y, y') \ge 0$ can be obtained via some manipulations and using the Lipschitz estimate on J. Indeed,

$$\sigma + \theta f_\sigma(y, y')$$
$$= \sigma\left(1 + \theta\left(\int_U J(x, u, x', u')\mathrm{d}\sigma'(u')\right.\right.$$
$$\left.\left. - \int_{U \times U} J(x, w, x', u')\mathrm{d}\sigma'(u')\mathrm{d}\sigma(w)\right)\right)$$
$$= \sigma\left(1 + \theta\int_U\left(J(x, u, x', u') - \int_U J(x, w, x', u')\mathrm{d}\sigma(w)\right)\mathrm{d}\sigma'(u')\right)$$
$$= \sigma\left(1 + \theta\int_{U \times U}(J(x, u, x', u') - J(x, w, x', u'))\mathrm{d}\sigma(w)\mathrm{d}\sigma'(u')\right)$$
$$\ge \sigma\left(1 - \theta L_J\int_U d_U(u, w)\mathrm{d}\sigma'(u')\right) \ge \sigma(1 - \theta L_J \operatorname{diam} U),$$

which is nonnegative as soon as $\theta \le (L_J \operatorname{diam} U)^{-1}$. By recalling that $f_\sigma(y, y') \in \mathscr{M}_0(U)$, we obtain that $\sigma + \theta f_\sigma(y, y') \in \mathscr{P}(U)$. $\quad\square$

4 Existence and Uniqueness of Lagrangian Solutions

4.1 Interaction systems in Banach spaces

Let us consider now a Banach space $(Y, \|\cdot\|)$ with a closed convex set C and a L-Lipschitz map

$$f\colon C \times C \to Y$$

satisfying the compatibility condition

$$(4.1) \quad \forall R > 0, \ \exists\theta > 0 \quad c, c' \in C \cap B_R(0) \quad \Rightarrow \quad c + \theta f(c, c') \in C.$$

Let us consider a continuous curve of measures $\Sigma \in \mathrm{C}([0, T]; (\mathscr{P}_1(C), W_1))$. Recalling that

$$(4.2) \qquad \|f(c, c')\| \leq \|f(c_0, c_0)\| + L(\|c - c_0\| + \|c' - c_0\|)$$

where c_0 is an arbitrary point in C, we can define the time-dependent vector field $b_\Sigma(t, \cdot): C \to Y$ by

$$(4.3) \qquad b_\Sigma(t, c) = b_{\Sigma_t}(c) := \int_C f(c, c') \mathrm{d}\Sigma_t(c'),$$

where the integral above can be interpreted in the strong sense, as a Bochner integral.

We are going to prove the following result, which provides (taking Proposition 3.6 into account) the proof of Theorem 3.3.

THEOREM 4.1. *Given* $\bar{\Sigma} \in \mathscr{P}_1(C)$ *there exists* $\Sigma \in \mathrm{C}^0([0, T]; (\mathscr{P}_1(C), W_1))$ *with* $\Sigma_0 = \bar{\Sigma}$ *such that the family of transition maps* $\mathbf{Y}_\Sigma(t, s, \cdot)$ *in* C *associated to the ODE*

$$\dot{y}_t = b_\Sigma(t, y_t), \quad y_s = y, \quad y \in \mathrm{C}^1([s, T]; Y), \quad y_t \in C \ \text{for every } t \in [s, T],$$

with the vector field b_Σ *given by* (4.3) *satisfying*

$$\Sigma_t = \mathbf{Y}_\Sigma(t, s, \cdot)_\# \Sigma_s \quad \text{for every } 0 \leq s \leq t \leq T.$$

In addition, one has the stability estimate

$$(4.4) \qquad W_1(\Sigma_t, \Sigma_t') \leq \mathrm{e}^{2L(t-s)} W_1(\Sigma_s, \Sigma_s') \quad \text{for every } 0 \leq s \leq t \leq T,$$

for the solutions Σ, Σ' *starting from* $\bar{\Sigma}$, $\bar{\Sigma}'$, *where* L *is the Lipschitz constant of* f.

4.2 Existence for the discrete problem

We first study the discrete problem for N particles evolving in Y, corresponding to the evolution of a discrete (atomic) measure. This case could be simply seen as a by-product of the more general "diffuse" measure well-posedness result; however, we include it both as a guideline to introduce the more general case and also as a constructive approximation (for N large; see also Remark 4.8 below), which could be useful for the purpose of numerical simulation. We consider the convex set C^N in Y^N with the norm

$$\|\boldsymbol{y}\|_{Y^N} := \frac{1}{N} \sum_{i=1}^{N} \|y_i\|, \quad \boldsymbol{y} = (y_1, \ldots, y_N) \in Y^N.$$

We define the map $\boldsymbol{f}^N = (f_1^N, \ldots, f_N^N): C^N \to Y^N$ by

$$f_i^N(\boldsymbol{y}) = \frac{1}{N} \sum_{j=1}^{N} f(y_i, y_j), \quad i = 1, \ldots, N.$$

We notice that $\boldsymbol{f}$ is Lipschitz, since

$$\|\boldsymbol{f}^N(\boldsymbol{y}) - \boldsymbol{f}^N(\boldsymbol{y}')\|_{Y^N} \leq \frac{1}{N} \sum_{i=1}^{N} \|f_i^N(\boldsymbol{y}) - f_i^N(\boldsymbol{y}')\|$$

$$\leq \frac{1}{N^2} \sum_{i,j=1}^{N} \|f(y_i, y_j) - f(y_i', y_j')\|$$

$$\leq \frac{L}{N^2} \sum_{i,j=1}^{N} (\|y_i - y_i'\| + \|y_j - y_j'\|)$$

$$\leq 2L\|\boldsymbol{y} - \boldsymbol{y}'\|_{Y^N}.$$

Let us now check that C^N satisfies the invariance properties with respect to $\boldsymbol{f}^N$: if $\boldsymbol{y} \in C^N$ with $\|\boldsymbol{y}\| \leq R$, then every component y_i belongs to C and $\|y_i\| \leq NR$. By (4.1) (applied to the constant NR) we may find a constant $\theta > 0$ such that

$$y_i + \theta f(y_i, y_j) \in C \quad \text{for every } i, j,$$

so that the convexity of C yields

$$y_i + \theta f_i^N(\boldsymbol{y}) = \frac{1}{N} \sum_{j=1}^{N} \left(y_i + \theta f(y_i, y_j) \right) \in C.$$

By applying Theorem B.1 with C^N and Y^N we obtain the following result.

COROLLARY 4.2. *For every $\bar{\boldsymbol{y}} \in C^N$, there exists a unique curve $\boldsymbol{y}$: $[0, +\infty) \to C^N$ of class C^1 such that*

$$\begin{cases} \dot{\boldsymbol{y}}(t) = \boldsymbol{f}^N(\boldsymbol{y}^N(t)), \\ \boldsymbol{y}(0) = \bar{\boldsymbol{y}}. \end{cases}$$

In particular, the family $\Sigma_t := \frac{1}{N} \sum_{i=1}^{N} \delta_{y_{i,t}}$ provides a solution to the existence part of Theorem 4.1 for the initial datum $\bar{\Sigma} := \frac{1}{N} \sum_{i=1}^{N} \delta_{\bar{y}_i}$.

4.3 Stability estimates

PROPOSITION 4.3 (Properties of b_t). *Let Λ, $\Lambda' \in \mathrm{C}([0, T]; \mathscr{P}_1(C))$, let b_Λ and $b_{\Lambda'}$ be defined as in (4.3) and let $f : C \times C \to Y$ be L-Lipschitz. Then*

(i) $\|b_\Lambda(t, y)\| \leq \|f(y_0, y_0)\| + L\|y - y_0\| + L \int_C \|y' - y_0\| \, \mathrm{d}\Lambda_t(y')$ *for all* $y_0 \in Y$.

(ii) $\|b_\Lambda(t, y) - b_\Lambda(t, z)\| \leq L\|y - z\|$.

(iii) $\|b_\Lambda(t, y) - b_\Lambda(s, y)\| \leq L W_1(\Lambda_t, \Lambda_s)$.

(iv) $\|b_\Lambda(t, y) - b_{\Lambda'}(t, y)\| \leq L W_1(\Lambda_t, \Lambda_t')$.

(v) *If there exists $\bar{R} > 0$ such that $\Lambda_t(C \setminus B_{\bar{R}}(0)) = 0$ for every $t \in [0, T]$, then for every $R > 0$ there exists $\theta > 0$ such that*

$$y \in C, \ \|y\| \leq R, \ t \in [0, T] \quad \Rightarrow \quad y + \theta b_\Lambda(t, y) \in C.$$

PROOF. Property (i) follows immediately from (4.2). To prove (ii), we notice that

$$\|b_\Lambda(t, y) - b_\Lambda(t, z)\| = \left\| \int_C \Big(f(y, y') - f(z, y') \Big) \mathrm{d}\Lambda_t(y') \right\| \le L \|y - z\|.$$

Estimate (iii) is a simple computation:

$$
\begin{aligned}
\|b_\Lambda(t, y) - b_\Lambda(s, y)\| &= \left\| \int_C f(y, y') \mathrm{d}(\Lambda_t - \Lambda_s)(y') \right\| \\
&= \sup_{z \in Y', \, \|z\|_{Y'} \le 1} \int_C \langle z, f(y, y') \rangle \, \mathrm{d}(\Lambda_t - \Lambda_s)(y') \\
&\le L W_1(\Lambda_t, \Lambda_s),
\end{aligned}
$$

where we have used that the map $y' \mapsto \langle z, f(y, y') \rangle$ is L-Lipschitz. The proof of (iv) is analogous.

Let us now consider the last statement (v); we may assume $R \ge \bar{R}$, and we can choose $\theta > 0$ such that (4.1) holds. Therefore

$$y + \theta b_\Lambda(t, y) = \int_C (y + \theta f(y, y')) \mathrm{d}\Lambda_t(y') \in C,$$

since C is convex and closed, and Λ_t is a probability measure. $\square$

COROLLARY 4.4. *Let* Λ, $\Lambda^i \in C([0, T]; \mathscr{P}_1(C))$, *let* b_Λ, b_{Λ^i} *be defined as in* (4.3), $y_0 \in Y$ *and let* $f \colon C \times C \to Y$ *be L-Lipschitz. Then*

(i) *for every* $y \in C$ *and* $s \in [0, T]$ *there exists a unique solution* $y_t = \mathbf{Y}_\Lambda(t, s, y)$ *from* $[s, T]$ *to* C *of class* C^1 *of the Cauchy problem*

(4.5) $$\dot{y}_r = b_\Lambda(r, y_r) = b_{\Lambda_r}(y_r), \quad y_s = y;$$

(ii) $\mathbf{Y}_\Lambda(t, 0, y)$ *satisfies the estimate*

(4.6) $$\|\mathbf{Y}_\Lambda(t, 0, y) - y_0\| \le \big(\|y - y_0\| + t B(\Lambda, t, y_0) \big) \mathrm{e}^{Lt}$$

   with

(4.7) $$B(\Lambda, t, y_0) := \|f(y_0, y_0)\| + L \max_{s \in [0, t]} \int_C \|y' - y_0\| \mathrm{d}\Lambda_s(y');$$

(iii) $\mathbf{Y}_\Lambda(\cdot, 0, y)$ *satisfies the estimate*

(4.8)
$$
\begin{aligned}
&\|\mathbf{Y}_\Lambda(t, 0, y) - \mathbf{Y}_\Lambda(t', 0, y)\| \\
&\qquad \le |t - t'| \big[B(\Lambda, T, y_0) + L \big(\|y - y_0\| + T B(\Lambda, T, y_0) \big) \mathrm{e}^{LT} \big];
\end{aligned}
$$

(iv) $\mathbf{Y}_\Lambda(t, s, \cdot)$ *satisfies the estimate*

$$\|\mathbf{Y}_\Lambda(t, s, y) - \mathbf{Y}_\Lambda(t, s, y')\| \le \mathrm{e}^{L(t-s)} \|y - y'\|, \quad 0 \le s \le t \le T;$$

(v) *more generally, $\mathbf{Y}_{\Lambda^1}$ and $\mathbf{Y}_{\Lambda^2}$ satisfy the estimate for $0 \le s \le t \le T$:*

$$\|\mathbf{Y}_{\Lambda^1}(t,s,y^1) - \mathbf{Y}_{\Lambda^2}(t,s,y^2)\| \le e^{L(t-s)}\|y^1 - y^2\|$$

(4.9)
$$+ L \int_s^t e^{L(t-\tau)} W_1(\Lambda_\tau^1, \Lambda_\tau^2)\mathrm{d}\tau.$$

PROOF. Let us first assume that Λ and Λ^i are concentrated on a ball of radius $\bar{R}$ in Y. Then statements (i) and (iv) immediately follow by Theorem B.1, thanks to the estimates of Proposition 4.3.

The proof of (ii) is a computation: recalling (3.8), Proposition 4.3(i) and the definition (4.7) of $B(\Lambda, t, y_0)$, the triangle inequality gives

$$\|\mathbf{Y}_\Lambda(t,0,y) - y_0\| \le \|y - y_0\| + \int_0^t \|b_{\Lambda_s}(y(s))\|\mathrm{d}s$$
$$\le \|y - y_0\| + tB(\Lambda,t,y_0) + L\int_0^t \|y(s) - y_0\|\mathrm{d}s,$$

which yields (4.6) by Gronwall's inequality. Estimate (4.8) follows from combining Proposition 4.3(i) with estimate (4.6), so that (iii) is proved.

Concerning (v), it is clearly sufficient to consider the case $s = 0$. Denoting by y_t^i, $t \in [0,T]$, the solutions to (4.5) with respect to the fields $b_t^i = b_{\Lambda_t^i}$ and the initial conditions $\bar{y}^i$, from $\dot{y}^i = b_t^i(y^i)$ we have, by Proposition 4.3(ii) and (iv)

$$\frac{\mathrm{d}}{\mathrm{d}t}\|y^1 - y^2\|(t) \le \|b_t^1(y^1) - b_t^1(y^2)\| + \|b_t^1(y^2) - b_t^2(y^2)\|$$
$$\le L\|y^1 - y^2\| + LW_1(\Lambda_t^1, \Lambda_t^2),$$

which gives by a simple comparison argument

$$\|y^1(t) - y^2(t)\| \le e^{Lt}\|y^1 - y^2\| + L\int_0^t e^{L(t-\tau)}W_1(\Lambda_\tau^1, \Lambda_\tau^2)\mathrm{d}\tau.$$

The general case when Λ may have unbounded support can be obtained by approximation, using once more Proposition 4.3, since the estimates are independent of $\bar{R}$. □

4.4 Contractivity and stability

We now fix $\bar{\Sigma} \in \mathscr{P}_1(C)$ and we consider the metric space

$$\mathscr{A} := \{\Lambda \in C([0,T]; (\mathscr{P}_1(C), W_1)) : \Lambda_0 = \bar{\Sigma}\},$$

complete when endowed with the usual sup distance (as a consequence of the completeness of $(\mathscr{P}_1(C), W_1)$). We define a map $\mathcal{T}: \mathscr{A} \to \mathscr{A}$ in the following way: given $\Lambda \in \mathscr{A}$ we first compute the flow map $\mathbf{Y}_\Lambda(t,s,\cdot)$ associated to b_Λ and then we define the curve $\mathcal{T}[\Lambda]: [0,T] \to \mathscr{P}_1(C)$ by

(4.10)
$$\mathcal{T}[\Lambda]_t := \mathbf{Y}_\Lambda(t,0,\cdot)_\# \bar{\Sigma}$$

It is immediate to check that $\mathcal{T}$ maps $\mathscr{A}$ to $\mathscr{A}$.

LEMMA 4.5. *For every* Λ, Λ^1, $\Lambda^2 \in \mathscr{A}$ *we have*

$$W_1(\mathcal{T}[\Lambda]_t, \mathcal{T}[\Lambda]_s)$$

(4.11)
$$\leq |t - s| \left[B(\Lambda, T, y_0) + L \left(\int_C \|y - y_0\| \, d\bar{\Sigma}(y) + T B(\Lambda, T, y_0) \right) e^{LT} \right],$$

(4.12)
$$W_1(\mathcal{T}[\Lambda^1]_t, \mathcal{T}[\Lambda^2]_t) \leq L \int_0^t e^{L(t-\tau)} W_1(\Lambda_\tau^1, \Lambda_\tau^2) d\tau,$$

where the constant $B(\Lambda, T, y_0)$ *is defined in* (4.7) *for* $t = T$.

PROOF. Estimate (4.11) follows immediately from the definition of $\mathcal{T}[\Lambda]$ in (4.10), estimate (4.8), and the fact that $\bar{\Sigma}$ is a probability measure. Estimate (4.12) is a direct consequence of (4.9), since

$$W_1(\mathcal{T}[\Lambda^1]_t, \mathcal{T}[\Lambda^2]_t) \leq \int_C \|\mathbf{Y}_{\Lambda^1}(t, 0, y) - \mathbf{Y}_{\Lambda^2}(t, 0, y)\| \, d\bar{\Sigma}(y)$$

and $\bar{\Sigma}$ is a probability measure. $\square$

COROLLARY 4.6. *The map* $\mathcal{T}$ *admits a unique fixed point, which provides the unique solution* Σ *in Theorem* 4.1.

PROOF. Let us fix a constant $L' > 2L$ so that $\ell := L/(L' - L) < 1$, and let us consider the equivalent distance in $\mathscr{A}$ given by

$$\mathsf{d}(\Lambda, \Lambda') := \max_{t \in [0,T]} e^{-L't} W_1(\Lambda_t, \Lambda_t').$$

Then, from (4.12) we immediately get

$$e^{-L't} W_1(\mathcal{T}[\Lambda]_t, \mathcal{T}[\Lambda']_t) \leq L \int_0^t e^{(L-L')(t-s)} ds \, W_1(\Lambda_r, \Lambda_r') \quad \forall t \in [0, T],$$

so that our choice of L' gives $\mathsf{d}(\mathcal{T}[\Lambda], \mathcal{T}[\Lambda']) \leq \ell \mathsf{d}(\Lambda, \Lambda')$. $\square$

We can slightly modify the previous argument in order to derive a stability estimate of the solution Σ_t in terms of the initial datum $\bar{\Sigma}$.

LEMMA 4.7. *Let* $\bar{\Sigma}^1$ *and* $\bar{\Sigma}^2$ *be initial data in* $\mathscr{P}_1(C)$, *and let* Σ_t^i *be the corresponding solutions. Then*

(4.13)
$$W_1(\Sigma_t^1, \Sigma_t^2) \leq e^{2Lt} W_1(\bar{\Sigma}^1, \bar{\Sigma}^2) \quad \forall t \in [0, T].$$

PROOF. Let us fix $0 \leq s < t \leq T$ and consider an optimal coupling Π_s between Σ_s^1 and Σ_s^2 so that using the fact that $(\mathbf{Y}_{\Sigma^1}(t, s, y^1), \mathbf{Y}_{\Sigma^2}(t, s, y^2))_\#\Pi_s$ is a coupling between Σ_t^1 and Σ_t^2 we can write

$$W_1(\Sigma_t^1, \Sigma_t^2) \leq \int_{C \times C} \|\mathbf{Y}_{\Sigma^1}(t, s, y^1) - \mathbf{Y}_{\Sigma^2}(t, s, y^2)\| \, d\Pi_s(y^1, y^2).$$

By (4.9) we get

$$W_1(\Sigma_t^1, \Sigma_t^2)$$

$$\leq e^{L(t-s)} \int_{C \times C} \|y^1 - y^2\| d\Pi_s(y^1, y^2) + L \int_s^t e^{L(t-\tau)} W_1(\Sigma_\tau^1, \Sigma_\tau^2) d\tau$$

$$= e^{L(t-s)} W_1(\Sigma_s^1, \Sigma_s^2) + L \int_s^t e^{L(t-\tau)} W_1(\Sigma_\tau^1, \Sigma_\tau^2) d\tau.$$

Choosing $t = s + h$ proves that the upper right derivative d/ds_+ of the map $s \mapsto W_1(\Sigma_s^1, \Sigma_s^2)$ satisfies

$$\frac{d}{ds_+} W_1(\Sigma_s^1, \Sigma_s^2) \leq 2L W_1(\Sigma_s^1, \Sigma_s^2)$$

and therefore $W_1(\Sigma_s^1, \Sigma_s^2) \leq e^{2Ls} W_1(\bar{\Sigma}^1, \bar{\Sigma}^2)$, which proves (4.13). $\square$

Remark 4.8 (Another existence proof). The following argument can provide an alternative strategy to constructing a solution starting from the discrete solutions of the previous section. In fact, one can use the contractivity to pass to the limit in the discrete problem. Choose $\bar{y}_i(\omega) \in C$ independent and identically distributed, with law $\bar{\Sigma}$, so that the random measures $\bar{\Sigma}^N(\omega) := \frac{1}{N} \sum_{i=1}^N \delta_{\bar{y}_i(\omega)}$ almost surely converge in $\mathscr{P}_1(C)$ to $\bar{\Sigma}$. We fix ω such that this happens, set $\bar{y}_i(\omega) = \bar{y}_i$, and let $y(t) = (y_1(t), \ldots, y_N(t))$ be the discrete evolution starting from $y(t) = (\bar{y}_1, \ldots, \bar{y}_N)$. Then, by contractivity, $\Sigma_t^N := \frac{1}{N} \sum_{i=1}^N \delta_{y_i(t)}$ converges weakly, and it is not hard to prove that $\Sigma_t = \lim_{N \to \infty} \Sigma_t^N$ provides a solution.

5 Uniqueness of Eulerian Solutions

In this section we address the uniqueness of Eulerian solutions, according to (3.10). Our first proof uses a classical duality argument, adapted to the infinite-dimensional space of measures and to the special structure

$$(5.1) \qquad b_\Sigma(t, c) = b_{\Sigma_t}(c) = \int_C f(c, c') d\Sigma_t(c')$$

of the vector field, with $f = (f_x, f_\sigma)$ as in (3.3) and (3.4). One of the advantages of the duality proof is that it provides uniqueness in the larger class of signed measures; the drawback is that, since we don't have at our disposal the mollification schemes of the finite-dimensional setting, we have to require C^1 regularity in place of Lipschitz regularity with respect to the x-variable of e and J. We use the special structure of the vector field, together with (3.16), also to make use of the flow map $\mathbf{Y}(s, t, x)$ backward in time, i.e., for times $s \leq t$; indeed, (3.16) yields that the abstract compatibility condition (4.1) holds also for $-b$, whose forward solutions correspond to backward solutions for b.

In this section we shall apply the abstract calculus tools of Appendices A.1 and B with $C = \mathbb{R}^d \times \mathscr{P}(U)$, $E = Y = \mathbb{R}^d \times F(U)$, the subspace $E_C = Y_C = \mathbb{R}^d \times \mathscr{M}_0(U)$, and its closure $\bar{E}_C = \mathbb{R}^d \times \{\sigma \in F(U) : \sigma(1) = 0\}$.

5.1 Uniqueness by duality

THEOREM 5.1. *Suppose that $J\colon (\mathbb{R}^d \times U)^2 \to \mathbb{R}$ and $e\colon \mathbb{R}^d \times U \to \mathbb{R}$ are Lipschitz maps, with $J(\cdot, u, x', u')$ of class C^1 for all $(u, x', u') \in U \times \mathbb{R}^d \times U$, and $e(\cdot, u)$ of class C^1 for all $u \in U$. Then, for all $\bar{\Sigma} \in \mathscr{M}(C)$ with $\int_C |x| \mathrm{d} |\bar{\Sigma}| < +\infty$, equation (3.10) admits a unique solution in the class of weakly continuous maps $t \in [0, T] \mapsto \Sigma_t \in \mathscr{M}(C)$ with $\sup_t \int_C (1 + |x|) \mathrm{d} |\Sigma_t| < +\infty$ and $\Sigma_0 = \bar{\Sigma}$.*

PROOF. Let us consider solutions Σ^1 and Σ^2 of (3.10) such that $\Sigma_0^1 = \Sigma_0^2$ and fix $h > 0$. Let us denote

$$\mathcal{R} := \left\{ r\colon [0, h] \to \mathrm{C}_b(Y)\colon r \text{ is Borel}, \ r_t \in \mathrm{C}_b^1(Y), \ \mathrm{Lip}(r_t) \leq 1 \text{ for all } t \in [0, h] \right\},$$

with the usual notation $r_t(y) = r(t, y)$.

The difference $\Sigma_t := \Sigma_t^1 - \Sigma_t^2$ solves

$$\frac{\mathrm{d}}{\mathrm{d}t} \Sigma_t + \mathrm{div}(b_{\Sigma_t^1} \Sigma_t) = -\mathrm{div}(b_t \Sigma_t^2)$$

in the weak sense of (3.10), with $b_t := b_{\Sigma_t^1} - b_{\Sigma_t^2}$. Let us stress that b_t has null first component, since it is the difference of the vector fields $b_{\Sigma_t^1}$ and $b_{\Sigma_t^2}$ that have the *same* first component (recall the notation (3.6)):

$$b_{\Sigma_t^i}(x, \sigma) = \left(\int_U e(x, u) \mathrm{d}\sigma(u), \left(J * \Sigma_t^i(x, \cdot) - \int_U J * \Sigma_t^i(x, u) \mathrm{d}\sigma(u) \right) \sigma \right),$$

for $i = 1, 2$. By linearity of the second components with respect to Σ we obtain the representation

$$b_t(x, \sigma) = \left(0, \left(J * \Sigma_t(x, \cdot) - \int_U J * \Sigma_t(x, u) \, \mathrm{d}\sigma(u) \right) \sigma \right).$$

We consider a bounded solution $g \in \mathrm{C}^1([0, h] \times C)$ of the backward transport equation with velocity field $b_{\Sigma_t^1}$, right-hand side r_t, and terminal condition $g_h = 0$ (with the usual notation $g_s(c) = g(s, c)$):

$$(5.2) \qquad \begin{cases} \partial_t g_t + \mathrm{D} g_t(b_{\Sigma_t^1}) = r_t & \text{in } [0, h] \times C, \\ g_h = 0 & \text{in } C. \end{cases}$$

We stress now that the differentiations acting on g_t with respect to c are all meant in the Fréchet sense (see Definition A.1) and that the pairing $\mathrm{D} g_t(b_{\Sigma_t^1})$ corresponds to the directional derivative of g_t along the vector $b_{\Sigma_t^1} \in Y_C$. By the classical method of characteristics, one can construct a solution g of (5.2) with the required C^1 regularity property by setting

$$(5.3) \qquad g_t(c) = -\int_t^h r_s(\mathbf{Y}_{\Sigma^1}(s, t, c)) \mathrm{d}s.$$

Indeed, the $\mathrm{C}_b^1(Y)$ regularity of r_t, the C^1 regularity of $\mathbf{Y}_{\Sigma^1}(s, t, \cdot)$ granted by Theorem B.2 (and the arguments below; see before formula (5.4)), and Theorem A.2 yield the C^1 regularity of the function g in (5.3), together with the exchange of

Fréchet differentiation with integration (the latter granted by (A.9)). Obviously $g_h = 0$, and one can check that g satisfies (5.2) with the following observation: since, thanks to (A.2),

$$\frac{\mathrm{d}}{\mathrm{d}t} g(t, c(t)) = \big(\partial_t g_t + \mathrm{D}g_t(b_{\Sigma_t^1})\big)(t, c(t))$$

along *any* C^1 solution $c(t)$ of the ODE

$$\frac{\mathrm{d}}{\mathrm{d}t} c(t) = b_{\Sigma_t^1}(c(t)),$$

if for any $c_0 \in C$ and $t_0 \in (0, h)$ we are able to find a C^1 solution $c(t)$ to the ODE above with $c(t_0) = c_0$ and $\frac{\mathrm{d}}{\mathrm{d}t} g(t, c(t)) = r_t(c(t))$ at $t = t_0$, we are done. Choosing $c(t) = \mathbf{Y}_{\Sigma^1}(t, 0, d)$ for some $d \in C$, from the semigroup property we get

$$g(t, c(t)) = -\int_t^h r_s(\mathbf{Y}_{\Sigma^1}(s, 0, d))\mathrm{d}s$$

so that we are able to check (5.2) at any (t_0, c_0) with $c_0 = \mathbf{Y}_{\Sigma^1}(t_0, 0, d)$. Choosing $d = \mathbf{Y}_{\Sigma^1}(0, t_0, c_0)$ (only at this point we are using the flow backwards in time), we obtain the global validity of (5.2).

As mentioned above, the C^1 regularity of $\mathbf{Y}_{\Sigma^1}(s, t, \cdot)$ follows from Theorem B.2 if we check that

$$(5.4) \qquad (t, c) \mapsto \mathrm{D}_{Y_C} A(t, \cdot)(c) \text{ is continuous from } [0, T] \times C \text{ to } \mathcal{L}(Y_C, \bar{Y}_C)$$

in the sense of (B.5), with

$$A(t, c) = b_{\Sigma_t^1}(c),$$

and we recall Proposition 4.3. It is at this stage that we need the extra C^1 assumption on $J(\cdot, u, x', u')$ and $e(\cdot, u)$. Indeed, thanks to the representation (5.1) of b_{Σ^1}, it is sufficient to check C^1 differentiability of $f(\cdot, c')$ for all $c' = (x', \sigma')$ in the direction $(v, \theta) \in Y_C$; the partial differential with respect to the x-variable at $c = (x, \sigma)$ is given by

$$v \in \mathbb{R}^d \mapsto \Bigg(\int_U \mathrm{D}_x e(x, u)(v)\mathrm{d}\sigma(u),$$

$$\Bigg(\int_U \mathrm{D}_x J(x, \cdot, x', u')(v)\mathrm{d}\sigma'(u')$$

$$- \int_{U \times U} \mathrm{D}_x J(x, w, x', u')(v)\mathrm{d}\sigma'(u')\mathrm{d}\sigma(w) \Bigg)\sigma \Bigg)$$

while a partial differential with respect to the σ-variable is given by

$$\theta \in \mathscr{M}_0(U) \mapsto \left(\int_U e(x,u)\mathrm{d}\theta(u), \left(\int_U J(x,\cdot,x',u')\mathrm{d}\sigma'(u') \right. \right.$$
$$\left. - \int_{U \times U} J(x,w,x',u')\mathrm{d}\sigma'(u')\mathrm{d}\sigma(w) \right)\theta$$
$$\left. - \left(\int_{U \times U} J(x,w,x',u')\mathrm{d}\sigma'(u')\mathrm{d}\theta(w) \right)\sigma \right),$$

so that the differential is continuous from $[0,T] \times C$ to $\mathcal{L}(Y_C, \bar{Y}_C)$.

Since r_t are 1-Lipschitz, and the Lipschitz constant of $b_{\Sigma_t^1}$ can be estimated from above by $L_* = L \sup_t |\Sigma_t^1|(C)$, with the L Lipschitz constant of the interaction term f in (5.1), from the Lipschitz estimate on $\mathbf{Y}_{\Sigma^1}(t,s,\cdot)$ granted by Theorem B.2, we get

$$(5.5) \qquad \|\mathrm{D}g_t(c)\|_{\mathcal{L}(Y_0,Y)} \le (h-t)\mathrm{e}^{L_* T} \quad \text{for all } (t,c) \in [0,h] \times C.$$

Since $g_t \Sigma_t$ vanishes at $t = 0$ and $t = h$, by the Leibniz rule and Theorem A.2 we get

$$0 = \int_0^h \frac{\mathrm{d}}{\mathrm{d}t} \int_C g_t \, \mathrm{d}\Sigma_t \, \mathrm{d}t$$
$$= \int_0^h \int_C (\partial_t g_t + \mathrm{D}g_t(b_{\Sigma_t^1}))\mathrm{d}\Sigma_t \, \mathrm{d}t + \int_0^h \int_C \mathrm{D}g_t(b_{\Sigma_t^1})\mathrm{d}\Sigma_t^2 \, \mathrm{d}t.$$

Using (5.2), we obtain that

$$(5.6) \qquad \sup_r \int_0^h \int_C r_t \, \mathrm{d}\Sigma_t \, \mathrm{d}t \le \int_0^h \int_C |\mathrm{D}g_t(b_t)| \mathrm{d}|\Sigma_t^2| \mathrm{d}t.$$

Motivated by this estimate, we employ the distance $\mathsf{d}_{\mathrm{C}^1}$ in $\mathscr{M}(C)$ of Lemma 2.1. Let us prove now that

$$(5.7) \qquad \sup_{r \in \mathcal{R}} \int_0^h \int_C r_t \, \mathrm{d}\Sigma_t \, \mathrm{d}t = \int_0^h \mathsf{d}_{\mathrm{C}^1}(\Sigma_t^1, \Sigma_t^2)\mathrm{d}t.$$

The inequality $\le$ is obvious, since r_t is an admissible function in the definition of $\mathsf{d}_{\mathrm{C}^1}$ for any $t \in [0,h]$. To prove the converse, we apply a measurable selection argument: since Y is separable, it is easily seen that

$$B := \left\{ v \in \mathrm{C}_b(Y): v \in \mathrm{C}_b^1(Y), \, \mathrm{Lip}(v) \le 1 \right\}$$

is a Borel and separable subset of $\mathrm{C}_b(Y)$, and that $t \mapsto \mathsf{d}_{\mathrm{C}^1}(\Sigma_t, \Sigma_t')$ is a Borel function. Then, for $\delta > 0$ fixed we consider the set

$$\Gamma := \left\{ (t,v) \in [0,h] \times \mathrm{C}_b(Y): v \in B, \, \mathsf{d}_{\mathrm{C}^1}(\Sigma_t^1, \Sigma_t^2) < \delta + \int_Y v \, \mathrm{d}\Sigma_t \right\},$$

which is measurable, thanks to the above-mentioned properties, with respect to the product of the Borel σ-algebras. Then, a measurable selection theorem [7, theorem 6.9.1] grants the existence of a Borel selection map $\bar{r} \colon [0, h] \to B$ satisfying $(t, \bar{r}_t) \in \Gamma$ for a.e. $t \in [0, h]$. Since, by construction, $\bar{r} \in \mathcal{R}$, it follows that

$$\sup_{r \in \mathcal{R}} \int_0^h \int_C r_t \, d\Sigma_t \, dt \geq \int_0^h \int_C \bar{r}_t \, d\Sigma_t \, dt$$

$$\geq -\delta \int_0^h |\Sigma_t|(C) dt + \int_0^h d_{C^1}(\Sigma_t^1, \Sigma_t^2) dt.$$

Since $\delta > 0$ is arbitrary, this proves (5.7).

Combining (5.6) and (5.7) we obtain

$$(5.8) \qquad \int_0^h d_{C^1}(\Sigma_t^1, \Sigma_t^2) dt \leq \int_0^h \int_C |Dg_t(b_t)| d|\Sigma_t^2| dt.$$

It remains to estimate from above the right-hand side in (5.8). Recalling that the norm on Y is given by $\|y\|_Y = \|(x, \sigma)\|_Y = |x| + \|\sigma\|_{\mathrm{BL}}$ and using (5.5), we have

$$\int_0^h \int_C |Dg_t(b_t)| d|\Sigma_t^2| dt$$

$$\leq \int_0^h \int_C \|Dg_t\|_{\mathcal{L}(Y_0, Y)} \|b_t(y)\|_Y \, d|\Sigma_t^2|(y) dt$$

$$\leq e^{L_*T} \int_0^h (h - t)$$

$$\cdot \int_C \left\| \left(0, \left(J * \Sigma_t(x, \cdot) - \int_U J * \Sigma_t(x, w) d\sigma(w) \right) \sigma \right) \right\|_Y d|\Sigma_t^2|(y) dt$$

$$\leq e^{L_*T} \int_0^h (h - t)$$

$$\cdot \int_C \left\| \left(J * \Sigma_t(x, \cdot) - \int_U J * \Sigma_t(x, w) d\sigma(w) \right) \sigma \right\|_{\mathrm{BL}} d|\Sigma_t^2|(y) dt.$$

Moreover, (2.8) gives

$$\left\| \left(J * \Sigma_t(x, \cdot) - \int_U J * \Sigma_t(x, w) d\sigma(w) \right) \sigma \right\|_{\mathrm{BL}} \leq 2 \int_U |J * \Sigma_t(x, u)| d\sigma(u),$$

and

$$|J * \Sigma_t(x, u)| \leq L_J d_{C^1}(\Sigma_t^1, \Sigma_t^2),$$

so that, with $S = \sup_t |\Sigma_t^2|(C)$, we get

$$(5.9) \qquad \int_0^h \int_C |Dg_t(b_t)| d|\Sigma_t^2| dt \leq 2S e^{L_*T} L_J h \int_0^h d_{C^1}(\Sigma_t^1, \Sigma_t^2) dt.$$

Combining (5.8) and (5.9) we obtain

$$\int_0^h \mathsf{d}_{\mathrm{C}^1}(\Sigma_t^1, \Sigma_t^2)\mathrm{d}t \le 2S\mathrm{e}^{L*T}L_J h \int_0^h \mathsf{d}_{\mathrm{C}^1}(\Sigma_t^1, \Sigma_t^2)\mathrm{d}t.$$

For $h > 0$ small enough such that $2S\mathrm{e}^{L*T}L_J h < 1$, one has

$$\mathsf{d}_{\mathrm{C}^1}(\Sigma_t^1, \Sigma_t^2) = 0 \quad \text{for all } 0 \le t \le h,$$

so that the curves Σ^1 and Σ^2 coincide in $[0, h]$ and the proof is achieved by repeating this argument finitely many times. $\qquad\square$

5.2 Uniqueness by superposition

In this section we prove uniqueness of Eulerian solutions, as defined in Definition 3.2, under the same assumptions of Theorem 3.3, dealing with Lagrangian solutions. In particular, we require the sole Lipschitz continuity of e and J, not requiring the C^1 smoothness of $J(\cdot, u, x', u')$ and $e(\cdot, u)$ of Theorem 5.1. The proof covers the more general setting of interaction systems in Banach spaces of Section 4.1; see Theorem 4.1 for the existence and stability of Lagrangian solutions.

Our main tool in the proof is the so-called superposition principle: it allows us to lift solutions to the continuity equation to probability measures on paths, thus recovering an (extended) Lagrangian representation; the principle, which remarkably works under no regularity assumption on the vector field, has by now many versions; see, for instance, [2, theorem 8.2.1] in Euclidean spaces and [39] in the context of the theory of currents. Here we consider the case when the state space is a separable Banach space.

THEOREM 5.2. *Let* $(Y, \|\cdot\|_Y)$ *be a separable Banach space, let* $b\colon (0, T)\times Y \to Y$ *be a Borel vector field, and let* $\mu_t \in \mathscr{P}(Y)$, $t \in [0, T]$, *be a continuous curve with*

(5.10) $$\int_0^T \int_Y \|b_t\|_Y \, \mathrm{d}\mu_t \, \mathrm{d}t < +\infty.$$

If

$$\frac{\mathrm{d}}{\mathrm{d}t}\mu_t + \mathrm{div}(b_t\mu_t) = 0$$

in duality with cylindrical functions $\phi \in \mathrm{C}_b^1(Y)$, *precisely of the form (here* $\langle \cdot, \cdot \rangle$ *denotes the duality map between* Y *and* Y'),

$$\varphi(\langle y, z_1' \rangle, \langle y, z_2' \rangle, \ldots, \langle y, z_N' \rangle)$$

with $\varphi \in \mathrm{C}_b^1(\mathbb{R}^N)$ *and* $z_1', \ldots, z_N' \in Y'$, *then there exists* $\eta \in \mathscr{P}(\mathrm{C}([0, T]; Y))$ *concentrated on absolutely continuous solutions to the ODE* $\dot{y} = b_t(y)$ *and with* $\mathrm{ev}_t(\eta) = \mu_t$ *for all* $t \in [0, T]$.

PROOF. We start from the version of the superposition principle in the space $\mathbb{R}^\infty$ of all sequences (x_i), $i \geq 1$, proved in [4, theorem 7.1] by finite-dimensional approximation. We endow $\mathbb{R}^\infty$ with the distance

$$\mathsf{d}_\infty(x, y) := 2^{-n} \min\{1, |x_n - y_n|\},$$

which makes it a complete and separable metric space. If $c\colon (0, T) \times \mathbb{R}^\infty \to \mathbb{R}^\infty$ is a vector field, with components c^i measurable with respect to the product of the Borel σ-algebras in the domain, and if $\nu_t \in \mathscr{P}(\mathbb{R}^\infty)$ satisfy $\int_0^T \int_{\mathbb{R}^\infty} |c_t^i| \mathrm{d}\nu_t\, \mathrm{d}t < +\infty$ for any i and solve the continuity equation

$$\frac{\mathrm{d}}{\mathrm{d}t} \nu_t + \mathrm{div}(c_t \nu_t) = 0$$

in duality with C_b^1 cylindrical function ϕ (i.e., dependent on finitely many coordinates x_i), there exists $\sigma \in \mathscr{P}(\mathrm{C}([0, T]; \mathbb{R}^\infty))$ such that

(1) σ is concentrated on continuous curves $\gamma\colon [0, T] \to \mathbb{R}^\infty$, with absolutely continuous components γ^i solving the infinite system of ODE $\dot{\gamma}^i = c_t^i(\gamma)$, $i \geq 1$;

(2) $\mathrm{ev}_t(\sigma) = \nu_t$ for all $t \in [0, T]$.

Given this basic result, if we strengthen the integrability assumption on c by requiring

(5.11) $$\int_0^T \int_{\mathbb{R}^\infty} \Psi(\sup_i |c_t^i|) \mathrm{d}\nu_t\, \mathrm{d}t < +\infty$$

with Ψ a nondecreasing function such that $\Psi(z)/z \to +\infty$ as $z \to \infty$, then it is immediately seen that σ is concentrated on a Borel set Θ (i.e., $\sigma(\mathrm{C}([0, T]; \mathbb{R}^\infty) \setminus \Theta) = 0$) made of curves γ absolutely continuous with respect to the norm $\|\cdot\|_\infty$. More precisely, one has

$$\sup_i |\gamma_s^i - \gamma_t^i| \leq \int_s^t \sup_i |c_\tau^i| \circ \gamma_\tau\, \mathrm{d}\tau \quad \text{for all } 0 \leq s \leq t \leq T \text{ and all } \gamma \in \Theta,$$

$$\int_0^T \Psi(\sup_i |c_\tau^i| \circ \gamma_\tau) \mathrm{d}\tau < +\infty \quad \text{for all } \gamma \in \Theta.$$

Now we turn to the case of a separable Banach space Y. Thanks to the Dunford-Pettis theorem when applied to the space-time measure $\mu_t\, \mathrm{d}t$, we can find a nondecreasing function $\Psi\colon [0, +\infty) \to [0, +\infty)$ with $\Psi(z)/z \to +\infty$ as $z \to +\infty$ such that (5.10) improves to

(5.12) $$\int_0^T \int_Y \Psi(\|b_t\|_Y) \mathrm{d}\mu_t\, \mathrm{d}t < +\infty.$$

Let $(z_i') \subset Y'$ as in the proof of Lemma 2.1 and let us consider the map $E\colon Y \to \ell_\infty \subset \mathbb{R}^\infty$ defined by

$$E(y) := (\langle y, z_1' \rangle, \langle y, z_2' \rangle, \dots).$$

It is immediately seen from (2.3) that the mapping E is an isometry of Y into ℓ_∞; hence, $E(Y)$ is closed and separable in ℓ_∞. In addition, [7, theorem 6.8.6] grants that E maps Borel sets of Y into Borel sets of $\mathbb{R}^\infty$; in particular, $E(Y)$ is a Borel set of $\mathbb{R}^\infty$. As a consequence, E^{-1} extended to 0 out of $E(Y)$ is μ-measurable for any $\mu \in \mathscr{P}(\mathbb{R}^\infty)$.

Let us consider the measures $\nu_t = E_\# \mu_t \in \mathscr{P}(\mathbb{R}^\infty)$; if we define

$$c_t^i := \langle b_t, z_i' \rangle \circ E^{-1} \quad \text{on } E(Y) \text{ for } i \geq 1$$

(and equal to 0 on $\mathbb{R}^\infty \setminus E(Y)$, this extension is irrelevant; since the measures ν_t are concentrated on $E(Y)$), we obtain by construction that the continuity equation holds, with the stronger integrability condition (5.11) coming from (5.12), since $\sup_i |c_t^i| \leq \|b_t\|_Y$ (recall that Ψ is nondecreasing).

Then, from the superposition theorem in $\mathbb{R}^\infty$ we obtain a probability measure $\sigma \in \mathscr{P}(C([0, T]; \mathbb{R}^\infty))$ concentrated on solutions of the ODE $\dot{\gamma}_t^i = c_t^i(\gamma)$ and absolutely continuous with respect to the norm $\|\cdot\|_\infty$. Moreover, since ν_t are concentrated on $E(Y)$, we obtain that, for any $t \in [0, T]$, $\gamma_t \in E(Y)$ for σ-a.e. γ. In particular, restricting t to a countable set, for σ-a.e. γ, one has

(5.13) $$\gamma_t \in E(Y) \quad \text{for any } t \in \mathbb{Q} \cap [0, T].$$

If we fix $\gamma \in \Theta$ (so that γ is absolutely continuous with respect to the $\|\cdot\|_\infty$ norm) and (5.13) holds, since our choice of the z_i' guarantees that the composition with E^{-1} is an isometry between the norm $\|\cdot\|_\infty$ and $\|\cdot\|_Y$, from $\gamma_t \in E(Y)$ for any $t \in \mathbb{Q} \cap [0, T]$, we deduce $\gamma([0, T]) \subset E(Y)$: indeed, if $t_n \in \mathbb{Q} \cap [0, T]$ and $t_n \to t$, $\gamma_{t_n} = E(y_n)$, then the absolute continuity of γ yields that (y_n) is a Cauchy sequence in Y; hence $y_n \to y$ for some $y \in Y$ and then $\gamma_t = E(y) \in E(Y)$. Therefore, we proved that for all $\gamma \in \Theta$ the property (5.13) improves to $\gamma([0, T]) \subset E(Y)$, and the same argument shows that the transformed curve $E^{-1}\gamma$ in Y is absolutely continuous.

The ODE $\dot{\gamma}_t^i = c_t^i(\gamma)$ becomes, for the transformed curve $y_t = E^{-1}\gamma_t$,

$$\frac{\mathrm{d}}{\mathrm{d}t} \langle y, z_i' \rangle = \langle b_t \circ y, z_i' \rangle \quad \text{for all } i \geq 1,$$

since the z_i' separate points in Y', which allows us to conclude that $\dot{y} = b_t \circ y$.

Let us now consider the map $\widetilde{E} \colon C([0, T]; Y) \to C([0, T]; \mathbb{R}^\infty)$, $\widetilde{E}(\gamma) := E \circ \gamma$, which naturally extends E to the corresponding complete and separable spaces of curves. Since $\widetilde{E}$ is continuous and injective, it maps Borel sets into Borel sets. It follows that the inverse map $E_*, \gamma \mapsto y = \widetilde{E}^{-1}\gamma$, arbitrarily defined to a constant on the Borel and σ-negligible set

$$C([0, T]; \mathbb{R}^\infty) \setminus \widetilde{E}\big(C([0, T]; Y)\big) \subset C([0, T]; \mathbb{R}^\infty) \setminus \Theta,$$

is Borel.

To conclude, having set $\eta = (E_*)_\# \sigma$, from $\mathrm{ev}_t(\sigma) = \nu_t$ one obtains $\mathrm{ev}_t(\eta) = \mu_t$ for all $t \in [0, T]$. The measure η then satisfies all stated properties. $\square$

THEOREM 5.3. *Let $\bar{\Sigma} \in \mathscr{P}_1(C)$, $f\colon C \times C \to Y$ be L-Lipschitz, and let b_Σ be defined as in* (5.1). *Then there is a unique $\Sigma \in C([0, T]; \mathscr{P}_1(C))$ with $\Sigma_0 = \bar{\Sigma}$, satisfying*

$$\int_C \phi(t, y)\,\mathrm{d}\Sigma_t(y) - \int_C \phi(0, y)\,\mathrm{d}\Sigma_0(y)$$
$$= \int_0^t \int_C \left(\partial_s \phi(s, y) + \mathrm{D}\phi(s, y)(b_{\Sigma_s}(y))\right)\mathrm{d}\Sigma_s(y)\,\mathrm{d}s$$

for all $\phi \in C_b^1([0, T] \times Y)$.

PROOF. Let us consider Eulerian solutions Σ^1 and Σ^2 starting from the same initial datum $\bar{\Sigma}$, and let us denote by $b^1(t, y)$, $b^2(t, y)$ the respective velocity fields and by $\mathbf{Y}^1(t, y)$, $\mathbf{Y}^2(t, y)$ the respective flow maps. By applying the superposition theorem to b^i (extended with the 0 value to $(0, T) \times (Y \setminus C)$), Σ^i we obtain that $\Sigma_t^i = \mathrm{ev}_t(\eta^i)$ for suitable $\eta^i \in \mathscr{P}(C([0, T]; Y))$ concentrated on absolutely-continuous-in-$[0, T]$ solutions to the Cauchy problem

$$\dot{y} = b_{\Sigma^i}(t, y) \quad \text{in } (0, T).$$

On the other hand, since from Theorem B.1 we know that the solution to the Cauchy problem is unique, the conditional probabilities η_y^i of η^i given the initial condition $y(0) = y$ have to be Dirac masses, precisely $\eta_y^i = \delta_{\mathbf{Y}^i(\cdot, y)}$. It follows that $\Sigma_t^i = \mathbf{Y}^i(t, \cdot)_\# \bar{\Sigma}$ for all $t \in [0, T]$. An application of the stability estimate (4.4) in Theorem 4.1 yields uniqueness. $\qquad\square$

Appendix A   Calculus in Banach Spaces

We adapt some basic calculus notions in Banach spaces to our framework, where the domain C of the functions we wish to differentiate is a convex subset of a normed space E but need not be open; in this case we denote by E_C the vector space $\mathbb{R}(C - C)$. For $c \in C$, we shall instead denote by E_c the convex cone of directions

$$E_c := \mathbb{R}_+(C - c) \subset E_C.$$

A.1 Differentiation

We first introduce a notion of multivalued Fréchet differential, adapted to functions defined on convex sets.

DEFINITION A.1 (Multivalued F-differential). Let E, F be normed vector spaces, $C \subset E$ convex, and let $f\colon C \to F$ be a map. We say that f is F-differentiable at $c \in C$ if there exists $L \in \mathcal{L}(E_C, F)$ such that

(A.1)
$$\lim_{C \ni c' \to c} \frac{\|f(c') - f(c) - L(c' - c)\|_F}{\|c' - c\|_E} = 0.$$

This notion is too strong for some applications, since the natural domain of L should be only the closure of the cone E_c, on which L is uniquely determined by (A.1). We denote the F-differential of f in $c \in C$ by

$$D f(c) = \{L \in \mathcal{L}(E_C, F) : L \text{ fulfills (A.1)}\},$$

and, if E_c is not dense in E_C, the map D is multivalued (as in the case of the subdifferential in convex analysis). By density, each $L \in D f(c)$ uniquely extends to an operator in $\mathcal{L}(\bar{E}_C, F)$. Hence, the F-differential $D f(c)$ is a closed convex subset of $\mathcal{L}(\bar{E}_C, F)$. For any $e \in E_C$ we denote

$$D f(c)e = D f(c)(e) = \{L(e) = Le : L \in D f(c)\}.$$

If $e \in \bar{E}_c$, then $D f(c)e$ is a singleton and in this case, with a slight abuse of notation, we may use $D f(c)e$ instead of Le for any $L \in D f(c)$.

A straightforward consequence of (A.1) is the chain rule for curves: if $t \mapsto c(t) \in C^1([0, T]; C)$ is a differentiable map, then $c'(t) \in \bar{E}_{c(t)} \subset \bar{E}_C$ and

$$(A.2) \qquad \left. \frac{d}{ds} f(c(s)) \right|_{s=t} = D f(c(t))(c'(t))$$

whenever f is F-differentiable at $c(t)$ (we remark on the abuse of notation mentioned above).

For F-differentiability we can also adopt the handy notation

$$f(c') - f(c) - D f(c)(c' - c) = o(\|c' - c\|_E).$$

If we consider the particular case when $E = F$ and $f : C \to C$, obviously any $L \in D f(c)$ is a linear operator from E_C to $\bar{E}_C$. With these notions, the proof of the following chain rule is standard.

THEOREM A.2 (Chain rule). *Suppose $f : C \to C$ and $g : C \to \mathbb{R}$ are F-differentiable, respectively, at $c \in C$ and at $d = f(c) \in C$. Then $g \circ f : C \to \mathbb{R}$ is also F-differentiable at c and*

$$D(g \circ f)(c) \subset D g(d) \circ D f(c),$$

where $D g(d) \circ D f(c) = \{M \circ L : M \in D g(d), L \in D f(c)\}$.

Notice that if $e \in E_c$, then $D f(c)(e)$ is unique (i.e., it does not depend on the choice of L) and belongs to $\bar{E}_d$. Therefore, also the choice of the element $M \in D g(d)$ is irrelevant.

PROOF. It suffices to write

$$g \circ f(c') - g \circ f(c) = g(d + D f(c)(c' - c) + R(c')) - g(d)$$

with $\|R(c')\|_E = o(\|c' - c\|_E)$ as $c' \to c$, and the expansion $g(d') - g(d) = D g(d)(d' - d) + o(\|d' - d\|_E)$ with $d' = d + D f(c)(c' - c) + R(c')$. $\square$

We say that f is of class C^1 and we write $f \in C^1(C; F)$ if f is F-differentiable at each $c \in C$, and there exists a selection $L(c) \in Df(c)$ for all $c \in C$ such that

(A.3) $c \mapsto L(c)$ is continuous from C to $\mathcal{L}(E_C, F)$

with $\mathcal{L}(E_C, F)$ endowed with the distance induced by the operator norm. To conform with the classical continuous Fréchet differentiability, below we may write with a slight abuse of notation that

(A.4) $c \mapsto Df(c)$ is continuous from C to $\mathcal{L}(E_C, F)$

to actually mean $f \in C^1(C; F)$ as in (A.3).

In the context of Theorem A.2, by choosing a continuous selection for $Df(c)$ and for $Dg(d)$, we have a continuous selection for the F-differential of the composition at c, $D(g \circ f)(c)$, thus granting the C^1 regularity of $g \circ f$.

DEFINITION A.3 (G-differentiation). Let E, F be normed vector spaces, $C \subset E$ convex, and let $f : C \to F$ be a map. We say that f is G-differentiable at $c \in C$ if the directional right derivatives

$$d f(c, e) := \lim_{h \to 0^+} \frac{f(c + he) - f(c)}{h} \text{ exists in } F \text{ for all } e \in E_c.$$

Of course, F-differentiability at c implies G-differentiability at c, with

$$d f(c, e) = Df(c)(e)$$

for all $e \in E_c$. In connection with the differentiability properties of the flow map, it is useful to establish the converse implication.

LEMMA A.4 (Criterion for C^1 regularity). *Let $f : C \to F$ be a continuous map and assume the existence of a continuous operator*

$$C \ni c \mapsto L(c) \in \mathcal{L}(E_C, F)$$

such that $d f(c, e) = L(c)e$ for all $c \in C$ and all $e \in E_c$. Then $L(c) \in Df(c)$ is an admissible choice in (A.1) *for all $c \in C$, so that $f \in C^1(C; F)$.*

PROOF. Set $c_t = c + t(c' - c)$ for $t \in [0, 1]$ and set $e = c' - c \in E_c$, and notice that $\pm e \in E_{c_t}$ for all $t \in (0, 1)$. Since $t \mapsto f(c_t)$ is continuous in $[0, 1]$ and differentiable in $(0, 1)$, we can write

(A.5) $f(c') - f(c) = \int_0^1 d f(c_t, e) dt = \int_0^1 L(c_t)e\, dt.$

Now we use that $L(c_t)e/\|e\|_E \to L(c)e/\|e\|_E$ uniformly in $[0, 1]$ as $c' \to c$ to obtain (A.1).   □

Remark A.5. Formula (A.5) also shows that a function $f \in C^1(C; F)$ with a uniformly bounded F-differential selection $L: c \mapsto L(c)$ as in (A.3), which means $L \in C_b(C; \mathcal{L}(E_C, F))$ satisfies a uniform bound

$$\|f(c)\|_F \le \|f(c_0)\|_F + \|c - c_0\|_E \sup_{c \in C} \|L(c)\|_{\mathcal{L}(E_C, F)},$$

where c_0 is a given point in C. In particular, f has linear growth

$$(A.6) \qquad \sup_{c \in C} \frac{\|f(c)\|_F}{1 + \|c - c_0\|_E} \leq \|f(c_0)\| + \sup_{c \in C} \|L_c\|_{\mathcal{L}(E_C, F)}.$$

The left-hand side of (A.6) defines a norm in the space $C_{lg}(C, F)$ of continuous functions with linear growth (the definition is in fact independent of c_0). The same argument of the proof of Lemma A.4 also shows that the graph of the multivalued operator D

$$(A.7) \qquad \begin{aligned} X := \big\{ (f, L) \in C^1(C, F) \times C_b(C, \mathcal{L}(E_C, F)): \\ L(c) \in Df(c) \text{ for every } c \in C \big\} \end{aligned}$$

between $C^1(C; F)$ and $C_b(C; \mathcal{L}(E_C, F))$ is closed, and thus a Banach space, with respect to the graph norm

$$\|(f, L)\|_X := \sup_{c \in C} \frac{\|f(c)\|_F}{1 + \|c - c_0\|_E} + \sup_{c \in C} \|L(c)\|_{\mathcal{L}(E_C, F)},$$

which in turn is equivalent to $\|f(c_0)\|_F + \sup_{c \in C} \|L(c)\|_{\mathcal{L}(E_C, F)}$.

The closedness of the graph X is equivalent to saying that if $c \mapsto L_n(c) \in Df_n(c)$ is a sequence of maps in $C_b(C; \mathcal{L}(E_C, F))$ uniformly converging to $L \in C_b(C; \mathcal{L}(E_C, F))$ and $\lim_{n \to \infty} f_n(c_0) = f(c_0)$ in F, then f_n is also converging uniformly on bounded sets to a function $f \in C^1(C; F)$ and $L(c) \in Df(c)$ for all $c \in C$.

A.2 Bochner integration

Let $(A, \mathcal{A}, \mu)$ be a σ-finite measure space, and let E be a Banach space. A μ-simple function $f \colon A \to E$ is representable as

$$f = \sum_{n=1}^{N} \chi_{A_n} e_n,$$

where $e_n \in E$ and $A_n \in \mathcal{A}$ with $\mu(A_n) < +\infty$.

DEFINITION A.6 (Bochner integral). A function $f \colon A \to E$ is μ-Bochner integrable if there exist simple functions $f_n \colon A \to E$ such that

    (i) $\lim_n f_n = f$ μ-a.e. (strong μ-measurability);
    (ii) $\lim_n \int_A \|f_n - f\|_E \, d\mu = 0$.

If f is μ-Bochner integrable, then

$$\int_A f \, d\mu = \lim_n \int_A f_n \, d\mu \in E$$

exists and is independent of the sequence (f_n). It is called the Bochner integral of f and satisfies

$$(A.8) \qquad \left\| \int_A f \, d\mu \right\|_E \leq \int_A \|f\|_E \, d\mu.$$

We shall use the fact that, in the case when (A, d_A) is a separable metric space and $\mathcal{A}$ is the Borel σ-algebra, any continuous function $f \colon A \to E$ is strongly μ-measurable. This is a consequence of the Pettis measurability theorem (see, for instance, [46, chap. 5]), since $x \mapsto \langle e', f(x) \rangle$ is continuous, hence μ-measurable, for any $e' \in E'$ (the so-called weak measurability property), and of the fact that the range of f is separable.

A simple criterion for Bochner integrability is the following:

PROPOSITION A.7 (Bochner integrability criterion). *A function that is strongly μ-measurable (as in* (i) *of Definition A.6)* $f \colon A \to E$ *is μ-Bochner integrable if and only if*

$$\int_A \|f\|_E \, \mathrm{d}\mu < +\infty.$$

The Bochner integral commutes in many ways with linear operators; let us illustrate these properties for duality operators, linear operators, and differentiation operators.

(1) For any $x' \in E'$ one has

$$\left\langle \int_A f \, \mathrm{d}\mu, x' \right\rangle_{E \times E'} = \int_A \langle f, x' \rangle_{E \times E'} \, \mathrm{d}\mu$$

and an analogous formula with the predual holds if E is a dual Banach space.

(2) If $f \colon A \to E$ is μ-Bochner integrable and $T \in \mathcal{L}(E, F)$, then $Tf \colon A \to F$ is also μ-Bochner integrable and

$$T \int_A f \, \mathrm{d}\mu = \int_A Tf \, \mathrm{d}\mu.$$

(3) Let E, F be normed spaces. If $(f, L) \colon A \to X$, where X in the Banach space defined in (A.7), and where $L(x, \cdot)$ is a continuous selection in $\mathrm{D}f(x, \cdot)$, is μ-Bochner integrable, then

(A.9) $$\mathrm{D} \int_A f(x, \cdot) \mathrm{d}\mu(x) \ni \int_A L(x, \cdot) \mathrm{d}\mu(x) \in \mathrm{C}_b(C; \mathcal{L}(E_C, F)).$$

This useful formula explains the correct exchange of F-differentiation and Bochner integration. In particular, it provides, in conjunction with item (1) above, the following simple extension of the fundamental theorem of calculus, for which in fact the simpler Riemann integral would be sufficient.

THEOREM A.8. *If $f \colon [a, b] \to F$ is continuous and differentiable in (a, b), and f' extends continuously to $[a, b]$, then*

$$f(b) - f(a) = \int_a^b f'(\xi) \mathrm{d}\xi.$$

Appendix B   Well-Posedness of ODEs in Banach Spaces and Linearization

First we recall Brezis' theorem [11, sec. I.3, theorem 1.4, cor. 1.1] on the well-posedness of ODEs in Banach spaces.

THEOREM B.1. *Let* $(E, \|\cdot\|_E)$ *be a Banach space,* C *a closed convex subset of* E, *and let* $A(t, \cdot): C \rightarrow E$, $t \in [0, T]$, *be a family of operators satisfying the following properties:*

 (i) *There exists a constant* $L \geq 0$ *such that*

(B.1)
$$\|A(t, c_1) - A(t, c_2)\|_E \leq L \|c_1 - c_2\|_E$$
$$\text{for every } c_1, c_2 \in C \text{ and } t \in [0, T].$$

 (ii) *For every* $c \in C$ *the map* $t \mapsto A(t, c)$ *is continuous in* $[0, T]$.
 (iii) *For every* $R > 0$ *there exists* $\theta > 0$ *such that*

(B.2)
$$c \in C, \ \|c\|_E \leq R \quad \Rightarrow \quad c + \theta A(t, c) \in C.$$

Then for every $\bar{c} \in C$ *there exists a unique curve* $c: [0, T] \rightarrow C$ *of class* C^1 *satisfying* $c_t \in C$ *for all* $t \in [0, T]$ *and*

$$\frac{d}{dt} c_t = A(t, c_t) \quad \text{in } [0, T], \ c_0 = \bar{c}.$$

Moreover, if c^1, c^2 *are the solutions starting from the initial data* $\bar{c}^1$, $\bar{c}^2 \in C$, *respectively, we have*

(B.3)
$$\left\| c_t^1 - c_t^2 \right\|_E \leq e^{Lt} \left\| \bar{c}^1 - \bar{c}^2 \right\|_E \quad \text{for every } t \in [0, T].$$

Assume now that $A(t, \cdot): C \rightarrow E$ is a family of operators as required by the assumptions of Theorem B.1, so that the flow map $\mathbf{Y}: [0, T] \times C \rightarrow C$ given by

(B.4)
$$\frac{d}{dt} \mathbf{Y}(t, c) = A(t, \mathbf{Y}(t, c)), \quad \mathbf{Y}(0, c) = c,$$

is well-defined. For simplicity we consider only the flow map starting from $t = 0$, but the same results hold for the full family of transition maps $\mathbf{Y}(t, s, c)$.

Thanks to (B.2), the operators A take their values in E_C. In this section we shall highlight additional conditions on A for the flow map $\mathbf{Y}(t, \cdot)$ to be continuously F-differentiable, uniformly in $t \in [0, T]$. More precisely, we assume that the operators $A(t, \cdot)$ satisfy

(B.5) $(t, c) \mapsto D_{E_C} A(t, \cdot)(c)$ is continuous from $[0, T] \times C$ to $\mathcal{L}(E_C, \bar{E}_C)$.

From (A.4), by (B.5) we mean that a selection $L_A(t, c) \in D_{E_C} A(t, \cdot)(c)$ exists such that $(t, c) \mapsto L_A(t, c)$ is continuous from $[0, T] \times C$ to $\mathcal{L}(E_C, \bar{E}_C)$.

Under this assumption, we prove by classical linearization and stability arguments the existence of the G-derivatives, and then their continuity as functions of the point of differentiation, using eventually Lemma A.4 to obtain the F-differentiability.

THEOREM B.2. *Under the assumptions of Theorem* B.1 *and* (B.5), *the map* $\mathbf{Y}(t, \cdot)$ *is continuously F-differentiable for all* $t \in [0, T]$, *with an F-differential* $D\mathbf{Y}$ *satisfying:*

(i) $(t, c) \mapsto D\mathbf{Y}(t, \cdot)(c)$ *continuous from* $[0, T] \times C$ *to* $\mathcal{L}(E_C, \bar{E}_C)$;

(ii) $\|D\mathbf{Y}(t, \cdot)(c)\|_{\mathcal{L}(E_C, \bar{E}_C)} \le e^{LT}$ *for all* $(t, c) \in [0, T] \times C$.

PROOF. For $c \in C$, $e \in E_c$, and $h > 0$ sufficiently small, let us consider the finite difference

$$\text{(B.6)} \qquad z_h(t, (c, e)) = \frac{\mathbf{Y}(t, c + he) - \mathbf{Y}(t, c)}{h}.$$

From (B.4) we obtain the equation

$$\text{(B.7)} \quad \frac{d}{dt} z_h(t, (c, e)) = \frac{A(t, \mathbf{Y}(t, c + he)) - A(t, \mathbf{Y}(t, c))}{h}, \quad z_h(0, (c, e)) = e.$$

In view of the well-posedness of (B.4), also (B.7) is well-posed. We highlight now some regularity properties of the functions $t \mapsto z_h(t, (c, e))$, which naturally come from the assumptions on A.

By definition of $z_h(t, (c, e))$ and (B.3) we observe that

$$\text{(B.8)} \qquad \sup_{t \in [0, T]} \|z_h(t, (c, e))\|_E \le \|e\|_E e^{LT},$$

where L is the constant of Theorem B.1(i). Hence $t \mapsto z_h(t, (c, e))$ are uniformly bounded in $[0, T]$ with respect to h. To compute their limit as $h \downarrow 0$, notice that an application of Theorem B.1 (or more classical results for linear ODEs) yields for all $f \in E_C$ the existence and the uniqueness of a C^1 map $t \mapsto z(t) \in \bar{E}_C$ solving the linear differential equation

$$\text{(B.9)} \qquad z'(t) = L_A(t, \mathbf{Y}(t, c))z(t), \quad z(0) = f,$$

for a continuous selection $L_A(t, \mathbf{Y}(t, c)) \in D_{E_C} A(t, \cdot)(\mathbf{Y}(t, c))$, whose existence is granted by (B.5). For later use we denote the solution $z = z_f$ to emphasize its initial datum.

From (B.6) and (B.7), with the notation $z^h(t) := z_h(t, (c, e))$, we have the identity

$$\frac{d}{dt}(z^h(t) - z_e(t))$$

$$= h^{-1}\Big[A(t, \mathbf{Y}(t, c) + hz^h(t)) - A(t, \mathbf{Y}(t, c))$$

$$\qquad - hL_A(t, \mathbf{Y}(t, c))z^h(t)\Big] + L_A(t, \mathbf{Y}(t, c))[z^h(t) - z_e(t)].$$

We recall now that for any $g \in C^1([0, T]; E)$ the map $t \mapsto \|g(t)\|_E$ is Lipschitz and satisfies (as a simple consequence of the fundamental theorem of calculus and (A.8))

$$\left| \frac{d}{dt} \|g(t)\|_E \right| \le \|g'(t)\|_E \quad \text{for almost every } t \in [0, T].$$

Then, we obtain the estimates

$$
\frac{\mathrm{d}}{\mathrm{d}t}\|z^h(t) - z_e(t)\|_E
$$

$$
\leq \left\| h^{-1}\big[A(t, \mathbf{Y}(t,c) + hz^h(t)) - A(t, \mathbf{Y}(t,c)) - hL_A(t, \mathbf{Y}(t,c))z^h(t)\big] \right.
$$

$$
\left. + L_A(t, \mathbf{Y}(t,c))[z^h(t) - z_e(t)] \right\|_E
$$

$$
\leq \left\| h^{-1}[A(t, \mathbf{Y}(t,c) + hz^h(t)) - A(t, \mathbf{Y}(t,c)) - hL_A(t, \mathbf{Y}(t,c))z^h(t)] \right\|_E
$$

$$
+ L\|z^h(t) - z_e(t)\|_E,
$$

since $\|L_A(t, \mathbf{Y}(t,c))\|_{\mathcal{L}(E_C, \bar{E}_C)} \leq L$, thanks to (B.1). By Gronwall's inequality and using that $z^h(0) = z_e(0) = e$, we obtain

$$
\|z^h(t) - z_e(t)\|_E
$$

$$
\leq \int_0^t \left\| h^{-1}[A(s, \mathbf{Y}(s,c) + hz^h(s)) - A(t, \mathbf{Y}(s,c)) \right.
$$

$$
\left. - hL_A(t, \mathbf{Y}(t,c))z^h(s)] \right\|_E e^{L(t-s)}\mathrm{d}s.
$$

By the pointwise limit

$$
\left\| h^{-1}[A(s, \mathbf{Y}(s,c) + hz^h(s)) - A(t, \mathbf{Y}(s,c)) - hL_A(t, \mathbf{Y}(t,c))z^h(s)] \right\|_E \to 0
$$
$$
\text{for } h \to 0,
$$

and by the dominated convergence ensured by assumption (B.1) and the uniform boundedness in (B.8), we conclude that

$$
\lim_{h \to 0} \sup_{t \in [0,T]} \|z^h(t) - z_e(t)\|_E = 0.
$$

Moreover, with a similar argument as above we have

$$
\frac{\mathrm{d}}{\mathrm{d}t}\|z_e(t)\|_E \leq L\|z_e(t)\|_E,
$$

and again by Gronwall's inequality and $z_e(0) = e$

$$
(\text{B.10}) \qquad\qquad \|z_e(t)\|_E \leq \|e\|_E e^{LT}.
$$

We shall denote now by $L(t,c)$ the linear operators induced by (B.9) by setting $L(t,c)f = z_f(t)$ for any $f \in E_C$ (not to be confused with $L_A(t,c) \in D_{E_C} A(t, \cdot)(c)$). Notice that in view of (B.10) the operators $L(t,c) \in \mathcal{L}(E_C, \bar{E}_C)$ are in fact uniformly bounded. With all of this we proved that $\mathbf{Y}(t, \cdot)$ is G-differentiable (uniformly in $t \in [0, T]$) and that

$$
\mathrm{d}\mathbf{Y}(t,c,e) = L(t,c)e \quad \text{for every } e \in E_c, \ t \in [0, T],
$$

In order to improve from G-differentiability to F-differentiability, we apply Lemma A.4 if we are able to show that the operator $(t,c) \mapsto L(t,c)$ is continuous from $[0, T] \times C$ to $\mathcal{L}(E_C, \bar{E}_C)$. Since uniform continuity with respect to t is obvious,

the continuity with respect to c can be again shown by considering, for c, $c' \in C$, the operator

$$Z(t, c, c') := L(t, c) - L(t, c').$$

Indeed, for c, c' fixed, we have in the space $\mathcal{L}(E_C, \bar{E}_C)$ the linear differential equation

$$\frac{\mathrm{d}}{\mathrm{d}t} Z(t, c, c') = L_A(t, \mathbf{Y}_t(c')) \circ Z(t, c, c')$$
$$+ [L_A(t, \mathbf{Y}_t(c')) - L_A(t, \mathbf{Y}_t(c))] \circ L(t, c),$$

with the initial condition $Z(0, c, c') = 0$; here we denoted $\mathbf{Y}_t(c) = \mathbf{Y}(t, c)$. Under the assumption (B.5) and by a similar argument based on the Gronwall inequality as above, since, for c fixed, $L_A(t, \mathbf{Y}_t(c')) \to L_A(t, \mathbf{Y}_t(c))$ as $\|c - c'\|_E \to 0$, we can conclude that $Z(t, c, c') \to 0$ as $\|c - c'\|_E \to 0$ uniformly in $[0, T]$, and therefore the continuity of $c \mapsto L(t, c)$. It follows that $\mathbf{Y}_t$ is of class C^1, uniformly with respect to $t \in [0, T]$, with a F-differential at c given by $L(t, c)$. $\qquad\square$

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