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## Measurement of the Low-Energy Antideuteron Inelastic Cross Section

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In this Letter, we report the first measurement of the inelastic cross section for antideuteron-nucleus interactions at low particle momenta, covering a range of  $0.3 \leq p < 4$  GeV/ $c$ . The measurement is carried out using  $p$ -Pb collisions at a center-of-mass energy per nucleon–nucleon pair of  $\sqrt{s_{NN}} = 5.02$  TeV, recorded with the ALICE detector at the CERN LHC and utilizing the detector material as an absorber for antideuterons and antiprotons. The extracted raw primary antiparticle-to-particle ratios are compared to the results from detailed ALICE simulations based on the GEANT4 toolkit for the propagation of (anti)particles through the detector material. The analysis of the raw primary (anti)proton spectra serves as a benchmark for this study, since their hadronic interaction cross sections are well constrained experimentally. The first measurement of the inelastic cross section for antideuteron-nucleus interactions averaged over the ALICE detector material with atomic mass numbers  $\langle A \rangle = 17.4$  and 31.8 is obtained. The measured inelastic cross section points to a possible excess with respect to the Glauber model parametrization used in GEANT4 in the lowest momentum interval of  $0.3 \leq p < 0.47$  GeV/ $c$  up to a factor 2.1. This result is relevant for the understanding of antimatter propagation and the contributions to antinuclei production from cosmic ray interactions within the interstellar medium. In addition, the momentum range covered by this measurement is of particular importance to evaluate signal predictions for indirect dark-matter searches.

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The possible presence of antinuclei in the Milky Way could be explained either by reactions of high-energy cosmic rays with the interstellar medium or by more exotic sources, such as dark-matter annihilation [1]. Some dark-matter models [2–6] predict that low-energy antideuterons are a promising probe for indirect dark-matter searches since the contributions from cosmic-ray interactions in the energy range below 1–2 GeV per nucleon [7–9] are expected to be rather small. For this reason, the search for antinuclei has been intensified in recent years with new satellite and balloon-borne experiments such as AMS-02 [10] and GAPS [11]. So far, only antiprotons have been detected in space [12], and no clear evidence of heavier antinuclei production has been found yet [13,14], but dedicated analyses searching for antideuteron and antihelium are currently ongoing [3,15].

In order to get a reliable baseline for antideuteron production at low energies, realistic models of cosmic-ray transport are necessary. In addition, also the predicted flux of antinuclei from dark-matter annihilation depends on the production mechanism and antinuclei transport properties

within the interstellar medium. There are three main relevant mechanisms that determine the signal and background rates: (i) the antideuteron production, either in  $p$ - $A$  and  $A$ - $A$  reactions between cosmic rays and the interstellar medium, depending on the element abundance or in dark-matter annihilation processes, (ii) the antideuteron propagation in the galaxy, the heliosphere and the Earth’s atmosphere, and (iii) inelastic processes such as nuclear breakup, charge exchange or annihilation that occur during propagation and in experiments inside the detectors. These three mechanisms must be measured as precisely as possible to interpret correctly any future measurement in satellite and balloon-borne experiments. While the propagation has been constrained by measuring different nuclei from primary and secondary cosmic rays [16–19], accelerator experiments can be used to study the production and the inelastic scattering cross sections.

Antimatter is copiously produced in high-energy collisions of protons and heavy ions [20,21]. This environment is hence well suited to study antinuclei properties. At RHIC, the STAR and PHENIX Collaborations have measured  $\bar{p}$ ,  $\bar{d}$ ,  ${}^3\bar{\text{He}}$  and  ${}^4\bar{\text{He}}$  [22–25] yields employing Au-Au collisions at center-of-mass energies per nucleon-nucleon pair of  $\sqrt{s_{NN}} = 130$  GeV and  $\sqrt{s_{NN}} = 200$  GeV. At the LHC, the ALICE Collaboration has studied  $\bar{p}$ ,  $\bar{d}$ ,  ${}^3\bar{\text{He}}$ , and  ${}^4\bar{\text{He}}$  production in  $pp$ ,  $p$ -Pb, and Pb-Pb collisions at center-of-mass energies per nucleon pair from 0.9 to 13 TeV [26–32], and the yields obtained for  $A \geq 2$  have

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been interpreted by means of coalescence or statistical hadronization models [33–36]. The LHC measurements combined with different coalescence models have been employed to estimate the antideuteron and antihelium flux from cosmic-ray interactions measurable by the AMS-02 and GAPS experiments [15,37–39]. Since the inelastic cross sections for antinuclei-nuclei interactions are measured precisely only for  $\bar{p}$  but barely known for heavier antinuclei, all the available calculations rely on poorly constrained parametrizations. For antideuterons, the inelastic cross sections have been measured for several materials only for two momentum values,  $p = 13.3$  GeV/ $c$  [40] and  $p = 25$  GeV/ $c$  [41]. However, the low-momentum range accessible by ALICE ( $p \leq 5$  GeV/ $c$ ) remains unexplored. For antihelium, no measurement of inelastic cross sections is available.

In this Letter, we present a method to evaluate the inelastic cross section of antinuclei based on the measurement of raw reconstructed antiparticle-to-particle ratios. Using ratios instead of individual particle yields allows us to extract the antideuteron and antiproton cross sections independently from their production cross sections and for a broad momentum range. We report the first measurement of the inelastic cross section for antideuteron-nucleus interactions in the momentum range of  $0.3 \leq p < 4$  GeV/ $c$ . The results presented are based on data collected during the 2016  $p$ -Pb LHC run at  $\sqrt{s_{NN}} = 5.02$  TeV. The performance of the ALICE detector and the description of its subsystems can be found in [42,43]. Collision events are selected by using the information from the V0 detector, which consists of two plastic scintillator arrays located on both sides of the interaction point at forward and backward pseudorapidities. A simultaneous signal in both arrays was used as a minimum-bias (MB) trigger. In total, about  $600 \times 10^6$  MB events are selected for further analysis, which correspond to an integrated luminosity of  $\mathcal{L}_{\text{int}}^{\text{MB}} = 287 \mu\text{b}^{-1}$ , with a relative uncertainty of 3.7% [44].

The charged-particle tracks are reconstructed in the ALICE central barrel with the inner tracking system (ITS) and the time projection chamber (TPC), which are located within a solenoid that provides a homogeneous magnetic field of 0.5 T in the direction of the beam axis. The ITS consists of six cylindrical layers of silicon detectors located at radial distances from the beam axis between 3.9 cm and 43 cm. The TPC extends radially from  $r = 85$  cm to  $r = 247$  cm, is 5 m long, and was filled with an Ar-CO<sub>2</sub> gas mixture during the 2016 data taking period. These two subsystems provide full azimuthal coverage for charged-particle trajectories in the pseudorapidity range  $|\eta_{\text{lab}}| < 0.8$ . The selected tracks must fulfill basic quality criteria established in antinuclei analyses in  $p$ -Pb collisions [31]. These criteria guarantee a resolution of about 2% on the momentum reconstructed at the primary vertex ( $p_{\text{primary}}$ ) in this analysis.

The TPC is also used for the particle identification (PID) of (anti)protons and (anti)deuterons via their specific energy loss  $dE/dx$  in the gas volume, with a resolution of about 5% [45]. The  $n(\sigma_i^{\text{TPC}})$  variable represents the PID response in the TPC expressed in terms of the deviation between the measured and expected  $dE/dx$  for a particle species  $i$ , normalized by the detector resolution  $\sigma$ . The expected  $dE/dx$  is computed with a parametrized Bethe-Bloch curve [43]. (Anti)protons and (anti)deuterons are selected by applying the selection criterion  $|n(\sigma_i^{\text{TPC}})| < 3$ . This selection is sufficient to obtain a purity close to 100% for (anti)protons and (anti)deuterons in the momentum range below 0.7 GeV/ $c$  and 1.4 GeV/ $c$ , respectively. For the momentum range above 0.7 GeV/ $c$  for (anti)protons and 0.9 GeV/ $c$  for (anti)deuterons, the PID is complemented by the time-of-flight (TOF) system, consisting of multigap resistive plate chambers. (Anti)proton and (anti)deuteron candidates selected in the TPC are matched to TOF hits, and fits to the squared-mass distributions are performed for different momentum intervals [31]. The PID purity in all momentum intervals is found to be higher than 88% and 47% for the (anti)proton and (anti)deuteron samples, respectively. The background is subtracted from the squared-mass spectra with a two-component fit [31].

The determination of the inelastic cross section requires precise knowledge of the ALICE detector material. The MC parametrization of the ALICE material budget up to the outer TPC vessel was validated with photon conversion analyses within a precision of  $\sim 4.5\%$  [43], and it is shown in the Supplemental Material [46]. The ALICE detector material from the primary interaction point up to the TOF has an average atomic number of  $\langle Z \rangle = 14.8$  and a mass number of  $\langle A \rangle = 31.8$ . For the detector material up to the middle of the TPC, these values amount to  $\langle Z \rangle = 8.5$  and  $\langle A \rangle = 17.4$ . These values have been obtained by weighing the contribution from different materials with their density times the length crossed by particles.

The selected (anti)proton and deuteron candidates include a substantial amount of background from secondary (anti)particles that originate from weak decays of hyperons or from spallation reactions in the detector material. Following the procedure described in [26,47,48], the contribution from secondary (anti)particles is subtracted by performing a fit to the distribution of the measured distance of closest approach (DCA) of track candidates to the primary vertex with templates from Monte Carlo (MC) simulations. In contrast to secondary particles, primary particles point back to the primary vertex; hence, a distinct structure peaked at zero in the DCA distribution characterizes the primary particles. Secondary particles correspond to a flat DCA distribution, and their contribution can therefore be separated [26,28]. The fraction of secondary (anti)protons is found to be around 20% in the lowest momentum interval analyzed ( $0.3 \leq p_{\text{primary}} < 0.4$  GeV/ $c$ ) and decreases monotonically

down to  $\sim 1.5\%$  at high momenta. The main contribution of secondary (anti)protons stems from weak decays. For deuterons, the dominant contribution of secondary particles comes from spallation processes in the detector material that lead to the ejection of fragments such as protons, neutrons, or deuterons. The fraction of secondary deuterons is found to be 23.5% in the lowest momentum interval ( $0.5 \leq p_{\text{primary}} < 0.6$  GeV/ $c$ ) and to decrease exponentially to negligible values at  $p_{\text{primary}} \sim 1.4$  GeV/ $c$ . For antiprotons and antideuterons, the contribution from spallation processes is absent. The feed down from weak decays of hyperons and hypernuclei has a negligible impact on the measured ratios [31,47,49]. Hence, the antideuteron sample is composed entirely from primaries. The total number of selected candidates amounts to  $7.57 \times 10^7$  protons,  $6.52 \times 10^7$  antiprotons,  $2.52 \times 10^5$  deuterons, and  $1.98 \times 10^5$  antideuterons. The momentum spectra are corrected for the background from secondary particles but not for the detector efficiency or losses of (anti)particles in the detector material, so they are referred to as raw primary spectra.

Figure 1 shows the  $\bar{p}/p$  and  $\bar{d}/d$  ratios as a function of  $p_{\text{primary}}$ . The systematic uncertainties due to tracking, particle identification, and contribution from secondaries are considered, and the total uncertainty is obtained as the quadratic sum of the individual contributions. It increases from 1% (2%) at low momentum up to 2% (6%) in the high-momentum region for  $\bar{p}/p$  ( $\bar{d}/d$ ). The uncertainty on the primordial antimatter-to-matter ratio produced in collisions is considered as a global uncertainty. The primordial  $\bar{p}/p$  ratio  $0.984 \pm 0.015$  is extrapolated from available measurements [47,48], and, under the assumption that the

(anti)deuteron yield is proportional to the squared yield of (anti)protons [50,51], the primary  $\bar{d}/d$  ratio amounts to  $0.968 \pm 0.030$ . These values are used as an input for detailed MC simulations based on the GEANT4 toolkit for the propagation of (anti)particles through the detector material [52]. For the description of antinucleus–nucleus inelastic cross sections, GEANT4 relies on a Glauber calculation convoluted with a MC averaging method [53]. Figure 1 shows that the GEANT4 -based simulations are able to describe the  $\bar{p}/p$  ratio and are in qualitative agreement with the data for the  $\bar{d}/d$  ratio.

The sensitivity of the antiparticle-to-particle ratios to the modifications of elastic and inelastic cross sections was benchmarked with the  $\bar{p}/p$  measurement. The (anti)proton cross sections have been measured by various experiments [54–60], and the results are described well by the GEANT4 parametrization. The blue boxes in Fig. 2 indicate the  $\pm 1\sigma$  limits for the measured  $\bar{p}/p$  ratio, where  $1\sigma$  corresponds to the quadratic sum of statistical, systematic, and global uncertainties. The green and magenta bands show the simulated ratios with a variation of  $\pm 25\%$  of the inelastic antiproton cross section along with the simulations using a default cross section (gray band). Only a variation of the total inelastic cross section has been carried out. The widths of the bands correspond to a quadratic sum of the contributions from two additional variations: (i) the elastic cross sections of protons and antiprotons are changed independently by  $\pm 20\%$ , which leads to  $\lesssim 1.5\%$  modification of the ratio, and (ii) the inelastic proton-nucleus cross section is varied by 3.5%, which is the uncertainty of the GEANT4 parametrizations obtained from fits of the experimental data for this cross section. This variation

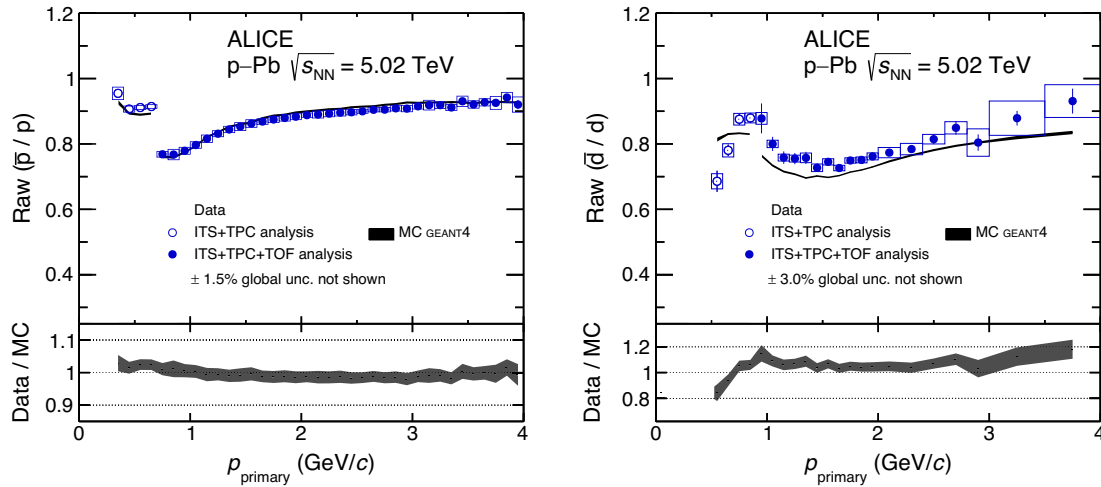


FIG. 1. Raw primary  $\bar{p}/p$  (left) and  $\bar{d}/d$  (right) ratios as a function of the momentum  $p_{\text{primary}}$ . Experimental data are shown in blue, and the statistical and systematic uncertainties are shown as vertical bars and boxes. The results from ALICE MC simulations based on GEANT4 using the FTFP\_INCLXX\_EMV physics list are shown in black. The width of the MC band represents the statistical uncertainty of the simulation. The global uncertainty due to the primordial ratio (1.5% for  $\bar{p}/p$  and 3% for  $\bar{d}/d$ ) is not shown in the top panels. The bottom panels display the ratios of experimental data to MC simulations with statistical, systematic and global uncertainties added in quadrature.

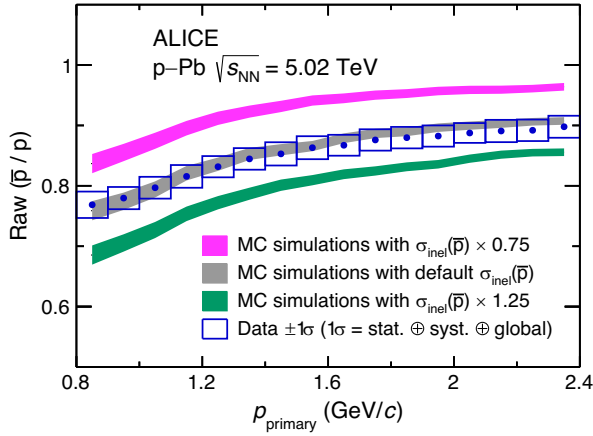


FIG. 2. Raw primary  $\bar{p}/p$  ratio as a function of momentum. Blue boxes indicate  $\pm 1\sigma$  experimental limits. The results from MC simulations with varied  $\sigma_{\text{inel}}(\bar{p})$  are shown as green and magenta bands, and the gray band corresponds to the results with default  $\sigma_{\text{inel}}(\bar{p})$ . The uncertainties on MC results include the variations of elastic cross sections and the variation of  $\sigma_{\text{inel}}(p)$ .

yields a modification of about 0.5% in the ratio. These systematic checks demonstrate that the antiparticle-to-particle ratio is mainly sensitive to the variation of the

inelastic cross sections and can therefore be used to measure the antideuteron inelastic cross section.

Extending this recipe, an iterative and momentum-dependent variation of  $\sigma_{\text{inel}}(\bar{p})$  within the GEANT4 simulations was carried out to obtain  $\bar{p}/p$  ratios that correspond to the  $\pm 1\sigma$  and  $\pm 2\sigma$  experimental limits. The resulting  $\pm 1\sigma$  and  $\pm 2\sigma$  limits for  $\sigma_{\text{inel}}(\bar{p})$  are presented in panels (a) and (b) of Fig. 3 together with standard GEANT4 parametrizations. Panel (a) refers to the ITS + TPC analysis and hence, corresponds to the inelastic interaction with nuclei that have average charge and mass number  $\langle Z \rangle = 8.5$  and  $\langle A \rangle = 17.4$ ; panel (b) refers to the analysis additionally employing the TOF and corresponds to  $\langle Z \rangle = 14.8$  and  $\langle A \rangle = 31.8$ . The inelastic cross sections shown in Fig. 3 are estimated as a function of the momentum  $p$  at which the inelastic interaction occurs. Because of the continuous energy loss of the particle inside the detector material, this momentum is lower than  $p_{\text{primary}}$  reconstructed at the primary vertex. The corresponding correction is estimated using MC simulations by looking at the average values of the annihilation momentum distribution in each  $p_{\text{primary}}$  interval. The RMS of the distributions is then propagated to the uncertainty of the cross section measurement. The minimum momentum reconstructed at the primary vertex

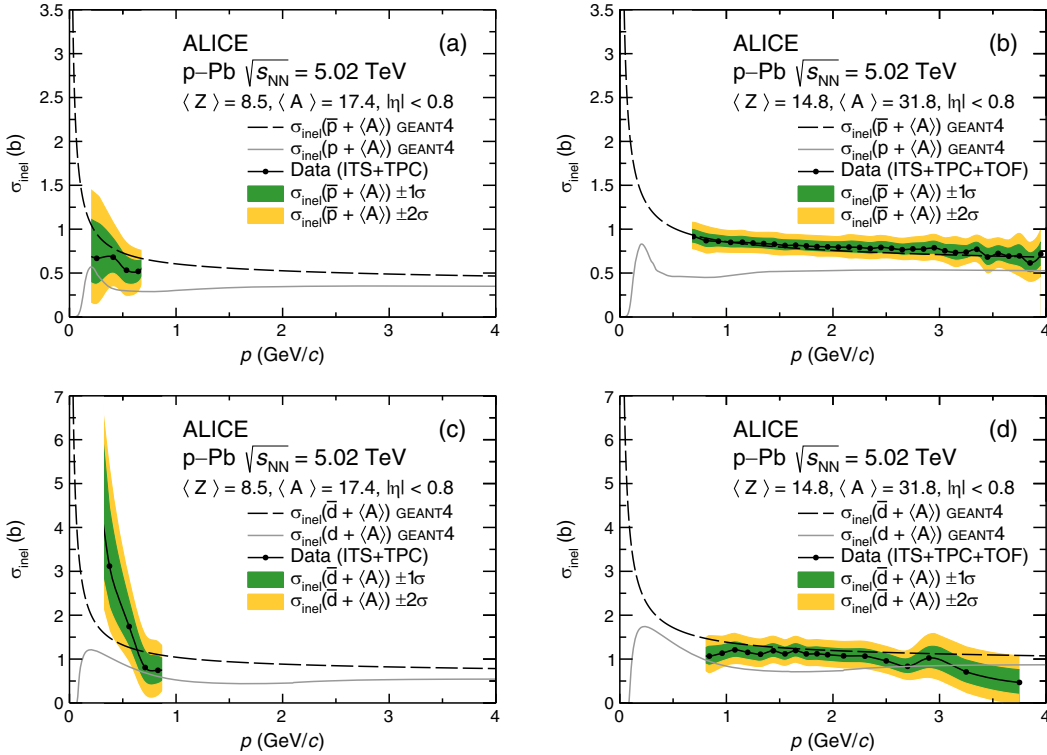


FIG. 3. Inelastic interaction cross section for antiprotons and antideuterons on an average material element of the ALICE detector as a function of the momentum  $p$ , at which the interaction occurs. The top row shows the results for antiprotons, the bottom row for antideuterons, and the results from the ITS + TPC (ITS + TPC + TOF) analysis are shown on the left (right). Dashed black lines represent the GEANT4 parametrizations for antinuclei, and full gray lines show the parametrizations for protons and deuterons. The experimental data points are shown connected by solid black lines, with green and orange bands corresponding to  $\pm 1$  and  $\pm 2\sigma$  constraints from the raw primary ratios.

amounts to  $p_{\text{primary}} = 0.3 \text{ GeV}/c$  for antiprotons and to  $p_{\text{primary}} = 0.5 \text{ GeV}/c$  for antideuterons, and the energy-loss correction transforms these values to  $p = 0.18 \text{ GeV}/c$  and  $p = 0.3 \text{ GeV}/c$ , correspondingly. For momenta  $p > 0.7 \text{ GeV}/c$ , the antiproton inelastic cross section is found to be in good agreement with the GEANT4 parametrizations, which, in turn, describe well the existing experimental data [53]. Thus, these results validate the analysis procedure, which then can be applied to (anti) deuterons.

In contrast to antideuterons, the deuteron inelastic cross section was measured on several materials at various momenta [61,62], and the data are well described by GEANT4 parametrizations. The antideuteron inelastic cross section can therefore be constrained via the comparison of the experimental  $\bar{d}/d$  ratio and the GEANT4 -based MC simulations with  $\sigma_{\text{inel}}(\bar{d})$  varied in a similar way as for antiprotons. For this purpose, the same uncertainties are considered: (i) the variation of elastic cross sections of (anti)deuterons by  $\pm 20\%$  that results in  $\lesssim 2\%$  deviation for the ratio, (ii) the variation of the inelastic deuteron cross section by  $7\%$  that corresponds to the precision of GEANT4 parametrizations ( $\lesssim 1\%$  uncertainty), and (iii) the uncertainty from the primordial  $\bar{d}/d$  ratio (3.0%).

The resulting upper and lower limits on  $\sigma_{\text{inel}}(\bar{d})$  for targets with  $\langle Z \rangle = 8.5$ ,  $\langle A \rangle = 17.4$ , and  $\langle Z \rangle = 14.8$ ,  $\langle A \rangle = 31.8$  are shown in panels (c) and (d) of Fig. 3, respectively. The extracted inelastic cross sections presented here include all inelastic antideuteron processes where the antideuteron is destroyed and represent the first measurement in this low-momentum range.

While the measured  $\sigma_{\text{inel}}(\bar{d})$  is found to be in agreement with the GEANT4 implementation within the  $0.9 \leq p < 4.0 \text{ GeV}/c$  momentum range, it rises faster than the simulated parametrization in the momentum range  $0.3 \leq p < 0.9 \text{ GeV}/c$ , reaching a maximal discrepancy of a factor 2.1 in the interval of  $0.3 \leq p < 0.47 \text{ GeV}/c$ .

These measurements can now help to better understand the antideuteron inelastic processes at low momenta and to improve the parametrization of the inelastic cross section used in GEANT4. Additionally, these results are now available for models of the propagation of antideuterons within the interstellar medium [3,7,38] and will impact the flux expectations at low momentum near Earth.

In summary, we have shown how the ALICE detector can be used as an absorber to study the antinuclei inelastic scattering cross section on detector material. The antiparticle-to-particle ratios method was validated using (anti)protons, and the sensitivity of the ratio to the variation of the inelastic cross section was demonstrated. In this way, the first measurement of the inelastic scattering cross section of antideuterons was performed on an effective target with mean charge number  $\langle Z \rangle = 8.5$  and mass number  $\langle A \rangle = 17.4$  in the momentum range  $0.3 \leq p < 0.9 \text{ GeV}/c$ , and with  $\langle Z \rangle = 14.8$  and

$\langle A \rangle = 31.8$  in  $0.9 \leq p < 4.0 \text{ GeV}/c$ . These cross sections can now be used in propagation models of antideuterons within the interstellar medium for dark-matter searches. Future studies of high-statistics  $pp$ ,  $p$ -Pb, and Pb-Pb data collected during the second (2015–2018) and third (scheduled to start in 2021) LHC run campaigns should allow the measurement of inelastic cross sections of heavier antinuclei such as  ${}^3\bar{\text{He}}$ , and  ${}^4\bar{\text{He}}$  in a similar way and the improvement of the current antideuteron results.

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