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Research Article

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Crystallization behavior of SiO₂-P₂O₅-CaO-MgO-Na₂O-K₂O bioactive glass powder

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Abstract: The crystallization process of a bioactive silicate glass with 47.5SiO₂-10Na₂O-10K₂O-10MgO-20CaO-2.5P₂O₅ molar composition was investigated by using nonisothermal differential thermal a nalvsis (DTA). The DTA plots recorded at different heating rates exhibited a single crystallization peak. The activation energy for crystallization was estimated by applying the equations proposed by Kissinger and Matusita-Sakka. The Johnson-Mehl-Avrami exponent (n) was assessed by using the Ozawa and Augis-Bennett methods. The analyses suggest that a surface crystallization mechanism with one-dimensional crystal growth is predominant. The activation energy for viscous flow was also assessed (176 kJ/mol) and was found lower than the activation energy for crystallization (271 kJ/mol). This confirms the stability of 47.5B against crystallization and its good sinterability, which is a highly attractive feature for producing glass products of biomedical interest, such as bioactive porous scaffolds for bone repair.

Keywords: Bioactive glass; Glass-ceramic; Thermal analysis; Crystallization

1 Introduction

Bioactive glasses were invented fifty years ago at the University of Florida by Prof. Larry Hench and his team, who developed the famous 45S5 composition ($45SiO_2-24.5CaO-24.5Na_2O-6P_2O_5$ wt.%) [1]. This glass, trade-named as

∂ Open Access. © 2019 E. Fiume *et al.*, published by De Gruyter. **(CC)** BY License Bioglass[®], was initially found able to bond both to bone and to soft collagenous tissues of the body [2]. Hence, it was soon considered an excellent biomaterial to repair bone tissue undergoing injury or dysfunctionality and, to date, it has been implanted in more than 1.5 million patients suffering from osseous defects in orthopedic and dental applications [3]. Recent studies carried out over the past two decades have also revealed additional highlyappealing properties of 45S5 Bioglass[®], including an inherent antibacterial effect due to the release of alkaline ions (Ca²⁺ and Na⁺) [4, 5] and the capability to stimulate angiogenesis, which is the key to accelerate wound healing process [6, 7].

The main drawback of 45S5 Bioglass® concerns its poor workability: in fact, the glass transition temperature (T_g) and the onset of crystallization (T_x) are very close, thus limiting the possibility of sintering by viscous flow [8]. As a result, 45S5 glass cannot be sintered without undergoing devitrification, which has an impact on the kinetics of apatite formation (bioactivity) in vitro and in vivo and, hence, on the bone-bonding capability and osteointegration rate. In this regard, Peitl et al. [9] showed that the onset time of hydroxyapatite formation on the surface of 45S5 Bioglass[®] increases from 8 to 25 h when the crystalline fraction increases from 0 to 60%; however, the bioactivity is not totally suppressed and the partially-crystallized material is still potentially suitable for biomedical applications [10-12]. A similar behavior was also found to occur in other silicate (e.g. S53P4, 53SiO₂-20CaO-23Na₂O-4P₂O₅ wt.% [13]) and phosphate glasses $(50P_2O_5 - (40 - x)CaO_5 - (40 - x)CaO_5$ $xSrO-10Na_2O$ with x = 0, 20, and 40 mol.%) [14]. Over the years, other glasses have been developed with a larger sintering window (e.g. 13-93, 53SiO₂-6Na₂O-12K₂O-5MgO-20CaO-4P₂O₅ wt.%) [15] and/or with the capability of retaining an excellent and fast apatite-forming ability also in partially-crystallized forms (e.g. CEL2, 43.8SiO₂-15.0Na₂O-6.1K₂O-4.6MgO-23.6CaO-6.9P₂O₅ wt.%) [16].

The fabrication of sintered bioactive glass products, such as three-dimensional macroporous scaffolds, is often a challenge because, on one hand, extensive densifica-

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tion is required to achieve enough strength for the glassderived product to be safely handled and, on the other hand, crystallization of the glass, which could negatively affect the bioactivity, may occur prior to significant densification.

Amorphous scaffolds made of sintered 13-93 glass were reported to have suitable mechanical properties for bone repair applications, but their apatite-forming ability was lower compared to 45S5 composition [17]. A comparative study between CEL2- and 45S5-derived glass-ceramic foams suggested that the apatite-forming kinetics of the former are faster, leading to a thicker surface layer of hydroxyapatite [18].

Verné et al. developed a silicate glass (47.5B) with an exceptionally large sinterability range $(T_x - T_g = 260^{\circ}C)$ and retaining in vitro bioactivity after partial devitrification at 800°C for 3 h [19]. These properties suggest the potential suitability of such a material for making bioactive porous scaffolds; however, the crystallization behavior should be determined first. In this regard, the present work aims at investigating the crystallization mechanism of 47.5B glass powder under non-isothermal conditions in order to assess the activation energies for crystallization and viscous flow. These parameters have not been determined for such glass composition so far. Over the past half century, a number of approaches have been developed to estimate the activation energy for crystallization [20], but there is a paucity of relevant studies about biomedical glasses. Of those, the vast majority focused on 45S5 Bioglass[®] [21, 22], while almost no information have been provided on other bioactive glass systems.

2 Materials and methods

2.1 Glass preparation

The analyzed material was a glass (47.5B) originally developed by Verné *et al.* [19] at Politecnico di Torino in the composition 47.5SiO₂-10Na₂O-10K₂O-10MgO-20CaO-2.5P₂O₅ (mol.%) and was produced using a standard melting method in a platinum crucible. The raw precursors (SiO₂, Na₂CO₃, K₂CO₃, (MgCO₃)₄·Mg(OH)₂·5H₂O, CaCO₃ and Ca₃(PO₄)₂ high-purity powders purchased from Sigma-Aldrich) were homogeneously mixed in the crucible and melted in air at 1500°C for 0.5 h; the melt was then quenched in deionized water to produce a frit that was ball milled (Pulverisette 0, Frtisch, Germany) and sieved below 32 µm by a stainless steel sieve (Giuliani Technologies Srl, Torino, Italy). According to the authors' previ-

ous experience on other glass compositions, this particle size range is very suitable to produce porous glass-derived scaffolds by various technologies, including foam replication and robocasting; this is the reason why the glass crystallization behavior was studied using 47.5B powder with such size.

2.2 Thermal properties

The glass transition temperature (T_g) and the peak crystallization temperature (T_p) were determined by differential thermal analysis (DTA; DTA404PC, Netzsch, Germany) at various heating rates β (10, 20, 30 and 40°C/min). The measurements were performed on 50 mg of glass that was heated to 1200°C in a platinum crucible under nitrogen flow (inert atmosphere) using Al₂O₃ powder as a reference material. The glass transition temperature was determined at the inflection point of the DTA curve, as obtained from the first derivative of the thermal plot. The peak crystallization temperature was found at the maximum of the exothermic peak. The accuracy of the measurements was $\pm 3^{\circ}$ C.

2.3 Analysis of the crystallization process

The crystallization process in glasses is usually described by two parameters, *i.e.* the Johnson-Mehl-Avrami (JMA) coefficient (n) – which depends on the mechanism of crystal nucleation and growth [23, 24] – and the activation energy for crystallization.

The JMA coefficient was quantified by suing two independent techniques. First, it was determined by the Ozawa equation [25]:

$$-n = \left. \frac{d \left(\ln \left(-\ln \left(1 - \chi \right) \right) \right)}{d (\ln \beta)} \right|_{T}$$
(1)

where χ is the volume fraction of crystallized phase at a fixed temperature T. The values of χ were estimated from the DTA plots as the ratio between the partial area under the exothermic peak (at a chosen temperature) and the total area under the peak. A plot of $\ln(-\ln(1-\chi))$ vs. $\ln\beta$, followed by linear interpolation, gives a straight line of slope n. The value of the JMA parameter n was calculated as mean ± SD of the slopes of the plots.

The coefficient n was also estimated by using the Augis-Bennet method [26]:

$$n = \frac{2.5RT_p^2}{\Delta T_{FWHM}E_c} \tag{2}$$

where ΔT_{FWHM} is the full width of the DTA exothermic peak at the half maximum. In this case, the JMA parameter

was calculated as mean \pm SD of the n-values determined at each heating rate (β).

The activation energy for crystallization can be assessed from non-isothermal DTA data by applying models that relate the variation of T_p to the heating rate β . Kissinger proposed the following equation to estimate the crystallization activation energy ($E_{c,K}$) [27]:

$$\ln \frac{\beta}{T_p^2} = -\frac{E_{c,K}}{RT_p} + constant \tag{3}$$

A plot of $\ln \frac{\beta}{T_p^2}$ vs. $\frac{1}{T_p}$, followed by linear interpolation, gives a straight line of slope $\frac{E_{c,K}}{R}$ (R = 8.314 J/mol K is the ideal gas constant), from which $E_{c,K}$ can be easily calculated.

However, Matusita *et al.* [28] demonstrated that Equation (3) is valid only when crystal growth occurs on a fixed number of nuclei. Hence, the Kissinger model yields an underestimated value of the activation energy for crystallization if most nuclei are formed during the DTA measurement, as the number of nuclei may vary with β . A more general equation deriving from the Kissinger model was then proposed by Matusita and Sakka [29]:

$$\ln \frac{\beta^n}{T_p^2} = -\frac{mE_c}{RT_p} + constant \tag{4}$$

where E_c is the correct activation energy for crystallization, n is the JMA exponent and m is a numerical factor which depends on the dimensionality of crystal growth, *i.e.* threedimensional sphere-like (m = 3), two-dimensional platelike (m = 2) or one-dimensional rod-like (m = 1) crystallization [30]. A plot of $\ln \frac{\beta^n}{T_p^2}$ vs. $\frac{1}{T_p}$, followed by linear interpolation, gives a straight line of slope $\frac{mE_c}{R}$, from which E_c can be calculated.

As a general rule, we have $m \approx 1$ when surface crystallization predominates, whereas $m \approx 3$ when bulk crystallization is predominant. The value of m is related to n as m = n - 1 when crystal nucleation occurs during DTA. When a surface crystallization mechanism is dominant, we have m = n = 1 regardless of whether nuclei are formed prior to or during thermal analysis: in this case, Matusita-Sakka equation (Equation (4)) essentially reduces to the Kissinger one (Equation (3), *i.e.* $E_c = E_{cK}$ [30].

The activation energy for viscous flow (E_{vf}) was estimated by following the approach proposed by Francis *et al.* [31] who suggested the application of a Kissinger-type model. Thus, a plot of $\ln \frac{\beta}{T_s^2}$ vs. $\frac{1}{T_g}$, followed by linear interpolation, gives a straight line of slope $\frac{E_{vf}}{R}$, from which E_{vf} can be easily calculated.

The validity of the JMA model was discussed by using the method proposed by Malek [24]. He showed that, for the model to be applicable, the crystallization rate $\frac{d\chi}{dt}$ must depend only on the fraction χ of the glass crystallized and the temperature T, *i.e.* the thermal history of the material must not affect the crystallization mechanism. This validity test is based on the analysis of a probe function $z(\chi) \propto \frac{d\chi}{dt}$, which in non-isothermal conditions is defined as:

$$z(\chi) = \phi T^2 \tag{5}$$

where ϕ is the specific heat flow.

According to the Equation (5), the function $z(\chi)$ does not depend on kinetic parameters, and the JMA model is valid if the maximum of $z(\chi)$ is located at $\chi = 0.63 \pm 0.02$ [32].

3 Results and discussion

Figure 1 shows the DTA plots of 47.5B glass powder at different heating rates. A single exothermic peak of crystallization was recorded, which is consistent with previous investigations on this glass system [19]. The glass transition and peak crystallization temperatures can be observed to increase as the heating rate β increases from 10 to 40°C/min. Specifically, T_g and T_p vary within 527-567°C and 765-848°C, respectively.

The JMA coefficient n was initially calculated from the Ozawa plots at five different temperatures (770, 780, 790, 800 and 810° C) (Figure 2) and found equal to 1.4 ± 0.03. Assessment of n by the Augis-Bennet method yielded the value 1.8 ± 0.2, which is comparable to the Ozawa's one.

Figure 3 shows the Kissinger plot for the calculation of $E_{c,K}$; the least squares fitting yielded an apparent activation energy for crystallization of 131 kJ/mol. However, since both the Ozawa model and the Augis-Bennet method vield a value of n > 1, the application of the Kissinger equation is not appropriate and could lead to an underestimation of the activation energy. As the crystal growth dimensionality m can be calculated as n - 1 in the case of a nonconstant number of nuclei, which happens during the DTA experiments while crystal nucleation occurs, we can assume that m is approximately equal to 1 [33, 34]. This value indicates that surface crystallization is dominant with onedimensional rod-like crystallization in 47.5B glass particles. Once the values of n and m have been calculated, it is possible to apply the equation proposed by Matusita and Sakka (Equation (4)) [29] to calculate the activation energy for crystallization in a more reliable way as compared to the Kissinger model. The relevant plot is reported in Figure 4 and yielded an energy value E_c of 251 kJ/mol (> $E_{c,K}$).



Figure 1: DTA thermographs of 47.5B glass powders collected at different heating rates (β).



Figure 2: Plots in accordance with the Ozawa equation to determine the JMA parameter (n).

Erol *et al.* [35] suggested a simplified approach to calculate E_c starting from the (incorrect) value of $E_{c,K}$ when $m \neq n \neq 1$. Equations (3) and (4) can be combined together to obtain the following formula:

$$E_c = \frac{n}{m} E_{c,K} - 2\frac{n-1}{m} RT_p \tag{6}$$

As the condition $E_c >> 20RT_p$ is verified in most oxide glass systems [36], the term $2\frac{n-1}{m}RT_p$ in Equation (6) can be eliminated introducing an error of less than 10% in the value of



Figure 3: Plot in accordance with the Kissinger equation to determine the apparent activation energy for crystallization (E_{cK}).



Figure 4: Plot in accordance with the Matusita and Sakka equation to determine the correct activation energy for crystallization (E_c).

E_c. Hence, the following simple relation can be obtained:

$$E_c \approx \frac{n}{m} E_{c,K} \tag{7}$$

The activation energy for crystallization obtained by Equation (7) is 236 kJ/mol, which actually differs by only 6.4 % from the value assessed by Matusita-Sakka interpolation (251 kJ/mol) (Figure 4).

It is interesting to observe that the value of E_c for 47.5B glass powder assessed in this study is lower than that observed for powder of classical 45S5 Bioglass® with the same particle size range below 32 µm (e.g. 311 kJ/mol in [37] and 381-399 kJ/mol in [22]); a comparison of the relevant parameters is shown in Table 1. This depends on the different glass formulations and, hence, on the bonding energies involved. Because of the high bonding strengths of Si-O and P-O bonds, it is unlikely that a significant number of these bonds is broken during devitrification; rather, as suggested by Clupper and Hench [21], crystallization is expected to involve the breaking of modifier-O bonds that may have bonding strengths lower than E_c . In particular, K–O bonds having a significantly low bonding energy [38] are present in 47.5B glass and absent in 45S5 Bioglass[®], which suggests that the former may require a compara-

Table 1: Thermal and kinetic parameters of 47.B and 45S5 Bioglass[®] particles.

Glass powder (<32 μm)	n	m	E_c (kJ/mol)	E _{vf} (kJ/mol)	K_{H}^{a}
47.5B	1.4-1.8	1	251	181	0.40
45S5 Bioglass®	1 [22, 37]	1 [22, 37]	311 [37], 381-399 [22]	338 [37]	0.066 [16]

 \overline{a} The Hruby parameter was determined using a heating rate $\beta = 20$ °C/min in the DTA experiment, with Tg = 542°C, Tx = 670°C and Tm = 990°C.

tively lower activation energy for crystallization. Furthermore, it is worth observing that the crystalline phases formed during devitrification of the two materials are different, and this could also justify the different values of E_c . In fact, 47.5B tends to crystallize to Na₂CaSi₃O₈ after thermal treatment at 800°C [19], whereas other phases were reported to develop in 45S5 Bioglass[®]. In this regard, phase assignment of crystallized 45S5 Bioglass[®] is not unanimous. Over the years, different research groups suggested different possible scenarios, i.e. crystallization of $Na_2Ca_2Si_3O_9$ [9, 21] or $Na_2CaSi_2O_6$ [8, 15] as the main phase, with the optional presence of $Na_2Ca_4(PO_4)_2SiO_4$ (silicorhenanite) as a secondary phase [39]. In the attempt to solve this controversy Bellucci et al. [40] hypothesized the coexistence of $Na_2Ca_2Si_3O_9$ and $Na_2CaSi_2O_6$, but this topic is still under debate.

It is also interesting to highlight that previous SEM observations on partially-devitrified 47.5B glass revealed the needle-like morphology of Na₂CaSi₃O₈ crystals [19], which is in agreement with the dimensionality of crystal growth (m = 1, one-dimensional rod-like) assessed in this work.

The Kissinger-type plot used for estimating the viscous flow activation energy is shown in Figure 5 and the least squares interpolation yields a value of 181 kJ/mol. E_{vf} is lower than E_c , showing that the energy needed for viscous flow to occur is lower than that required for diffusion to progress the crystallization. This confirms that 47.5B glass exhibits a higher stability against crystallization as compared to 45S5 Bioglass[®], which conversely has $E_c < E_{vf}$ (Table 1) and a rapid tendency to crystallize just above T_g [37].

This behavior can be further quantified by calculating the Hruby parameter that is defined as follows [41]:

$$K_H = \frac{T_x - T_g}{T_m - T_x} \tag{8}$$

where T_x is the onset of crystallization and T_m the melting temperature determined from DTA plot.

The Hruby parameter K_H serves as a measure of the glass stability against crystallization: the larger K_H of a certain glass, the greater its stability against crystallization upon heating [42]. As shown in Table 1, the Hruby parameter of 47.5B glass (0.40) is six times higher than that cal-



Figure 5: Kissinger-type plot for the determination of the activation energy for the viscous flow ($E_{\nu f}$).



Figure 6: Plots of the normalized $z(\chi)$ functions (Malek test).

culated for 45S5 Bioglass[®] powder (0.066) [16] under the same experimental conditions.

The JMA model validity was assessed by applying the method proposed by Malek [24]. The normalized $z(\chi)$ functions, reported in Figure 6, have a typical bell-type shape and their maximum varies from 0.40 to 0.48 as the hating rate β increases from 10 to 40°C/min. Therefore, the crystallization process of 47.5B bioactive glass can be just approximately described by the JMA model, which is valid only if the maximum of the $z(\chi)$ function is located at $\chi = 0.63 \pm 0.02$ [32]. This suggests that the crystallization mechanism in 47.5B glass is more complex than a simple process of crystal nucleation and growth and may also change as a function of the heating rate. Future investigations could

be addressed to elucidate whether, for example, glass-inglass phase separation takes place during the crystallization of 47.5B, as observed in other bioactive glass systems such as 45S5 Bioglass[®] [8, 22, 43].

4 Conclusions

This work investigated the crystallization process of a SiO₂-P₂O₅-CaO-MgO-Na₂O-K₂O bioactive glass under non-isothermal conditions. The analyses gave evidence of a complex crystallization process where surface crystallization associated to the development of rod-like crystals was dominant. The activation energy for viscous flow was found lower than that required for crystallization, which suggests the capability of this glass to achieve significant densification before crystallization begins when it is thermally-treated above T_g . From a methodological viewpoint, this study also demonstrates that using multiple and complementary methods, i.e. the Ozawa and Augis-Bennet methods for determining the JMA exponent n as well as the Kissinger and Matusita-Sakka equations for assessing the crystallization activation energy, is fundamental to accurately describe the crystallization parameters.

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