POLITECNICO DI TORINO
Repository ISTITUZIONALE

Fourier integral operators and the index of symplectomorphisms on manifolds with boundary

Availability:

This version is available at: 11583/2624976 since: 2015-12-04T17:18:21Z

Publisher:
Academic Press Inc.

Published
DOI:10.1016/j.jfa.2015.06.001

Terms of use:

This article is made available under terms and conditions as specified in the corresponding bibliographic description in the repository

Publisher copyright

(Article begins on next page)

03 May 2024

FOURIER INTEGRAL OPERATORS OF
BOUTET DE MONVEL TYPE

U. BATTISTI, S. CORIASCO AND E. SCHROHE

Abstract. Given two compact manifolds with boundary X, Y, and a boundary preserving symplectomorphism $\chi : T^*Y \setminus 0 \to T^*X \setminus 0$, which is one-homogeneous in the fibers and satisfies the transmission condition, we introduce Fourier integral operators of Boutet de Monvel type associated with χ. We study their mapping properties between Sobolev spaces, develop a calculus and prove a Egorov type theorem. We also introduce a notion of ellipticity which implies the Fredholm property. Finally, we show how – in the spirit of a classical construction by A. Weinstein – a Fredholm operator of this type can be associated with χ and a section of the Maslov bundle. If $\dim Y > 2$ or the Maslov bundle is trivial, the index is independent of the section and thus an invariant of the symplectomorphism.

Introduction

We develop a calculus of Fourier integral operators (FIOs) on compact manifolds with boundary, which extends the calculus of pseudodifferential boundary value problems defined by Boutet de Monvel [3]. Given two compact manifolds with boundary, Y and X, we base our operators on symplectomorphisms $\chi : T^*Y \setminus 0 \to T^*X \setminus 0$, which are positively homogeneous of degree 1 in the fibers. In case $X = Y$ and $\chi = id$ we recover the Boutet de Monvel calculus.

Apart from the general interest in operators of this type, our main objective is to provide the analytic framework for an index problem in the spirit of A. Weinstein [32]. Weinstein considered two *closed* manifolds and a corresponding symplectomorphism between the cotangent bundles with the zero section removed. He proved that this symplectomorphism defines, in a natural way, a FIO F with the Fredholm property and that its index is a remarkable quantity: Let X and Y be additionally riemannian with Laplacians Δ_X and Δ_Y and suppose that the principal symbols satisfy $\sigma(\Delta_Y) = \sigma(\Delta_X) \circ \chi$. Denote by $\lambda_j(X)$ and $\lambda_j(Y)$ their sequences of eigenvalues. Then, under a mild additional assumption (the Maslov class of χ has to be $\equiv 0 \bmod 4$) the sequence $|\lambda_j(Y) - \lambda_{j-\mathrm{ind}\, F}(X)|$ is bounded. Weinstein then asked for an expression of this index in topological terms. See [33] for on overview of the problem.

In this article we show that, under suitable natural assumptions on the symplectomorphism χ, we can associate with χ a FIO in the calculus we develop, and establish the Fredholm property.

In the boundaryless case, C. Epstein and R. Melrose [10] solved Weinstein's problem under the assumption that both manifolds coincide, relying on previous results by V. Guillemin [17], L. Boutet de Monvel [4], and S. Zelditch [34]. They reduced the task to the computation of the index of a Dirac operator on a closed manifold, constructed from the data, and thus to the Atiyah-Singer index theorem. This construction has been refined by C. Epstein in [11], [12], [13]. The general case, where X and Y are closed but possibly different, was treated by E. Leichtnam,

R. Nest and B. Tsygan [22] in the framework of deformation quantization. A generalization to symplectomorphisms on a manifold with conical singularities has been studied by V. E. Nazaĭkinskiĭ, B. Sternin and B.-W. Schulze [26], [25], relying on the work of Epstein and Melrose. In the present article, however, we shall not tackle the problem of the computation of the index.

A second motivation for considering this class of FIOs is a theorem of J.J. Duistermaat and I. Singer [9]. They showed that – under a mild topological condition – every order-preserving isomorphism $i : L_{cl}(X) \to L_{cl}(Y)$ between the algebras of classical pseudodifferential operators on closed manifolds X and Y, respectively, is of the form $i(A) = F^{-1}AF$, where F is a FIO associated with a symplectomorphism as above. Recently, V. Mathai and R. Melrose [23] found a proof which avoids the topological condition. An analog of the Duistermaat-Singer theorem in the semiclassical setting has been given by H. Christianson [6].

We show in this article that conjugation with a FIO in our class provides an order-preserving isomorphism of Boutet de Monvel's algebra and we expect these to be all.

The paper is organized as follows. In Section 1, we define the class of symplectomorphisms we work with. A variant of Moser's trick shows that a symplectomorphism χ as above can always be extended to a symplectomorphism $\tilde{\chi} : T^*\widetilde{Y} \setminus 0 \to T^*\widetilde{X} \setminus 0$, where $\widetilde{Y}$ and $\widetilde{X}$ are neighborhoods of X and Y, respectively, in larger closed manifolds. The homogeneity, together with the fact that it preserves the boundary, implies that χ induces a symplectomorphism $\chi_\partial : T^*\partial Y \setminus 0 \to T^*\partial X \setminus 0$, which is the lift of a diffeomorphism $b : \partial Y \to \partial X$.

We then analyze operators of the form $r^+A^\chi e^+$, where A^χ is a (FIO) associated with $\tilde{\chi}$, r^+ is the restriction operator to int X and e^+ is the extension-by-zero operator on functions in Y.

As χ preserves the boundary, $r^+A^\chi e^+$ maps $C^\infty(Y)$ to $C^\infty(\text{int } X)$. We require in addition that each component of χ satisfies the transmission condition. This implies the continuity of $r^+A^\chi e^+ : C^\infty(Y) \to C^\infty(X)$ and $r^+(A^\chi)^*e^+ : C^\infty(X) \to C^\infty(Y)$. In fact, it is also necessary, as we will show in a forthcoming paper [2].

The assumptions on χ place the analysis here in a framework which is in a sense complementary to that considered by A. Hirschowitz and A. Piriou in [18]. They studied the transmission property for Fourier distributions conormal to hypersurfaces in $T^*X \setminus 0$.

In Section 2, we establish the continuity properties of the above truncated FIOs $r^+A^\chi e^+$, relying in a crucial way on the technique of operator-valued symbols. We prove that, for a symbol $a \in S^m(\mathbb{R}^n \times \mathbb{R}^n)$ satisfying the transmission condition,

$$r^+\text{Op}_n^\psi(a)e^+ : u \mapsto r^+ \int e^{i\psi(x', x_n, \xi', \xi_n) - i\psi_\partial(x', \xi')}a(x', x_n, \xi', \xi_n)\widehat{e^+u}(\xi_n)\,d\!\!\!{}^-\xi_n$$

is an operator-valued symbol in $S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+)\right)$. Here ψ is a phase function which locally represents χ, while the phase ψ_∂ represents the symplectomorphism χ_∂. A key point is the analysis of $r^+\text{Op}_n^\psi(a)\delta_0^{(j)}$, where $\delta_0^{(j)}$ is the j-th derivative of the Dirac distribution at the origin, see Theorem 2.6. In contrast to the corresponding result in Boutet de Monvel's calculus, however, it is not true that $r^+\text{Op}_n^\psi(a)e^+$ belongs to $S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; H^s(\mathbb{R}_+), H^{s-m}(\mathbb{R}_+)\right)$ for each $s \in \mathbb{R}$ as we show in Remark 2.12. The section ends with the proof of the continuity of $r^+A^\chi e^+$ in the scale of Sobolev spaces.

In Section 3 we complement the above truncated FIOs to matrices of Boutet de Monvel type operators of the form

$$\mathcal{A} := \begin{pmatrix} r^+ A^{\chi} e^+ + G^{\chi_\partial} & K^{\chi_\partial} \\ T^{\chi_\partial} & S^{\chi_\partial} \end{pmatrix} : \begin{matrix} C^\infty(Y, E_1) \\ \oplus \\ C^\infty(\partial Y, F_1) \end{matrix} \rightarrow \begin{matrix} C^\infty(X, E_2) \\ \oplus \\ C^\infty(\partial X, F_2), \end{matrix}$$

acting between sections of vector bundles E_1 over Y, E_2 over X, F_1 over ∂Y and F_2 over ∂X. Here, G^{χ_∂}, K^{χ_∂}, T^{χ_∂}, S^{χ_∂} are FIOs with Lagrangian submanifold defined by the graph of χ_∂ and, respectively, a singular Green symbol g of order m and type d, a potential symbol k of order m, a trace symbol t of order m and type d, a symbol $s \in S^m(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1})$.

The set of such operators $\mathcal{A}$ is denoted by $\mathscr{B}_\chi^{m,d}(X \times Y)$. We then develop the local version of a calculus which is – under the usual restrictions – closed under composition, that is

$$\mathscr{B}_{\chi_1}^{m_1,d_1} \circ \mathscr{B}_{\chi_2}^{m_2,d_2} \subseteq \mathscr{B}_{\chi_1 \circ \chi_2}^{m_1+m_2,d}, \quad d = \max\{m_2 + d_1, d_2\}.$$

We conclude the section with the proof of a Egorov type theorem for this class of operators, see Theorem 3.8.

In Section 4, the principal symbols of operators in $\mathscr{B}_\chi^{m,d}(X \times Y)$ are introduced. Similarly as in Boutet de Monvel's calculus, the interior principal symbol $\sigma(\mathcal{A})$ is defined as the principal symbol of A^χ, restricted to $T^*Y \setminus 0$. The technique of $\mathscr{L}(\mathscr{S}(\mathbb{R}_+))$-valued symbols then enables us to identify also a homogeneous operator-valued boundary principal symbol $\sigma_\partial(\mathcal{A})$. Ellipticity, defined as the invertibility of both, then allows the construction of a parametrix in the calculus.

In Section 5, we extend the above local calculus to compact manifolds with boundary. One result we obtain is the following: Whenever $\mathcal{A} \in \mathscr{B}_\chi^{0,0}(X \times Y)$ is invertible,

(1) $$j: \mathscr{B}^{m,d}(X) \rightarrow \mathscr{B}^{m,d}(Y): \mathcal{P} \mapsto \mathcal{A}^{-1} \circ \mathcal{P} \circ \mathcal{A}$$

is defined for all m and d, hence extends to an isomorphism between algebras of Boutet de Monvel operators preserving the order, in the spirit of [9].

In the last Section 6 we show how an index can be associated with an admissible symplectomorphism χ.

To this end we first reduce to the case of a one by one matrix. Following the approach of A. Weinstein in [32], the natural candidate for a Fredholm operator associated with χ is

$$\mathcal{U} = (r^+ U^\chi e^+)$$

where U^χ is a FIO defined by χ with a principal symbol s which is a unitary section of the associated Maslov bundle. We then establish the Fredholm property of this operator by showing the invertibility of $\sigma(\mathcal{A})$ and $\sigma_\partial(\mathcal{A})$. At this point the analysis is more subtle than in the case of closed manifolds. A priori, it is not clear why the boundary symbol should be invertible. In order to show this, we use a deformation of the phase function via a scaling of the normal variable. In the limit, we obtain an invertible operator. However, as the phase is, in general, discontinuous at the zero section, this is not a continuous deformation on $L^2(\mathbb{R}_+)$. Instead, we work on the weighted space $L^2(\mathbb{R}_+, (1 + x^2)^{-1}dx)$, where Schur's lemma implies the desired norm continuity.

It turns out that the index of $\mathcal{U}$ is independent of the choice of the section s whenever the Maslov bundle is trivial or $\dim Y > 2$.

For the case $X = Y$, examples of admissible symplectomorphisms can be constructed by deforming the identity by means of a Hamiltonian flow. Of course, the index of the associated operator $\mathcal{U}$ will then be zero. In this respect, the

situation is similar to the case of closed manifolds, where explicit examples of symplectomorphisms with nonzero index are lacking.

1. Admissible Phase Functions

The cotangent bundle T^*Y of a manifold with boundary $(Y, \partial Y)$ is a symplectic manifold with boundary $T^*_{\partial Y} Y$. In this article we consider two compact n-dimensional manifolds with boundary $(Y, \partial Y)$ and $(X, \partial X)$ and a symplectomorphism

$$(2) \qquad \chi : T^*Y \setminus 0 \to T^*X \setminus 0$$

which is positively homogeneous of degree one in the fibers. We require that χ preserves the boundary, that is

$$(3) \qquad \chi\,(\partial T^*Y \setminus 0) = \partial T^*X \setminus 0.$$

The following lemma, which is proven in [24], analyzes symplectomorphisms of this type.

Lemma 1.1. *Under assumptions* (2) *and* (3), χ *induces a symplectomorphism*

$$(4) \qquad \chi_\partial : T^*\partial Y \setminus 0 \to T^*\partial X \setminus 0,$$

positively homogeneous of degree one in the fibers, such that the following diagram commutes:

$$
\begin{CD}
T^*_{\partial Y} Y \setminus N^*\partial Y @>\chi>> T^*_{\partial X} X \setminus N^*\partial X \\
@V{i^*_Y}VV @VV{i^*_X}V \\
T^*\partial Y \setminus 0 @>\chi_\partial>> T^*\partial X \setminus 0.
\end{CD}
$$

Remark 1.2. *In Lemma 1.1 we have considered the induced symplectomorphism χ_∂ outside the zero section. Actually, since χ is smooth on $\partial T^*Y \setminus 0$, the induced symplectomorphism χ_∂ is also smooth at the zero section. Since χ_∂ is positively homogeneous of degree one in the fibers, the smoothness at the zero section implies that χ_∂ is trivial in the fibers. That is, χ_∂ is the lift of a diffeomorphism*

$$(5) \qquad b : \partial Y \to \partial X$$

of the boundaries (see [5]*).*

The manifolds X and Y embed into closed manifolds of the same dimension. Moser's trick, see [5, Ch. 7], then allows us to extend χ to a symplectomorphism

$$(6) \qquad \widetilde{\chi} : T^*\widetilde{Y} \setminus 0 \to T^*\widetilde{X} \setminus 0,$$

positively homogeneous of degree one in the fibers over neighborhoods $\widetilde{X}$ of X and $\widetilde{Y}$ of Y in these closed manifolds.

4

It will be important to understand the form of the Jacobian of χ in a collar neighborhood of the boundary. We write

$$\chi(y', y_n, \eta', \eta_n)$$

(7)
$$= (x'^*(y', y_n, \eta', \eta_n), x_n^*(y', y_n, \eta', \eta_n), \xi'^*(y', y_n, \eta', \eta_n), \xi_n^*(y', y_n, \eta', \eta_n)).$$

We suppose that the coordinates (y', y_n, η', η_n), $(x'^*, x_n^*, \xi'^*, \xi_n^*)$ determine a collar neighborhood of the boundary, y_n, x_n^* being boundary defining functions. Since the boundary is preserved, $x_n^*(y', 0, \eta', \eta_n) = 0$ for all (y', η', η_n). Hence, $\partial_{y'} x_n^*$, $\partial_{\eta'} x_n^*$, $\partial_{\eta_n} x_n^*$ are identically zero at $y_n = 0$. Moreover, Lemma 1.1 implies that the restrictions $x_\partial'^*$ and $\xi_\partial'^*$ of x'^* and ξ'^* to $y_n = 0$ locally define the symplectomorphism χ_∂. These restrictions are then independent of the conormal direction, that is $\partial_{\eta_n} x'^*$ and $\partial_{\eta_n} \xi'^*$ are identically zero at the boundary. Hence, we find that the Jacobian of χ at the boundary has the form

(8)
$$J(\chi)_{T_{\partial Y}^* Y} = \left. \left(\begin{array}{cccc} \partial_{y'} x'^* & \partial_{\eta'} x'^* & \partial_{y_n} x'^* & 0 \\ \partial_{y'} \xi'^* & \partial_{\eta'} \xi'^* & \partial_{y_n} \xi'^* & 0 \\ 0 & 0 & \partial_{y_n} x_n^* & 0 \\ \partial_{y'} \xi_n^* & \partial_{\eta'} \xi_n^* & \partial_{y_n} \xi_n^* & \partial_{\eta_n} \xi_n^* \end{array} \right) \right|_{y_n=0}.$$

Moreover, since

(9)
$$J_{\chi_\partial} = \left(\begin{array}{cc} \partial_{y'} x_\partial'^* & \partial_{\eta'} x_\partial'^* \\ \partial_{y'} \xi_\partial'^* & \partial_{\eta'} \xi_\partial'^* \end{array} \right)$$

is a symplectic matrix, it has determinant 1. Clearly also $J(\chi)_{T_{\partial Y}^* Y}$ has determinant equal to 1, because χ is a symplectomorphism. This implies that $\partial_{y_n} x_n^*|_{y_n=0} \cdot \partial_{\eta_n} \xi_n^*|_{y_n=0} = 1$. In particular, $\partial_{y_n} x_n^*|_{y_n=0}$ and $\partial_{\eta_n} \xi_n^*|_{y_n=0}$ can never vanish. As the boundary is compact, they are bounded away from zero. We even see that $\partial_{y_n} x_n^*|_{y_n=0} > 0$, since x_n^* and y_n are boundary defining functions.

We recall a well-known property of Lagrangian subspaces, which extends to the case of manifolds with boundary.

Proposition 1.3. *Let $\Lambda \subset T^*Z \setminus 0$ be a conic Lagrangian submanifold. Then, for each $\lambda_0 = (z_0, \theta_0) \in \Lambda$, there exists a neighborhood U of z_0 and a phase function ϕ defined in a conic neighborhood $U \times \Gamma \subseteq U \times \mathbb{R}^N$ – N large enough – such that ϕ parametrizes Λ in a conic neighborhood of λ. That is*

$$C_\phi = \{(z, \theta) \mid \phi'_\theta(z, \theta) = 0\} \to T^*Z$$

$$(z, \theta) \mapsto (z, \phi'_z(z, \theta))$$

induces a diffeomorphism in $U \times \Gamma$.

If Λ is locally defined by the graph of a symplectomorphism, we have a splitting of the variable z as $z = (x, y)$, and $\Lambda \subseteq T^*X \times T^*Y$. In this case, we can choose a phase function of the particular form $\phi(x, y, \theta) = \psi(x, \theta) - y \cdot \theta$, with $\phi \in C^\infty(\Omega_{x_0} \times \Omega_{y_0} \times \Gamma)$, where Ω_{x_0} and Ω_{y_0} are neighborhoods of x_0 and y_0, respectively, Γ is a cone in $\mathbb{R}^n \setminus 0$, $2n$ is the dimension of Λ.

Remark 1.4. *We can apply Proposition 1.3 also to the symplectomorphism χ_∂ in order to obtain a phase function $\phi_\partial(x', y', \theta') = \psi_\partial(x', \theta') - y' \cdot \theta'$ which represents χ_∂. Since χ_∂ is the lift of a diffeomorphism, the phase function $\psi_\partial(x', \theta')$ is smooth at $\theta' = 0$ and therefore linear in θ'.*

We will not recall here the notion of Maslov bundle, see [19] for its precise description.

Lemma 1.5. *Under the above assumptions the Maslov bundle of*

(10)
$$\Lambda = \mathrm{Graph}(\chi)' = \{(x, \xi), (y, -\eta) \mid \chi(y, \eta) = (x, \xi)\} \subseteq T^*X \setminus 0 \times T^*Y \setminus 0$$

is trivial in a collar neighborhood of $\partial\Lambda = (\partial T^*X \times \partial T^*Y) \cap \Lambda$.

Proof. Let $\lambda_0 = (x_0, y_0, \xi_0, \eta_0) \in \Lambda$, $U_{\lambda_0}^\Lambda$ a conic neighborhood of λ_0 and $\phi(x, y, \theta) = \psi(x, \theta) - y \cdot \theta$ a phase function representing Λ on $U_{\lambda_0}^\Lambda$. Then the Maslov bundle is trivialized on $U_{\lambda_0}^\Lambda$ by $e^{\frac{i}{4}\operatorname{sgn}\phi_\theta''}$. By Remark 1.4, χ_∂ is linear. This implies that $\phi_{\theta'}''$ vanishes identically at the boundary for each covector θ' tangent to the boundary. Moreover, $\partial_{\theta_n}\phi$ is constant at the boundary, hence $\partial_{\theta_n}^2\phi$ vanishes identically at the boundary. Therefore, $e^{\frac{i}{4}\operatorname{sgn}\phi_\theta''} = 1$ identically in a small neighborhood of λ_0. As this holds for all $\lambda_0 \in \partial\Lambda$, the result follows from the compactness of the boundary. $\quad\square$

In general it is not possible to find a global phase function defining the whole Lagrangian submanifold. In fact, this is impossible whenever the Maslov bundle is not trivial, see [21] for the precise statement. In our setting, the triviality of the Maslov bundle implies the following:

Proposition 1.6. *For every $x_0 \in \partial X$ there exist neighborhoods U_{x_0} and U_{y_0}, of x_0 and $y_0 = b^{-1}(x_0)$, respectively, such that $\chi(T^*U_{y_0}) \subseteq T^*U_{x_0}$. Moreover it is possible to define phase functions*

$$
\begin{aligned}
\phi_L(x, y, \eta) &= \psi_L(x, \eta) - y \cdot \eta \\
\phi_R(x, y, \xi) &= x \cdot \xi - \psi_R(y, \xi)
\end{aligned}
\tag{11}
$$

which parametrize χ in $\Lambda \cap \left(T^*U_{x_0} \times T^*U_{y_0}\right)$, where Λ is given by (10).

Proof. We write

$$
\chi : (y, \eta) \mapsto (x^*(y, \eta), \xi^*(y, \eta))
$$

and

$$
\chi^{-1} : (x, \xi) \mapsto (y^*(x, \xi), \eta^*(x, \xi)).
$$

Since the symplectomorphism preserves the boundary it is possible to find neighborhoods U_{x_0}, U_{y_0} such that $\chi(T^*U_{y_0}) \subseteq T^*U_{x_0}$. Here we write $(x, \xi) = (x', x_n, \xi', \xi_n)$, $(y, \eta) = (y', y_n, \eta', \eta_n)$ with boundary defining functions x_n and y_n. In view of the considerations around (8), we can suppose - possibly restricting U_{y_0} - that

$$
\det\left(\partial_\eta \xi^*(y, \eta)\right) \neq 0 \text{ on } T^*U_{y_0},
\tag{12}
$$

and

$$
\det\left(\partial_\xi \eta^*(x, \xi)\right) \neq 0 \text{ on } T^*U_{x_0}.
\tag{13}
$$

Following the idea of [21] we introduce

$$
\begin{aligned}
\widetilde{\psi}_L(x, y, \xi) &= (y^*(x, \xi) - y) \cdot \eta^*(x, \xi) \\
\widetilde{\psi}_R(x, y, \eta) &= (x - x^*(y, \eta)) \cdot \xi^*(y, \eta).
\end{aligned}
$$

Since χ and χ^{-1} preserve the canonical 1-form we have that

$$
\begin{aligned}
\xi^* \cdot x_{\eta_k}^* &= 0, & \eta^* \cdot y_{\xi_k}^* &= 0, \\
\xi^* \cdot x_{y_k}^* &= \eta_k, & \eta^* \cdot y_{x_k}^* &= \xi_k.
\end{aligned}
$$

The above relations together with the non degeneracy conditions (12), (13) imply that $\widetilde{\psi}_L$ and $\widetilde{\psi}_R$ are phase functions representing $\Lambda \cap \left(T^*U_{x_0} \times T^*U_{y_0}\right)$. In order to have the phase function as in (11), we use the inverse mapping theorem. In fact, by (13) and (12) it is possible to invert $\xi^*(y, \eta)$ and $\eta^*(x, \xi)$ in $T^*U_{y_0}$ and $T^*U_{x_0}$

respectively. We denote by $\tilde{\eta}(y, \xi)$ and $\tilde{\xi}(x, \eta)$ the inverse functions of $\xi^*(y, \eta)$ and $\eta^*(x, \xi)$, respectively. We then set

$$\phi_L(x, y, \eta) = (y^*(x, \tilde{\xi}(x, \eta)) - y) \cdot \eta$$
$$\phi_R(x, y, \xi) = (x - x^*(y, \tilde{\eta}(y, \xi))) \cdot \xi$$

and we obtain the assertion. $\square$

In order to define a suitable calculus for FIOs on manifolds with boundary, we need to introduce the transmission condition, see, e.g., [3, 14, 16, 27, 29]. Consider the function spaces:

$$H^+ = \{\mathscr{F}(e^+ u) \mid u \in \mathscr{S}(\mathbb{R}_+)\} \quad \text{and} \quad H_0^- = \{\mathscr{F}(e^- u) \mid u \in \mathscr{S}(\mathbb{R}_-)\},$$

where $\mathscr{S}(\mathbb{R}_\pm) = r^\pm \mathscr{S}(\mathbb{R})$ is the restriction of Schwartz functions on $\mathbb{R}$ to the right (left) half line, and $e^\pm$ is the extension by zero of a function defined on $\mathbb{R}_\pm$. It is easy to prove that the functions in H^+ and H_0^- decay to first order at infinity. Moreover, we denote by H' the set of all polynomials in one variable. Then we define

$$H = H^+ \oplus H_0^- \oplus H'.$$

Definition 1.7. *Let $a \in S^m(\mathbb{R}^n \times \mathbb{R}^n \times \mathbb{R}^n)$. Then a satisfies the transmission condition at $x_n = y_n = 0$ provided that, for all k, l,*

$$\partial_{y_n}^k \partial_{x_n}^l a(x', 0, y', 0, \xi', \langle \xi' \rangle \xi_n) \in S^m(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1} \times \mathbb{R}^{n-1}) \hat{\otimes}_\pi H_{\xi_n}.$$

We denote by $S_{tr}^m(\mathbb{R}^n \times \mathbb{R}^n \times \mathbb{R}^n)$ the subset of symbols satisfying the transmission condition.

For symbols positively homogeneous of degree m in ξ for large $|\xi|$, Definition 1.7 is equivalent to

$$(14) \qquad \partial_{x_n}^k \partial_{y_n}^l \partial_\xi^\alpha \partial_{x'}^\beta a(x', 0, y', 0, 0, 1) = (-1)^{m-|\alpha|} \partial_{x_n}^k \partial_{y_n}^l \partial_\xi^\alpha \partial_{x'}^\beta a(x', 0, y', 0, 0, -1)$$

for all $k, l \in \mathbb{N}, \alpha, \beta \in \mathbb{N}^{n-1}$. The above condition is often called symmetry condition. The proof of the equivalence can be found in [27].

Definition 1.8 (Admissible symplectomorphism). *We call a symplectomorphism χ as above admissible, if all the components of χ locally satisfy the transmission condition at the boundary. A phase function that represents an admissible symplectomorphism will be called admissible.*

Remark 1.9. *Definition 1.8 has an invariant meaning, because a change of coordinates in the cotangent bundle, induced by a change of coordinates in the base manifold, is linear with respect to the fibers. Hence, if the transmission condition is satisfied in one local chart then it is satisfied also in any other.*

Example 1.10. *The easiest way to construct an admissible symplectomorphism is to consider $X = Y$ and a Hamiltonian flow generated by a function $f \in C^\infty(T^*X \setminus 0)$ which is positively homogeneous of degree 1 in the fibers, satisfies the transmission conditions with respect to ∂X and has vanishing normal derivative at $T_{\partial X}^* X \setminus 0$.*

2. Oscillatory Integrals

In this section we will analyze the continuity properties of oscillatory integrals arising from FIOs associated with Lagrangian submanifolds obtained from admissible symplectomorphisms as in Definition 1.8. We will use the concept of operator-valued pseudodifferential operators acting on weighted Sobolev spaces over $\mathbb{R}_+$; see the Appendix for basic definitions and results.

Let us consider $\widetilde{A} \in I^m_{comp}(\widetilde{X} \times \widetilde{Y}, \widetilde{\Lambda})$, where $\widetilde{\Lambda} = \text{Graph}(\widetilde{\chi})'$. The definition of FIOs implies that for all $(x_0, y_0, \xi_0, \eta_0) = \lambda_0 \in \widetilde{\Lambda}$, the operator is microlocally given by a kernel of the type

$$(15) \qquad \int e^{i\phi(x', x_n, y', y_n, \xi', \xi_n)} a(x', x_n, y', y_n, \xi', \xi_n) \, d\!\!\!/\xi' \, d\!\!\!/\xi_n,$$

up to smooth kernels. We will focus on the situation where boundary points are involved, so we suppose that, at $(x_0, y_0) \in \widetilde{X} \times \widetilde{Y}$, the local coordinates $(x_1, \ldots, x_n, y_1, \ldots, y_n) = (x', x_n, y', y_n)$ are chosen so that x_n, y_n are boundary defining functions. We also identify the chart domains on $\widetilde{X}$ and $\widetilde{Y}$ with the corresponding open subsets $\Omega_x, \Omega_y \subset \mathbb{R}^n$. As the Lagrangian submanifold $\widetilde{\Lambda}$ is defined by the graph of a symplectomorphism, we can always assume that

$$\phi(x', x_n, y', y_n, \xi', \xi_n) = \psi(x', x_n, \xi', \xi_n) - iy' \cdot \xi' - iy_n \cdot \xi_n.$$

The phase function ϕ is (initially) defined in an open conic neighborhood Γ in $\Omega_x \times \Omega_y \times (\mathbb{R}^n \setminus \{0\})$, and the symbol $a(x, y, \xi)$ has support contained in Γ. We set

$$\Omega_x^\partial = \Omega_x \cap \{x_n = 0\}, \quad \Omega_y^\partial = \Omega_y \cap \{y_n = 0\},$$

and

$$(16) \qquad \Omega_x^+ = \Omega_x \cap \{x_n \geq 0\}, \quad \Omega_y^+ = \Omega_y \cap \{y_n \geq 0\}.$$

We also recall that, since the FIO is associated with a symplectomorphism, we can rely on representations both by left and right quantization, see [20, Ch. 25].

For convenience, we will proceed under the following technical assumptions.

Assumptions 1.
- *The amplitude a satisfies the transmission condition w.r.t. $x_n = 0$, $y_n = 0$. As χ is admissible, this is preserved under changes of coordinates.*
- *ψ is defined on $\mathbb{R}^n \times (\mathbb{R}^n \setminus \{0\})$. In fact, since a vanishes outside Γ, we can choose any good extension for ψ; see [1] for the extension that we will use below.*
- *Since the kernel (15) represents an operator on a compact manifold, a can be assumed to vanish unless $|x_n|$ and $|y_n|$ are small, or x', y' lie outside suitable compact subsets of $\mathbb{R}^{n-1}$. Moreover, since the kernels (15) are given modulo smoothing operators, it is no restriction to assume also that a vanishes for ξ in a neighborhood of the origin. Otherwise, we can insert a 0-excision function in the amplitude, which changes (15) by a smooth kernel.*

Assumptions 1 allow us to focus on oscillatory integrals of the type

$$(17) \qquad \int e^{i\psi(x', x_n, \xi', \xi_n) - iy' \cdot \xi' - iy_n \cdot \xi_n} a(x', x_n, y', y_n, \xi', \xi_n) \, d\!\!\!/\xi' \, d\!\!\!/\xi_n,$$

with ψ and a as above.

We will next analyze the action in the normal direction of an operator with kernel as in (17). Before, however, we introduce a class of functions which will be useful in the sequel.

Definition 2.1. *A function $a \in C^\infty(\mathbb{R}^{n-1}_{x'} \times \mathbb{R}_{x_n} \times \mathbb{R}^{n-1}_{\xi'} \times \mathbb{R}_{\xi_n})$ belongs to the set $BS^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; S^l(\mathbb{R})\right)$ if, for all $\alpha, \beta \in \mathbb{N}^{n-1}$,*

$$(x_n, \xi_n) \mapsto \partial_{\xi'}^\alpha \partial_{x'}^\beta a\left(x', \frac{x_n}{\langle \xi' \rangle}, \xi', \xi_n \langle \xi' \rangle\right) \in S^l(\mathbb{R} \times \mathbb{R})$$

and each seminorm can be estimated uniformly by $\langle \xi' \rangle^{m - |\alpha|}$. That is, for all $\gamma, \delta \in \mathbb{N}$ there exists a constant $C_{\gamma, \delta}$ such that

$$\left| \partial_{\xi_n}^\gamma \partial_{x_n}^\delta \left[\partial_{\xi'}^\alpha \partial_{x'}^\beta a\left(x', \frac{x_n}{\langle \xi' \rangle}, \xi', \xi_n \langle \xi' \rangle\right) \right] \right| \leq C_{\gamma, \delta} \langle \xi_n \rangle^{l - |\gamma|} \langle \xi' \rangle^{m - |\alpha|}, \quad x, \xi \in \mathbb{R}^n.$$

Definition 2.1 implies that $a \in BS^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}, S^m(\mathbb{R})\right)$, if $a \in S^m(\mathbb{R}^n \times \mathbb{R}^n)$. This a consequence of the fact that $\partial^\alpha_{\xi'}\partial^\beta_{x'} a \in S^{m-|\alpha|}(\mathbb{R}^n \times \mathbb{R}^n)$ and the estimate

$$\left| \partial^\gamma_{\xi_n} \partial^\delta_{x_n} \left[a\left(x', \frac{x_n}{\langle \xi' \rangle}, \xi', \xi_n \langle \xi' \rangle\right)\right] \right| \leq C \langle \xi' \rangle^m \langle \xi_n \rangle^{m-|\gamma|}, \quad x, \xi \in \mathbb{R}^n.$$

Moreover, it is clear that BS-spaces satisfy a multiplicative property, that is

$$BS^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; S^l(\mathbb{R})\right) \cdot BS^{m'}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; S^{l'}(\mathbb{R})\right)$$
$$\subseteq BS^{m+m'}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}, S^{l+l'}(\mathbb{R})\right).$$

The proof of the Lemma 2.2 and Theorem 2.4, below, can be found in [1].

Lemma 2.2. *Let $a \in S^m(\mathbb{R}^n \times \mathbb{R}^n)$ and ψ satisfy Assumptions 1. Then*

(18)
$$\partial^\alpha_{\xi'}\partial^\beta_{x'} \left(e^{i\psi(x', x_n, \xi', \xi_n) - i\psi_\partial(x', \xi')} a\left(x', x_n, \xi', \xi_n\right) \right)$$
$$= e^{i\psi(x', x_n, \xi', \xi_n) - i\psi_\partial(x', \xi')} \tilde{a}\left(x', x_n, \xi', \xi_n\right),$$

where $\tilde{a} \in BS^{m-|\alpha|}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; S^{m+|\beta|}(\mathbb{R})\right)$.

Remark 2.3. *Since χ preserves the boundary, $\partial_{\xi_n}\psi(x', 0, \xi', \xi_n)$ is identically equal to zero, hence $\psi(x', 0, \xi', \xi_n)$ is independent of ξ_n. We set*

(19)
$$\psi\left(x', 0, \xi', \xi_n\right) = \psi_\partial\left(x', \xi'\right)$$

and notice that ψ_∂ represents the symplectomorphism at the boundary χ_∂, described in Remark 1.2.

Theorem 2.4. *Let $a \in S^m_{tr}(\mathbb{R}^n \times \mathbb{R}^n)$ and ψ be as in Assumptions 1. Then, the operator*

$$\mathrm{Op}^\psi_n(a) : u \mapsto \int e^{i\psi(x', x_n, \xi', \xi_n) - i\psi_\partial(x', \xi')} a(x', x_n, \xi', \xi_n) \hat{u}(\xi_n) đ\xi_n$$

belongs to $S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}), \mathscr{S}(\mathbb{R})\right)$.

Remark 2.5. *Theorem 2.4 is also valid for $\left(\mathrm{Op}^\psi_n(a)\right)^t$, hence $\mathrm{Op}^\psi_{x_n}(a)$ can be extended to an operator-valued symbol in $S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}), \mathscr{S}'(\mathbb{R})\right)$. As $\iota : L^2(\mathbb{R}) \to \mathscr{S}'(\mathbb{R})$ and $e^+ : L^2(\mathbb{R}_+) \to L^2(\mathbb{R})$ can both be interpreted as operator-valued symbols of order 0, we also have*

(20)
$$\mathrm{Op}^\psi_n(a)e^+ \in S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; L^2(\mathbb{R}_+), \mathscr{S}'(\mathbb{R})\right).$$

Moreover, since the proof of Theorem 2.4 is based on the theory of SG FIOs, see [7, 8], it is also possible to prove that for every (x', ξ')

(21)
$$\mathrm{Op}^\psi_n(a)e^+ \in \mathcal{L}(L^2(\mathbb{R}_+), H^{-m}(\mathbb{R})).$$

In general, however, we can not replace $\mathscr{S}'(\mathbb{R}_+)$ by $H^{-m}(\mathbb{R}_+)$ in (20), see Remark 2.12.

The next step is to consider the action of an oscillatory integral as in (17) on derivatives of Dirac distributions at the origin.

Theorem 2.6. *Let $a \in S^m_{tr}(\mathbb{R}^n \times \mathbb{R}^n)$ and ψ be as in Assumptions 1. Then*

$$k_j(x', \xi') = r^+ \mathrm{Op}^\psi_n(a)\delta^{(j)}_0$$
$$= r^+ \int e^{i\psi(x', x_n, \xi', \xi_n) - i\psi_\partial(x', \xi')} a(x', x_n, \xi', \xi_n)\hat{\delta}^{(j)}_0(\xi_n) đ\xi_n$$

defines an operator-valued symbol in $S^{m+\frac{1}{2}+j}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathbb{C}, \mathscr{S}(\mathbb{R}_+)\right)$. Here, $\delta^{(j)}_0$ is the j-th derivative of the Dirac distribution at 0.

Proof. We start by considering the operator $\mathrm{Op}^\psi(a)$ acting on smooth functions defined on the whole of $\mathbb{R}^n$.

$$\mathrm{Op}^\psi(a): \ C_c^\infty(\mathbb{R}^n) \to C^\infty(\mathbb{R}^n)$$

$$u \mapsto \int e^{i\psi(x',x_n,\xi',\xi_n)} a(x',x_n,\xi',\xi_n)\hat{u}(\xi',\xi_n)\, đ\xi'\, đ\xi_n$$

$$= \int e^{i\psi_\partial(x',\xi')} \int e^{ir(x',x_n,\xi',\xi_n)} a(x',x_n,\xi',\xi_n)\hat{u}(\xi',\xi_n)\, đ\xi'\, đ\xi_n,$$

where $r(x',x_n,\xi',\xi_n) = \psi(x',x_n,\xi',\xi_n) - \psi_\partial(x',\xi')$. First let $j = 0$. Going over to the right quantization, we find

$$\mathrm{Op}^\psi(a)(\phi \otimes \delta_0)(x_n)$$

$$= \int e^{i\psi_\partial(x',\xi')} \int e^{i\psi(x',x_n,\xi',\xi_n)-\psi_\partial(x',\xi')} a(x',x_n,\xi',\xi_n)\hat{\delta}_0(\xi_n)\, đ\xi_n \hat{\phi}(\xi')đ\xi'$$

$$= \int e^{ix'\cdot\xi'+ix_n\cdot\xi_n-i\psi^{-1}(y',y_n,\xi',\xi_n)} a_R(y',y_n,\xi',\xi_n)\phi(y') \otimes \delta_0(y_n)\, dy'\, dy_n đ\xi'\, đ\xi_n$$

$$= \int e^{ix'\cdot\xi'-i\psi_\partial^{-1}(y',\xi')} \int e^{ix_n\cdot\xi_n-ir^{-1}(y',y_n,\xi',\xi_n)} a_R(y',y_n,\xi',\xi_n) \cdot$$

(22)
$$\cdot\phi(y') \otimes \delta_0(y_n)\, dy'\, dy_n đ\xi'\, đ\xi_n,$$

where the equality is modulo operators with smooth kernel. In (22),

$$r^{-1}(y',y_n,\xi',\xi_n) = \psi^{-1}(y',y_n,\xi',\xi_n) - \psi_\partial^{-1}(y',\xi'),$$

ψ^{-1} is the phase function representing the symplectomorphism χ^{-1}, and $\psi_\partial^{-1} = \psi^{-1}|_{y_n=0}$. Now, we focus on the action in the normal direction, namely, the expression

$$B\left(y',\xi',x_n\right)(\delta_0) = \int e^{ix_n\cdot\xi_n} \int e^{-ir^{-1}(y',y_n,\xi',\xi_n)} a_R\left(y',y_n,\xi',\xi_n\right) \delta_0(y_n)\, dy_n đ\xi_n.$$

The symbol a_R satisfies the transmission condition, so we can write

$$a_R\left(y',0,\xi',\xi_n\langle\xi'\rangle\right) = \sum_{k=0}^m s_k^R\left(y',\xi'\right)\xi_n^k\langle\xi'\rangle^k + \sum_{l=0}^\infty \lambda_l b_l^R\left(y',\xi'\right)\hat{h}_l(\xi_n),$$

where $s_k^R \in S^{m-k}\left(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1}\right)$, $\{\lambda_l\}_{l\in\mathbb{N}} \in l^1$, $\left\{b_l^R\right\}_{l\in\mathbb{N}} \subset S^m\left(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1}\right)$ is a null sequence, $\{h_l\}_{l\in\mathbb{N}}$ a null sequence in $\mathscr{S}(\mathbb{R}_+) \oplus \mathscr{S}(\mathbb{R}_-)$. By the definition of operators on distributions, we have for all $u \in C_c^\infty(\mathbb{R})$

$$\langle \kappa_{\langle\xi'\rangle^{-1}} B\delta_0, u\rangle = \langle \delta_0, B^t\left(\kappa_{\langle\xi'\rangle} u\right)\rangle$$

$$= \langle \delta_0, \langle\xi'\rangle^{\frac{1}{2}} \int e^{-ir^{-1}(y',x_n,\xi',\xi_n)+iy_n\cdot\xi_n} a_R(y',x_n,\xi',\xi_n)u(\langle\xi'\rangle y_n) dy_n đ\xi_n\rangle$$

$$= \langle\xi'\rangle^{\frac{1}{2}} \int a_R\left(y',0,\xi',\xi_n\langle\xi'\rangle\right)\hat{u}(-\xi_n)đ\xi_n$$

(23)
$$= \langle\xi'\rangle^{\frac{1}{2}} \sum_{k=0}^m s_k^R(y',\xi') \int \xi_n^k\langle\xi'\rangle^k \hat{u}(-\xi_n)đ\xi_n$$

$$+\langle\xi'\rangle^{\frac{1}{2}} \sum_{l=0}^\infty \lambda_l b_l^R\left(y',\xi'\right) \int \hat{h}_l(\xi_n)\hat{u}(-\xi_n)đ\xi_n.$$

Using the properties of the Fourier transform,

$$\langle \kappa_{\langle \xi' \rangle^{-1}} B\delta_0, u \rangle = \langle \xi' \rangle^{\frac{1}{2}} \sum_{k=0}^{m} s_k^R(y', \xi') \langle \xi' \rangle^k (-1)^k i^k \langle \delta_0^{(k)}, u \rangle$$

$$(24) \qquad\qquad + \langle \xi' \rangle^{\frac{1}{2}} \sum_{l=0}^{\infty} \lambda_l b_l^R(y', \xi') \int h_l(x_n) u(x_n) dx_n.$$

Applying the restriction operator r^+, all terms that depend on $\delta_0^{(k)}$ vanish, so we get

$$(25) \qquad \kappa_{\langle \xi' \rangle^{-1}} r^+ (B\delta_0)(y', x_n, \xi') = \langle \xi' \rangle^{\frac{1}{2}} \sum_{l=0}^{\infty} \lambda_l b_l^R(y', \xi') r^+ h_l(x_n).$$

Derivatives w.r.t. (x', ξ') can be treated in the same way. Hence $(r^+ B\delta_0)(y', \xi') \in S^{m+\frac{1}{2}}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathbb{C}, \mathscr{S}(\mathbb{R}_+))$. Inserting (25) into (22), we obtain

$$r^+ \mathrm{Op}^{\psi}(a)(\phi \otimes \delta_0)(x_n) = \int e^{ix' \cdot \xi' - i\psi_{\partial}^{-1}(y', \xi')} (r^+ B\delta_0)(y', \xi') \phi(y') dy' \, d\!\!\!\!/\xi'$$

$$= \langle \xi' \rangle \sum_{l=0}^{\infty} r^+ h_l(\langle \xi' \rangle x_n) \int e^{ix' \cdot \xi' - i\psi_{\partial}^{-1}(y', \xi')} \lambda_l b_l^R(y', \xi') \phi(y') dy' \, d\!\!\!\!/\xi'.$$

Switching back to the left quantization, we obtain, modulo smoothing operators, the symbol-kernel

$$\kappa_{\langle \xi' \rangle} r^+ \mathrm{Op}_n^{\psi}(a)\delta_0 = \langle \xi' \rangle^{\frac{1}{2}} \sum_{l=0}^{\infty} \lambda_l b_l(x', \xi') r^+ h_l(x_n),$$

with a suitable null sequence $\{b_l\} \subset S^m(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1})$. This implies the assertion for $j = 0$. The proof for $j > 0$ is similar. In fact, it is enough to notice that

$$(26) \qquad\qquad \delta_0^{(j)}(\xi_n) = (i\xi_n)^j \delta_0,$$

so we can follow the same steps, but with a symbol of order $m + j$.

Finally, we have to take into account all the seminorms of $\mathscr{S}(\mathbb{R}_+)$, hence to consider derivatives with respect to the x_n-variable. Lemma 2.2 implies that this step can be obtained from the previous one using a different symbol, which still satisfies Assumptions 1, since the phase is admissible. $\qquad\square$

Remark 2.7. *With the same notation as in Theorem 2.6, we obtain that*

$$r^- \int e^{i\psi(x', x_n, \xi', \xi_n) - i\psi_{\partial}(x', \xi')} a(x', x_n, \xi', \xi_n) \delta_0^{(j)} \, d\!\!\!\!/\xi_n$$

is a symbol in $S^{m+\frac{1}{2}+j}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathbb{C}, \mathscr{S}(\mathbb{R}_-))$.

Remark 2.8. *The proof of Theorem 2.6, in particular (23), shows more. Suppose that the symbol a vanishes at $x_n = 0$ to order $\geq m_+ + 1$, where $m_+ = \max\{m, 0\}$. Then $a_R(x', 0, \xi)$ has no polynomial part. Hence*

$$(27) \qquad\qquad e^+ r^+ \mathrm{Op}^{\psi}(a)\delta_0 = \chi_{\mathbb{R}_+} \mathrm{Op}^{\psi}(a)\delta_0,$$

where $\chi_{\mathbb{R}_+}$ is the characteristic function of $\mathbb{R}_+$.

If a vanishes even to order $\geq m_+ + N + 1$ for some $N \in \mathbb{N}$, then Equation (27) also holds with δ_0 replaced by $\delta_0^{(j)}$ for $j = 0, \dots, N$.

Remark 2.9. *We noted in the proof of Theorem 2.6 that $r^+ \mathrm{Op}_n^{\psi}(a)\delta_0^{(j)} = r^+ \mathrm{Op}_n^{\psi}(a\xi_n^j)\delta_0$, $j > 0$. If we then switch to the right quantization, we obtain, instead of the powers of ξ_n, derivatives of a_R with respect to y_n and factors of y_n-derivatives of ψ^{-1}. The top order term will involve only the j-th power of $\partial_{y_n} \psi^{-1}$, and no terms containing higher order*

derivatives of the phase. Finally, evaluating this at the boundary, we find an expression depending only on the right symbol at the boundary, multiplied by $(\partial_{y_n}\psi^{-1}|_{y_n=0})^j$. This remark will be crucial in the definition of the principal symbol in Section 4, more precisely in Theorem (4.2).

The action on the Dirac distributions at the origin is the key point to prove the continuity of the operator $r^+\mathrm{Op}_n^\psi(a)e^+$ in $\mathscr{S}(\mathbb{R}_+)$. Before, we need a technical lemma, whose proof we omit, since it is straightforward.

Lemma 2.10. *Let ζ be an excision function of the origin, that is $0 \le \zeta \le 1$, $\zeta(t) = 0$ for $|t| \le 1$ and $\zeta(t) = 1$ for $|t| \ge 2$. Then, the function $\xi_n \mapsto \zeta\left(\frac{\xi_n}{\langle\xi'\rangle}\right)\xi_n^{-l}$ belongs to $S^{-l}(\mathbb{R}\times\mathbb{R})$.*

Theorem 2.11. *Let a and ψ satisfy Assumptions 1. Then*

$$r^+\mathrm{Op}_n^\psi(a)e^+ : \mathscr{S}(\mathbb{R}_+) \to \mathscr{S}(\mathbb{R}_+)$$

$$u \mapsto r^+ \iint e^{i\psi(x',x_n,\xi',\xi_n)-i\psi_\partial(x',\xi')-iy_n\cdot\xi_n}a(x',x_n,\xi',\xi_n)e^+u(y_n)dy_n\,đ\xi_n$$

is an operator-valued symbol in $S^m\left(\mathbb{R}^{n-1},\mathbb{R}^{n-1};\mathscr{S}(\mathbb{R}_+),\mathscr{S}(\mathbb{R}_+)\right)$. Corresponding results hold for $r^-\mathrm{Op}_n^\psi(a)e^+$ and $r^+\mathrm{Op}_n^\psi(a)e^-$.

Proof. Choose ζ as in Lemma 2.10 and write

$$r^+\mathrm{Op}_n^\psi(a)(e^+u) = r^+A_1(e^+u) + r^+A_2(e^+u),$$

with

$$A_1 = \mathrm{Op}_n^\psi\left(\zeta\left(\frac{\xi_n}{\langle\xi'\rangle}\right)a\right) \text{ and } A_2 = \mathrm{Op}_n^\psi\left(\left(1-\zeta\left(\frac{\xi_n}{\langle\xi'\rangle}\right)\right)a\right).$$

Lemma 2.10 implies that both A_1 and A_2 satisfy the hypotheses of Theorem 2.6. Moreover, the kernel of the operator A_2 is smooth with Schwartz decay, due to the compact support of the symbol in the ξ_n variable and in the x_n variable. Each $\mathscr{S}(\mathbb{R}_+)$-seminorm of $r^+A_2e^+u$ is bounded by $\langle\xi'\rangle^m$, because the symbol belongs to $S^m(\mathbb{R}^n\times\mathbb{R}^n)$. Therefore $r^+A_2e^+$ belongs to $S^m\left(\mathbb{R}^{n-1},\mathbb{R}^{n-1};\mathscr{S}(\mathbb{R}_+),\mathscr{S}(\mathbb{R}_+)\right)$. We can write

$$(28) \quad r^+A_1e^+u = \int e^{i\psi(x',x_n,\xi',\xi_n)-i\psi_\partial(x',\xi')}\zeta\left(\frac{\xi_n}{\langle\xi'\rangle}\right)a(x',x_n,\xi',\xi_n)\frac{1}{(i\xi_n)^l}\widehat{(i\xi_n)^le^+u}(\xi_n)đ\xi_n.$$

The properties of the Fourier transform assure that $\widehat{i\xi_ne^+u}(\xi_n) = \widehat{e^+\partial u} + u(0)\widehat{\delta_0}$. Theorem 2.6 implies that the part of the integral depending on derivatives of Dirac distributions satisfies the desired bound. To analyze the other part and the derivatives with respect to x',ξ',x_n, we note that, by taking l large enough, we are left with an integral operator in the normal direction, for which the sup norm of the kernel can be estimated in terms of the seminorms of a times $\langle\xi'\rangle^m$. The details are left to the reader. $\square$

Remark 2.12. *The proof of Theorem 2.11 indicates that the order of the operator in the normal direction may increase if we differentiate w.r.t. the x'-variable. In contrast to the usual Boutet de Monvel calculus, it is not true in general that*

$$r^+\mathrm{Op}_n^\psi(a)e^+ \in S^m\left(\mathbb{R}^{n-1},\mathbb{R}^{n-1};H^s(\mathbb{R}_+),H^{s-m}(\mathbb{R}_+)\right).$$

This can be seen explicitly by means of the following example. Define

$$A : \mathscr{S}(\mathbb{R}^n) \to \mathscr{S}(\mathbb{R}^n)$$

$$(29) \qquad u \mapsto \iint e^{i[(x'-y')\cdot\eta'+(f(x')x_n-y_n)\cdot\eta_n]}u(y',y_n)dy'dy_n đ\eta' đ\eta_n,$$

12

where f is a strictly positive function. The phase function of the FIO A in (29) represents a symplectomorphism χ of $T^*\overline{\mathbb{R}^n}_+ = \overline{\mathbb{R}^{2n}_+}$ onto itself of the form

$$\chi : \overline{\mathbb{R}^{2n}}_+ \to \overline{\mathbb{R}^{2n}}_+$$

(30)
$$(y', y_n, \eta', \eta_n) \mapsto \left(y', f(y')^{-1}y_n, \eta' + f'(y')y_n\frac{\eta_n}{f(y')}, f(y')\eta_n\right).$$

The symplectomorphism χ is admissible since it preserves the boundary:

$$\chi(y', y_n, \eta', \eta_n) \in \partial\overline{\mathbb{R}^n}_+ \Leftrightarrow (y', y_n, \eta', \eta_n) \in \partial\overline{\mathbb{R}^n}_+, \text{ that is } y_n = 0.$$

It is linear in the fibers, therefore all components have the transmission property. Looking at the action along the normal direction, we see that (29) cannot be extended to an operator-valued symbol in $S^0\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; H^s\left(\mathbb{R}\right), H^s\left(\mathbb{R}\right)\right)$. Indeed,

$$\kappa_{\langle\eta'\rangle^{-1}}\partial_{x'_j}\iint e^{i(f(x')x_n - y_n)\cdot\eta_n}\kappa_{\langle\eta'\rangle}u(y_n)dy_n\,đ\eta_n$$

$$= \kappa_{\langle\eta'\rangle^{-1}}\left(\langle\eta'\rangle^{\frac{1}{2}}\iint e^{i(f(x')x_n - y_n)\cdot\eta_n}i(\partial_{x'_j}f)(x')x_n\eta_nu(\langle\eta'\rangle y_n)dy_n\,đ\eta_n\right)$$

$$= \kappa_{\langle\eta'\rangle^{-1}}\left(\langle\eta'\rangle^{-\frac{1}{2}}\iint e^{i(\langle\eta'\rangle f(x')x_n - z_n)\cdot\frac{\eta_n}{\langle\eta'\rangle}}i(\partial_{x'_j}f)(x')x_n\eta_nu(z_n)dz_n\,đ\eta_n\right)$$

$$= \kappa_{\langle\eta'\rangle^{-1}}\left(\langle\eta'\rangle^{-\frac{1}{2}}\int e^{if(x')x_n\cdot\eta_n}i(\partial_{x'_j}f)(x')x_n\eta_n\hat{u}\left(\frac{\eta_n}{\langle\eta'\rangle}\right)đ\eta_n\right)$$

$$= \kappa_{\langle\eta'\rangle^{-1}}\left(\langle\eta'\rangle^{\frac{1}{2}}\int e^{i\langle\eta'\rangle f(x')x_n\cdot\theta_n}i(\partial_{x'_j}f)(x')x_n\theta_n\langle\eta'\rangle\hat{u}(\theta_n)đ\theta_n\right)$$

$$= \kappa_{\langle\eta'\rangle^{-1}}\left(\langle\eta'\rangle^{\frac{3}{2}}\int e^{i\langle\eta'\rangle f(x')x_n\cdot\theta_n}(\partial_{x'_j}f)(x')x_n\widehat{\partial_{x_n}u}(\theta_n)đ\theta_n\right)$$

$$= \kappa_{\langle\eta'\rangle^{-1}}\left(\langle\eta'\rangle^{\frac{3}{2}}(\partial_{x'_j}f)(x')x_n\left(\partial_{x_n}u\right)(\langle\eta'\rangle f(x')x_n)\right)$$

$$= \left(\partial_{x'_j}f\right)(x')x_n\left(\partial_{x_n}u\right)(f(x')x_n).$$

Hence we lose one derivative and the operator-valued symbol does not send $H^s(\mathbb{R})$ into $H^s(\mathbb{R})$.

Now, we recall a technical lemma, proven in [27], p. 122.

Lemma 2.13. Let $a \in S^m(\mathbb{R}^n \times \mathbb{R}^n)$ be a symbol with the transmission property. Then there exists a symbol $a_1 \in S^m(\mathbb{R}^n \times \mathbb{R}^n)$ having the transmission property for all hyperplanes $x_n = \epsilon, \epsilon \geq 0$, such that

$$\partial_{x_n}^k\left(a\left(x', x_n, \xi', \xi_n\right) - a_1\left(x', x_n, \xi', \xi_n\right)\right)|_{x_n=0} = 0$$

for all $k \in \mathbb{N}$ and for all x', ξ', ξ_n. The symbol a_1 has the following expression

(31)
$$a_1(x', x_n, \xi', \xi_n) = \sum_{j=0}^{\infty}\frac{x_n^j}{j!}\partial_{x_n}^j a(x', 0, \xi', \xi_n)\phi(t_j x_n),$$

where ϕ is a cut-off function at the origin and $\{t_j\}$ is a sequence in $\mathbb{R}_+$ such that the series in (31) converges in $S^m\left(\mathbb{R}^n \times \mathbb{R}^n\right)$.

Proposition 2.14. Let a and ψ satisfy Assumptions 1. Then, it is possible to write

$$r^+\text{Op}_n^{\psi}(a)e^+ = r^+\text{Op}_n^{\psi}(a_d)e^+ + r^+\text{Op}_n^{\psi}(a_0)e^+,$$

with a_0 such that

(32)
$$\left(\text{Op}_n^{\psi}(a_0)e^+\right)u \in L^2(\mathbb{R}), \quad u \in \mathscr{S}(\mathbb{R}_+),$$

and $a_d \in S^m_{tr}(\mathbb{R}^n \times \mathbb{R}^n)$ is a polynomial in ξ_n.

Proof. The proof follows from Remark 2.8 and an observation in the proof of Theorem 2.11. Choose a, a_1 as in Lemma 2.13 and set $b = a - a_1$. In view of the transmission property of a we can write

$$\partial^j_{x_n} a(x', 0, \xi', \xi_n) = \sum_{k=0}^m a_{k,j}(x', \xi')\xi_n^k + \sum_{k=0}^\infty \lambda_{k,j} b_{k,j}(x', \xi') h_{k,j}\left(\frac{\xi_n}{\langle \xi' \rangle}\right),$$

with $a_{k,j} \in S^{m-k}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1})$, $b_{k,j} \in S^m(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1})$, $h_{k,j} \in H^+ \oplus H_0^-$, $(\lambda_{k,j})_k \in l^1$, and

$$b_{k,j}, h_{k,j} \to 0, \quad k \to \infty.$$

Then, we set

$$a_d(x', x_n, \xi', \xi_n) = \sum_{j=0}^\infty \sum_{k=0}^m \frac{x_n^j}{j!} a_{k,j}(x', \xi')\xi_n^k \phi(t_j x_n),$$

$$a_0^1(x', x_n, \xi', \xi_n) = \sum_{j=0}^\infty \sum_{k=0}^\infty \frac{x_n^j}{j!} \lambda_{k,j} b_{k,j}(x', \xi') h_{k,j}\left(\frac{\xi_n}{\langle \xi' \rangle}\right) \phi(t_j x_n),$$

$$a_0 = a_0^1 + b.$$

By construction, $a = a_d + a_0$. Notice that a_0^1 is a symbol of order zero w.r.t. the ξ_n-variable. We then conclude from (21) with $m = 0$ that $\mathrm{Op}^\psi_n(a_0^1)e^+u$ belongs to $L^2(\mathbb{R})$. To analyze $\mathrm{Op}^\psi_n(b)e^+u$ we use the scheme in Theorem 2.11. We split the operator into two parts: one smoothing and the other with a symbol vanishing to infinite order at $\xi_n = 0$. For the smoothing part, (32) holds. In the other part, we divide and multiply by ξ_n^l as in (28), with l arbitrary. We obtain a symbol which is of order $m - l$ in the ξ_n variable and derivatives of Dirac's delta up to the order $l - 1$. To handle these terms, we use Remark 2.8, valid for symbols vanishing to infinite order at the origin. Namely,

$$\tag{33} \mathrm{Op}^\psi_n\left(b(x', x_n, \xi', \xi_n)\zeta\left(\frac{\xi_n}{\langle \xi' \rangle}\right)(i\xi_n)^{-l}\right)\delta_0^{(j)} \in L^2(\mathbb{R}), \quad j \in \mathbb{N},$$

concluding the proof. $\qquad\square$

Remark 2.15. *We have proven that $r^+\mathrm{Op}^\psi_n(a)e^+$ is a continuous operator from $\mathscr{S}(\mathbb{R}_+)$ to itself, so it is possible to define the transposed operator*

$$\left(r^+\mathrm{Op}^\psi_n(a)e^+\right)^t.$$

It is important to stress that, in general,

$$\tag{34} \left(r^+\mathrm{Op}^\psi_n(a)e^+\right)^t u \neq r^+\left(\mathrm{Op}^\psi_n(a)\right)^t e^+ u, \quad u \in \mathscr{S}(\mathbb{R}_+).$$

A simple counterexample is the operator $(r^+\partial e^+)^t$, since

$$(r^+\partial_{x_n} e^+)^t u = -r^+\partial_{x_n} e^+ u - u(0)\delta_0.$$

Nevertheless, if a_0 is as in Proposition 2.14 and therefore (32) is fulfilled, then equality holds in (34). Indeed, for $u, f \in \mathscr{S}(\mathbb{R}_+)$,

$$\langle (r^+\mathrm{Op}^\psi_n(a_0)e^+)^t u, f \rangle = \langle u, r^+\mathrm{Op}^\psi_n(a_0)e^+ f \rangle = \langle e^+ u, \mathrm{Op}^\psi_n(a_0)e^+ f \rangle$$

$$\tag{35} = \langle \mathrm{Op}^\psi_n(a_0)^t e^+ u, e^+ f \rangle = \langle r^+(\mathrm{Op}^\psi_n a_0)^t e^+ u, f \rangle.$$

Theorem 2.11 then implies that $r^+\mathrm{Op}_n^\psi(a_0)e^+ \in S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}'(\mathbb{R}_+)\right)$. Since $C_c^\infty(\mathbb{R}_+)$ is dense in $\mathscr{S}'(\mathbb{R}_+)$, we can also define the action on $\mathscr{S}'(\mathbb{R}_+)$ by

$$(r^+\mathrm{Op}_n^\psi(a_0)e^+)u = \lim_{k\to\infty} r^+\mathrm{Op}_n^\psi(a_0)e^+u_k, \quad u_k \to u \text{ in } \mathscr{S}'(\mathbb{R}_+).$$

Lemma 2.16. *Let $u \in \mathscr{S}'(\mathbb{R})$ be smooth on $\mathbb{R}_+$. Then the following statements are equivalent:*

i) *For all $j \in \mathbb{N}$*

$$\lim_{x\to 0^+} \partial^j u(x) = c^j, \quad c^j \in \mathbb{C}.$$

In particular the function r^+u can be extended smoothly up to zero.

ii) *For all $j \in \mathbb{N}$ and for all sequences $\left\{\psi_m^j\right\}_{m\in\mathbb{N}} \subseteq C_c^\infty(\mathbb{R}_+)$ such that*

(36) $$\psi_m^j \to (-1)^j \delta_0^{(j)} \quad \text{in } \mathscr{S}'(\mathbb{R}),$$

we have

$$\lim_{m\to\infty} \langle u, \psi_m^j \rangle = c^j, \quad c^j \in \mathbb{C}.$$

There is a trivial continuous inclusion $i : C_c^\infty(\mathbb{R}_+) \to C_c^\infty(\mathbb{R})$ given by the extension by zero, so the limit (36) is well defined.

The proof is left to the reader. The implication *i)* $\Rightarrow$ *ii)* is almost trivial. For the converse one can argue by contradiction.

Theorem 2.17. *Let ψ and a satisfy Assumptions 1. By Proposition 2.14 we can write $a = a_d + a_0$, where $\mathrm{Op}_n^\psi(a_0)e^+$ maps $\mathscr{S}(\mathbb{R}_+)$ to $L^2(\mathbb{R})$. Hence $(e^+r^+ - 1)\mathrm{Op}_n^\psi(a_0)e^+ = -e^-r^-\mathrm{Op}_n^\psi(a_0)e^+$. Moreover, $r^-\mathrm{Op}_n^\psi(a_0)e^+$ extends to an operator*

(37) $$r^-\mathrm{Op}_n^\psi(a_0)e^+ : \mathscr{S}'(\mathbb{R}_+) \to \mathscr{S}(\mathbb{R}_-)$$

and defines a symbol in $S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_-)\right)$.

Notice that e^+ in (37) is not defined on $\mathscr{S}'(\mathbb{R}_+)$. The operator $r^-\mathrm{Op}_n^\psi(a_0)e^+$ is defined as the dual of $r^+\mathrm{Op}_n^\psi(a_0)^t e^- : \mathscr{S}(\mathbb{R}_-) \to \mathscr{S}(\mathbb{R}_+)$.

Proof. We want to prove that for all s_1, s_2, for all δ, γ, we have

$$p_{\delta,\gamma}\left(\kappa_{\langle\xi'\rangle^{-1}} r^-\mathrm{Op}_n^\psi(a_0)e^+ \kappa_{\langle\xi'\rangle} u\right) \prec \langle\xi'\rangle^m \|u\|_{H_0^{-s_1, -s_2}}, \quad u \in C_c^\infty(\mathbb{R}_+),$$

$\left\{p_{\delta,\gamma}\right\}$ being the seminorms of $\mathscr{S}(\mathbb{R}_-)$. By Assumptions 1, a vanishes for $|x_n| > \epsilon$, ϵ small. As $\partial_{\xi_n}\psi(x', x_n, \xi) \neq 0$ if $x_n \neq 0$, the phase function has no critical points on $\mathrm{supp}\, a_0$ outside $\{x_n = 0\}$. An integration by parts argument shows that $\mathrm{singsupp}\,\mathrm{Op}_n^\psi(a_0)u \subseteq \{x_n = 0\}$. Hence, we only need to consider the behavior as $x_n \to 0^-$. From Theorem 2.11 and Proposition 2.14, we recall that the following maps are continuous:

(38) $$r^-\mathrm{Op}_n^\psi(a_0)e^+ : \mathscr{S}(\mathbb{R}_+) \to \mathscr{S}(\mathbb{R}_-)$$

and

(39) $$r^-\mathrm{Op}_n^\psi(a_0)e^+ : \mathscr{S}'(\mathbb{R}_+) \to \mathscr{S}'(\mathbb{R}_-).$$

In order to prove that $r^-\mathrm{Op}_n^\psi(a_0)e^+$ belongs to $S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_-)\right)$, we have to analyze

(40) $$\lim_{x_n\to 0^-} \partial_{x_n}^k \kappa_{\langle\xi'\rangle^{-1}}\left(\partial_{x'}^\beta \partial_\xi^\alpha r^-\mathrm{Op}_n^\psi(a_0)e^+ \kappa_{\langle\xi'\rangle} u\right)(x_n).$$

We start with the case $|\alpha| = |\beta| = 0$. By definition, we have to prove that, for all s_1, s_2, k we have

$$\left| \lim_{x_n \to 0^-} \partial_{x_n}^k \kappa_{\langle \xi' \rangle^{-1}} \left(r^- \mathrm{Op}_n^{\psi}(a_0) e^+ \kappa_{\langle \xi' \rangle} u \right)(x_n) \right| \prec \langle \xi' \rangle^m \|u\|_{H_0^{-s_1, -s_2}(\mathbb{R})}, \quad u \in C_c^{\infty}(\mathbb{R}_+).$$

Using the idea of Lemma 2.16, we consider

$$\langle \kappa_{\langle \xi' \rangle^{-1}} r^- \mathrm{Op}_m^{\psi}(a_0) e^+ \kappa_{\langle \xi' \rangle} u, (-1)^k \partial_{x_n}^k \psi_l \rangle,$$

where $\{\psi_l\}_{l \in \mathbb{N}} \subseteq C_c^{\infty}(\mathbb{R}_-)$ is a sequence such that

(41) $$\psi_l \to \delta_0, \quad \text{in } \mathscr{S}'(\mathbb{R}).$$

Notice that (41) implies that $\kappa_{\langle \xi' \rangle} \psi_l$ converges to $\langle \xi' \rangle^{-\frac{1}{2}} \delta_0$. By Remark 2.15 and Proposition 2.14 we have

$$\langle \kappa_{\langle \xi' \rangle^{-1}} r^- \mathrm{Op}_n^{\psi}(a_0) \kappa_{\langle \xi' \rangle} e^+ u, \psi_l \rangle = \langle u, \kappa_{\langle \xi' \rangle^{-1}} r^+ \left(\mathrm{Op}_n^{\psi}(a_0)^t \kappa_{\langle \xi' \rangle} \right) e^- \psi_l \rangle.$$

By (39) and (41) we get

$$\lim_{l \to \infty} \langle u, \kappa_{\langle \xi' \rangle^{-1}} r^+ \mathrm{Op}_n^{\psi}(a_0)^t \kappa_{\langle \xi' \rangle} e^- \psi_l \rangle = \langle \kappa_{\langle \xi' \rangle^{-1}} u, r^+ \mathrm{Op}_n^{\psi}(a_0)^t \kappa_{\langle \xi' \rangle} \delta_0 \rangle.$$

By Theorem 2.6 we know that

$$r^+ \left(\mathrm{Op}_n^{\psi}(a_0)^t \right) \delta_0 \in S^{m+\frac{1}{2}} \left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathbb{C}, \mathscr{S}(\mathbb{R}_+) \right),$$

so, finally, since $\mathscr{S}(\mathbb{R}_-) = \text{proj-lim}_s H^s(\mathbb{R}_-)$,

$$\lim_{l \to \infty} \left| \langle u, \kappa_{\langle \xi' \rangle^{-1}} r^+ \mathrm{Op}_n^{\psi}(a_0)^t \kappa_{\langle \xi' \rangle} e^- \psi_l \rangle \right|$$

$$\leq \|\kappa_{\langle \xi' \rangle^{-1}} r^+ \mathrm{Op}_n^{\psi}(a_0)^t \kappa_{\langle \xi' \rangle} \delta_0\|_{H^{s_1, s_2}(\mathbb{R}_+)} \|u\|_{H_0^{-s_1, -s_2}(\overline{\mathbb{R}}_-)} \leq C \langle \xi' \rangle^m \|u\|_{H_0^{-s_1, -s_2}(\overline{\mathbb{R}}_-)}.$$

For the derivatives w.r.t. ξ' and x' of orders α and β, we use Lemma 2.2 and a slight variation of Theorem 2.6 in the setting of symbols belonging to the class $BS^{m-|\alpha|} \left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; S^{m+|\beta|}(\mathbb{R}) \right)$. □

Theorem 2.18. *Let a and ψ satisfy Assumptions 1. Then*

$$r^+ \mathrm{Op}^{\psi}(a) e^+ : H^s(\mathbb{R}_+^n) \to H^{s-m}(\mathbb{R}_+^n), \quad s > -\frac{1}{2},$$

continuously.

Proof. For $s \leq 0$ the result follows from the continuity of $e^+ : H^s(\mathbb{R}_+^n) \to H^s(\mathbb{R}^n)$, $r^+ : H^s(\mathbb{R}^n) \to H^s(\mathbb{R}_+^n)$ and the continuity properties of FIOs with homogeneous phase. It remains to consider the case $s > 0$. Using interpolation we may assume $s \in \mathbb{N}$. We write

$$r^+ \mathrm{Op}^{\psi}(a) e^+ = r^+ \mathrm{Op}^{\psi}(a) e^+ \circ \Lambda_+^{-s} \circ \Lambda_+^s,$$

where $\Lambda_+^s = r^+ \Lambda^s e^+$ is a truncated pseudodifferential operator in the sense of Boutet de Monvel such that $\Lambda_+^s : H^s(\mathbb{R}_+^n) \cong L^2(\mathbb{R}_+^n)$, and Λ_+^{-s} is the inverse of Λ_+^s. So, we only need to prove that $r^+ \mathrm{Op}^{\psi}(a) e^+ \circ \Lambda_+^{-s} : L^2(\mathbb{R}_+^n) \to H^{s-m}(\mathbb{R}_+^n)$ is continuous. We observe that

(42) $$r^+ \mathrm{Op}^{\psi}(a) e^+ \circ r^+ \Lambda_+^{-s} e^+ = r^+ \mathrm{Op}^{\psi}(a) \circ \Lambda^{-s} e^+ - r^+ \mathrm{Op}^{\psi}(a) \left(e^+ r^+ - 1 \right) \Lambda_+^{-s} e^+.$$

The operator $\mathrm{Op}^{\psi}(a) \circ \Lambda^{-s}$, by the properties of FIOs is, modulo operators with smoothing kernel, a FIO of order $m - s$ with phase ψ. Thus, $r^+ \mathrm{Op}^{\psi}(a) \circ \Lambda^{-s} e^+ : L^2(\mathbb{R}_+^n) \to H^{s-m}(\mathbb{R}_+^n)$ is continuous, since e^+ is continuous on L^2. Now, we have to analyze the second term of (42). We treat it as a FIO defined on the boundary with operator-valued symbol. Notice that Λ_+^{-s} is of negative order, and the differential part of the decomposition in Proposition 2.14 vanishes, so

$$r^+ \mathrm{Op}_n^{\psi}(a)(e^+ r^+ - 1)\Lambda^{-s} e^+ u = -r^+ \mathrm{Op}_n^{\psi}(a) e^- r^- \Lambda^{-s} e^+ u, \quad u \in C_c^{\infty}(\mathbb{R}_+).$$

According to the general theory of Boutet de Monvel's calculus, $r^-\Lambda_+^{-s}e^+$ extends to a symbol in $S^{-s}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_-)\right)$; by Theorem 2.11, we know that

$$r^+\mathrm{Op}_n^\psi(a)e^- \in S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{-1}; \mathscr{S}(\mathbb{R}_-), \mathscr{S}(\mathbb{R}_+)\right).$$

So, $r^+\mathrm{Op}_n^\psi(a)e^- r^-\Lambda_+^{-s}e^-$ is a symbol in $S^{m-s}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$. We can therefore write $r^+\mathrm{Op}_n^\psi(a)(e^+r^+ - 1)\Lambda_+^{-s}e^+$ as an operator-valued FIO defined on the boundary with phase function ψ_∂ and an amplitude belonging to $S^{m-s}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$. The continuity of operator-valued pseudodifferential operators on wedge Sobolev spaces implies that

$$
\begin{array}{ccc}
L^2(\mathbb{R}_+^n) & \lhook\joinrel\longrightarrow & \mathscr{W}^0(\mathbb{R}^{n-1}; L^2(\mathbb{R}_+)) \\
& & \Big\downarrow{\scriptstyle r^+\mathrm{Op}^\psi(e^+r^+-1)\Lambda^{-s}r^+} \\
H^{s-m}(\mathbb{R}_+^n) & \longleftarrow\joinrel\rhook & \mathscr{W}^{s-m}(\mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+)),
\end{array}
$$

where $\mathscr{W}^s(\mathbb{R}^{n-1}; E)$ denotes the wedge Sobolev space of order s with values in the topological vector space E, see the Appendix 7. $\qquad\square$

3. Fourier Integral Operators of Boutet de Monvel Type

We recall the definition of three symbol classes in the Boutet de Monvel calculus.

i) A *potential symbol* of order m is an element of

$$S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathbb{C}, \mathscr{S}(\mathbb{R}_+)) = \text{proj-lim}_\mathbf{s} S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathbb{C}, H^\mathbf{s}(\mathbb{R}_+)).$$

ii) A *trace symbol* of order m and type zero is an element of the set

$$S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathbb{C}) = \text{proj-lim}_\mathbf{s} S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; H_0^\mathbf{s}(\overline{\mathbb{R}}_+), \mathbb{C});$$

it also defines a symbol in $S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; H^{s_1, s_2}(\mathbb{R}_+), \mathbb{C})$, if $s_1 > -\frac{1}{2}$.

   A *trace symbol* of type d is a sum of the form

$$t = \sum_{j=0}^d t_j \partial_+^j, \quad t_j \in S^{m-j}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathbb{C}),$$

where t is in $S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; H^{s_1, s_2}(\mathbb{R}_+), \mathbb{C})$ and ∂_+ is the derivative in the normal direction,

$$\partial_+ = r^+\partial_{x_n}e^+ \in S^1(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; H^\mathbf{s}(\mathbb{R}_+), H^{\mathbf{s}-(1,0)}(\mathbb{R}_+)).$$

iii) A *singular Green symbol* of order m and type zero is an element of

$$S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}), \mathscr{S}(\mathbb{R}_+)) =$$
$$\text{proj-lim}_\mathbf{s} S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; H_0^{-\mathbf{s}}(\overline{\mathbb{R}}_+), H^\mathbf{s}(\mathbb{R}_+));$$

this also is a symbol in $S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; H^{s_1, s_2}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$, provided $s_1 > -\frac{1}{2}$.

   A *singular Green symbol* of order m and type d is a sum of the form

$$g = \sum_{j=0}^d g_j \partial_+^j, \quad g_j \in S^{m-j}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+)).$$

Obviously, g then is in $S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; H^{s_1, s_2}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$, $s_1 > d - \frac{1}{2}$.

Remark 3.1. *The trace operator γ_j is a trace symbol of order $j + \frac{1}{2}$ and type $j + 1$, see* [29].

In Definition 3.2, below, we employ the notation introduced in (16).

Definition 3.2. *Let χ be an admissible symplectomorphism, and let χ_∂ be the induced symplectomorphism at the boundary. For $m \in \mathbb{R}$ and $d \leq \max\{m, 0\}$ we define*

$$\mathcal{A} := \begin{pmatrix} r^+ A^\chi e^+ + G^{\chi_\partial} & K^{\chi_\partial} \\ T^{\chi_\partial} & S^{\chi_\partial} \end{pmatrix},$$

and we write $\mathcal{A} \in \mathscr{B}^{m,d}_\chi\left(\Omega^+_x \times \Omega^+_y\right)$, if $A^\chi \in I^m_{comp}(\Omega_x \times \Omega_y, \widetilde{\Lambda})$, the symbol of A^χ satisfies the transmission condition, and G^{χ_∂}, K^{χ_∂}, T^{χ_∂} and S^{χ_∂} are FIOs defined on the boundary, with Lagrangian submanifolds defined by χ_∂. Their respective symbols are

- (i) *a singular Green symbol of order m and type d;*
- (ii) *a potential symbol of order m;*
- (iii) *a trace symbol of order m and type d;*
- (iv) *a usual pseudodifferential symbol of order m.*

Remark 3.3. *In the above definition, we consider FIOs with operator-valued symbol. By Remark 1.2, χ_∂ is the lift to the cotangent bundle of a diffeomorphism, locally represented as $b : \Omega^\partial_y \to \Omega^\partial_x$. One can extend b to*

$$(43) \qquad \tilde{b} : \Omega^\partial_y \times [0,1) \to \Omega^\partial_x \times [0,1) : (y', y_n) \mapsto (b(y'), y_n).$$

Then K^{χ_∂}, T^{χ_∂}, G^{χ_∂} can be considered as the pullbacks of pseudodifferential operators with operator-valued symbol by the diffeomorphism $\tilde{b}$. For example, $G^{\chi_\partial}u = \tilde{b}^* G_1 u = G_2(\tilde{b}^* u)$, where G_1, G_2 are usual singular Green operators, and the equality holds modulo operators with smooth kernels. Similarly for K^{χ_∂} and T^{χ_∂}.

Here we think of functions with compact support as extended by zero to $\mathbb{R}^n$. If the symplectomorphism χ is the identity, and therefore $\Omega^+_x = \Omega^+_y$, the class $\mathscr{B}^{m,d}_\chi(\Omega^+_x \times \Omega^+_x)$ of FIOs coincides with the class $\mathscr{B}^{m,d}(\Omega^+_x)$ of operators of order m and class d in the Boutet de Monvel calculus, see [3, 15, 27, 29]. As a consequence of Theorem 2.18 and of the Sobolev continuity of FIOs defined through operator-valued symbols, we obtain the following theorem:

Theorem 3.4. *Every $\mathcal{A} \in \mathscr{B}^{m,d}_\chi(\Omega^+_x \times \Omega^+_y)$ induces a continuous operator*

$$\mathcal{A} : H^s_{\mathrm{comp}}(\Omega^+_y) \oplus H^s_{\mathrm{comp}}(\Omega^\partial_y) \to H^{s-m}_{\mathrm{loc}}(\Omega^+_x) \oplus H^{s-m}_{\mathrm{loc}}(\Omega^\partial_x),$$

provided $s > d - \frac{1}{2}$.

Now, we analyze the composition of FIOs of Boutet de Monvel type. Recall that we assume the involved symbols to have compact support w.r.t. the space variable.

Theorem 3.5. *Let $\mathcal{B} \in \mathscr{B}^{m_\mathcal{B}, d_\mathcal{B}}_{\chi'}(\Omega^+_x \times \Omega^+_y)$ and $\mathcal{A} \in \mathscr{B}^{m_\mathcal{A}, d_\mathcal{A}}_\chi(\Omega^+_y \times \Omega^+_z)$ be FIOs of Boutet de Monvel type associated with the symplectomorphisms χ, χ'. Then $\mathcal{B}\mathcal{A}$ is a FIO of Boutet de Monvel type of order $m = m_\mathcal{B} + m_\mathcal{A}$ and type $d = \max\{(m_\mathcal{A} + d_\mathcal{B}), d_\mathcal{A}\}$ defined by the symplectomorphism $\chi' \circ \chi$, that is $\mathcal{B}\mathcal{A} \in \mathscr{B}^{m,d}_{\chi' \circ \chi}(\Omega^+_x \times \Omega^+_z)$.*

Proof. By definition we can write

$$\mathcal{B}\mathcal{A} = \begin{pmatrix} r^+ B^{\chi'} e^+ + G^{\chi'_\partial}_\mathcal{B} & K^{\chi'_\partial}_\mathcal{B} \\ T^{\chi'_\partial}_\mathcal{B} & S^{\chi'_\partial}_\mathcal{B} \end{pmatrix} \circ \begin{pmatrix} r^+ A^\chi e^+ + G^{\chi_\partial}_\mathcal{A} & K^{\chi_\partial}_\mathcal{A} \\ T^{\chi_\partial}_\mathcal{A} & S^{\chi_\partial}_\mathcal{A} \end{pmatrix}.$$

We start with the composition of the elements in the upper left corner. We can write

$$r^+ B^{\chi'} e^+ r^+ A^\chi e^+ = r^+ B^{\chi'} A^\chi e^+ + r^+ B^{\chi'}(e^+ r^+ - 1) A^\chi e^+.$$

By the general theory of FIOs, $B^{\chi'} A^\chi$ is a FIO of order $m_\mathcal{B} + m_\mathcal{A}$ with canonical transformation $\chi' \circ \chi$. We prove next that the operator $r^+ B^{\chi'}(e^+ r^+ - 1) A^\chi e^+$ is a FIO

on the boundary, associated with the canonical transformation $\chi'_\partial \circ \chi_\partial$ and with a Green symbol of order $m_\mathcal{B} + m_\mathcal{A}$ and type $(m_\mathcal{A})_+ = \max\{m_\mathcal{A}, 0\}$. As both the symplectomorphisms are admissible, the boundary is preserved. B^χ' and A^χ are operator-valued FIO at the boundary with symbols b and a satisfying Assumptions 1 and Lagrangian submanifold induced by χ'_∂ and χ_∂, respectively. So, we consider the composition

$$r^+\mathrm{Op}^{\psi'}(b)(e^+r^+ - 1)\mathrm{Op}^\psi(a)e^+,$$

ψ', ψ being general phase functions associated with χ' and χ, respectively. We start by studying the composition in the normal direction. We decompose the symbol $a = a_d + a_0$ as in Proposition 2.14. First, we analyze the differential part

$$(44) \qquad r^+\mathrm{Op}_n^{\psi'}(b)\,(e^+r^+ - 1)\,\mathrm{Op}_n^\psi\left(\sum_{j=1}^{m_\mathcal{A}} a_j(x', x_n, \xi')\,\xi_n^j\right)e^+u,$$

where $a_j(x', x_n, \xi') \in S^{m_\mathcal{A}-j}(\mathbb{R}^n \times \mathbb{R}^{n-1})$. Since, on $\mathscr{S}(\mathbb{R}_+)$

$$(45) \qquad \xi_n\widehat{e^+u(y_n)}(\xi_n) = -i\widehat{e^+\partial_{y_n}u}(\xi_n) - iu(0)\widehat{\delta_0},$$

induction shows

$$(46) \qquad \xi_n^j\widehat{e^+u(y_n)} = (-i)^j\left(\widehat{e^+\partial_{x_n}^j u} + \sum_{l=0}^{j-1} u^{(l)}(0)\widehat{\delta_0^{(j-l-1)}}\right).$$

Now

$$\mathrm{Op}_n^\psi\left(\sum_{j=1}^{m_\mathcal{A}}(-i)^j a_j(x', x_n, \xi')\sum_{l=0}^{j-1} u^{(l)}(0)\widehat{\delta_0^{(j-l-1)}}\right)$$

$$(47) \qquad = \sum_{j=1}^{m_\mathcal{A}}(-i)^j a_j(x', x_n, \xi')\sum_{l=0}^{j-1} u^{(l)}(0)\mathrm{Op}_n^\psi\left(\xi^{j-l-1}\right)(\delta_0).$$

Following the scheme of the proof of Theorem 2.6, one obtains

$$(48) \qquad (e^+r^+ - 1)\mathrm{Op}_n^\psi\left(\xi_n^{j-l-1}\right)(\delta_0) = \sum_{k=0}^{j-l-1} d_k(x', \xi')\delta_0^{(k)} + c_{j,l}(x', x_n, \xi')$$

where $d_k \in S^{j-l-1-k}(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1})$ and $c_{j,l} \in S^{j-l-\frac{1}{2}}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathbb{C}, \mathscr{S}(\mathbb{R}_-)\right)$. Observing that $\gamma_l : u \mapsto u^{(l)}(0)$ is a trace symbol of order $l + \frac{1}{2}$ and type l, in view of (48) and Theorem 2.6, we obtain that $r^+\mathrm{Op}_n^{\psi'}(b)\,(e^+r^+ - 1)$ applied to (47) is a Green symbol of order $m_\mathcal{A}$ and type $(m_\mathcal{A})_+$. Next, we analyze

$$r^+\mathrm{Op}_n^{\psi'}(b)(e^+r^+ - 1)\sum_{j=1}^{m_\mathcal{A}}\mathrm{Op}_n^\psi\left(a_j(x', x_n, \xi')\widehat{e^+\partial_{y_n}^j u}(\xi_n)\right)$$

$$(49)$$

$$= r^+\mathrm{Op}_n^{\psi'}(b)(e^+r^+ - 1)\sum_{j=1}^{m_\mathcal{A}} a_j\mathrm{Op}_n^\psi(1)e^+\partial_+^j u.$$

Recall that $e^+r^+ - 1 = -e^-r^-$ on regular distributions. By Theorem 2.17 we find that $r^-a_j\mathrm{Op}_n^\psi(1)e^+$ is a symbol in $S^{m_\mathcal{A}-j}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_-)\right)$. Moreover, by Theorem 2.11 we know that $r^+\mathrm{Op}_n^{\psi'}(b)e^- \in S^{m_\mathcal{B}}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_-), \mathscr{S}(\mathbb{R}_+)\right)$. Hence, the symbol in (49) is a Green symbol of order $m_\mathcal{B} + m_\mathcal{A}$ and type $(m_\mathcal{A})_+$. In view of the decomposition $a = a_d + a_0$, we now have to consider

$$(50) \qquad r^+\mathrm{Op}_n^{\psi'}(b)(e^+r^+ - 1)\mathrm{Op}_n^\psi(a_0)e^+.$$

Theorem 2.17 implies that $r^- \mathrm{Op}_n^\psi(a_0)e^+ \in S^{m_{\mathcal{A}}}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_-)\right)$. Observing that $r^+ \mathrm{Op}_n^{\psi'}(b)e^-$ is an element of $S^{m_{\mathcal{B}}}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_-), \mathscr{S}(\mathbb{R}_+)\right)$, we see that (50) defines a Green symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$ and type zero.

The other compositions can be analyzed in a similar way. We omit most of the details. They are all FIOs at the boundary with a Lagrangian distribution defined by $\chi'_\partial \circ \chi_\partial$.

(1) $\left(r^+ B^{\chi'} e^+\right) G_{\mathcal{A}}^{\chi_\partial}$ has a singular Green symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$ and type $d_{\mathcal{A}}$.

(2) $G_{\mathcal{B}}^{\chi'_\partial}(r^+ A^\chi e^+)$ has a singular Green symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$ and of type $(m_{\mathcal{A}} + d_{\mathcal{B}})_+ = \max\{m_{\mathcal{A}} + d_{\mathcal{B}}, 0\}$.

(3) $G_{\mathcal{B}}^{\chi'_\partial} G_{\mathcal{A}}^\chi$ has a Green symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$ and type $d_{\mathcal{A}}$.

(4) $\left(r^+ B^{\chi'} e^+\right) K_{\mathcal{A}}^{\chi_\partial}$ has a potential symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$.

(5) $G_{\mathcal{B}}^{\chi'_\partial} K_{\mathcal{A}}^{\chi_\partial}$ has a potential symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$.

(6) $K_{\mathcal{B}}^{\chi'_\partial} T_{\mathcal{A}}^{\chi_\partial}$ has a Green symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$ and type $d_{\mathcal{A}}$.

(7) $K_{\mathcal{B}}^{\chi'_\partial} S_{\mathcal{A}}^{\chi_\partial}$ has a potential symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$.

(8) $T_{\mathcal{B}}^{\chi'_\partial} r^+ A^\chi e^+$ has a trace symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$ and type $(m_{\mathcal{A}} + d_{\mathcal{B}})_+$.

(9) $T_{\mathcal{B}}^{\chi'_\partial} G_{\mathcal{A}}^{\chi_\partial}$ has a trace symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$ and type $d_{\mathcal{A}}$.

(10) $S_{\mathcal{B}}^{\chi'_\partial} T_{\mathcal{A}}^{\chi_\partial}$ has a trace symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$ and type $d_{\mathcal{A}}$.

(11) $T_{\mathcal{B}}^{\chi'_\partial} K_{\mathcal{A}}^{\chi_\partial}$ has a symbol in $S^{m_{\mathcal{B}}+m_{\mathcal{A}}}(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1})$.

(12) $S_{\mathcal{B}}^{\chi'_\partial} S_{\mathcal{A}}^{\chi_\partial}$ has a symbol in $S^{m_{\mathcal{B}}+m_{\mathcal{A}}}(\mathbb{R}^{n-1} \times \mathbb{R}^{n-1})$.

The composition in 1) follows from Remark 3.3 and the composition laws of operator-valued pseudodifferential operators. The compositions in 3), 4), 5), 6), 7), 9), 10), 11), 12) can be treated similarly, exploiting the properties of operator-valued symbols. The compositions in 2) and 8) are slightly more delicate. Let us analyze the composition in 2). We first suppose $d_{\mathcal{B}} = 0$. By Proposition 2.14, a can be written in the form $a = a_d + a_0$. Using (45) and (46), we can write

$$\left(r^+ \mathrm{Op}_n^\psi(a_d)e^+\right)u =$$

$$\sum_{j=1}^{m_{\mathcal{A}}}(-i)^j r^+ a_j(x', x_n, \xi')\left(\mathrm{Op}_n^\psi(1)e^+\partial_{y_n}^j u + \sum_{l=0}^{j-1}\mathrm{Op}_n^\psi(1)\left(\delta_0^{(j-l-1)}\right)\gamma_l(u)\right).$$

By Theorem 2.6, Remark 2.15 and the properties of trace operators, the sum in j, l can be written as

$$r^+ \mathrm{Op}_n^\psi(a_d)e^+ = \sum_{j=1}^{m_{\mathcal{A}}} \tilde{a}_j(x', \xi')\partial_+^j + \sum_{l=0}^{j-1} k_l(x', \xi')\gamma_l(u),$$

$\tilde{a}_j \in S^{m_{\mathcal{A}}-j}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}, \mathscr{S}'(\mathbb{R}_+), \mathscr{S}'(\mathbb{R}_+)\right)$, $k_l \in S^{m_{\mathcal{A}}-l-\frac{1}{2}}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}, \mathbb{C}, \mathscr{S}(\mathbb{R}_+)\right)$. By the properties of trace operators and the definition of Green symbols of type zero, we then obtain that $g_{\mathcal{B}}\left(r^+ \mathrm{Op}_n^\psi(a_d)e^+\right)$ is a Green symbol of order $m_{\mathcal{B}} + m_{\mathcal{A}}$ and type $(m_{\mathcal{A}})_+$. To prove the same result for a_0, we notice that $r^+ \mathrm{Op}_n^\psi(a_0)e^+$ extends to a symbol in $S^{m_{\mathcal{A}}}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}'(\mathbb{R}_+)\right)$, and the assertion follows from the definition of Green symbols. If $d_{\mathcal{B}} \neq 0$, we see that

$$\partial_+ \int e^{i\psi(x', x_n, \xi', \xi_n) - i\psi_\partial(x', \xi')}a(x', x_n, \xi', \xi_n)\hat{u}(\xi_n)đ\xi_n$$

$$= r^+ \int e^{i\psi(x', x_n, \xi', \xi_n) - i\psi_\partial(x', \xi')}\tilde{a}(x', x_n, \xi', \xi_n)\hat{u}(\xi_n)đ\xi_n,$$

20

where $\tilde{a} = \partial_{x_n} a + (i \partial_{x_n} \psi) a$, which implies $\tilde{a} \in S^{m_\mathcal{A}+1}(\mathbb{R}^n \times \mathbb{R}^n)$. Iteratively we can reduce to the case $d_\mathcal{B} = 0$, raising the order from $m_\mathcal{A}$ to $m_\mathcal{A} + d_\mathcal{B}$. To handle the composition 8), we proceed similarly. $\qquad \square$

We will now show a Egorov type theorem for operators in $\mathscr{B}_\chi^{m,d}\left(\Omega_x^+ \times \Omega_y^+\right)$. In analogy with the usual calculus of FIOs on closed manifolds, we analyze the adjoint of the operator $\mathcal{A}$.

Theorem 3.6. *The formal adjoint $\mathcal{A}^*$ of an operator $\mathcal{A} \in \mathscr{B}_\chi^{m,0}(\Omega_x^+ \times \Omega_y^+)$, $m \leq 0$, is a FIO of Boutet de Monvel type, namely $\mathcal{A}^* \in \mathscr{B}_{\chi^{-1}}^{m,0}\left(\Omega_y^+ \times \Omega_x^+\right)$. Moreover,*

$$(51) \qquad \mathcal{A}^* = \begin{pmatrix} r^+ \left(A^\chi\right)^* e^+ + (G^{\chi_\partial})^* & (T^{\chi_\partial})^* \\ (K^{\chi_\partial})^* & (S^{\chi_\partial})^* \end{pmatrix},$$

where $(A^\chi)^*$ is the formal adjoint of A^χ, and the Lagrangian submanifold is defined by the graph of χ^{-1}. The operators $(G^{\chi_\partial})^*$, $(K^{\chi_\partial})^*$, $(T^{\chi_\partial})^*$, $(S^{\chi_\partial})^*$ appearing in (51) are the adjoints of G^{χ_∂}, K^{χ_∂}, T^{χ_∂}, S^{χ_∂}, respectively, that is, they are FIOs with Lagragian submanifold given by the graph of χ_∂^{-1}.

Proof. Since $m \leq 0$, A^χ is continuous from $L^2(\mathbb{R}^n)$ to itself. Moreover, $e^+ : L^2(\mathbb{R}_+^n) \to L^2(\mathbb{R}^n)$ is continuous and its adjoint is r^+. So we can write

$$(r^+ A^\chi e^+ u, v)_{L^2(\mathbb{R}_+^n)} = (A^\chi e^+ u, e^+ v)_{L^2(\mathbb{R}^n)}$$
$$= (e^+ u, (A^\chi)^* e^+ v)_{L^2(\mathbb{R}^n)} = (u, r^+ (A^\chi)^* e^+ v)_{L^2(\mathbb{R}_+^n)}.$$

For the other components of $\mathcal{A}^*$, we apply Remark 3.3 and recall that the adjoint of a Green operator of order m and type 0 is a Green operator of the same order and type, the adjoint of a potential operator of order m is a trace operator of order m and type 0 and the adjoint of a trace operator of order m and type 0 is a potential operator of order m, see [29]. $\qquad \square$

Definition 3.7. *For every $m \in \mathbb{Z}$ we define the operator*

$$[\Lambda_+^m] := \begin{pmatrix} r^+ \Lambda^m e^+ & 0 \\ 0 & \mathrm{Op}(\langle \xi' \rangle^m) \end{pmatrix}$$

where $r^+ \Lambda^m e^+ : H^m(\mathbb{R}_+^n) \to L^2(\mathbb{R}_+^n)$ is an isomorphism in Boutet de Monvel's calculus. The operator $[\Lambda_+^m]$ is an element of $\mathscr{B}^{m,0}(\overline{\mathbb{R}^n}_+)$ and it is invertible.

Theorem 3.8. *Let $\mathcal{A} \in \mathscr{B}_\chi^{m,d}(\Omega_x^+ \times \Omega_y^+)$, $d \leq m_+ = \max\{m, 0\}$, $m \in \mathbb{Z}$, be a FIO of Boutet de Monvel type. Then,*

  i) *If $m \leq 0$ and $d = 0$, $\mathcal{A}\mathcal{A}^*$ is a an element of $\mathscr{B}^{2m,0}(\Omega_x^+)$;*
  ii) *If $m > 0$ and $\omega \in C_c^\infty(\Omega_y)$, $(\mathcal{A}\omega[\Lambda_+^{-m}])(\mathcal{A}\omega[\Lambda_+^{-m}])^*$ is an element of $\mathscr{B}^{0,0}(\Omega_x^+)$.*

Proof. The proof of part *i)* essentially follows from Theorem 3.6 and from Egorov's theorem for FIOs on closed manifolds. The second part follows from the first, noticing that $\mathcal{A}\omega[\Lambda_+^{-m}]$ belongs to $\mathscr{B}_\chi^{0,0}(\Omega_x^+ \times \overline{\mathbb{R}^n}_+)$ by Theorem 3.5. $\qquad \square$

In general, one cannot expect a Egorov type theorem for FIOs of Boutet de Monvel type of all orders and types. In fact, the formal adjoint $\mathcal{P}^*$ of an operator $\mathcal{P} \in \mathscr{B}^{m,d}(\overline{\mathbb{R}^n}_+)$, with $m > 0$ or $d > 0$, in general is not even a Boutet de Monvel operator. However, by means of Theorems 3.5 and 3.8, it is possible to prove that

Theorem 3.9. *If $\mathcal{P} \in \mathscr{B}^{m',d'}(\Omega_y^+)$ and $\mathcal{A} \in \mathscr{B}_\chi^{m,0}(\Omega_x^+ \times \Omega_y^+)$, $m \leq 0$, then $\mathcal{A}\mathcal{P}\mathcal{A}^*$ belongs to $\mathscr{B}^{m',d'}(\Omega_x^+)$.*

4. Principal Symbols

We next define the boundary principal symbol $\sigma_\partial(\mathcal{A})$ of a FIO of Boutet de Monvel type

$$\mathcal{A} := \begin{pmatrix} r^+ A^\chi e^+ + G^{\chi_\partial} & K^{\chi_\partial} \\ T^{\chi_\partial} & S^{\chi_\partial} \end{pmatrix} \in \mathscr{B}_\chi^{m,d}(\Omega_x \times \Omega_y).$$

We have shown that $\mathcal{A}$ can be seen as an operator-valued FIO defined on the boundary with a symbol belonging to $S^m\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+) \oplus \mathbb{C}, \mathscr{S}(\mathbb{R}_+) \oplus \mathbb{C}\right)$. We will now switch to *classical* operator-valued symbols, see the Appendix, and show that in this setting we can define

$$(52) \qquad \sigma_\partial(\mathcal{A}) = \sigma_\partial \begin{pmatrix} r^+ A^\chi e^+ & 0 \\ 0 & 0 \end{pmatrix} + \sigma_\partial \begin{pmatrix} G^{\chi_\partial} & K^{\chi_\partial} \\ T^{\chi_\partial} & S^{\chi_\partial} \end{pmatrix}.$$

In view of the fact that $G^{\chi_\partial} = \tilde{b}^* G_1$, where $\tilde{b}$ is the diffeomorphism in (43) and G_1 is a usual Green operator, we already have a suitable principal symbol for G^{χ_∂}, namely the pullback of the principal symbol of G_1 under χ_∂. Similar arguments apply to $K^{\chi_\partial}, T^{\chi_\partial}, S^{\chi_\partial}$. Hence we can focus on finding a principal boundary symbol $r^+ A_\partial^\chi e^+$ for $r^+ A^\chi e^+$.

The natural candidate for the boundary principal symbol of $r^+ A^\chi e^+$ is

$$(53) \quad r^+ A_\partial^\chi e^+ (x', \eta') : u \mapsto r^+ \int e^{i x_n \partial_{x_n} \psi(x',0,\eta',\eta_n) - i y_n \eta_n} a_m(x', 0, \eta', \eta_n) e^+(u)\, dy_n \, đ\eta_n,$$

$$(x', \eta') \in \Omega_x^\partial \times T^*_{(b^{-1}(x'),0)} \Omega_y \setminus 0.$$

For simplicity this is written for $x' \in \Omega_x^\partial$ and $\eta' \in T^*_{(b^{-1}(x'),0)} \Omega_y \setminus 0$. To have a consistent definition of principal symbol, we have to express η' in terms of (x', ξ'). The symbol a_m is a section of the Maslov bundle. Since, in our case, this bundle is trivial near the boundary, it is defined as a section of the cotangent bundle restricted to the boundary. Notice that the operator-valued symbol in (53) is homogeneous of degree m in the sense of Definition 7.1 when a_m is the homogeneous principal part of $a \in S^m_{\mathrm{cl}}(\mathbb{R}^n \times \mathbb{R}^n) \cap S^m_{\mathrm{tr}}(\mathbb{R}^n \times \mathbb{R}^n)$.

Lemma 4.1. *The operator-valued function $r^+ A_\partial^\chi e^+$, defined in (53), belongs to the symbol class $S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$.*

Proof. One can divide the operator in two parts: one with compact support in x_n and one with symbol vanishing to infinite order at $x_n = 0$. For the part with compact support we repeat the arguments in the proofs of Theorems 2.6 and 2.11. Using essentially integration by parts, one proves that also the second operator belongs to $S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$. $\qquad\square$

In view of Proposition 1.6 and Assumptions 1 we shall write in the sequel $A_n^\chi = \mathrm{Op}_n^\psi(a)$.

Theorem 4.2. *Let a and ψ satisfy the Assumptions 1. Moreover, assume that a is a classical symbol. Then, $r^+ A_\partial^\chi e^+$, defined in (53), satisfies*

$$r^+ A_n^\chi e^+ - r^+ A_\partial^\chi e^+ \in S^{m-1}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+)\right).$$

Proof. Let $a_m^0(x', \eta) = a_m(x', 0, \eta)$ and $A_{n,0}^\chi = \mathrm{Op}_n^\psi(a_m^0)$. We claim that

$$(54) \qquad r^+ A_n^\chi e^+ - r^+ A_{n,0}^\chi e^+ \in S^{m-1}\left(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+)\right).$$

In the above equation, we consider the phase function ψ extended to $\mathbb{R}^n$ such that the extension coincides, outside a compact subset of $\mathbb{R}^n \times \mathbb{R}^n$, with the standard pseudodifferential phase $x \cdot \eta$, see [1] for details. We claim that

(55)
$$r^+ A_{n,0}^\chi e^+ \in S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+)).$$

To see this we split $r^+ A_{n,0}^\chi e^+$ into two parts. The first one has a symbol vanishing near $\{x_n = 0\}$ and a standard pseudodifferential phase, thus belongs to $S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$. The same is true for the second whose symbol is compactly supported w.r.t. the x_n-variable, so that we can apply Theorem 2.11. Next we note that $a - a_m \in S^{m-1}(\mathbb{R}^n \times \mathbb{R}^n)$ and write

$$a_m(x', x_n, \eta', \eta_n) - a_m(x', 0, \eta', \eta_n) = x_n b(x', x_n, \eta', \eta_n),$$

where $b \in S^m(\mathbb{R}^n \times \mathbb{R}^n)$. Since the multiplication by x_n is an element in $S^{-1}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$, it is enough to check that

$$r^+ A_{n,0}^\chi e^+ - r^+ A_\partial^\chi e^+ \in S^{m-1}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+)).$$

Explicitly, we have to consider the seminorms of

(56)
$$\kappa_{\langle\eta'\rangle^{-1}} r^+ \left(A_{n,0}^\chi - A_\partial^\chi \right) e^+ \kappa_{\langle\eta'\rangle}.$$

Evaluating the semigroup actions, we obtain

$$\kappa_{\langle\eta'\rangle^{-1}} r^+ \left(A_{n,0}^\chi - A_\partial^\chi \right) e^+ \kappa_{\langle\eta'\rangle} u(x) =$$
$$r^+ \int e^{i\psi\left(x', \frac{x_n}{\langle\eta'\rangle}, \eta', \eta_n\langle\eta'\rangle\right) - i\psi_\partial(x', \eta')} a_m(x', 0, \eta', \eta_n\langle\eta'\rangle) \widehat{e^+ u}(\eta_n) \, đ\eta_n -$$
$$r^+ \int e^{i\frac{x_n}{\langle\eta'\rangle} \partial_{x_n} \psi(x', 0, \eta', \eta_n\langle\eta'\rangle)} a_m(x', 0, \eta', \eta_n\langle\eta'\rangle) \widehat{e^+ u}(\eta_n) \, đ\eta_n$$
$$= \int e^{i\frac{x_n}{\langle\eta'\rangle} \partial_{x_n} \psi(x', 0, \eta', \eta_n\langle\eta'\rangle)} \left(e^{i\psi\left(x', \frac{x_n}{\langle\eta'\rangle}, \eta', \eta_n\langle\eta'\rangle\right) - i\psi_\partial(x', \eta') - i\frac{x_n}{\langle\eta'\rangle} \partial_{x_n} \psi(x', 0, \eta', \eta_n\langle\eta'\rangle)} - 1 \right)$$

(57)
$$\cdot a_m(x', 0, \eta', \eta_n\langle\eta'\rangle) \widehat{e^+ u}(\eta_n) \, đ\eta_n.$$

Arguing now as in the proof of Theorem 2.11, we can use an excision function and split (57) in

  i) a part which is smoothing w.r.t. the η_n-variable;
  ii) one which involves operators of order $m - j - 1$ acting on derivatives of order j of the Dirac distribution;
  iii) a part with an arbitrary decay rate w.r.t. the η_n-variable.

For term *ii)* we observe that the top order terms of the operators arising from $r^+ A_{n,0}^\chi e^+$ and $r^+ A_n^\chi e^+$, applied to $\delta_0^{(j)}$, both agree. This is a consequence of Remark 2.9, since they both involve the same power of the y_n-derivative of the phase function associated with χ^{-1}, and the expression is then evaluated at $y_n = 0$.

For the terms *i)* and *iii)* we notice that the Taylor expansion at $x_n = 0$ of

$$e^{i\psi\left(x', \frac{x_n}{\langle\eta'\rangle}, \eta', \eta_n\langle\eta'\rangle\right) - i\psi_\partial(x', \eta') - i\frac{x_n}{\langle\eta'\rangle} \partial_{x_n} \psi(x', 0, \eta', \eta_n\langle\eta'\rangle)} - 1$$

is simply given by

$$x_n^2 b(x', x_n, \eta', \eta_n\langle\eta'\rangle),$$

where b is a symbol order one. As the multiplication by x_n^2 is an operator-valued symbol of order -2, we gain one order of decay in η'. To check all the other seminorms, we take advantage of the fact that the symbol can be assumed to have arbitrarily high decay w.r.t. η_n. $\qquad\square$

The above formulation of the principal symbol is invariant since, with the notation in (7), (9), and the fact that χ_∂ is the lift of a diffeomorphism of the boundary,

$$\partial_{x_n}\psi(x',0,\eta',\eta_n) = (\xi_n)^*(y',0,\eta',\eta_n), \quad x' = x'^*(y',0,\eta',\eta_n) = x_\partial'^*(y',\eta') = x_\partial'^*(y').$$

So we can write

Definition 4.3. *Let A^χ be a FIO with phase function and symbol satisfying Assumptions 1, and let its principal symbol be denoted by $\sigma(A)$. Then, its boundary principal symbol is defined, on $T^*\Omega_x^\partial$, as*

(58)
$$r^+ A_\partial^\chi e^+(x',\xi') : u \mapsto$$
$$r^+ \int e^{ix_n(\xi_n)^*(x',0,(\xi_\partial'^*)^{-1}(\xi'),\eta_n)-i\eta_n y_n}\sigma(A)(x',0,(\xi_\partial'^*)^{-1}(\xi'),\eta_n)e^+(u)dy_n\,đ\eta_n.$$

This is the homogeneous principal symbol in the sense of Definition 7.1 in the Appendix.

Definition 4.4. *Let $\mathcal{A}$, χ, a, G^{χ_∂}, T^{χ_∂}, S^{χ_∂} be as in Definition 3.2. We define, on $T^*\Omega_x^\partial$,*

(59) $\sigma_\partial(\mathcal{A})(x',\xi') = \begin{pmatrix} r^+ A_\partial^\chi e^+ + \sigma_\partial(G^{\chi_\partial}) & \sigma_\partial(K^{\chi_\partial}) \\ \sigma_\partial(T^{\chi_\partial}) & \sigma_\partial(S^{\chi_\partial}) \end{pmatrix}$

$$\in S^m\left(\mathbb{R}^{n-1},\mathbb{R}^{n-1};\mathscr{S}(\mathbb{R}_+)\oplus\mathbb{C},\mathscr{S}(\mathbb{R}_+)\oplus\mathbb{C}\right),$$

where $\sigma_\partial(G^{\chi_\partial}),\sigma_\partial(K^{\chi_\partial}),\sigma_\partial(T^{\chi_\partial}),\sigma_\partial(S^{\chi_\partial})$ are the homogeneous principal symbols of the corresponding operators. The symbol $\sigma_\partial(\mathcal{A})$ is called the boundary principal symbol of $\mathcal{A}$.

Proposition 4.5. *If $\mathcal{A} \in \mathcal{B}_{\chi_1}^{m_1,d_1}(\Omega_x \times \Omega_z)$ and $\mathcal{B} \in \mathcal{B}_{\chi_2}^{m_2,d_2}(\Omega_z \times \Omega_y)$ then $\sigma_\partial(\mathcal{AB}) = \sigma_\partial(\mathcal{A})\,\sigma_\partial(\mathcal{B})$.*

Proof. Let us write

$$\mathcal{A} = \begin{pmatrix} r^+ A^{\chi_1} e^+ + G_1^{(\chi_1)_\partial} & K_1^{(\chi_1)_\partial} \\ T_1^{(\chi_1)_\partial} & S_1^{(\chi_1)_\partial} \end{pmatrix}$$

and

$$\mathcal{B} = \begin{pmatrix} r^+ B^{\chi_2} e^+ + G_2^{(\chi_2)_\partial} & K_2^{(\chi_2)_\partial} \\ T_2^{(\chi_2)_\partial} & S_2^{(\chi_2)_\partial} \end{pmatrix}.$$

It is clear that one has to consider only the FIO part; the Green, potential, trace and pseudodifferential terms in the lower right corner are simple compositions. The definition of boundary principal symbols implies that

$$r^+ A_\partial^{\chi_1} e^+ - r^+ A_n^{\chi_1} e^+ \in S^{m_1-1}(\mathbb{R}^{n-1},\mathbb{R}^{n-1};\mathscr{S}(\mathbb{R}_+),\mathscr{S}(\mathbb{R}_+))$$
$$r^+ B_\partial^{\chi_2} e^+ - r^+ B_n^{\chi_2} e^+ \in S^{m_2-1}(\mathbb{R}^{n-1},\mathbb{R}^{n-1};\mathscr{S}(\mathbb{R}_+),\mathscr{S}(\mathbb{R}_+)).$$

Therefore, we can write

$$r^+ A_n^{\chi_1} e^+ r^+ B_n^{\chi_2} e^+ = r^+ A_\partial^{\chi_1} e^+ r^+ B_\partial^{\chi_2} e^+ + r$$

where $r \in S^{m_1+m_2-1}(\mathbb{R}^{n-1},\mathbb{R}^{n-1};\mathscr{S}(\mathbb{R}_+),\mathscr{S}(\mathbb{R}_+))$. □

Theorem 3.6 describes precisely $\mathcal{A}^*$ for $\mathcal{A} \in \mathcal{B}_\chi^{0,0}(\Omega_x \times \Omega_y)$. One can then easily prove the next proposition.

Proposition 4.6. *For any $\mathcal{A} \in \mathcal{B}_\chi^{0,0}(\Omega_x \times \Omega_y)$ we have $\sigma_\partial(\mathcal{A}^*) = [\sigma_\partial(\mathcal{A})]^*$.*

Proposition 4.7. *Let $\mathcal{A}$ be an operator belonging to $\mathcal{B}_\chi^{0,0}(\Omega_x \times \Omega_y)$ whose boundary principal symbol $\sigma_\partial(\mathcal{A})$ is invertible on $T^*\Omega_x^\partial$. Then $[\sigma_\partial(\mathcal{A})]^{-1}$ is a boundary principal symbol of an operator belonging to $\mathcal{B}_{\chi^{-1}}^{0,0}(\Omega_y \times \Omega_x)$.*

Proof. Proposition 4.6 implies that $\sigma_\partial(\mathcal{A}^*) = [\sigma_\partial(\mathcal{A})]^*$. From Theorem 3.8 we know that $\mathcal{A}\mathcal{A}^* = \mathcal{P}$ is a usual Boutet de Monvel operator. Proposition 4.5 shows that $\sigma_\partial(\mathcal{A}\mathcal{A}^*) = \sigma_\partial(\mathcal{A})\,\sigma_\partial(\mathcal{A}^*)$ is invertible. Hence, by the general theory of Boutet de Monvel's calculus, $[\sigma_\partial(\mathcal{A}\mathcal{A}^*)]^{-1}$ is the boundary principal symbol of a Boutet de Monvel operator $\mathcal{P}'$. The inverse of $\sigma_\partial(\mathcal{A})$ can now be written as

$$[\sigma_\partial(\mathcal{A})]^{-1} = \sigma_\partial(\mathcal{A}^*)\,\sigma_\partial(\mathcal{P}') = \sigma_\partial(\mathcal{A}^*\mathcal{P}').$$

Hence $[\sigma_\partial(\mathcal{A})]^{-1}$ is the boundary principal symbol of $\mathcal{A}^*\mathcal{P}' \in \mathscr{B}_{\chi^{-1}}^{0,0}\big(\Omega_x \times \Omega_y\big)$, by Theorem 3.5. $\qquad\square$

Remark 4.8. *The proof of Proposition 4.7 implies that the FIO part in the upper left corner of $[\sigma_\partial(\mathcal{A})]^{-1}$ is $r^+ B_\partial^{\chi^{-1}} e^+$, where $B^{\chi^{-1}}$ is a parametrix of A^χ, hence a FIO associated with the symplectomorphism χ^{-1}.*

Definition 4.9. *An operator $\mathcal{A} \in \mathscr{B}_\chi^{m,d}(\Omega_x \times \Omega_y)$, $d \le m_+$, is called elliptic, if*

(i) *The interior symbol $\sigma(\mathcal{A})$ is elliptic. That is, there exists a FIO $B^{\chi^{-1}}$ of order $-m$ such that $A^\chi B^{\chi^{-1}} - \mathrm{Id}$ and $B^{\chi^{-1}} A^\chi - \mathrm{Id}$ are FIOs of order -1 on $\Omega_x \times \Omega_x$ and $\Omega_y \times \Omega_y$ respectively.*

(ii) *The boundary principal symbol*

$$\sigma_\partial(\mathcal{A})(x', \xi') : \mathscr{S}(\mathbb{R}_+) \oplus \mathbb{C} \to \mathscr{S}(\mathbb{R}_+) \oplus \mathbb{C}$$

is an invertible operator for any $(x', \xi') \in T^*\Omega_x \setminus 0$.

Theorem 4.10. *If $\mathcal{A} \in \mathscr{B}_\chi^{m,d}(\Omega_x \times \Omega_y)$, $d \le m_+$, is elliptic, then there exists an operator $\mathcal{B} \in \mathscr{B}_{\chi^{-1}}^{-m,d_-}(\Omega_y \times \Omega_x)$, $d_- = \max\{-m, 0\}$, such that $\mathcal{A}\mathcal{B}$ and $\mathcal{B}\mathcal{A}$ are both equal to the identity modulo lower order operators.*

Proof. Using an order reduction it is sufficient to consider the case $m = d = 0$. Since $\mathcal{A}$ is elliptic, there exists a parametrix $B^{\chi^{-1}}$ of A^χ such that $B^{\chi^{-1}} A^\chi$ and $A^\chi B^{\chi^{-1}}$ are equal to the identity modulo lower order operators. Let

$$[\sigma_\partial(\mathcal{A})]^{-1} = \begin{pmatrix} r^+ B_\partial^{\chi^{-1}} e^+ + g' & k' \\ t' & s' \end{pmatrix}.$$

We set

$$\mathcal{B} = \begin{pmatrix} r^+ B^{\chi^{-1}} e^+ + \mathrm{Op}^{\chi_\partial^{-1}}(g') & \mathrm{Op}^{\chi_\partial^{-1}}(k') \\ \mathrm{Op}^{\chi_\partial^{-1}}(t') & \mathrm{Op}^{\chi_\partial^{-1}}(s') \end{pmatrix}.$$

By construction, the composition $\mathcal{A}\mathcal{B}$ is equal, modulo lower order terms, to the identity outside a neighborhood of the boundary.

The definition of the boundary principal symbol and Proposition 4.7 imply that

$$r^+ A_n^\chi e^+ - r^+ A_\partial^\chi e^+ = r_1 \in S^{-1}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$$

(60) $$\qquad r^+ B_n^{\chi^{-1}} e^+ - r^+ B_\partial^{\chi^{-1}} e^+ = r_2 \in S^{-1}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+)).$$

Since the Maslov bundle is trivial for both χ and χ^{-1}, $\mathcal{A}$ and $\mathcal{B}$ can be considered as operator-valued FIOs at the boundary, with a Lagrangian submanifold defined by the graphs of χ_∂ and χ_∂^{-1} respectively. Then we can write

$$r^+ A_n^\chi e^+ r^+ B_n^{\chi^{-1}} e^+ = (r_1 + r^+ A_\partial^\chi e^+)(r_2 + r^+ B_\partial^{\chi^{-1}} e^+).$$

Since $r^+ A_\partial^\chi e^+$ and $r^+ B_\partial^{\chi^{-1}} e^+$ are elements of $S^0(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+))$ the relations (60) and the composition laws of Theorem 3.5 imply that $r^+ A_n^\chi e^+ \, r^+ B_n^{\chi^{-1}} e^+$ is equal to the identity up to a symbol of order -1. Hence $\mathcal{A}\mathcal{B} = \mathrm{Id}$ modulo lower order operators. The argument for $\mathcal{B}\mathcal{A}$ is analogous. $\qquad\square$

5. FIOs of Boutet de Monvel Type on Manifolds

Let X and Y be two compact manifolds with boundary, and χ be an admissible symplectomorphism $\chi\colon T^*Y \setminus 0 \to T^*X \setminus 0$. We recall that χ extends to a symplectomorphism $\tilde{\chi}\colon T^*\widetilde{Y} \setminus 0 \to T^*\widetilde{X} \setminus 0$ of open neighborhoods of Y and X, respectively. We denote by $\widetilde{\Lambda}$ the graph of $\tilde{\chi}$. As pointed out in Section 1, the restriction χ_∂ of χ is the lift of a diffeomorphism $b\colon \partial Y \to \partial X$. As in (43), we can extend b to a diffeomorphism of collar neighbourhoods $\tilde{b}\colon \partial Y \times [0,1) \to \partial X \times [0,1)$ by setting $\tilde{b}(y', y_n) = (b(y'), y_n)$. We will now define FIOs of Boutet de Monvel type

$$
(61) \qquad \mathcal{A} = \begin{pmatrix} r^+ A^X e^+ + G^{\chi_\partial} & K^{\chi_\partial} \\ T^{\chi_\partial} & S^\partial \end{pmatrix} : \begin{matrix} C^\infty(Y, E_1) \\ \oplus \\ C^\infty(\partial Y, F_1) \end{matrix} \to \begin{matrix} C^\infty(X, E_2) \\ \oplus \\ C^\infty(\partial X, F_2) \end{matrix},
$$

acting between vector bundles over X, Y and their boundaries.

Definition 5.1. *The linear operator $\mathcal{A}$ in (61) belongs to $\mathscr{B}_\chi^{m,d}(X \times Y)$, if*

(i) *A^X is a FIO defined by $\widetilde{\Lambda}$ and its local symbols satisfy the transmission condition;*
(ii) *in a collar neighborhood of the boundary, G^{χ_∂}, K^{χ_∂}, T^{χ_∂} and S^{χ_∂} are, modulo operators with smooth kernel, the pullbacks of standard singular Green, potential, trace, and pseudodifferential operators by $\tilde{b}$.*

The following theorem is obvious from Theorem 3.4.

Theorem 5.2. *Let $\mathcal{A} \in \mathscr{B}_\chi^{m,d}(X \times Y)$, $d \le m_+$, and $s \in \mathbb{R}$ with $s > d - 1/2$. Then $\mathcal{A}$ extends to a bounded operator*

$$
(62) \qquad \mathcal{A} : \begin{matrix} H^s(Y, E_1) \\ \oplus \\ H^s(\partial Y, F_1) \end{matrix} \to \begin{matrix} H^s(X, E_2) \\ \oplus \\ H^s(\partial X, F_2). \end{matrix}
$$

Definition 5.3. *Let $\mathcal{A} \in \mathscr{B}_\chi^{m,d}(X \times Y)$. We define*

(i) *the interior principal symbol $\sigma(\mathcal{A})$, as the principal symbol $\sigma(A^X)$ of A^X, restricted to the Lagrangian submanifold associated with the graph of χ;*
(ii) *the boundary principal symbol $\sigma_\partial(\mathcal{A})$*

$$
\sigma_\partial(\mathcal{A})(x', \xi') = \begin{pmatrix} r^+ A_\partial^X(x', \xi') e^+ + \sigma_\partial(G^{\chi_\partial})(x', \xi') & \sigma_\partial(K^{\chi_\partial})(x', \xi') \\ \sigma_\partial(T^{\chi_\partial})(x', \xi') & \sigma_\partial(S^{\chi_\partial})(x', \xi') \end{pmatrix},
$$

for $(x', \xi') \in T^*\partial X \setminus 0$. Here $\sigma_\partial(G^{\chi_\partial})$, $\sigma_\partial(K^{\chi_\partial})$, $\sigma_\partial(T^{\chi_\partial})$, $\sigma_\partial(S^{\chi_\partial})$ are the principal symbols of the corresponding operators. We consider $\sigma_\partial(\mathcal{A})(x', \xi')$ as an operator

$$
\sigma_\partial(\mathcal{A})(x', \xi') : \begin{matrix} \mathscr{S}(\mathbb{R}_+, E_1^*)_{(y', \eta')} \\ \oplus \\ F_{1,(y', \eta')}^* \end{matrix} \to \begin{matrix} \mathscr{S}(\mathbb{R}_+, E_2^*)_{(x', \xi')} \\ \oplus \\ F_{2,(x', \xi')}^* \end{matrix}
$$

Here, $\chi_\partial(y', \eta') = (x', \xi')$, E_1^* is the pullback of $E_1|_{\partial Y}$ to $T^*\partial Y$ with $\mathbb{R}_+$ identified with the normal to ∂Y at y', F_1^* is the pullback of F_1 to $T^*\partial Y$, and similarly for E_2 and F_2.

The boundary principal symbol defined above has an invariant meaning. Indeed, the principal symbol of A^X is invariantly defined, and since the Maslov bundle can be trivialized in a neighborhood of the boundary, the operator-valued symbol (53) is defined on $T^*\partial X \setminus 0$ as we have noted after (53). The other symbols involved also have an invariant meaning since the corresponding operators are compositions of the diffeomorphism $\tilde{b}$ and usual operator-valued pseudodifferential operators.

In view of Theorem 3.5, Proposition 4.5 and the well known properties of composition of FIOs associated with a canonical transformation we obtain the following theorem.

Theorem 5.4. *Let $\mathcal{B} \in \mathscr{B}_{\chi'}^{m_\mathcal{B}, d_\mathcal{B}}(X \times Y)$ and $\mathcal{A} \in \mathscr{B}_{\chi}^{m_\mathcal{A}, d_\mathcal{A}}(Y \times Z)$ then $\mathcal{B}\mathcal{A}$ belongs to $\mathscr{B}_{\chi' \circ \chi}^{m_\mathcal{B} + m_\mathcal{A}, d}(X \times Z)$, where $d = \max\{(m_\mathcal{A} + d_\mathcal{B}), d_\mathcal{A}\}$. Moreover,*

$$\sigma_\partial(\mathcal{B}\mathcal{A}) = \sigma_\partial(\mathcal{B})\sigma_\partial(\mathcal{B}).$$

Remark 5.5. *In view of Definition 5.3, the principal interior symbol of the composition $\mathcal{B}\mathcal{A}$ is given by the usual formula, see Hörmander [20], Theorem 25.2.3 and Section 25.3. The formula for the boundary symbol of the composition appears simpler. Note, however, that the symplectomorphism χ_∂ enters also in Definition 4.3.*

Theorem 5.6. *Let $\mathcal{A} \in \mathscr{B}_{\chi}^{m,0}(X \times Y)$ with $m \le 0$. Then, $\mathcal{A}^*$ belongs to $\mathscr{B}_{\chi^{-1}}^{m,0}(Y \times X)$, and*

$$\sigma_\partial(\mathcal{A}^*) = \sigma_\partial(\mathcal{A})^*.$$

Definition 5.7. *An operator $\mathcal{A} \in \mathscr{B}_{\chi}^{m,d}(X \times Y)$, $d \le m_+$, is elliptic if $\sigma(\mathcal{A})$ and $\sigma_\partial(\mathcal{A})$ are invertible.*

Theorem 5.8. *Let $\mathcal{A} \in \mathscr{B}_{\chi}^{m,d}(X \times Y)$, $d \le m_+$, be elliptic. Then there exists a parametrix $\mathcal{B} \in \mathscr{B}_{\chi^{-1}}^{-m,d_-}(Y \times X)$, $d_- = \max\{-m, 0\}$, such that $\mathcal{A}\mathcal{B}$ and $\mathcal{B}\mathcal{A}$ are equal to the identity modulo operators with smooth integral kernels. In particular, $\mathcal{A}$ in (62) is a Fredholm operator.*

Proof. Using an order reduction we may assume $m = d = 0$. By the composition properties of Theorem 5.4, the operator $\mathcal{P} = \mathcal{A}\mathcal{A}^*$ is an elliptic Boutet de Monvel operator, hence it admits a parametrix $\mathcal{C}$ such that

$$\mathcal{A}\mathcal{A}^*\mathcal{C} = \mathcal{P}\mathcal{C} = \mathrm{Id} + \mathcal{K}, \quad \mathcal{K} \text{ a smoothing operator.}$$

Therefore, $\mathcal{A}^*\mathcal{C}$ is a right parametrix of $\mathcal{A}$. In the same way we obtain a left parametrix. □

From Theorem 5.4 we also obtain the following Egorov type theorem.

Theorem 5.9. *If $\mathcal{P} \in \mathscr{B}^{m', d'}(Y)$ and $\mathcal{A} \in \mathscr{B}_{\chi}^{m,0}(X \times Y)$, $m \le 0$, then $\mathcal{A}\mathcal{P}\mathcal{A}^*$ belongs to $\mathscr{B}^{m', d'}(X)$.*

6. Index of Admissible Symplectomorphisms

Following the idea of Weinstein [32], we will now associate a FIO of Boutet de Monvel type with an admissible symplectomorphism $\chi : T^*Y \setminus 0 \to T^*X \setminus 0$ and a unitary section of the Maslov bundle u. Choose on Y and X the trivial line bundle and on ∂Y and ∂X the zero bundle. We consider the FIO of Boutet de Monvel type defined as

(63)
$$\mathcal{U}^\chi = \begin{pmatrix} r^+ U^\chi e^+ \end{pmatrix} : C^\infty(Y) \to C^\infty(X),$$

where U^χ has principal symbol u.

Theorem 6.1. *The operator $\mathcal{U}$, defined in (63), is elliptic, that is the interior and the boundary symbol are both invertible.*

Proof. The interior symbol is invertible by construction. So we have to show the invertibility of the boundary symbol

$$\sigma_\partial(\mathcal{U}^\chi)(x', \xi') : \mathscr{S}(\mathbb{R}_+) \to \mathscr{S}(\mathbb{R}_+)$$

(64)
$$u \mapsto r^+ \int e^{ix_n \partial_{x_n} \psi(x', 0, (\xi'^*)^{-1}(\xi'), \eta_n)} \widehat{e^+ u}(\eta_n) \, \djmath\eta_n$$

for each $(x', \xi') \in T^*\partial X \setminus 0$. As we have noticed in Equation (21), $\sigma_\partial(\mathcal{U}^X)(x', \xi') \in \mathscr{L}\left(L^2(\mathbb{R}_+)\right)$. In the sequel it will be easier to consider the operator

$$\sigma_\partial(\mathcal{U}^X)(x', \eta'): L^2(\mathbb{R}_+) \to L^2(\mathbb{R}_+)$$

(65)
$$u \mapsto r^+ \int e^{ix_n \partial_{x_n} \psi(x', 0, \eta', \eta_n)} \widehat{e^+ u}(\eta_n) \, d\eta_n.$$

In the following, we will first show that $\sigma_\partial(\mathcal{U}^X)(x', \eta')$ is invertible on a weighted L^2-space. From this we will infer the invertibility on $L^2(\mathbb{R}_+)$ and on $\mathscr{S}(R_+)$.

First notice that for each $\lambda \in \mathbb{R}$

$$\kappa_\lambda: L^2(\mathbb{R}_+) \to L^2(\mathbb{R}_+): u(x) \mapsto \lambda^{\frac{1}{2}} u(\lambda x)$$

is a unitary. Hence, the invertibility of (65) is equivalent to that of

$$\kappa_{\lambda^{-1}} \circ \sigma_\partial(\mathcal{U}^X)(x', \eta') \circ \kappa_\lambda u(x_n) = r^+ \int e^{ix_n \partial_{x_n} \psi(x', 0, \lambda^{-1}\eta', \eta_n)} \widehat{e^+ u}(\eta_n) \, d\eta_n.$$

That is, we can always suppose that $|\eta'|$ is as small as necessary. We next consider the larger space $L^2_{w\pm} := L^2(\mathbb{R}_\pm, w)$, $w(x) = (1 + |x|)^{-2}$, which contains the constants and all bounded continuous functions. We set $L^2_\omega = L^2(\mathbb{R}, \omega)$. For $0 \le t \le 1$, let

$$A(t) = A_n^\psi(t)(x', \eta') : u \mapsto \int e^{ix_n \partial_{x_n} \psi(x', 0, t\eta', \eta_n)} \widehat{u}(\eta_n) \, d\eta_n$$

$$= \kappa_t \int e^{ix_n \partial_{x_n} \psi(x', 0, \eta', \eta_n)} \widehat{\kappa_{t^{-1}} u}(\eta_n) \, d\eta_n.$$

In view of the fact that $\partial_{x_n} \psi(x', 0, 0, \eta_n)$ is homogeneous of degree 1 in η_n, we have $\partial_{x_n} \psi(x', 0, 0, \eta_n) = c(x')\eta_n$ for a smooth function c. Moreover, $c(x') > 0$, since $\partial_{\eta_n} \partial_{x_n} \psi(x', 0, 0, \eta_n) > 0$, as noticed after (9). We have

$$r^+ A(0)e^+ u(x_n) = r^+ A_n^\psi(0)e^+ u(x_n) = r^+ \int e^{ix_n \partial_{x_n} \psi(x', 0, 0, \eta_n)} \widehat{e^+ u}(\eta_n) \, d\eta_n$$

$$= r^+ \int e^{ix_n c(x')\eta_n} \widehat{e^+ u}(\eta_n) \, d\eta_n = u(c(x')x_n).$$

Being a dilation in the variable by a positive factor, $r^+ A(0)e^+$ is clearly invertible.

Let us now consider on L^2_{w+} the operator family $P(t)$, $t \in [0, 1]$, given by

(66)
$$P(t) = r^+ A(t)e^+ (r^+ A(t)e^+)^* = r^+ A(t)e^+ r^+ A(t)^* e^+.$$

Since $P(0)$ is invertible on L^2_{w+}, invertibility of $P(t)$ for small t will follow, if we prove that $P(t)$ is continuous at $t = 0$. In this case, $r^+ A(t)e^+$ will be surjective on L^2_{w+}. A similar argument for $(r^+ A^*(t)e^+)(r^+ A(t)e^+)$ will imply injectivity for small t. Hence, for t close to zero, $r^+ A(t)e^+$ will be invertible on L^2_{w+}.

So our task is to prove the continuity of $P(t)$ in $t = 0$. Notice that

(67)
$$P(t) = r^+ A(t)e^+ r^+ A(t)^* e^+ = r^+ A(t)A(t)^* e^+ + r^+ A(t)e^- r^- A^*(t)r^+.$$

The operator $A(t)A^*(t)$ can be written explicitly

$$(A(t)A^*(t))u(x_n) = \int e^{i(x_n - y_n)\partial_{x_n} \psi(x', 0, t\eta', \eta_n)} u(y_n) \, dy_n \, d\eta_n.$$

As observed after (9), $\partial_{x_n} \partial_{\eta_n} \psi(x', 0, t\eta', \eta_n) \ge \delta > 0$, hence it is possible to invert $\eta_n \mapsto \partial_{x_n} \psi(x', 0, t\eta', \eta_n)$. Letting $\xi_n = \partial_{x_n} \psi(x', 0, t\eta', \eta_n)$, $\eta_n = \eta_n(x', t\eta', \xi_n)$, we obtain the pseudodifferential operator

(68) $$(A(t)A^*(t))u(x_n) = \int e^{i(x_n - y_n)\xi_n} \partial_{x_n} \partial_{\eta_n} \psi(x', 0, t\eta', \eta_n(x', t\eta', \xi_n))^{-1} u(y_n) \, dy_n \, d\xi_n.$$

28

Notice that $\partial_{\eta_n}\partial_{x_n}\psi(x',0,0,\eta_n(x',0,\xi_n))^{-1} = c(x')^{-1}$. Therefore

$$A(0)A^*(0)u(x_n) = \int e^{i(x_n-y_n)\xi_n}c(x')^{-1}u(y_n)dy_n đ\xi_n = c(x')^{-1}u(x_n),$$

so that $A(0)A^*(0)$ is invertible. For each fixed t, the operator in (68) is an SG-pseudodifferential operator of order $(0,0)$, hence continuous on L_w^2.

We next study the limit as $t \to 0^+$. In view of the fact that $\partial_{x_n}\partial_{\eta_n}\psi$ is in general not continuously extendable to $\eta = 0$, we choose a zero excision function ζ with $\zeta(\xi_n) = 0$ for $|\xi_n| < 1$ and $\zeta(\xi_n) = 1$ for $|\xi| > 2$ and consider the two operators

$$(P_1(t)u)(x_n) = \int e^{i(x_n-y_n)\xi_n}\partial_{x_n}\partial_{\eta_n}\psi(x',0,t\eta',\eta_n(x',t\eta',\xi_n))^{-1}\zeta(\xi_n)u(y_n)dy_n đ\xi_n,$$

$$(P_2(t)u)(x_n) = \int e^{i(x_n-y_n)\xi_n}\partial_{x_n}\partial_{\eta_n}\psi(x',0,t\eta',\eta_n(x',t\eta',\xi_n))^{-1}(1-\zeta(\xi_n))u(y_n)dy_n đ\xi_n.$$

Then $P_1(t)$ is an SG-pseudodifferential operator of order $(0,0)$ whose symbol has uniformly bounded seminorms w.r.t. to t. Furthermore,

$$\lim_{t\to 0+}\zeta(\xi_n)\left(\partial_{x_n}\partial_{\eta_n}\psi(x',0,t\eta',\eta_n(x',t\eta',\xi_n))^{-1} - c(x')^{-1}\right) = 0$$

uniformly in the topology of SG-symbols on $\mathbb{R}$ and therefore

$$\lim_{t\to 0}P_1(t) = P_1(0) : u \mapsto \int e^{i(x_n-y_n)\xi_n}c(x')^{-1}\zeta(\xi_n)u(y_n)dy_n đ\xi_n$$

in the norm topology of $\mathscr{L}\left(L_w^2\right)$.

The kernel of the operator $P_2(t)$ is

$$K_{P_2(t)}(x_n,y_n) = \int e^{i(x_n-y_n)\xi_n}\partial_{x_n}\partial_{\eta_n}\psi(x',0,t\eta',\eta_n(x',t\eta',\xi_n))^{-1}(1-\zeta(\xi_n))đ\xi_n.$$

The function $\partial_{x_n}\partial_{\eta_n}\psi(x',0,t\eta',\eta_n(x',t\eta',\xi_n))^{-1}$ is positively homogeneous of degree zero in (η',η_n), hence bounded. As $t \mapsto 0^+$, Lebesgue's theorem of dominated convergence implies that

$$(1-\zeta(\xi_n))\left(\partial_{x_n}\partial_{\eta_n}\psi(x',0,t\eta',\eta_n(x',t\eta',\xi_n)^{-1} - c(x'))^{-1}\right) \to 0 \quad \text{in } L^1(\mathbb{R}_{\xi_n}),$$

and hence

$$\lim_{t\to 0}K_{P_2(t)}(x_n,y_n) = \int e^{i(x_n-y_n)\xi_n}c(x')^{-1}(1-\zeta(\xi_n))đ\xi_n = K_{P_2(0)}(x_n,y_n),$$

uniformly on $\mathbb{R}_+ \times \mathbb{R}_+$.

Schur's Lemma then implies that $P_2(t)$ tends to $P_2(0)$ in the norm of $\mathscr{L}(L_{w+}^2)$, since

$$\sup_{x_n\in\mathbb{R}_+}\int_{\mathbb{R}_+}\left|K_{P_2(t)}(x_n,y_n) - K_{P_2(0)}(x_n,y_n)\right|(1+|y_n|)^{-2}dy_n$$

$$\leq \sup_{\mathbb{R}_+\times\mathbb{R}_+}\left|K_{P_2(t)}(x_n,y_n) - K_{P_2(0)}(x_n,y_n)\right|\int_{\mathbb{R}_+}(1+|y_n|)^{-2}dy_n,$$

$$\sup_{y_n\in\mathbb{R}_+}\int_{\mathbb{R}_+}\left|K_{P_2(t)}(x_n,y_n) - K_{P_2(0)}(x_n,y_n)\right|(1+|x_n|)^{-2}dx_n$$

$$\leq \sup_{\mathbb{R}_+\times\mathbb{R}_+}\left|K_{P_2(t)}(x_n,y_n) - K_{P_2(0)}(x_n,y_n)\right|\int_{\mathbb{R}_+}(1+|x_n|)^{-2}dx_n.$$

This implies the continuity of $t \mapsto A(t)A^*(t)$ at $t = 0$ in $\mathscr{L}(L_\omega^2)$. In view of the fact that extension by zero and restriction define bounded maps, we also see that $t \mapsto r^+A(t)A^*(t)e^+$ is continuous at $t = 0$ in L_{w+}^2. Hence, we have proven the continuity of the first operator in the right hand side of (67).

Next consider $G(t) = r^+ A(t)e^- r^- A^*(t)r^+$. In view of the fact that $r^+ A(0)e^- = 0$, we want to prove that $G(t) \to 0$ in L^2_{w+}. Notice that

$$\left\| r^+ A(t)e^- r^- A^*(t)e^+ - r^+ A(0)e^- r^- A^*(0)e^+ \right\|_{\mathscr{L}(L^2_{w+})}$$

$$= \left\| r^+ A(t)e^- \left(r^- A^*(t)e^+ - r^- A^*(0)e^+\right) + \left(r^+ A(t)e^- - r^+ A(0)e^-\right) r^- A^*(0)e^+ \right\|_{\mathscr{L}(L^2_{w+})}$$

$$\leq \max\left\{ \sup_t \left\| r^+ A(t)e^- \right\|_{\mathscr{L}(L^2_{w-}, L^2_{w+})}, \left\| r^- A^*(0)e^+ \right\|_{\mathscr{L}(L^2_{w+}, L^2_{w-})} \right\}$$

$$(69) \qquad \cdot \left(\left\| r^- (A^*(t) - A^*(0))e^+ \right\|_{\mathscr{L}(L^2_{w+}, L^2_{w-})} + \left\| r^+ (A(t) - A(0))e^- \right\|_{\mathscr{L}(L^2_{w-}, L^2_{w+})} \right).$$

We start by considering $\|r^+ (A(t) - A(0))e^-\|_{\mathscr{L}(L^2_{w-}, L^2_{w+})}$, that is, the operator

$$u \mapsto r^+ \iint_{\mathbb{R}_+} \left(e^{i x_n \partial_{x_n} \psi(x', 0, t\eta', \eta_n) - i y_n \eta_n} - e^{i x_n c(x')\eta_n - i y_n \eta_n} \right) e^- u(y_n) dy_n đ\eta_n.$$

The associated kernel is

$$k(t; x_n, y_n) = \int \left(e^{i x_n \partial_{x_n} \psi(x', 0, t\eta', \eta_n) - i y_n \eta_n} - e^{i x_n c(x')\eta_n - i y_n \eta_n} \right) đ\eta_n, \quad x_n > 0, y_n < 0.$$

For convenience, we invert the sign of the y_n-variable and consider

$$\tilde{k}(t; x_n, y_n) = \int \left(e^{i x_n \partial_{x_n} \psi(x', 0, t\eta', \eta_n) + i y_n \eta_n} - e^{i x_n c(x')\eta_n + i y_n \eta_n} \right) đ\eta_n$$

$$= \int e^{i x_n c(x')\eta_n + i y_n \eta_n} \left(e^{i x_n \left(\partial_{x_n} \psi(x', 0, t\eta', \eta_n) - c(x')\eta_n\right)} - 1 \right) đ\eta_n, \quad x_n > 0, y_n > 0.$$

Choose a zero excision function ζ as above and let

$$(70) \qquad k_1(t; x_n, y_n) = \int e^{i x_n c(x')\eta_n + i y_n \eta_n} \left(e^{i x_n \left(\partial_{x_n} \psi(x', 0, t\eta', \eta_n) - c(x')\eta_n\right)} - 1 \right) \zeta(\eta_n) đ\eta_n$$

$$(71) \qquad k_2(t; x_n, y_n) = \int e^{i x_n c(x')\eta_n + i y_n \eta_n} \left(e^{i x_n \left(\partial_{x_n} \psi(x', 0, t\eta', \eta_n) - c(x')\eta_n\right)} - 1 \right) (1 - \zeta(\eta_n)) đ\eta_n.$$

First we analyze k_1. For an application of Schur's Lemma, on L^2_{w+} it will be enough to check that

$$(72) \qquad \sup_{x_n, y_n} |k_1(t; x_n, y_n)| \to 0 \quad \text{as } t \to 0^+.$$

We recall that the phase function $\partial_{x_n} \psi(x', 0, t\eta', \eta_n)$ satisfies the transmission condition, so that

$$\partial_{x_n} \psi(x', 0, t\eta', \eta_n) - c(x')\eta_n = \sum_{|\alpha|=1} \partial_{x_n} \partial_{\eta'}^\alpha \psi(x', 0, 0, \eta_n)(t\eta')^\alpha + r_1(x', t\eta', \eta_n)$$

$$(73) \qquad = \sum_{|\alpha|=1} \partial_{x_n} \partial_{\eta'}^\alpha \psi(x', 0, 0, 1)(t\eta')^\alpha + \sum_{|\alpha|=2} (t\eta')^\alpha r_{1,\alpha}(x', \sigma(t)\eta', \eta_n), \quad \sigma(t) \in (0, t).$$

Using integration by parts, we can write

$$|k_1(t; x_n, y_n)| = \left| \int e^{ix_n c(x')\eta_n + iy_n\eta_n} \left(e^{ix_n\left(\partial_{x_n}\psi(x',0,t\eta',\eta_n) - c(x')\eta_n\right)} - 1 \right) \zeta(\eta_n)\,d\eta_n \right|$$

$$\leq \left| \int e^{ix_n c(x')\eta_n + iy_n\eta_n} \frac{x_n \partial_{\eta_n} r_1(x', t\eta', \eta_n)}{(x_n c(x') + y_n)} e^{ix_n\left(\partial_{x_n}\psi(x',0,t\eta',\eta_n) - c(x')\eta_n\right)} \zeta(\eta_n)\,d\eta_n \right|$$

$$+ \left| \int e^{ix_n c(x')\eta_n + iy_n\eta_n} \frac{1}{(x_n c(x') + y_n)} \left(e^{ix_n\left(\partial_{x_n}\psi(x',0,t\eta',\eta_n) - c(x')\eta_n\right)} - 1 \right) \partial_{\eta_n}\zeta(\eta_n)\,d\eta_n \right|$$

$$\leq \frac{x_n}{x_n c(x') + y_n} \int \left| \partial_{\eta_n} r_1(x', t\eta', \eta_n) \right| \zeta(\eta_n)\,d\eta_n$$

$$(74) \qquad + \frac{1}{x_n c(x') + y_n} \int \left| \left(e^{ix_n\left(\partial_{x_n}\psi(x',0,t\eta',\eta_n) - c(x')\eta_n\right)} - 1 \right) \right| |\partial_{\eta_n}\zeta(\eta_n)|\,d\eta_n.$$

Notice that $r_1(x', t\eta', \eta_n)$ is in general not continuous in $(x', 0, 0)$. Nevertheless, as a function in η_n, it satisfies estimates as a symbol in $S^{-1}(\mathbb{R} \times \mathbb{R})$ for $|\eta_n| > 1$. Hence, $\zeta(\eta_n)\partial_{\eta_n} r_1(x', t\eta, \eta_n)$ is integrable and uniformly bounded for $t \in (0, 1]$. Moreover,

$$\lim_{t \to 0} \zeta(\eta_n)(t\eta')^\alpha \partial_{\eta_n} r_{1,\alpha}(x', \sigma(t)\eta, \eta_n) = 0, \quad |\alpha| = 2.$$

Notice also that

$$\frac{x_n}{c(x')x_n + y_n} \leq c(x')^{-1},$$

is uniformly bounded for $x_n > 0$, $y_n > 0$. In (74), $\partial_{\eta_n}\zeta(\eta_n)$ is a function with compact support and, for $\eta \neq 0$,

$$\left| e^{ix_n\left(\partial_{x_n}\psi(x',0,t\eta',\eta_n) - c(x')\eta_n\right)} - 1 \right| \leq \left| x_n \left(\partial_{x_n}\psi(x', 0, t\eta', \eta_n) - c(x')\eta_n \right) \right|$$

$$\leq x_n t |\eta'| \left| \sum_{|\alpha'|=1} \int_0^1 \partial_{x_n}\partial_\xi^{\alpha'} \psi(x', 0, st\eta', \eta_n)\,ds \right| \leq M x_n t |\eta'|,$$

since $\partial_{x_n}\partial_\xi^{\alpha'} \psi(x', 0, st\eta', \eta_n)$ is uniformly bounded for $|\alpha'| = 1$, $\eta_n \neq 0$, $(x', \eta') \in T^*\partial X \setminus 0$, $t \in (0, 1]$. Lebesgue's theorem on dominated convergence therefore implies that (72) holds and thus the operator associated with k_1 tends to zero as $t \to 0^+$ by Schur's lemma. A similar, but simpler, argument holds for $k_2(t; x_n, y_n)$. So we have proven that $r^+ A(t)e^- \to 0$ in the norm topology.

A similar argument shows $r^- A^*(t)e^+ \to 0$. Hence, it is possible to find a constant $C > 0$ such that for small $\bar{t} > 0$

$$\left\| r^+ A(t)e^- \right\|_{\mathscr{L}\left(L_{\omega-}^2, L_{\omega+}^2\right)} \leq C, \quad \left\| r^- A^*(t)e^+ \right\|_{\mathscr{L}\left(L_{\omega+}^2, L_{\omega-}^2\right)} \leq C, \quad t \in [0, \bar{t}].$$

In view of the above argument and Equation (69), we find

$$\lim_{t \to 0} \|G(t)\|_{\mathscr{L}\left(L_{\omega+}^2\right)} \leq C \left(\lim_{t \to 0} \left\| r^- A(t)e^+ \right\|_{\mathscr{L}\left(L_{\omega+}^2, L_{\omega-}^2\right)} + \left\| r^+ A(t)e^- \right\|_{\mathscr{L}\left(L_{\omega-}^2, L_{\omega+}^2\right)} \right)$$

$$= 0.$$

Therefore, we have proven that

$$\lim_{t \to 0} r^+ A(t)e^+ \left(r^+ A(t)e^+ \right)^* = r^+ A(0)e^+ r^+ A^*(0)e^+ = c(x')^{-1}\mathrm{Id}$$

in the norm topology of $\mathscr{L}(L_{w+}^2)$. It follows that $r^+ A(t)e^+ \left(r^+ A(t)e^+ \right)^*$ is invertible for, say, $t < t_1$. This implies that $r^+ A(t)e^+$ is surjective on L_{w+}^2. We can apply the same argument to $t \mapsto \left(r^+ A(t)e^+ \right)^* r^+ A(t)e^+$ and find that it is also continuous at $t = 0$ and hence invertible for small t, say $t < t_2$. In particular, $r^+ A(t)e^+$ is injective on L_{w+}^2. Hence $r^+ A(t)e^+$ in invertible on L_{w+}^2 for $t < \min\{t_1, t_2\}$, and the same is true for $\left(r^+ A(t)e^+ \right)^*$.

In particular, both $r^+A(t)e^+$ and its $L^2(\mathbb{R}_+)$-adjoint $(r^+A(t)e^+)^*$ map $L^2(\mathbb{R}_+)$ to itself and are injective. Therefore, $r^+A(t)e^+$ is invertible on $L^2(\mathbb{R}_+)$ for small t.

Finally, consider $r^+A(t)e^+$ for small t on $\mathscr{S}(\mathbb{R}_+)$. It clearly is injective, since it is so on $L^2(\mathbb{R}_+)$. The inverse is $((r^+A(t)e^+)^*r^+A(t)e^+)^{-1}(r^+A(t)e^+)^*$. As $(r^+A(t)e^+)^*r^+A(t)e^+$ is an operator of order and type zero in Boutet de Monvel's calculus, so is its inverse by spectral invariance, cf. [28]. In particular, the inverse to $r^+A(t)e^+$ maps $\mathscr{S}(\mathbb{R}_+)$ to itself. Hence $r^+A(t)e^+$ also is surjective on $\mathscr{S}(\mathbb{R}_+)$ and therefore invertible. □

In view of Theorem 6.1, we can define the index of an admissible symplecto-morphism.

Definition 6.2. *Let X and Y be compact manifolds with boundary, and $\chi : T^*Y \setminus 0 \to T^*X \setminus 0$ an admissible symplectomorphism. Then, for each unitary section u of the Maslov bundle associated with χ, we define*

$$\mathrm{ind}(\chi, u) = \mathrm{ind}\,\mathcal{U}^\chi$$

with $\mathcal{U}^\chi$ as in (63).

As in the case of closed manifolds, the index $\mathrm{ind}(\chi, u)$ in general might depend on the chosen unitary section u. If the Maslov bundle is trivial, or the dimension of X and Y is at least three, then it turns out that $\mathrm{ind}(\chi, u)$ is independent of u. The argument is standard, and we shortly recall it for the sake of completeness. Consider two unitary sections u_1, u_2 and the corresponding operators $\mathcal{U}^\chi_{u_1}$ and $\mathcal{U}^\chi_{u_2}$. Theorem 3.8 implies that $\mathcal{U}^\chi_{u_1}(\mathcal{U}^\chi_{u_2})^*$ is an element of $\mathscr{B}^{0,0}(X)$, whose symbol, modulo lower order terms, coincides with $u_1\overline{u_2}$. If the Maslov bundle is trivial, or the dimension of X and Y is at least three, so that the cospheres are simply connected, there is a homotopy deforming $u_1\overline{u_2}$ to 1 and thus $\mathcal{U}^\chi_{u_1}(\mathcal{U}^\chi_{u_2})^*$ to the identity. Then,

$$0 = \mathrm{ind}[\mathcal{U}^\chi_{u_1}(\mathcal{U}^\chi_{u_2})^*] = \mathrm{ind}\,\mathcal{U}^\chi_{u_1} + \mathrm{ind}(\mathcal{U}^\chi_{u_2})^* = \mathrm{ind}\,\mathcal{U}^\chi_{u_1} - \mathrm{ind}\,\mathcal{U}^\chi_{u_2},$$

as claimed. In this case, we can actually define

$$\mathrm{ind}\,\chi := \mathrm{ind}(\chi, u).$$

7. Appendix: Operator-Valued Symbols

We denote by $H^s(\mathbb{R}^n)$, $s \in \mathbb{R}$, the usual Sobolev space on $\mathbb{R}^n$ and by $H^{s_1,s_2}(\mathbb{R}^n)$ the weighted space $\langle x \rangle^{-s_2}H^{s_1}(\mathbb{R}^n)$. We often write $\mathbf{s} = (s_1, s_2)$. Moreover, we let $H^{\mathbf{s}}(\mathbb{R}^n_+) = \{u|_{\mathbb{R}^n_+} : u \in H^{\mathbf{s}}(\mathbb{R}^n)\}$ and $H^{\mathbf{s}}_0(\overline{\mathbb{R}}^n_+) = \{u \in H^{\mathbf{s}}(\mathbb{R}^n) : \mathrm{supp}\,u \subseteq \overline{\mathbb{R}}^n_+\}$.

Next we recall a few facts on operator-valued symbols. For details see [29] or [30]. Let E and F be Banach spaces with strongly continuous group actions κ^E and κ^F of $\mathbb{R}_+$, that is, $\kappa^E : \mathbb{R}_+ \to \mathscr{L}(E)$ is strongly continuous and $\kappa^E(\lambda_1\lambda_2) = \kappa^E(\lambda_1)\kappa^E(\lambda_2)$ for $\lambda_1, \lambda_2 > 0$. We will write κ^E_λ instead of $\kappa^E(\lambda)$ for $\lambda > 0$. The corresponding notation is used for κ^F.

On the spaces $H^{\mathbf{s}}(\mathbb{R}^n), H^{\mathbf{s}}(\mathbb{R}^n_+)$ and $H^{\mathbf{s}}_0(\overline{\mathbb{R}}^n_+)$, $\mathbf{s} \in \mathbb{R}^2$, we always use the group action induced by the unitary action on $L^2(\mathbb{R}^n)$ given by

$$(\kappa_\lambda u)(x) = \lambda^{-n/2}u(\lambda x), \quad x \in \mathbb{R}^n.$$

Definition 7.1. *A smooth family $a(y, \eta)$, $y, \eta \in \mathbb{R}^q$, of operators in $\mathscr{L}(E, F)$ is called an operator-valued symbol of order $m \in \mathbb{R}$, if for all multi-indices α, β,*

$$\|\kappa^F_{\langle\eta\rangle^{-1}}D^\alpha_\eta D^\beta_y a(y, \eta)\kappa^F_{\langle\eta\rangle}\|_{\mathscr{L}(E,F)} = O(\langle\eta\rangle^{m-|\alpha|}).$$

We denote the space of all these symbols by $S^m(\mathbb{R}^q, \mathbb{R}^q; E, F)$.

The symbol a is called homogeneous of degree m, provided that

$$\kappa^F_{\lambda^{-1}}a(y, \lambda\eta)\kappa^E_\lambda = \lambda^m a(y, \eta), \quad \lambda > 0, \eta \neq 0,$$

and classical, if it has an asymptotic expansion into homogeneous terms.

This concept includes the usual pseudodifferential operators, by choosing $E = F = \mathbb{C}$ with the trivial group action.

The definition extends to more general spaces E and F. Noting that $\mathscr{S}(\mathbb{R}_+) = \text{proj-lim}_{\mathbf{s} \in \mathbb{R}^2} H^{\mathbf{s}}(\mathbb{R}_+)$, and $\mathscr{S}'(\mathbb{R}_+) = \text{ind-lim}_{\mathbf{s} \in \mathbb{R}^2} H_0^{\mathbf{s}}(\overline{\mathbb{R}}_+)$ we let

$$
\begin{array}{rcl}
S^m(\mathbb{R}^q, \mathbb{R}^q; \mathbb{C}, \mathscr{S}(\mathbb{R}_+)) & = & \text{proj-lim}_{\mathbf{s} \in \mathbb{R}^2} S^m(\mathbb{R}^q, \mathbb{R}^q; \mathbb{C}, H^{\mathbf{s}}(\mathbb{R}_+)) \\[4pt]
S^m(\mathbb{R}^q, \mathbb{R}^q; \mathscr{S}'(\mathbb{R}_+), \mathbb{C}) & = & \text{proj-lim}_{\mathbf{s} \in \mathbb{R}^2} S^m(\mathbb{R}^q, \mathbb{R}^q; H_0^{\mathbf{s}}(\overline{\mathbb{R}}_+), \mathbb{C}) \\[4pt]
S^m(\mathbb{R}^q, \mathbb{R}^q; \mathscr{S}'(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+)) & = & \text{proj-lim}_{\mathbf{s} \in \mathbb{R}^2} S^m(\mathbb{R}^q, \mathbb{R}^q; H_0^{-\mathbf{s}}(\overline{\mathbb{R}}_+), H^{\mathbf{s}}(\mathbb{R}_+)) \\[4pt]
S^m(\mathbb{R}^q, \mathbb{R}^q; \mathscr{S}(\mathbb{R}_+), \mathscr{S}(\mathbb{R}_+)) & = & \text{ind-lim}_{\mathbf{t} \in \mathbb{R}^2} \text{proj-lim}_{\mathbf{s} \in \mathbb{R}^2} S^m(\mathbb{R}^q, \mathbb{R}^q; H_0^{-\mathbf{t}}(\overline{\mathbb{R}}_+), H^{\mathbf{s}}(\mathbb{R}_+)).
\end{array}
$$

Lemma 7.2. *Let $a \in S^m(\mathbb{R}^n \times \mathbb{R}^n)$. For fixed $(x', \xi') \in \mathbb{R}^{n-1} \times \mathbb{R}^{n-1}$*

$$(x_n, \xi_n) \mapsto a(x', x_n, \xi', \xi_n)$$

defines a symbol in $S^m(\mathbb{R} \times \mathbb{R})$. For $u \in \mathscr{S}(\mathbb{R})$ we let

$$\left(\text{Op}_n(a)(x', \xi')\right) u(x_n) = \int e^{i x_n \xi_n} a(x', x_n, \xi', \xi_n) \hat{u}(\xi_n) \đ\xi_n.$$

Then $\text{Op}_n(a)$ is an operator-valued symbol in $S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}, H^{s_1, s_2}(\mathbb{R}_+), H^{s_1 - m, s_2}(\mathbb{R}_+))$ for all $s_1, s_2 \in \mathbb{R}$.

As $\kappa_{\langle \xi' \rangle}^{-1} \text{Op}_n(a) \kappa_{\langle \xi' \rangle}$ has the symbol $a(x', x_n / \langle \xi' \rangle, \xi', \xi_n \langle \xi' \rangle)$, the assertion follows from the usual symbol estimates. The same consideration shows

Lemma 7.3. *For $s_1, s_2 \in \mathbb{R}$ multiplication by x_n defines an operator-valued symbol in $S^{-1}(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}, H^{s_1, s_2}(\mathbb{R}), H^{s_1, s_2 + 1}(\mathbb{R}))$.*

With $a \in S^m(\mathbb{R}^q, \mathbb{R}^q; E, F)$ we associate the pseudodifferential operator $\text{Op}(a)$, defined on the space $\mathscr{S}(\mathbb{R}^q, E)$ of rapidly decreasing E-valued functions by

$$\text{Op}(a) u(y) = \int e^{i y \eta} a(y, \eta) \hat{u}(\eta) \, \đ\eta.$$

It maps $\mathscr{S}(\mathbb{R}^q, E)$ continuously to $\mathscr{S}(\mathbb{R}^q, F)$. Here $\hat{u}$ is the vector-valued Fourier transform

$$\hat{u}(\eta) = \int e^{i y \eta} u(y) \, \đy.$$

Theorem 7.4. *Let $a_1 \in S^{m_1}(\mathbb{R}^q, \mathbb{R}^q; F, G)$ and $a_2 \in S^{m_2}(\mathbb{R}^q, \mathbb{R}^q; E, F)$ for Banach spaces E, F and G. Then $\text{Op}(a_1) \text{Op}(a_2) = \text{Op}(a)$ for an operator-valued symbol $a \in S^{m_1 + m_2}(\mathbb{R}^q, \mathbb{R}^q; E, G)$ with the asymptotic expansion*

$$a(y, \eta) \sim \sum_\alpha \frac{1}{\alpha!} \partial_\eta^\alpha a_1(y, \eta) D_y^\alpha a_2(y, \eta).$$

The proof is as in the standard case.

Definition 7.5. *Let $s \in \mathbb{R}$ and E be a Banach space. The wedge Sobolev space $\mathscr{W}^s(\mathbb{R}^q, E)$ is the completion of $C_c^\infty(\mathbb{R}^q, E)$ with respect to the norm*

$$\|u\|_{\mathscr{W}^s(\mathbb{R}^q, E)}^2 = \int \langle \eta \rangle^{2s} \|\kappa_{\langle \eta \rangle}^{-1} \hat{u}(\eta)\|_E^2 \, \đ\eta.$$

The definition extends to projective limits as above. A direct computation shows the first identity in the following lemma; the second then is immediate.

Lemma 7.6. $\mathscr{W}^s(\mathbb{R}^{n-1}, H^s(\mathbb{R})) = H^s(\mathbb{R}^n)$ *and* $\mathscr{W}^s(\mathbb{R}^{n-1}, H^s(\mathbb{R}_+)) = H^s(\mathbb{R}_+^n)$, $s \in \mathbb{R}$.

A proof of the following statement can be found in [31].

Theorem 7.7. *Let $a \in S^m(\mathbb{R}^q, \mathbb{R}^q; E, F)$ for Hilbert spaces E and F. Then*

$$\mathrm{Op}(a) : \mathscr{W}^s(\mathbb{R}^q, E) \to \mathscr{W}^{s-m}(\mathbb{R}^q, F)$$

is bounded, $s \in \mathbb{R}$.

As $\mathscr{W}^0(\mathbb{R}^{n-1}, L^2(\mathbb{R}_+)) = L^2(\mathbb{R}_+^n)$, we can also consider the formal adjoint $\mathrm{Op}(a)^*$ of $\mathrm{Op}(a)$ with respect to the $L^2(\mathbb{R}^n)$ scalar product, when $a \in S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; E, F)$ and E, F are in the scales $H^s(\mathbb{R}_+)$ or $H_0^s(\overline{\mathbb{R}}_+)$, $s \in \mathbb{R}$. It is easy to see that $\mathrm{Op}(a)^* = \mathrm{Op}(a^{(*)})$ for an $a^{(*)} \in S^m(\mathbb{R}^{n-1}, \mathbb{R}^{n-1}; F', E')$ with the asymptotic expansion

$$a^{(*)}(y, \eta) \sim \sum_\alpha \frac{1}{\alpha!} \partial_\eta^\alpha D_y^\alpha a(y, \eta)^*.$$

Here $a(y, \eta)^* \in \mathscr{L}(F', E')$ is the adjoint of $a(y, \eta)$.

References

[1] U. Battisti, S. Coriasco, and E. Schrohe. On a Class of Fourier Integral Operators on Manifolds with Boundary. arXiv:1406.0636.

[2] U. Battisti, S. Coriasco, and E. Schrohe. On the Transmission Property for Fourier Integral Operators of Boutet de Monvel Type. In preparation.

[3] L. Boutet de Monvel. Boundary problems for pseudo-differential operators. *Acta Math.*, 126(1-2):11–51, 1971.

[4] L. Boutet de Monvel and V. Guillemin. *The spectral theory of Toeplitz operators*, volume 99 of *Annals of Mathematics Studies*. Princeton University Press, Princeton, NJ, 1981.

[5] A. Cannas da Silva. *Lectures on symplectic geometry*, volume 1764 of *Lecture Notes in Mathematics*. Springer-Verlag, Berlin, 2001.

[6] H. Christianson. Classification of order preserving isomorphisms between algebras of semiclassical operators. *Proc. Amer. Math. Soc.*, 139(2):499–510, 2011.

[7] S. Coriasco. Fourier integral operators in SG classes II: Application to SG hyperbolic Cauchy problems. *Ann. Univ. Ferrara, Sez. VII - Sc. Mat.*, 44:81–122, (1998), 1999.

[8] S. Coriasco. Fourier integral operators in SG classes. I. Composition theorems and action on SG Sobolev spaces. *Rend. Sem. Mat. Univ. Politec. Torino*, 57(4):249–302 (1999), 2002.

[9] J. J. Duistermaat and I. M. Singer. Order-preserving isomorphisms between algebras of pseudo-differential operators. *Comm. Pure Appl. Math.*, 29(1):39–47, 1976.

[10] C. Epstein and R. Melrose. Contact degree and the index of Fourier integral operators. *Math. Res. Lett.*, 5(3):363–381, 1998.

[11] C. L. Epstein. Subelliptic Spin$_{\mathbb{C}}$ Dirac operators. I. *Ann. of Math. (2)*, 166(1):183–214, 2007.

[12] C. L. Epstein. Subelliptic Spin$_{\mathbb{C}}$ Dirac operators. II. Basic estimates. *Ann. of Math. (2)*, 166(3):723–777, 2007.

[13] C. L. Epstein. Subelliptic Spin$_{\mathbb{C}}$ Dirac operators. III. The Atiyah-Weinstein conjecture. *Ann. of Math. (2)*, 168(1):299–365, 2008.

[14] G. Grubb. Complex powers of pseudodifferential boundary value problems with the transmission property. In *Pseudodifferential operators (Oberwolfach, 1986)*, volume 1256 of *Lecture Notes in Math.*, pages 169–191. Springer, Berlin, 1987.

[15] G. Grubb. *Functional Calculus for Boundary Value Problems*. Birkahäuser, Boston, 1996. Reprint of the 1982 edition.

[16] G. Grubb and L. Hörmander. The transmission property. *Math. Scand.*, 67(2):273–289, 1990.

[17] V. Guillemin. Toeplitz operators in n dimensions. *Integral Equations Operator Theory*, 7(2):145–205, 1984.

[18] A. Hirschowitz and A. Piriou. La propriété de transmission pour les distributions de Fourier; application aux lacunes. In *Séminaire Goulaouic-Schwartz (1976/1977), Équations aux dérivées partielles et analyse fonctionnelle, Exp. No. 14*, page 19. Centre Math., École Polytech., Palaiseau, 1977.

[19] L. Hörmander. *The analysis of linear partial differential operators III. Pseudo-differential operators*. Classics in Mathematics. Springer, Berlin, 2007. Reprint of the 1994 edition.

[20] L. Hörmander. *The analysis of linear partial differential operators IV*. Classics in Mathematics. Springer, Berlin, 2007. Reprint of the 1994 edition.

[21] A. Laptev, Y. Safarov, and D. Vassiliev. On global representation of Lagrangian distributions and solutions of hyperbolic equations. *Comm. Pure Appl. Math.*, 47(11):1411–1456, 1994.

[22] E. Leichtnam, R. Nest, and B. Tsygan. Local formula for the index of a Fourier integral operator. *J. Differential Geom.*, 59(2):269–300, 2001.

[23] V. Mathai and R. B. Melrose. Geometry of pseudodifferential algebras bundles and Fourier integral operators. arXiv:1210.0990.

[24] R. B. Melrose. Transformation of boundary problems. *Acta Math.*, 147(3-4):149–236, 1981.

[25] V. E. Nazaikinskii and B. Y. Sternin. Surgery and the relative index in elliptic theory. *Abstr. Appl. Anal.*, pages Art. ID 98081, 16, 2006.

[26] V. E. Nazaĭkinskiĭ, B. Y. Sternin, and B.-V. Shul'tse. The index of Fourier integral operators on manifolds with conical singularities. *Izv. Ross. Akad. Nauk Ser. Mat.*, 65(2):127–154, 2001.

[27] S. Rempel and B.-W. Schulze. *Index theory of elliptic boundary problems*. North Oxford Academic Publishing Co. Ltd., London, 1985. Reprint of the 1982 edition.

[28] E. Schrohe. Fréchet algebra techniques for boundary value problems on noncompact manifolds: Fredholm criteria and functional calculus via spectral invariance. *Math. Nachr.*, 199:145–185, 1999.

[29] E. Schrohe. A short introduction to Boutet de Monvel's calculus. In *Approaches to singular analysis (Berlin, 1999)*, volume 125 of *Oper. Theory Adv. Appl.*, pages 85–116. Birkhäuser, Basel, 2001.

[30] B.-W. Schulze. *Pseudo-differential operators on manifolds with singularities*, volume 24 of *Studies in Mathematics and its Applications*. North-Holland Publishing Co., Amsterdam, 1991.

[31] J. Seiler. Continuity of edge and corner pseudodifferential operators. *Math. Nachr.*, 205:163–182, 1999.

[32] A. Weinstein. Fourier integral operators, quantization, and the spectra of Riemannian manifolds. In *Géométrie symplectique et physique mathématique (Colloq. Internat. CNRS, No. 237, Aix-en-Provence, 1974)*, pages 289–298. Éditions Centre Nat. Recherche Sci., Paris, 1975. With questions by W. Klingenberg and K. Bleuler and replies by the author.

[33] A. Weinstein. Some questions about the index of quantized contact transformations. *Sūrikaisekikenkyūsho Kōkyūroku*, (1014):1–14, 1997. Geometric methods in asymptotic analysis (Japanese) (Kyoto, 1997).

[34] S. Zelditch. Index and dynamics of quantized contact transformations. *Ann. Inst. Fourier (Grenoble)*, 47(1):305–363, 1997.