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Case study: A GPR survey on a morainic lake in northern Italy for bathymetry, water volume and sediment characterization.

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Abstract

We carried out an extensive waterborne GPR survey consisting of 50 profiles with a total length of nearly 37 km on the morainic lake of Candia northerly Turin (Italy). Our aim was to test the capability of GPR to estimate the bathymetry, the water volume and the sediment type. We enhanced and controlled the GPR data processing and interpretation with bathymetry acquired with an acoustic echo sounder and measured conductivity and temperature profile of the water column with a multiparametric probe. We also analysed the diffraction hyperbola that originated within the sediments in order to estimate the velocity and relative permittivity. With the permittivity and dielectric mixing rules, we estimated the porosity of the sediments above the diffracting objects and drew a map of the bottom lake porosity.
1. Introduction

Here we discuss an extensive waterborne GPR survey of a morainic lake in northern Italy that we performed along with an echo-sounder and spotted measurements along the water column with a multiparametric sonde. The survey aimed to evaluate the capability of a non-seismic tool as the GPR to estimate the water volume and the bottom sediment type.

The use of non-seismic methods to investigate sediments beneath shallow freshwater has increased in the last twenty years (Sambuelli and Butler, 2009; Butler, 2009). Monitoring river erosion, lake filling and understanding the connection between surface water and underground water are critical environmental issues. Many publications related to the application of GPR to shallow-water environments can be found in literature. Annan and Davis (1977) and Kovacs (1978) conducted some early applications of GPR in low-conductivity melting water in arctic areas. In such low-conductivity environments a noticeable subbottom penetration in water depths exceeding 25 m can be achieved (Delaney et al., 1992). Bathymetric maps and sediments morphology of ice-covered lakes (Moorman and Michel, 1997; Schwamborn et al., 2002) rivers and reservoirs (Arcone et al., 1992; Best et al., 2005; Hunter et al., 2003) have been reported as well. In temperate areas GPR has been tested for monitoring stream discharge (Haeni et al., 2000; Melcher et al., 2002; Cheng et al., 2004; Costa et al., 2006), studying bottom deposits (Arcone et al., 2006; Buynevich and Fitzgerald, 2003; Fuchs et al., 2004; Shields et al., 2004; Moorman et al., 2001, Sambuelli et al., 2009), mapping bathymetry (Powers et al., 1999) and subbottom strata (Arcone et al., 2010).

GPR could provide complementary information to seismic methods, especially in very shallow water where reverberation can prevent the interpretation of subbottom reflectors or when gas pockets in sediments rich in organic matter block seismic signal penetration (Delaney et al., 1992; Schwamborn et al., 2002; Arcone et al., 2006; Mellet, 1995). Waterborne GPR can be used in many different setups: case histories report the use of antennas directly coupled to water from the surface (Mellet, 1995; Sellmann et al., 1992), noncontact systems such as helicopter mounts (Melcher et al., 2002) or hanging slings (Haeni et al., 2000; Cheng et al., 2004; Costa et al., 2000), and antennas mounted on the bottom of non-metallic boats (Sambuelli et al., 2009; Sellmann et al., 1992, Bradford et al., 2005; Jol and Albrecht, 2004; Porsani et al., 2004; Park et al., 2004). The potential of GPR to allow estimates bottom sediment types has been mentioned by Ulriksen (1982), who suggested that fine sediments can be identified from smooth reflectors, rough surfaces such as moraine can be identified from speckled and weak signals, while boulders may produce hyperbolic diffractions. As a rule of thumb the coarser is the sediment then the higher are the amplitudes of the reflected GPR signals. The same qualitative approach to the analysis of basin-bottom characteristics has been tested positively by other authors (Mellet, 1995, Beres and Haeni, 1991). In particular,
Dudley and Giffen (1999) found optimal agreement between GPR results and direct sampling of sediments. Recently some more quantitative approaches to the discrimination of sediment type using the amplitude of the GPR signal reflected from the water bottom have been proposed (Sambuelli et al., 2009; Lin et al., 2010).

Estimating the bathymetry and discriminating sediment types by analysing the time delay and amplitude of reflections mainly assumes that the water is homogeneous. This assumption is valid in shallow rivers or lakes where there is little or no thermocline and therefore no reflectors before the bottom, but must be questioned for deeper lakes (Bradford et al., 2007).

2. Site description
The lake of Candia (Turin, Italy) is an example of inter-moraine lake belonging to the morainic amphitheatre of Ivrea, northerly the town of Turin (Figure 1). The lake (226 m a.s.l.) has a surface of about 1.52 km², a perimeter of about 5.5 km and an estimated volume of about 7.2 Mm³. The amphitheatre was built up by the Balteo Glacier which flowed through the Aosta Valley during the Quaternary period. The Balteo Glacier had an average width of 3 km, a thickness of around 600 m and was spread over 300 km through the Alps. At its maximum expansion it went down to 25 km into the Po plain with a width of about 20 km, reaching today’s northern border of the town of Turin. According to recent studies the lake of Candia was left after the recession of the Riss Glaciation (about 150000 years ago). The lake of Candia then settled in a plain that was eroded by the advancing glacier and then refilled to a far lower extent by the glacial deposits of the glacier retreat. According to Gianotti et al. (2008), the Candia lake lies on the Piverone alloformation (Late Pleistocene) above the Serra alloformation (end of the Middle Pleistocene) each one referable to a different glacial episode.

A borehole (Alice Superiore borehole) drilled about 20 km NNW in the same geological environment gave the stratigraphic log reported in Figure 2. It is evident the dominance of more or less dense gravels with boulders that may be included in the mud. A palustrine gyttja layer made by peaty brown clay represents the interstadial episode between the two glacial events. Until about 50 years ago, before an eutrophication process took place, it was possible to see the white gravels of the lake bottom and the fresh water columns rising up to the surface of the lake from underwater springs. Today, the silts and the vegetation prevent the view of the lake bottom. The lake has neither affluent nor effluent but a small ditch used sometimes for irrigation; the water recharge is from rains and underwater springs.

3. Methods
3.1 Data acquisition
We acquired the GPR data with a K2 Ids radar and a Subecho-70 monostatic antenna with peak frequency in air at 89 MHz. We placed the antenna at the bottom of a fiber-reinforced plastic boat and the routes of the profiles were planned and tracked with a Garmin GPS60. We acquired 50 GPR profiles for a total length of about 37 km (Figure 3) at a sampling frequency of 1.7 GHz, trace length of 600 ns and a traces interval of 4 cm. Neither filters nor gains were used during the acquisition.

We also acquired bathymetry with a 170kHz Airmar DT800 echo sounder. The dominant wavelength in water roughly equalled 1 cm, and our average sampling interval was 1 m. The echo sounder probe was fixed at the side of the boat and approximately aligned with the centre of the GPR antenna. We also used a Hydrolab Datasonde 4 Multi-Parametric Probe (MPP) to measure temperature, T, and conductivity, \( \sigma \), along the water column at three points on the lake (Figure 3).

3.2 Data processing
We needed to estimate the water conductivity to evaluate the GPR signal attenuation and the water permittivity to evaluate the GPR pulse velocity. We also estimated the dominant frequency of the GPR pulse by doing a spectral analysis of windowed reflected pulses from different water depths. The results shown a dominant frequency range from 66 to 80 MHz with a mode around 70 MHz which was the value selected for further calculations. As a solution to Maxwell’s equations the amplitude at a distance \( z \) and time \( t \) of the electric field \( E(z, t, \omega) \) associated with a monochromatic plane wave of angular frequency \( \omega \) is

\[
E(z, t, \omega) = E(z_0, t_0, \omega)e^{i(\omega z + k_\omega z + \omega \varepsilon_0 z + \varepsilon_0 t)},
\]

where \( E(z_0, t_0, \omega) \) is the amplitude at reference time and distance \( t_0 \) and \( z_0 \). The quantity

\[
k_w = i\alpha_w + \frac{\omega}{v_w}
\]

is the complex wave number (subscript \( w \) refers to water) where

\[
\alpha_w = \omega \sqrt{\frac{\mu_\omega \mu_0}{\varepsilon_\omega \varepsilon_0}} \left[ \frac{1}{\varepsilon_\omega \varepsilon_0} \right]^{\frac{1}{2}} \left[ \frac{1}{\varepsilon_\omega \varepsilon_0 (\omega)} \right]^{\frac{1}{2}}
\]

and

\[
v_w = \frac{1}{\sqrt{\frac{\mu_\omega \mu_0}{\varepsilon_\omega \varepsilon_0}} \left[ \frac{1}{\varepsilon_\omega \varepsilon_0} \right]^{\frac{1}{2}} \left[ 1 + \left( \frac{\sigma_\omega}{\varepsilon_\omega \varepsilon_0} \right)^2 \right]^{\frac{1}{2}}}
\]
are respectively, the attenuation factor due to dissipative phenomena and the phase velocity. Given 
\( \mu_0 = 4\pi \times 10^{-7} \) [H/m], \( \mu_w = \mu_0 \), \( \epsilon_0 = (36\pi)^{-1} \times 10^{-9} \) [F/m] and \( \omega = 2\pi \times 70 \times 10^6 \) [rad/s], we need \( \varepsilon_w \) and \( \sigma_w \) to evaluate \( \alpha_w \) and \( v_w \).

The MPP gives the T and \( \sigma_w \) so that we calculated the permittivity using the formula given by Nyshadam et al. (1992):

\[
\varepsilon_w = 87.74 - 4.0008 \times 10^{-1} T - 0.09398 \times 10^{-2} T^2 - 0.00141 \times 10^{-3} T^3 \pm 0.009
\]

The average salinity of the lake water was about 0.06 ppt (part per thousand) and so the correction to apply to eq. 5 for such values of salinity is negligible. We plotted the T and the \( \sigma_w \) along the water columns according to the MPP data (Figure 4 a,b). The small temperature gradient below the first meter was so small and showed such small variations (Figure 4 c,d,e) that we rejected the hypothesis of a thermo cline. We then used the following values of velocity and attenuation in the GPR processing: \( v_w = 0.033 \) [m/ns] and \( \alpha_w = 2.24 \) [dB/m].

We processed the GPR profiles with the Reflexw® software. We corrected for the direct arrival delay and eliminated the very low frequency trend. We band-pass filtered from 20 to 190 MHz with a Butterworth filter to attenuate both low and high frequency noise and then resampled the traces to obtain a trace interval of 0.2 m. We removed the steady background horizontal continuous signals with a spatial high-pass filter. We applied firstly a divergence compensation, using \( v_w \), to recover the amplitude attenuation due to the geometrical spreading, and then secondly an exponential gain, using \( \alpha_w \), to recover the amplitude decrease due to the dissipation related to the water conductivity. As far as the wavelength is concerned, the dominant wavelength of the GPR signal bottom reflection ranged from about 0.4 m in the more shallow waters to about 0.5 m in the deepest. To enhance the continuity of the reflectors we then applied a horizontal running average over 11 traces that, giving the new spatial sampling interval, roughly corresponds to a low pass filter that attenuates the wavelengths shorter than 2 m. We then picked the first break of the first bottom reflection to estimate the bathymetry of the lake and to calculate the volume of water. Finally we applied a muting above the picked times to cancel out residual noise in water. In the following some processed radargrams will be presented.

To assess the bathymetry given by the GPR we compared the water depths to those obtained with the acoustic sounder. On average the difference between the two bathymetric data was about 0.17 m with a standard deviation of 0.14 m. This discrepancy could be caused by different reflecting boundaries because the acoustic signal should have no penetration whereas the GPR signal is sensitive to subtle near surface layering. This difference could also be enhanced by the
largely different wavelengths so that the leading edges of the pulses could be differently estimated. It could be also due to the fact that the echo sounder sensor was roughly connected to the boat side and we could not account for the exact depth of the sensor below the water surface nor its rolling and pitching with respect to the boat.

To estimate the bottom sediment porosity we used the diffraction events that later than the bottom lake reflection. The diffractions are caused by point-like objects within the sediments. The $t$ and $x$ coordinates of these events obey the following equations

\[
t = \frac{2h_1}{v_1 \cos(i)} + \frac{2h_2}{v_2 \cos(r)}
\]

and

\[
x = h_1 \tan(i) + h_2 \tan(r),
\]

where $h_1$ is the water depth above the vertex of the diffraction pattern, $h_2$ is the thickness of sediment above the vertex, $i$ is the incident angle, $r$ is the refracted angle and $v_1$ and $v_2$ are, respectively, the velocity of GPR wave in water and in sediment (Figure 5). The incident and the refracted angles $i$ and $r$, $v_1$ and $v_2$ are related by Snell’s law

\[
\frac{\sin(i)}{\sin(r)} = \frac{v_1}{v_2}
\]

For each one of the diffraction events within the radargrams we can read $n$ experimental data pairs $(t_e, x_e)$, $2h_1/v_1$ and $2h_2/v_2$. These latter are respectively, the two-way-traveltime (twt) of the reflection from water-sediment interface above the event vertex (i.e., $t_1$) and the twt relative to the path within the sediments from the above mentioned point to the vertex of the event (i.e., $t_2$). However, the two equations [6, 7] cannot be solved for a unique solution in the form $t=t(x)$. Therefore we set up a procedure to determine the unknowns $h_2$ and $v_2$ which consists of fitting the $(t_e, x_e)$ with $(t, x)$ obtained from equations [6, 7] by spanning all possible incident angles with values of $v_2$ ranging from the velocity in water to the velocity in air. We then found the sediment velocity $v_2$ by searching for the minimum of a misfit function between calculated and experimental traveltimes related to each diffraction event. To have a more robust estimate, less sensitive to outliers, the selected function was the mean absolute difference (mad) that is:
\[ mad = \frac{\sum |t_e - t|}{N}, \]  

where \( t_e \) are the experimental picked twt in \( x_e \), \( t \) are the calculated twt in \( x_e \) and \( N \) is the number of twt’s considered for each diffraction event. A simplified script referring to the routine implemented for the procedure is reported in Appendix A.

We assessed the effectiveness of this procedure by performing a synthetic test with the finite difference modeller contained in the Reflexw® software. We built a model consisting of a point-like scattering target within sediments under a water layer. The synthetic GPR pulse had a dominant frequency of 100 MHz; the scatter source had a diameter \( D = 0.16 \) m and the dominant wavelength of the GPR signal was \( \lambda = 0.32 \) m. In Figure 5 a sketch of the geometry implied in equations 6, 7, 8 with the main physical and geometrical parameters used for the synthetic model is shown. In Figure 6a, b the model and the synthetic radargram are respectively shown. The fitting procedure results were \( \epsilon_2 = 25.06 \) and \( h_2 = 1.72 \) m, with a relative percent error with respect to the expected values respectively of 0.24% and 1.16%.

We then selected the clearest diffraction events along the profiles and, according to the proposed procedure, estimated the local porosities. In Figures 7 and 8a, 8b an example of the aforementioned procedure is shown. Figure 7 shows one of the selected diffraction events observed beneath the bottom sediment. Figure 8a shows the trend of the misfit function between calculated and observed data. Figure 8b shows the fitting between experimental and calculated (twt, x) pairs. Over the 51 diffraction events we analysed we had the following statistical figures: mean(min(mad)) = 0.95 ns, st.dev.(min(mad)) = 1.16 ns. We considered these statistical results positively. In fact the sampling interval during acquisition was 0.588 ns so that the mean of the selected mad’s (mean(min(mad))) was less than twice the sampling interval of the radar traces.

We then inserted the \( v_2 \) values, converted in the dielectric permittivity of sediments \( \epsilon_s \), in a mixing rule that, given the dielectric permittivity of water and solid particles, allowed the local sediment porosity (Sambuelli, 2009) to be estimated. Among the many mixing formulas available we selected the Maxwell (Equation 10) (Maxwell, 1892) and the so called CRIM (Birchak et al., 1974) which are

\[ \phi = 3\epsilon_w \frac{\epsilon_s - \epsilon_w}{(\epsilon + 2\epsilon_w)(\epsilon_w - \epsilon_s)} , \]  

and

\[ \phi = \frac{\sqrt{\epsilon_s} - \sqrt{\epsilon_{ss}}}{\sqrt{\epsilon_w} - \sqrt{\epsilon_{ss}}} \]
respectively, where $\phi$ is the porosity, $\varepsilon_r$ the relative permittivity of the sediments calculated from $v_2$, $\varepsilon_w=82$ (the relative permittivity of the water taken from the MPP data), and $\varepsilon_{ss}=6$ is an average relative permittivity of the solid grains taken from the literature (Reynolds, 1997; Parkhomenko, 1967). We then averaged the porosities.

4. Results and discussion

The GPR survey shows that the bottom of Candia lake consists of two main types of sediment. One is gravel and pebbles, which shows a high reflectivity and therefore is not penetrated by GPR signal. The other is organic-rich silt, with a lower reflectivity, more transparent, and allows the underlying bedding to be detected. In Figure 9 we show the radargrams recorded along the red lines in Figure 3. From the top to the bottom of the figure the profiles are from East to West. The gravels are clearly visible as is their progressive cover by the silts going westward. The top radargram, in particular, shows the high reflectivity of the gravels underlain by a multiple reflection. On the other hand the transparency of the silt is clearly depicted by the evident reflection of the gravels beneath the silt. The gravels are likely moraine while the silt (younger) is likely later sedimentation always superimposed on the gravels. We confirmed the outcropping of the gravels and pebbles of the moraine by sampling with a Van Veen grab bucket (Figure 10a). We actually found respectively organic rich slime (Figure 10b) and pebbles (Figure 10c) where the radargrams showed weak and strong reflections.

By converting the twt’s of the bottom reflections in water depth and using the velocity of the GPR pulse in water obtained with the MPP, we plotted the bathymetry of the lake (Figure 11). The water volume we calculated agreed with the value (about 7 Mm$^3$) referred to in literature (Comino et al., 2005).

In Figure 12 we show a porosity map of the lake bottom (Figure 12). We mapped porosity 0.15 up 0.65. The lowest values are compatible with poorly sorted gravel with sand and silt filling the spaces between the largest grains. The highest value are compatible with silt or fine sand rich in organic material. These unusually high porosity values have also been found in the sea bottom (Ullman and Aller, 1982), and in pond (Avnimelech et al., 2001) and lake bottoms (Hartke and Hill, 1974). Even if the sampled points were few and oddly distributed it seems evident that significant areas of the bottom sediments have porosity larger than 0.55. These areas have a great role in the biological and hydro-geochemical processes developing within the lake.

5. Conclusion
Our three main results are: the identification of an area of the lake bottom made by the coarse materials of the moraine, the bathymetry and water volume and the estimated porosity of the bottom sediments.

The waterborne GPR can be effective to estimate bathymetry of shallow freshwater basins, together with the bedding and porosity of the bottom sediments, provided that the water resistivity, using a GPR central frequency ranging from 50 to 100 MHz, is higher than 40 $\Omega$ m to reach 3 to 10 m depths. The depth can be even greater if very low conductivity water derives from ice melting. Profiling of the subbottom bedding largely depends on the bottom sediment type. The finer are the bottom sediments the lower is their reflectivity the higher is the probability to image strata beneath the bottom surface. An estimation of the bottom sediment porosity can also be achieved if point-like scattering targets within the sediments are found. Then our diffraction matching procedure allows an estimation of the GPR pulse velocity in the sediment and hence, with suitable mixing rules, a porosity range for the sediment above the scatterers can be derived.

Appendix A

We present a scheme of the script we wrote to implement the procedure for the determination of the velocity of sediments $v_2$. The knowns quantities are: $t_1$, $v_1$, $h_1$, $t_2(=h_2/2v_2)$, $\{t_e\}$, $\{x_e\}$. In the script: $v_w =$ velocity in water; $v_{air} =$ velocity in air; $t_e =$ experimental times picked on the diffraction hyperbola; $x_e =$ coordinate related to $t_e$; $r =$ refraction angle; $v_2 =$ velocity in sediments. The line “spline $(t,x) \rightarrow t(x | v_2)$” means that a cubic spline is used to interpolate the calculated $(t,x)$ and calculate the value of $t$ corresponding to each $x_e$ coordinate.

for $v_2 = v_w$ to $v_{air}$ [200 steps]
  calculate $h_2$
  for $r = 0^\circ$ to $75^\circ$ [25 steps]
    calculate $i$ (Eq.8)
    calculate $t$ (Eq.6)
    calculate $x$ (Eq.7)
  next $r$
  spline $(t, x) \rightarrow t(x | v_2)$
  calculate $t(x_e | v_2)$ % calculate the $t$’s after spline interpolation
  read $(t_e)$ at $x_e$ % read the experimental data
  calculate mad (Eq.9) % using $t(x_e | v_2)$ and $t_e$
  next $v_2$

plot $(v_2, mad)$ (Figure 8a) % plot mad versus velocity
find $v^*_2 = v_2(\text{min(mad)})$ % find the velocity corresponding to the minimum mad
plot $t(x | v_2)$ and $(x_e, t_e)$ (Figure 8b) % plot the calculated hyperbola and the experimental data

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References


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Caption of Figures

Figure 1: The location of the Candia Lake in northern Italy.

Figure 2: Stratigraphic log of the Alice Superiore borehole (redrawn from Gianotti et al. [34]).

Figure 3: The violet lines show the GPR profiles. The four yellow lines refer to the profiles (1, 4, 6, 7 from East to West) that are more deeply illustrated in the paper. Yellow pins refer to the locations of the Multi-Parametric Probe measurements.

Figure 4: The vertical profiles of: measured temperatures (a); measured conductivities (b); calculated relative permittivity (c); calculated GPR pulse velocity (d); calculated GPR pulse attenuation (e).

Figure 5: Sketch of the geometrical model used to calculate the GPR pulse velocity in sediments using the diffraction events from objects within the sediments (Equations 6, 7, 8). The electromagnetic and geometrical parameters used in the finite difference synthetic model are also shown.

Figure 6: The synthetic model (a); the synthetic radargram (b).

Figure 7: An experimental diffraction event within the sediments.

Figure 8: Analysis of the event in figure 7. The search for the velocity giving the minimum mad (a); the fitting of the calculated (continuous line) and experimental (circles) twt,x pairs (b).

Figure 9: The processed radargrams corresponding to the four profiles shown as yellow lines in Figure 3. All the radargrams runs from South to North. On the left (near the South shore) the gravels, and the related multiples, are visible as brighten reflections in all the radargrams. In profiles 6 and 7 the gravels are covered by the silt that fills the great part of the lake bottom.

Figure 10: The direct sampling of the lake bottom. The Van Veen Grab Bucket (a); the silt (b); the pebbles (c).

Figure 11: The bathymetry map of the lake.

Figure 12: The porosity map of the lake sediments. The black crosses refer to the points were we found the analyzed diffraction events within the sediments.
Figure 5

Water $\varepsilon = 82$
$v_1$ $\rho_{\text{w}} = 80 \ \Omega \text{m}$
$h_1 = 2.7 \text{ m}$

Sediments $\varepsilon_s = 25$
$v_2$ $\rho_s = 100 \ \Omega \text{m}$
$h_2 = 1.7 \text{ m}$

Scatter $\varepsilon_{\text{sc}} = 6$
Source $\rho_{\text{sc}} = 500 \ \Omega \text{m}$

Figure 6
Figure 7

Figure 8
Figure 9
Figure 10

Figure 11
Figure 12