

Identification of BeO and BeOxDy in melted zones of the JET Be limiter tiles: Raman study using comparison with laboratory samples

*Original*

Identification of BeO and BeOxDy in melted zones of the JET Be limiter tiles: Raman study using comparison with laboratory samples / Kumar, M.; Makepeace, C.; Pardanaud, C.; Ferro, Y.; Hodille, E.; Martin, C.; Roubin, P.; Widdowson, A.; Dittmar, T.; Linsmeier, C. h.; Lungu, C. P.; Porosnicu, C.; Jepu, I.; Dinca, P.; Lungu, M.; Pompilian, O. G.; Subba, F.; Jet, Contributors. - In: NUCLEAR MATERIALS AND ENERGY. - ISSN 2352-1791. - 17:(2018), pp. 295-301. [10.1016/j.nme.2018.11.008]

*Availability:*

This version is available at: 11583/2986749 since: 2024-03-11T10:13:48Z

*Publisher:*

ELSEVIER

*Published*

DOI:10.1016/j.nme.2018.11.008

*Terms of use:*

This article is made available under terms and conditions as specified in the corresponding bibliographic description in the repository

*Publisher copyright*

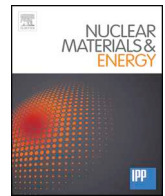
(Article begins on next page)



ELSEVIER

Contents lists available at ScienceDirect

## Nuclear Materials and Energy

journal homepage: [www.elsevier.com/locate/nme](http://www.elsevier.com/locate/nme)

# Identification of BeO and BeOxD<sub>y</sub> in melted zones of the JET Be limiter tiles: Raman study using comparison with laboratory samples

M. Kumar<sup>a</sup>, C. Makepeace<sup>b,f</sup>, C. Pardanaud<sup>a,\*</sup>, Y. Ferro<sup>a</sup>, E. Hodille<sup>a</sup>, C. Martin<sup>a</sup>, P. Roubin<sup>a</sup>, A. Widdowson<sup>b</sup>, T. Dittmar<sup>c</sup>, C.h. Linsmeier<sup>c</sup>, C.P. Lungu<sup>d</sup>, C. Porosnicu<sup>d</sup>, I. Jecu<sup>b,d</sup>, P. Dinca<sup>d</sup>, M. Lungu<sup>d</sup>, O.G. Pompilian<sup>d</sup>, JET contributors<sup>e,\*\*</sup>

<sup>a</sup> Aix Marseille Univ, CNRS, PIIM, Marseille, France

<sup>b</sup> Department of Materials, University of Oxford, Parks Road, Oxford OX1 3PH, UK

<sup>c</sup> Forschungszentrum Jülich GmbH, Institut für Energie- und Klimaforschung - Plasmaphysik, Jülich 52425, Germany

<sup>d</sup> National Institute for Laser, Plasma and Radiation Physics, Magurele-Bucharest 077125, Romania

<sup>e</sup> EUROfusion Consortium, JET, Culham Science Centre, Abingdon OX14 3DB, UK

<sup>f</sup> Culham Centre for Fusion Energy, Culham Science Centre, Abingdon, OX14 3DB, UK

## ARTICLE INFO

## Keywords:

Beryllium oxide  
Beryllium deuterioxide  
D trapping  
Melting

## ABSTRACT

Beryllium oxide (BeO) and deuterioxide (BeO<sub>x</sub>D<sub>y</sub>) have been found on the melted zone of a beryllium tile extracted from the upper dump plate of JET-ILW (2011–2012 campaign). Results have been obtained using Raman microscopy, which is sensitive to both the chemical bond and crystal structure, with a micrometric lateral resolution. BeO is found with a wurtzite crystal structure. BeO<sub>x</sub>D<sub>y</sub> is found as three different types which are not the β-phase but behaves as molecular species like Be(OD)<sub>2</sub>, O(Be-D)<sub>2</sub> and DBEOD. The presence of a small amount of trapped D<sub>2</sub>O is also suspected. Our results therefore strongly suggest that D trapping occurs after melting through the formation of deuterioxides. The temperature increase favors the formation of crystal BeO which favors deuterium trapping through OD bonding.

## 1. Introduction

Plasma facing components (PFCs) are first wall made of tungsten and/or beryllium in nowadays tokamaks and experience plasma-wall interactions [1], such as in the JET tokamak which is equipped with a ITER-like wall (ILW), since 2011 [2]. There, the PFCs are subject to high heat loads and ion irradiation that modify material physical properties, particularly in leading to erosion and deposition processes [3]. Predicting these modifications is of concern for long time operation and safety issues [4]. Some key points under study are: the migration of eroded elements in the tokamak [5,6], the melting of PFC materials [7,8], the production of dust [9–11], the hydrogen isotope retention [12,13], the formation of oxides because of the presence of both impurity and hot surfaces [14], and finally the presence of hydrogen inside oxides [15], with hydroxides formation hypothesized [16].

The Raman microscopy technique is sensitive to the near surface (few tens of nanometers) with a micrometric lateral resolution. It probes atomic vibrations, being thus able to distinguish chemical bonds, and has been used in the fusion community for many years,

especially for carbon material studies [17]. In a series of recent works, we have benchmarked this technique for beryllium materials on laboratory synthesized samples in order to use it for the post-mortem analysis of JET samples [18–20].

We report here on the specific Raman analysis of two tiles which occupied the same place on the JET upper dump plate, but in two different campaigns, one which led to melting zones and a second one without melting zones and we focus on oxide and hydrogen behaviors. This comparative work is complementary to a wider work that includes other post-mortem analyses on these two tiles and that deals with the effect of beryllium oxide on retention in JET-ILW. The paper is organized as follows. Details on the tiles and the campaigns are given together with the Raman technique in Section 2. Results are given in Section 3 and are discussed in Section 4. Finally, a conclusion is given in Section 5.

## 2. Experimental section

A first tile, situated in the upper dump plate of JET-ILW was

\* Corresponding author.

E-mail address: [cedric.pardanaud@univ-amu.fr](mailto:cedric.pardanaud@univ-amu.fr) (C. Pardanaud).

\*\* See the author list of “X. Litaudon et al 2017 Nucl. Fusion 57 102001.

extracted in 2012 from the limiter, after the 2011–2012 campaign. The plasma of that campaign was composed of 6 h of limiter phase and 13 h of divertor phase, with an input energy per tile of 150 GJ [21]. This tile will be referred as “2011–2012 campaign tile” in the following document. A second tile, situated in the same location as the first, was extracted in 2014, from the limiter, after the 2013–2014 campaign. The plasma of that campaign was composed of 5.2 h of limiter phase and 14.2 h of divertor phase, with an input energy per tile of 201 GJ. This tile will be referred as “2013–2014 campaign tile” in the following document. Both tiles were analyzed in between the end of 2017 and the beginning of 2018 by means of Raman microscopy.

A codeposit laboratory sample was produced in the facilities of the National Institute for Laser, Plasma and Radiation Physics (NILPRP) in Romania, using thermionic vacuum (TVA), [22]. Be<sub>2</sub>O and C were deposited in a D<sub>2</sub> atmosphere, to mimick the composition of deposits found on JET divertor tile 1 [23].

Raman spectra were recorded in back-scattering geometry in a Raman mapping mode using a Horiba-Jobin-Yvon HR LabRAM HR800 apparatus (with  $\times 100$  objective, with a numerical aperture of 0.9 and sometimes  $\times 50$  objective with a numerical aperture of 0.5 for mappings, both with  $\lambda_L = 514.5$  nm). The laser power was kept at  $\sim 1$  mW  $\mu\text{m}^{-2}$  to prevent any damage. The grating used (600 lines/mm) leads to a spectral resolution of  $1\text{ cm}^{-1}$ . The lateral definition of the maps was most of the time chosen to be few microns, except for a zone which was investigated in deeper details, with  $0.5\text{ }\mu\text{m}$ .

### 3. Results

Raman spectroscopy was performed on many locations of an upper dump plate extracted from the 2011–2012 and 2013–2014 campaigns. The signal was found to be heterogeneous and Fig. 1 is obtained by averaging on a  $\approx 1000\text{ }\mu\text{m}^2$  zone. This figure therefore displays the comparison between typical averaged Raman spectra of these two upper dump plate tiles and those of some reference samples: defective Be, beryllium deuteride, oxide and carbide. These two latter components could in principle be found in the walls as oxygen and carbon are major pollutants. We first describe the Raman spectroscopy of the reference samples (bullets) and then we make first conclusions for JET samples.

- Crystalline Be spectrum displays a unique Raman active mode close to  $455\text{ cm}^{-1}$ . Introducing defects gives rise to two groups of bands called PDOS and 2PDOS (PDOS stands for phonon density of states), with an empty spectral region between these two groups from  $650$  to  $750\text{ cm}^{-1}$ . These PDOS bands were never found to exceed 25% of the height of the  $455\text{ cm}^{-1}$  band. More details on the Be band can be found in [18], and later on, in Fig. 4.
- Crystalline Be deuteride (BeD<sub>2</sub>) can exist under several forms [24], and we display here the  $\alpha$  phase spectrum, extracted from [20]. Modes are distributed from  $250$  to  $1000\text{ cm}^{-1}$ , then there is a gap from  $1000$  to  $1400\text{ cm}^{-1}$  and again modes from  $1400$ – $1550\text{ cm}^{-1}$  (see [20] for the higher frequency modes at  $1700\text{ cm}^{-1}$ ). The main difference with the other reported deuteride phases, according to calculations, is that, for the other phases, only the low-lying frequency modes exist and there is no vibrational modes at a frequency higher than  $1240\text{ cm}^{-1}$ . Then, the presence of bands falling in the  $1400$ – $1550\text{ cm}^{-1}$  spectral region are a strong indicator of the presence of  $\alpha$ -BeD<sub>2</sub> phase. In Fig. 1, the stars correspond to values found in [25] for  $\alpha$ -BeH<sub>2</sub>, modified to take account of the isotopic shift (modification is done by using multiplying by a scaling factor that takes into account the ratio of the reduced masses of the Be-H and Be-D vibrators).
- Crystalline Be oxide (BeO) can exist in different phases [26], but the most stable is the wurtzite [27] (and references therein). Its spectrum is mainly composed of a strong band at  $\approx 680\text{ cm}^{-1}$ , a medium band at  $723\text{ cm}^{-1}$  and two weak bands at  $337$  and

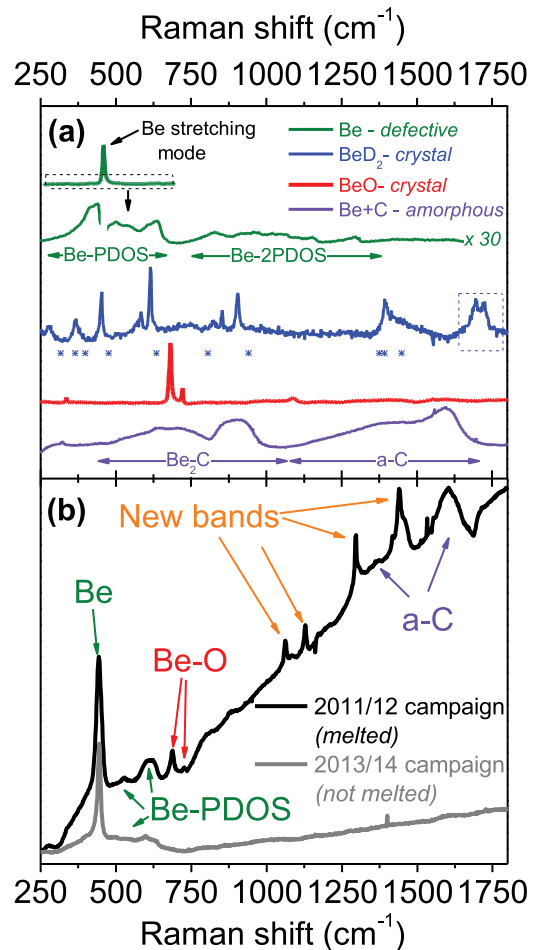
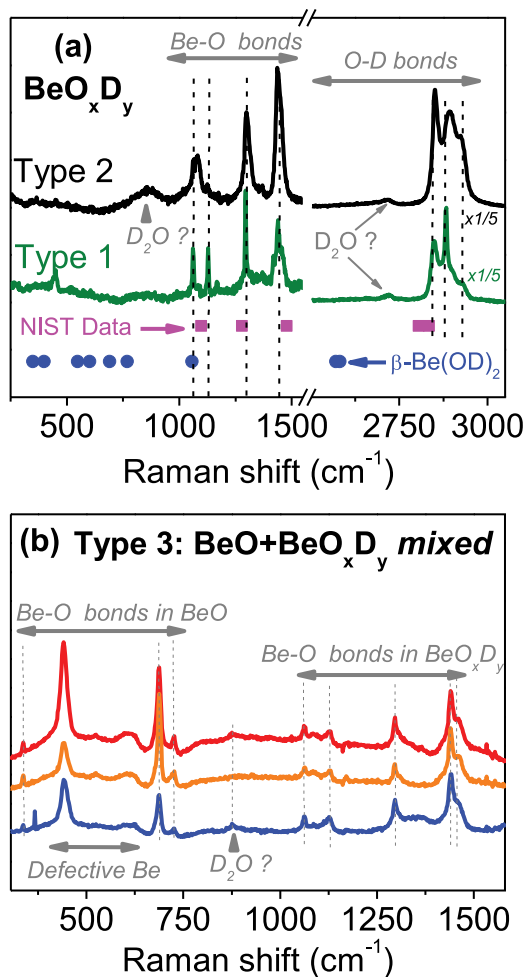


Fig. 1. Comparison of reference sample and JET-ILW sample Raman spectra. (a) Raman spectra from reference samples are defective beryllium (extracted from [18]), beryllium deuteride (extracted from [20]), beryllium oxide and beryllium carbide (extracted from [18]). Stars are from beryllium hydrides data [25] isotopically shifted for the purpose. (b) Raman spectra from two JET samples: 2011–2012 and 2013–2014 campaign tiles.

$1088\text{ cm}^{-1}$  [28]. Recent calculations have predicted the existence of a BeO<sub>2</sub> phase (beryllium peroxide) [29], with vibrational modes falling in the range  $480$ – $530$  and  $680$ – $850\text{ cm}^{-1}$ , plus another degenerate mode at  $1125\text{ cm}^{-1}$ .

- The first Raman spectra of amorphous beryllium carbide were obtained in [18]. Bands fall in the range  $500$ – $1000\text{ cm}^{-1}$ . Be<sub>2</sub>C was found to coexist with amorphous carbon, whose bands fall in the range  $1000$ – $1600\text{ cm}^{-1}$ . Note that hydrogenated amorphous carbon can also give rise to a linear background with a positive slope [30] (and references therein). Linear backgrounds when dealing with carbon materials are most of the time due to some photoluminescence mechanism involving hydrogen atoms in competition with the Raman mechanism.

The 2011–2012 and 2013–2014 campaign tile spectra display the crystalline Be band and they also display the PDOS and 2PDOS bands: this shows that both samples correspond to defective crystalline Be. Spectra display a positive slope which is much higher in the case of the 2011–2012 campaign tile than in the case of the 2013–2014 campaign tile. As mentioned previously, this slope can be associated to the presence of hydrogenated amorphous carbon that is thus most probably present in the 2011–2012 campaign tile while it is less clear for the 2013–2014 campaign tile. Beryllium oxide under the wurtzite form is detected only in the 2011–2012 campaign tile. Again, in the case of the



**Fig. 2.** Raman spectra of the 2011–2012 campaign tile in various zones. (a) Identification of Type 1 and Type 2 deuterioxides. The intensity corresponding to the 2700–3000  $\text{cm}^{-1}$  spectral range has been divided by 5. Blue points are from [31] and magenta points are from [32] (b) Identification of Type 3 deuterioxide, coexisting with the oxide. A background due to carbon contamination was removed in (a) and (b). (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

2011–2012 campaign tile only, four strong new bands, lying at 1062, 1129, 1296 and 1440  $\text{cm}^{-1}$ , are detected. Instinctively, one thinks to attribute these bands to the presence of beryllium deuteride phase. Even if the 1440  $\text{cm}^{-1}$  band falls in the 1400–1500  $\text{cm}^{-1}$  spectral region corresponding to the  $\alpha\text{-BeD}_2$  phase, and could in a first look be attributed to a beryllium deuterated phase, this hypothesis is ruled out because:

- the three other bands fall in the 1000–1400  $\text{cm}^{-1}$  spectral region, where there is no band related to beryllium deuteride phase,
- there are no bands at low frequencies, whereas there should be present for beryllium deuteride phase.

This is then not possible to attribute the bands at 1062, 1129, 1296 and 1440  $\text{cm}^{-1}$  to beryllium deuteride phase. The question of the origin of these new bands remains open.

Another candidate to explain these four bands in the 2011–2012 campaign tile is beryllium deuterioxide. However, the phase diagram of such material is still lacking. Nevertheless, whatever the deuterioxide phase is, there should exist some O–D bond signals that generally falls in the range 2200–2800  $\text{cm}^{-1}$ . Fig. 2a displays Raman spectra from two different spots of the 2011–2012 campaign tile which are spectrally

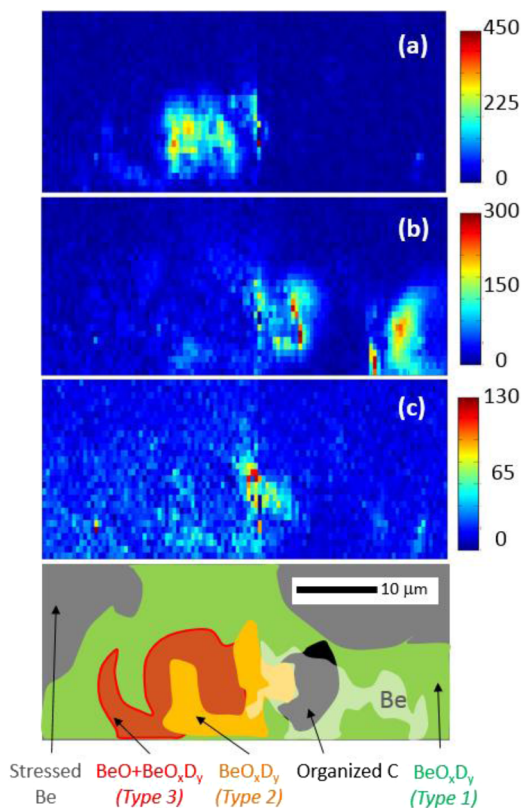
extended up to 3000  $\text{cm}^{-1}$ . Contrary to Fig. 1 for which data were averaged on a large zone, these spectra correspond to homogenous responses of the sample (spot size of typically a few  $\mu\text{m}^2$ ) and two types (named Type 1 and Type 2) have been identified here. For these two types, three very intense overlapping bands at 2850, 2885 and 2928  $\text{cm}^{-1}$  together with a weak asymmetric band at 2720  $\text{cm}^{-1}$  are measured. These band positions are consistent with the presence of O–D bonds. Raman spectra of the most documented deuterioxide phase (the  $\beta\text{-Be(OD)}_2$  phase, called behoite) is composed of weak bands at 349, 400, 549, 602, 690 and 769  $\text{cm}^{-1}$  together with more intense bands at 1056, 2570 and 2582  $\text{cm}^{-1}$  [31]. As there is no agreement between the frequencies of the JET sample spectrum and those of the  $\beta\text{-Be(OD)}_2$  reference spectrum, we conclude that the deuterioxide detected in JET is not the behoite phase. The other reference materials we found in the literature are also displayed in the NIST webbook chemistry database and comes from [32]. Be–O bonds were found at 1098, 1279, and 1477  $\text{cm}^{-1}$ , Be–D bonds were found at 1643 and 1663  $\text{cm}^{-1}$ , and O–D bonds were found at 2805 and 2834  $\text{cm}^{-1}$ . There is a good agreement between the bands displayed in Fig. 2a for the JET samples (Type 1 and Type 2) and for these NIST data. The bands at 1643 and 1663  $\text{cm}^{-1}$  are not observed in JET samples, but it could be explained because the technique used in [32] is infrared spectroscopy, which can change the intensities according to quantum selection rules. It can also be due to the underlying broad band due to hydrogenated amorphous carbon which peaks at 1600  $\text{cm}^{-1}$  and which can hinder the Be–D band detection, as these bands are broad (typically  $\approx 100 \text{ cm}^{-1}$ ) and not well-defined. We thus conclude that these Type 1 and Type 2 spectra most probably correspond to deuterioxides.

Type 1, 2, 3 spectra displayed in Fig. 2 are representative of the three distinct physico-chemical environments found in this investigation. Type 1 and 2 can be observed with or without the presence of Be or wurtzite BeO bands while Type 3 (Fig. 2b) exists only when wurtzite BeO bands (at 680 and 723  $\text{cm}^{-1}$ ) are present. Type 1 and 2 spectra only slightly differ. As for band widths, they are smaller for Type 1 than for Type 2 bands (typically 7  $\text{cm}^{-1}$  and 30  $\text{cm}^{-1}$ , respectively), suggesting a more organized environment in the case of Type 1 than in the case of Type 2. As for band positions:

- at  $\sim 1100 \text{ cm}^{-1}$ , there is a doublet for Type 1 (1060 and 1127  $\text{cm}^{-1}$ ) and a singlet for Type 2 (1077  $\text{cm}^{-1}$ )
- at  $\sim 1440 \text{ cm}^{-1}$ , there is an overlapping triplet for Type 1 (1420, 1437 and 1454  $\text{cm}^{-1}$ ) and a singlet for Type 2 (1438  $\text{cm}^{-1}$ )
- two weak bands are observed at 1167 and 1368  $\text{cm}^{-1}$  for Type 1 and only the latter is observed for Type 2
- the bands at 1295  $\text{cm}^{-1}$  and 2882  $\text{cm}^{-1}$  for Type 1 are upshifted at 1299  $\text{cm}^{-1}$  and 2893  $\text{cm}^{-1}$  for Type 2, respectively
- there is a broad band at 855  $\text{cm}^{-1}$  which is more intense for Type 2 than for Type 1

As for the intensities of the main bands in the Be–O (O–D) region, the intensity at 1295 (2882)  $\text{cm}^{-1}$  is larger than that at 1437 (2848)  $\text{cm}^{-1}$  for Type 1 ( $I_{1295}/I_{1437}$  is  $\sim 1.35$ ) while it is the contrary for Type 2 ( $I_{1295}/I_{1437}$  is  $\sim 0.71$ ). The relative ratios  $I_{2882}/I_{1295}$  and  $I_{2848}/I_{1437}$  stay roughly uniform over the sample zone investigated, meaning that they correspond to the same species. Therefore, even though Type 1 and Type 2 are probably not composed with a single deuterioxide species, we can characterize Type 1 and Type 2 with these two main couples of bands, (1295 and 2882  $\text{cm}^{-1}$ ) and (1437 and 2848  $\text{cm}^{-1}$ ), respectively.

Type 3 spectra are very similar to those of Type 1 (see Fig. 2b with three examples of Type 3 spectra are displayed, and a Type 1 spectrum is reported from Fig. 2a in green). The two bands at 680 and 723  $\text{cm}^{-1}$  from the wurtzite BeO are always observed in the case of Type 3 spectra, with various relative intensities, the lowest value of  $I_{723}/I_{680}$  being  $\sim 0.35$ . We also report a weak and broad band at 878  $\text{cm}^{-1}$  which origin could be similar to that at 855  $\text{cm}^{-1}$  in the case of Type 1

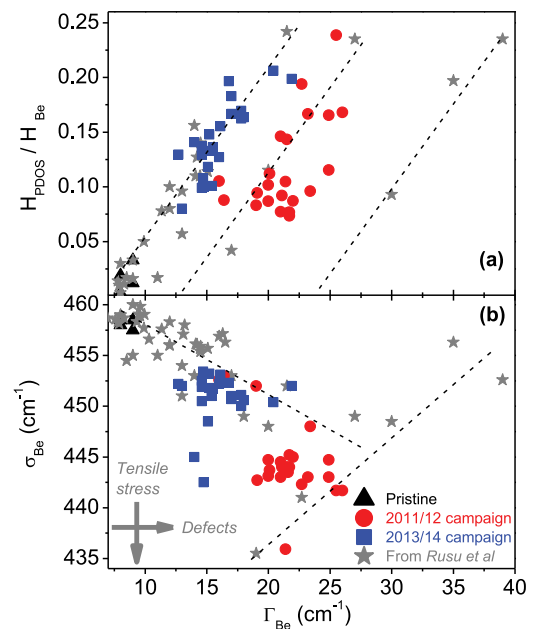


**Fig. 3.** Raman mapping of a zone of the 2011–2012 campaign tile. (a) BeO band intensity. (b) BeO<sub>x</sub>D<sub>y</sub> Type 1 band intensity. (c) BeO<sub>x</sub>D<sub>y</sub> Type 2 band intensity. (d) Spatial repartition of the different phases. This cartoon also takes into account other maps of the band positions and intensities (not displayed here). Note that amorphous carbon is found everywhere.

and 2. Remarkably, these three types of BeO<sub>x</sub>D<sub>y</sub> are found only in melted zones of the 2011–2012 campaign tile and not in the 2013–2014 campaign tile which was not melted, nor in the non-melted zones of the 2011–2012 campaign tile.

Maps of the intensities and frequencies of the most relevant Raman bands have been recorded for a melted zone of the 2011–2012 campaign tile and Fig. 3abc displays three of these maps while Fig. 3d displays a scheme of the spatial distribution of the various phases, Be, BeO, BeO<sub>x</sub>D<sub>y</sub> and C, that could be deduced from the ensemble of maps. The spatial distribution of BeO has been obtained by displaying the 680 cm<sup>-1</sup> band intensity (Fig. 3a) which correlates very well with the frequency map (not shown). The spatial distribution of BeO<sub>x</sub>D<sub>y</sub> Type 1 and 2 have been obtained by displaying the 1295 and 855 cm<sup>-1</sup> band intensities, (Fig. 3b and c, respectively). The corresponding frequency maps (not shown) correlates only partially, due to complex radiative transfer processes at play in heterogeneous media. The main points are listed below:

- two types of beryllium structures are found: stressed and not stressed. The band frequency of the stressed Be is ~ 10 cm<sup>-1</sup> downshifted compared to normal Be.
- BeO was found at the lower middle part of the investigated zone, in a ≈ 10 μm nucleus
- BeO<sub>x</sub>D<sub>y</sub> Type 3 is found mixed with BeO
- BeO<sub>x</sub>D<sub>y</sub> Type 2 is found only close to BeO and interfacing with Type 1
- BeO<sub>x</sub>D<sub>y</sub> Type 1 is the main phase found in the zone, outside the Type 3 and Type 2 locations, and with the maximum intensity in the lower left corner of the zone (Fig. 3b).
- hydrogenated amorphous carbon is found everywhere (easily



**Fig. 4.** Raman parameter plots characterizing the Be phase. (a)  $H_{\text{PDOS}}/H_{\text{Be}}$  as a function of  $\Gamma_{\text{Be}}$ . (b)  $\sigma_{\text{Be}}$  as a function of  $\Gamma_{\text{Be}}$ . Data in grey and black from [18] have been added for comparison. Lines are guides for the eyes. (For interpretation of the references to colour in this figure legend, the reader is referred to the web version of this article.)

observed because of its high Raman cross-section)

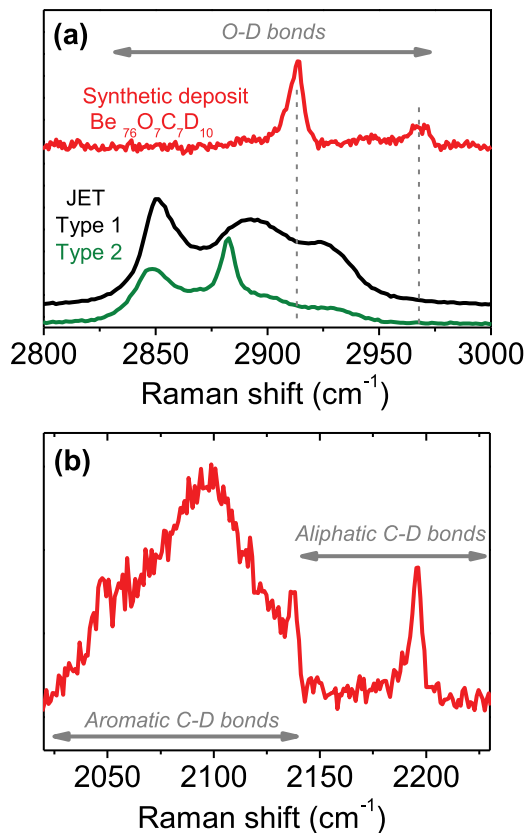
- a zone of organized carbon is found at the left of the BeO nucleus

This Raman map of the various phases suggests that BeO<sub>x</sub>D<sub>y</sub> Types 1, 2 and 3 are roughly concentrically distributed according to the nucleus containing BeO. More precisely, Type 3 is mixed in the BeO nucleus, Type 2 is found at the neighboring of the BeO nucleus whereas Type 1 is found elsewhere.

Fig. 4 focuses on the Be phases and displays spectroscopic parameters of the 455 cm<sup>-1</sup> Be band for spectra recorded at various locations separated by millimeter distances on the 2011–2012 campaign tile (in red) and the 2013–2014 campaign tile (in blue). The spectroscopic parameters are the relative band height ratio formed with the PDOS related band height and the Be Raman active mode related band height:  $H_{\text{PDOS}}/H_{\text{Be}}$ , the band position related to the Be Raman active mode:  $\sigma_{\text{Be}}$ , and its full width at half maximum:  $\Gamma_{\text{Be}}$ .  $H_{\text{PDOS}}/H_{\text{Be}}$  (Fig. 4a) and  $\sigma_{\text{Be}}$  (Fig. 4b) are plotted in function of  $\Gamma_{\text{Be}}$ . It was shown that these parameters can be correlated and reflect the existence of defects and tensile stress in the material. Previous data from [18] obtained for various implanted and deposited mixed Be based samples synthesized in laboratory are displayed for comparison (in grey, see also Figure 7 of [18]). Crystalline Be corresponds to  $H_{\text{PDOS}}/H_{\text{Be}} \approx 0$ ,  $\sigma_{\text{Be}} \approx 457$  cm<sup>-1</sup> and  $\Gamma_{\text{Be}} \approx 9$  cm<sup>-1</sup> while averaged parameters for the 2013–2014 campaign tile (the 2011–2012 campaign tile) are  $H_{\text{PDOS}}/H_{\text{Be}} \approx 0.14$  (0.16),  $\sigma_{\text{Be}} \approx 451$  (445) cm<sup>-1</sup> and  $\Gamma_{\text{Be}} \approx 16$  (22) cm<sup>-1</sup>. Fig. 4 shows that data from both JET and laboratory samples are not randomly spread and that they aligned similarly. JET results can thus be interpreted in the light of this previous analysis and we can argue that beryllium is more tensile stressed in the 2011–2012 campaign tile than in the 2013–2014 campaign tile. In particular, red points (2011–2012 campaign tile) are clearly at lower frequencies and higher widths than blue points (2013–2014 campaign tile), and they compare well with stressed beryllium carbide points measured in [18].

Fig. 5 displays the Raman spectrum of a laboratory sample that was synthesized by codepositing Be<sub>2</sub>O and C in a D<sub>2</sub> atmosphere, to mimicking the composition of deposits found on JET divertor tile 1. A photoluminescence background, with an origin that has to be





**Fig. 5.** Comparison of a sample mimicking JET tile 1 divertor deposit (non-melted) and the 2011–2012 campaign tile (melted). (a) O–D spectral region. (b) C–D spectral region. Note that a very intense background, not yet understood, has been subtracted from the synthetic deposit spectrum.

understood, has been subtracted. Bands are found at 2047, 2098, 2137, 2196, 2913 and 2968  $\text{cm}^{-1}$ . They are due to aromatic C–D bonds for the 2047, 2098 and 2137  $\text{cm}^{-1}$  bands, aliphatic C–D bonds for the 2196  $\text{cm}^{-1}$  band and Be–O–D bonds for the 2913 and 2968  $\text{cm}^{-1}$  bands. This last band may also be due to  $\text{D}_2$  trapped in vacancies, but it needs more investigations [15]. D was then found to be bonded with both O and C but we did not detect Be–D bonds.

#### 4. Discussion

Deuterioxides are found by detecting Be–O and O–D vibrational modes in melted zones of 2011–2012 campaign JET tile. These modes are detected neither in non-melted zones of the 2011–2012 campaign tile nor in the 2013–2014 campaign tile. We have identified three types of deuterioxides that seem approximately concentrically organized around a  $\approx 10 \mu\text{m}$  nucleus containing BeO. Type 3 is found only mixed with BeO in the nucleus while Type 2 is found at its surrounding and Type 1, which actually covers the largest area in Fig. 3d, is found in between Type 2 and a tensile stressed Be phase. We have observed that the couple of bands at 1295 and 2882  $\text{cm}^{-1}$  is representative of Type 1 whereas the couple of bands at 1437 and 2848  $\text{cm}^{-1}$  is representative of Type 2. In [32] deuterioxides were obtained by laser ablating a beryllium target in deuterated water and collecting the molecules formed in a cryogenic matrix and three molecules were obtained:  $\text{D}^{\ominus}\text{Be-O-Be-D}$ ,  $\text{Be(OD)}_2$  and  $\text{D-Be-OD}$ . According to band frequencies collected in [30], our Type 1 phase may be composed of D bonded mainly like in  $\text{D-Be-O-Be-D}$  and  $\text{D-Be-OD}$  whereas the Type 2 phase may be composed of D bonded mainly like in  $\text{Be(OD)}_2$  bonds. Type 3 is similar to Type 1, but with the presence of BeO wurtzite nanocrystallites in its vicinity (let us recall that the lateral resolution in

Raman is micrometric). D is then trapped in various physico-chemical environments. The stoichiometry of Type 1 may be estimated in the range  $\text{Be}_{1.2}\text{O}_1\text{D}_2$ , for Type 1 and  $\text{BeO}_2\text{D}_2$  for Type 2. Remarkably, beryllium hydride was not detected. Note that its low Raman cross section however makes its detection difficult.

The nucleus containing BeO schematized in Fig. 4 is about  $10 \mu\text{m}$  wide and TEM measurements show that it extends down to  $2 \mu\text{m}$  in depth. This corresponds to a thickness much larger than that of native oxide layer found at the surface of Be samples due to air contamination. A thin layer would actually not lead to intense signal. The presence of such a nucleus probably reveals that, when melting occurs, there is a massive diffusion and concentration of the oxygen deeply trapped in the bulk of the pristine Be material which leads to the formation of beryllium oxide nuclei and also, thanks to the concomitant ion bombardment, to the formation of deuterioxides. Note that the source of oxygen contamination that leads to these oxides and hydroxides cannot be inferred to sample air exposure after being extracted from the tokamak but comes from residual vacuum and leaks. The zone analyzed in Fig. 3, for example, is  $\approx 2 \mu\text{m}$  thick, far from the typical nanometer scale contamination that can occur after venting.

In both Type 1 and 2, a weak asymmetric band at 2720  $\text{cm}^{-1}$  is recorded, which might be attributed to trapped  $\text{D}_2\text{O}$  molecules. The vibrational frequencies of trapped  $\text{D}_2\text{O}$  depend strongly on the host material because of possible strong non covalent interactions:  $\nu_1$  and  $\nu_3$  modes of  $\text{D}_2\text{O}$  trapped in inert matrices are at 2655 and 2765  $\text{cm}^{-1}$ , respectively [33], whereas they lie in the range 2200–2679  $\text{cm}^{-1}$  for crystalline deuterated hydrates [34]. The asymmetry of the band at 2720  $\text{cm}^{-1}$  could thus be due to the overlapping of modes  $\nu_1$  and  $\nu_3$ . Trapped  $\text{H}_2\text{O}$  molecules display a broad band at about 1200  $\text{cm}^{-1}$  [35], which, taking into account an isotopic shift, would correspond to a frequency of 870  $\text{cm}^{-1}$  for trapped  $\text{D}_2\text{O}$ . This value is close to the weak and broad bands observed at 855  $\text{cm}^{-1}$  in Type 2 and 878  $\text{cm}^{-1}$  in Type 3. This analysis shows that a small amount of  $\text{D}_2\text{O}$  may be trapped together with deuterioxides.

Pletnev et al found in [36] that behoite  $\beta\text{-Be(OH)}_2$  incompletely dehydrates when heated to 770 K, forming water and two phases: a first phase with BeO coordinated with OH groups and a second phase of amorphous oxyhydroxide, vanishing at 1470 K. In our study, the sample was melted, meaning the temperature rose above 1470 K. Our observation of the various deuterioxide types with the presence of beryllium oxide, together with the absence of the behoite phase, is thus consistent with these results from [36]. Note that the temperature at which the behoite phase decomposes is close to the 680 K thermal desorption peak of  $\text{D}_2$  reported in [16].

The presence of tensile stress in the beryllium phase neighboring Type 1 deuterioxide is attested in Figs. 3d and 4 by the downshift of the 455  $\text{cm}^{-1}$  Be band. There are strong similarities in the Be band frequency versus width behavior with that of reference sample of  $\text{Be}_2\text{C}$  [18] where Be neighboring Be carbide was stressed. These results then suggest that here Be is stressed by the surrounding hydroxide phase or hydrogenated amorphous carbon phase present almost everywhere, stress being induced by the volume occupied by species with CD or OD bonds.

Analyzing the 3D structure is necessary to get a deeper insight on these processes and Raman information should be crossed with the EELS (Energy Electron Loss Spectroscopy) analysis performed on a TEM (Transmission Electron Microscopy) cross section of a FIB (Focused Ion Beam) slice cut in this zone. From the absolute Raman band intensities, we cannot yet derive quantitative data because of the 3D structure of the material which affects light propagation and because the Raman cross sections for the involved compounds are not yet all available.

#### 5. Conclusion

We have analyzed with a micrometric lateral resolution two tiles occupying the same place in the limiter upper dump plate of the JET-

ILW tokamak during the two campaigns, 2011–12 and 2013–14, using Raman mapping. In particular, we have analyzed melted zones of the 2011–12 campaign tile where we have detected the wurtzite beryllium oxide BeO phase and various species attributed to deuterioxides BeO<sub>x</sub>D<sub>y</sub>. Three types of deuterioxides were observed, corresponding to different stoichiometries, none of them being the well-documented crystalline phase, the behoite β-phase. This is in agreement with what is known about the thermal stability of this phase. The presence of a small amount of D<sub>2</sub>O is also suspected. On the contrary, no oxides nor deuterioxides were detected in the non-melted 2013–2014 campaign tile. Our results therefore suggest that D trapping occurs after melting through the formation of deuterioxides. The temperature increase favors the migration of oxygen impurities which favors deuterium trapping through OD bonding. For the laboratory sample mimicking the composition of deposits found on JET divertor tile 1, D was found to bond with both O and C but we did not detect Be-D bonds. This confirms that the formation of deuterioxides is probably an efficient way of trapping D. Note that it can also affect droplet formation and composition. To check the importance of melting in these processes, we have to compare our results with additional data, especially from deposited zones in JET. Impact on retention is not in the scope of this paper as Raman cross section of these compounds are not yet known.

### Conflict of interest

The authors declare no conflict of interest.

### Acknowledgements

This work has been carried out within the framework of the EUROfusion Consortium and has received funding from the Euratom research and training programme 2014–2018 under grant agreement No 633053. The views and opinions expressed herein do not necessarily reflect those of the European Commission.

### Bibliography

- [1] S. Brezinsek, J.W. Coenen, T. Schwarz-Selinger, K. Schmid, A. Kirschner, A. Hakola, F.L. Tabares, H.J. van der Meiden, M.L. Mayoral, M. Reinhart, E. Tsitrone, T. Ahlgren, M. Aints, M. Airila, S. Almaviva, E. Alves, T. Angot, V. Anita, R.A. Parra, F. Aumayr, M. Balden, J. Bauer, M. Ben Yaala, B.M. Berger, R. Bisson, C. Bjorkas, I.B. Radovic, D. Borodin, J. Bucalossi, J. Butikova, B. Butoi, I. Cadez, R. Caniello, L. Caneve, G. Cartry, N. Catarino, M. Cekada, G. Ciraolo, L. Ciupinski, F. Colao, Y. Corre, C. Costin, T. Craciunescu, A. Cremona, M. De Angeli, A. de Castro, R. Dejarnac, D. Dellasega, P. Dinca, T. Dittmar, C. Dobrea, P. Hansen, A. Drenik, T. Eich, S. Elgeti, D. Falie, N. Fedorczak, Y. Ferro, T. Fornal, E. Fortuna-Zalesna, L. Gao, P. Gasior, M. Gherendi, F. Ghezzi, Z. Gosar, H. Greuner, E. Grigore, C. Grisolia, M. Groth, M. Gruca, J. Grzonka, J.P. Gunn, K. Hassouni, K. Heinola, T. Eich, S. Huber, W. Jacob, I. Jezu, X. Jiang, I. Jogi, A. Kaiser, J. Karhunen, M. Kelemen, M. Koppen, H.R. Koslowski, A. Kreter, M. Kubkowska, M. Laan, L. Laguardia, A. Lahtinen, A. Lasa, V. Lasic, N. Lemahieu, J. Likonen, J. Linke, A. Litnovsky, C. Linsmeier, T. Loewenhoff, C. Lungu, M. Lungu, G. Maddaluno, H. Maier, T. Makkonen, A. Manhard, Y. Marandet, S. Markelj, L. Marot, C. Martin, A.B. Martin-Rojo, Y. Martynova, R. Mateus, D. Matveev, M. Mayer, G. Meisl, N. Mellet, A. Michau, J. Miettinen, S. Moller, T.W. Morgan, J. Mougenot, M. Mozetic, V. Nemanic, R. Neu, K. Nordlund, M. Oberkofler, E. Oyarzabal, M. Panjan, C. Pardanaud, P. Paris, M. Passoni, B. Pegourie, P. Pelicon, P. Petersson, K. Piip, G. Pintsuk, G.O. Pompilian, G. Popa, C. Porosnicu, G. Primec, M. Probst, J. Raisanen, M. Rasinski, S. Ratynskaia, D. Reiser, D. Ricci, M. Richou, J. Riesch, G. Riva, M. Rosinski, P. Roubin, M. Rubel, C. Ruset, E. Safi, G. Sergienko, Z. Siketic, A. Sima, B. Spilker, R. Stadlmayr, I. Stuedel, P. Strom, T. Tadic, D. Tafalla, I. Tale, D. Terentyev, A. Terra, V. Tiron, I. Tiseanu, P. Toliaas, D. Tskhakaya, A. Uccello, B. Unterberg, I. Uytendhove, E. Vassallo, P. Vavpetic, P. Veis, L.L. Velicu, J.W.M. Vernimmen, A. Voitikans, U. von Toussaint, A. Weckmann, M. Wirtz, A. Zoloznik, R. Zaplotnik, W.P. Contributors, Plasma-wall interaction studies within the EUROfusion consortium: progress on plasma-facing components development and qualification, *Nucl. Fusion* 57 (2017).
- [2] G.F. Matthews, M. Beurskens, S. Brezinsek, M. Groth, E. Joffrin, A. Loving, M. Kear, M.L. Mayoral, R. Neu, P. Prior, V. Riccardo, F. Rimini, M. Rubel, G. Sips, E. Villedieu, P. de Vries, M.L. Watkins, E.-J. Contributors, JET ITER-like wall-overview and experimental programme, *Phys. Scr.* T145 (2011).
- [3] M. Mayer, S. Krat, A. Baron-Wiechec, Y. Gasparyan, K. Heinola, S. Koivuranta, J. Likonen, C. Ruset, G. de Saint-Aubin, A. Widdowson, J.E.T. Contributors, Erosion and deposition in the JET divertor during the second ITER-like wall campaign, *Phys. Scr.* T170 (2017).
- [4] R. Neu, A.U. Team, E.P. Taskforce, J.E. Contributors, Preparing the scientific basis for an all metal ITER, *Plasma Phys. Control. Fusion* 53 (2011).
- [5] S. Brezinsek, A. Widdowson, M. Mayer, V. Philipps, P. Baron-Wiechec, J.W. Coenen, K. Heinola, A. Huber, J. Likonen, P. Petersson, M. Rubel, M.F. Stamp, D. Borodin, J.P. Coad, A.G. Carrasco, A. Kirschner, S. Krat, K. Krieger, B. Lipschultz, C. Linsmeier, G.F. Matthews, K. Schmid, J.E.T. Contributors, Beryllium migration in JET ITER-like wall plasmas, *Nucl. Fusion* 55 (2015).
- [6] I. Bykov, H. Bergsaker, G. Possnert, Y. Zhou, K. Heinola, J. Pettersson, S. Conroy, J. Likonen, P. Petersson, A. Widdowson, J.E.T. Contributors, Studies of Be migration in the JET tokamak using AMS with Be-10 marker, *Nucl. Instrum. Methods Phys. Res. Sect. B-Beam Interact. Mater. Atoms* 371 (2016) 370–375.
- [7] K. Krieger, S. Brezinsek, M. Reinelt, S.W. Lisgo, J.W. Coenen, S. Jachmich, S. Marsen, A. Meigs, G. van Rooij, M. Stamp, O. van Hoey, M.F. Ivanova, T. Loarer, V. Philipps, J.E.T.E.F.D.A. Contributors, Beryllium migration and evolution of first wall surface composition in the JET ILW configuration, *J. Nucl. Mater.* 438 (2013) S262–S266.
- [8] K. Schmid, M. Reinelt, K. Krieger, An integrated model of impurity migration and wall composition dynamics for tokamaks, *J. Nucl. Mater.* 415 (2011) S284–S288.
- [9] E. Fortuna-Zalesna, J. Grzonka, S. Moon, M. Rubel, P. Petersson, A. Widdowson, J.E.T. Contributors, Fine metal dust particles on the wall probes from JET-ILW, *Phys. Scr.* T170 (2017).
- [10] E. Fortuna-Zalesna, J. Grzonka, M. Rubel, A. Garcia-Carrasco, A. Widdowson, A. Baron-Wiechec, L. Ciupinski, J.E.T. Contributors, Studies of dust from JET with the ITER-Like Wall: Composition and internal structure, *Nucl. Mater. Energy* 12 (2017) 582–587.
- [11] S. Masuzakii, M. Tokitani, T. Otsuka, Y. Oya, Y. Hatan, M. Miyamoto, R. Sakamoto, N. Ashikawa, S. Sakurada, Y. Uemura, K. Azuma, K. Yumizurus, M. Oyaizu, T. Suzuki, H. Kurotaki, D. Hamaguchi, K. Isobel, N. Asakura, A. Widdowson, K. Heinola, S. Jachmich, M. Rubel, J.E.T. contributors, Analyses of microstructure, composition and retention of hydrogen isotopes in divertor tiles of JET with the ITER-like wall, *Phys. Scr.* T170 (2017).
- [12] K. Heinola, A. Widdowson, J. Likonen, E. Alves, A. Baron-Wiechec, N. Barradas, S. Brezinsek, N. Catarino, P. Coad, S. Koivuranta, Long-term fuel retention in JET ITER-like wall, *Phys. Scr.* T167 (2016) 014075.
- [13] J. Likonen, K. Heinola, A. De Backer, S. Koivuranta, A. Hakola, C.F. Ayres, A. Baron-Wiechec, P. Coad, G.F. Matthews, M. Mayer, A. Widdowson, J. Contributors, Deuterium trapping and release in JET ITER-like wall divertor tiles, *Phys. Scr.* T167 (2016) 014074.
- [14] E.A. Hodille, F. Ghorghiu, Y. Addab, A. Zoloznik, M. Minissale, Z. Piazza, C. Martin, T. Angot, L. Gallais, M.F. Barthe, C.S. Becquart, S. Markelj, J. Mougenot, C. Grisolia, R. Bisson, Retention and release of hydrogen isotopes in tungsten plasma-facing components: the role of grain boundaries and the native oxide layer from a joint experiment-simulation integrated approach, *Nucl. Fusion* 57 (2017).
- [15] E.A. Hodille, Y. Ferro, Z.A. Piazza, C. Pardanaud, Hydrogen in beryllium oxide investigated by DFT: on the relative stability of charged-state atomic versus molecular hydrogen, *J. Phys. Condens. Matter* 30 (2018).
- [16] M. Reinelt, C. Linsmeier, Ion implanted deuterium retention and release from clean and oxidized beryllium, *J. Nucl. Mater.* (2009) 568–571 390-91.
- [17] C. pardanaud, C. Martin, P. Roubin, Raman Microscopy: A Suitable Tool for Characterizing Surfaces in Interaction with Plasmas in the Field of Nuclear Fusion, in: K. Maaz (Ed.), *Raman Spectroscopy and Applications*, InTech, 2017, pp. 3–30.
- [18] M.I. Rusu, C. Pardanaud, Y. Ferro, G. Giacometti, C. Martin, Y. Addab, P. Roubin, M. Minissale, L. Ferri, F. Virot, M. Barrachin, C.P. Lungu, C. Porosnicu, P. Dinca, M. Lungu, M. Koppen, P. Hansen, C. Linsmeier, Preparing the future post-mortem analysis of beryllium-based JET and ITER samples by multi-wavelengths Raman spectroscopy on implanted Be, and co-deposited Be, *Nucl. Fusion* (2017) 57.
- [19] C. Pardanaud, M.I. Rusu, G. Giacometti, C. Martin, Y. Addab, P. Roubin, C.P. Lungu, C. Porosnicu, I. Jezu, P. Dinca, M. Lungu, O.G. Pompilian, R. Mateus, E. Alves, M. Rubel, J. contributors, Raman microscopy investigation of beryllium materials, *Phys. Scr.* T167 (2016) 014027.
- [20] C. Pardanaud, M.I. Rusu, C. Martin, G. Giacometti, P. Roubin, Y. Ferro, A. Allouche, M. Oberkofler, M. Köppen, T. Dittmar, C. Linsmeier, Hydrogen retention in beryllium: concentration effect on crystalline growth, *J. Phys. Condens. Matter* 27 (2015) 475401.
- [21] K. Heinola, A. Widdowson, J. Likonen, T. Ahlgren, E. Alves, C.F. Ayres, A. Baron-Wiechec, N. Barradas, S. Brezinsek, N. Catarino, P. Coad, C. Guillemat, I. Jezu, S. Krat, A. Lahtinen, G.F. Matthews, M. Mayer, J.E.T. Contributors, Experience on divertor fuel retention after two ITER-Like Wall campaigns, *Phys. Scr.* T170 (2017).
- [22] C.P. Lungu, C. Porosnicu, I. Jezu, M. Lungu, A. Marcu, C. Luculescu, C. Ticos, A. Marin, C.E.A. Grigorescu, The behavior of W, Be and C layers in interaction with plasma produced by terawatt laser beam pulses, *Vacuum* 110 (2014) 207–212.
- [23] A. Hakola, K. Heinola, J. Likonen, C. Lungu, C. Porosnicu, E. Alves, R. Mateus, I. Bogdanovic Radovic, Z. Siketic, V. Nemanic, C. Pardanaud, a.E.W.P. Contributors, Production of JET-ILW relevant Be-containing laboratory samples for fuel retention investigations, *Nucl. Mater. Energy* (2018) This volume.
- [24] X.Y. An, T.X. Zeng, W.Y. Ren, Structural, electronic and optical properties of BeH<sub>2</sub>: A density functional theory study, *Mater. Res. Express* 4 (2017).
- [25] C. Pépin, P. Loubeyre, Layered structure and re-entrant disproportionation observed in crystalline BeH<sub>2</sub> under pressure, *Phys. Rev. B* 93 (2016) 224104.
- [26] U.D. Wdowik, Structural stability and thermal properties of BeO from the quasi-harmonic approximation, *J. Phys. Condens. Matter* 22 (2010).
- [27] F. Luo, Z.C. Guo, X.L. Zhang, C.Y. Yuan, C.A. Liu, L.C. Cai, Pressure effects on structural and elastic properties of BeO from first-principles calculations, *Phys. Status Solidi B-Basic Solid State Phys.* 252 (2015) 212–218.
- [28] A. Bosak, K. Schmalzl, M. Krisch, W. van Beek, V. Kolobanov, Lattice dynamics of

- beryllium oxide: inelastic x-ray scattering and ab initio calculations, *Phys. Rev. B* 77 (2008).
- [29] S.T. Zhang, F. Li, H.Y. Xu, G.C. Yang, Pressure-induced stable beryllium peroxide, *Inorg. Chem.* 56 (2017) 5233–5238.
- [30] C. Pardanaud, C. Martin, P. Roubin, G. Giacometti, C. Hopf, T. Schwarz-Selinger, W. Jacob, Raman spectroscopy investigation of the H content of heated hard amorphous carbon layers, *Diam. Relat. Mater.* 34 (2013) 100–104.
- [31] H.D. Lutz, C. Jung, R. Mortel, H. Jacobs, R. Stahl, Hydrogen bonding in solid hydroxides with strongly polarising metal ions, beta-Be(OH)(2) and epsilon-Zn(OH)(2), *Spectrochim. Acta Part a-Mol. Biomol. Spectrosc.* 54 (1998) 893–901.
- [32] C.A. Thompson, L. Andrews, Reactions of laser ablated Be atoms with H<sub>2</sub>O: Infrared spectra and density functional calculations of HOBeOH, HBeOH, and HBeOBeH, *J. Phys. Chem.* 100 (1996) 12214–12221.
- [33] S. Coussan, P. Roubin, J.P. Perchard, Infrared induced isomerizations of water polymers trapped in nitrogen matrix, *Chem. Phys.* 324 (2006) 527–540.
- [34] W. Mikenda, S. Steinbock, Stretching frequency versus bond distance correlation of O-D(H)...F hydrogen-bonds in solid hydrates, *J. Mol. Struct.* 326 (1994) 123–130.
- [35] W.W. Rudolph, D. Fischer, G. Irmer, C.C. Pye, Hydration of beryllium(II) in aqueous solutions of common inorganic salts. a combined vibrational spectroscopic and ab initio molecular orbital study, *Dalton Trans.* (2009) 6513–6527.
- [36] R.N. Pletnev, V.S. Kiiko, Y.N. Makurin, A.A. Nepryakhin, Proton magnetic resonance and the state of hydrogen in beryllium hydroxide, *Refract. Ind. Ceram.* 46 (2005) 273–275.