

THE NORM OF TIME-FREQUENCY AND WAVELET LOCALIZATION OPERATORS

Original

THE NORM OF TIME-FREQUENCY AND WAVELET LOCALIZATION OPERATORS / Nicola, F.; Tilli, P.. - In: TRANSACTIONS OF THE AMERICAN MATHEMATICAL SOCIETY. - ISSN 0002-9947. - STAMPA. - 376:10(2023), pp. 7353-7375. [10.1090/tran/8990]

Availability:

This version is available at: 11583/2982558 since: 2023-09-28T14:55:42Z

Publisher:

American Mathematical Society

Published

DOI:10.1090/tran/8990

Terms of use:

This article is made available under terms and conditions as specified in the corresponding bibliographic description in the repository

Publisher copyright

(Article begins on next page)

THE NORM OF TIME-FREQUENCY AND WAVELET LOCALIZATION OPERATORS

FABIO NICOLA AND PAOLO TILLI

ABSTRACT. Time-frequency localization operators (with Gaussian window) $L_F : L^2(\mathbb{R}^d) \rightarrow L^2(\mathbb{R}^d)$, where F is a weight in \mathbb{R}^{2d} , were introduced in signal processing by I. Daubechies in 1988, inaugurating a new, geometric, phase-space perspective. Sharp upper bounds for the norm (and the singular values) of such operators turn out to be a challenging issue with deep applications in signal recovery, quantum physics and the study of uncertainty principles.

In this note we provide optimal upper bounds for the operator norm $\|L_F\|_{L^2 \rightarrow L^2}$, assuming $F \in L^p(\mathbb{R}^{2d})$, $1 < p < \infty$ or $F \in L^p(\mathbb{R}^{2d}) \cap L^\infty(\mathbb{R}^{2d})$, $1 \leq p < \infty$. It turns out that two regimes arise, depending on whether the quantity $\|F\|_{L^p}/\|F\|_{L^\infty}$ is less or greater than a certain critical value. In the first regime the extremal weights F , for which equality occurs in the estimates, are certain Gaussians, whereas in the second regime they are proved to be Gaussians *truncated above*, degenerating into a multiple of a characteristic function of a ball for $p = 1$. This phase transition through Gaussians truncated above appears to be a new phenomenon in time-frequency concentration problems. For the analogous problem for wavelet localization operators -where the Cauchy wavelet plays the role of the above Gaussian window- a complete solution is also provided.

1. INTRODUCTION

The uncertainty principle is an ubiquitous theme in mathematics and represents an endless source of challenging and inspirational problems. The literature in this connection is enormous. For a deep introduction to the topic, we address the reader to the classical contributions [18, 25, 28] and the recent account [55].

Roughly speaking, the uncertainty principle states that a function and its Fourier transform cannot both be too concentrated, or equivalently that a time-frequency distribution cannot be too concentrated in the time-frequency space. To better develop the latter point of view, we recall the definition of a time-frequency distribution known in harmonic analysis and signal processing as short-time, or windowed, Fourier transform [23, 31, 47, 48], and in mathematical physics as coherent state transform [43].

2010 *Mathematics Subject Classification.* 42B10, 49Q20, 49R05, 81S30, 94A12.

Key words and phrases. Short-time Fourier transform, wavelet transform, localization operator, Toeplitz operator, uncertainty principle.

Let φ be the ‘‘Gaussian window’’

$$(1.1) \quad \varphi(x) = 2^{d/4} e^{-\pi|x|^2}, \quad x \in \mathbb{R}^d,$$

normalized in such way that $\|\varphi\|_{L^2} = 1$. The short-time Fourier transform with Gaussian window of a function $f \in L^2(\mathbb{R}^d)$, is defined as

$$\mathcal{V}f(x, \omega) = \int_{\mathbb{R}^d} e^{-2\pi iy \cdot \omega} f(y) \varphi(x - y) dy, \quad x, \omega \in \mathbb{R}^d.$$

It turns out that $\mathcal{V} : L^2(\mathbb{R}^d) \rightarrow L^2(\mathbb{R}^{2d})$ is an isometry, so that the function $|\mathcal{V}f(x, \omega)|^2$ can be interpreted as the time-frequency energy density of f . Consequently, the integral $\int_{\Omega} |\mathcal{V}f(x, \omega)|^2 dx d\omega$, over a measurable subset $\Omega \subset \mathbb{R}^{2d}$, will be the fraction of its energy trapped in Ω (see e.g. [1, 5, 6, 27, 49, 53, 50]). Sharp upper bounds for this quantity were recently proved in [50]; the relevant estimate, that we state here in dimension $d = 1$ for simplicity, reads

$$(1.2) \quad \int_{\Omega} |\mathcal{V}f(x, \omega)|^2 dx d\omega \leq (1 - e^{-|\Omega|}) \|f\|_{L^2}^2$$

(see Theorem 2.3 for the general statement), where $|\Omega|$ is the Lebesgue measure of $\Omega \subset \mathbb{R}^2$. Since when $|\Omega|$ is finite the constant on the right-hand side is strictly less than one, this result can be regarded as a manifestation of the uncertainty principle. Upper bounds in the same spirit, for certain Cantor-type rotationally invariant subsets $\Omega \subset \mathbb{R}^2$ have recently been obtained in [3, 35, 36, 37] in connection with the fractal uncertainty principle. We also address to [13, 24, 30, 32, 51] for other forms of the uncertainty principle involving the short-time Fourier transform.

While the localization in the time-frequency plane by means of the characteristic function of a subset Ω , as in (1.2), is certainly natural, following a practice which dates back at least to [18] and fully promoted in [19], one can similarly localize using other *weight* functions $F(x, \omega) \geq 0$ (satisfying possibly other constraints such as $F(x, \omega) \leq 1$, to avoid an amplification) and ask for similar estimates for the weighted energy

$$\langle L_F f, f \rangle = \int_{\mathbb{R}^{2d}} F(x, \omega) |\mathcal{V}f(x, \omega)|^2 dx d\omega.$$

Here we introduced the so-called *time-frequency localization operator* L_F associated with the weight F , defined as

$$(1.3) \quad L_F = \mathcal{V}^* F \mathcal{V},$$

or weakly as

$$(1.4) \quad \langle L_F f, g \rangle_{L^2(\mathbb{R}^d)} = \langle F, \overline{\mathcal{V}f} \mathcal{V}g \rangle_{L^2(\mathbb{R}^{2d})}.$$

Indeed, if $F \in L^p(\mathbb{R}^{2d})$, $1 \leq p \leq \infty$, is any complex-valued function, it is known that L_F is a bounded operator in $L^2(\mathbb{R}^d)$ (see e.g. [56]). This class of operators was introduced by I. Daubechies [19] as a joint time-frequency version of the celebrated Landau-Pollack-Slepian operator [39, 40, 54] from

signal theory, and by F. A. Berezin [10], motivated by the quantization problem in quantum mechanics. Since then, they have been an object of intensive studies, especially regarding boundedness, compactness, Schatten properties, and asymptotics of the eigenvalues, even for more general weights and function spaces; see the classical references [11, 20, 21, 56], the more recent contributions [2, 4, 9, 16, 17, 26, 45, 46] and the references therein.

In this note we confine ourselves to the basic case of weights in Lebesgue spaces, and we prove optimal estimates of the kind

$$(1.5) \quad \|L_F\|_{L^2(\mathbb{R}^2) \rightarrow L^2(\mathbb{R}^d)} \leq C = C(p, A, B)$$

for the norm of the localization operator L_F , for any function $F \in L^p(\mathbb{R}^{2d})$ satisfying the double constraint

$$(1.6) \quad \|F\|_{L^\infty} \leq A \quad \text{and} \quad \|F\|_{L^p} \leq B.$$

Here and throughout the paper, we assume that $p \in [1, \infty)$, $A \in (0, \infty)$ and $B \in (0, \infty)$, with the further condition that $A < \infty$ when $p = 1$. Moreover, with some abuse of notation, the symbol $\|F\|_{L^\infty}$ is always meant in the broader sense of $\text{ess sup } |F|$, i.e. $+\infty$ when $F \notin L^\infty(\mathbb{R}^{2d})$. With this proviso, the case where $A = \infty$ (and $p > 1$) corresponds to dropping the L^∞ constraint: in this case, optimal bounds can be deduced from Lieb's uncertainty inequality [41] (see below) and the optimal functions, i.e. those functions F that achieve equality in (1.5), are certain Gaussians. Hence in estimating $\|L_F\|$, the L^∞ constraint in (1.6) is only relevant when A is less than the L^∞ norm of these Gaussians.

We are also able to characterize in all cases the optimal functions F , as well as all pairs of functions f and g (normalized in $L^2(\mathbb{R}^d)$) such that

$$(1.7) \quad \|L_F\|_{L^2 \rightarrow L^2} = |\langle L_F f, g \rangle|,$$

as Gaussian functions of the kind

$$(1.8) \quad x \mapsto ce^{2\pi i x \cdot \omega_0} \varphi(x - x_0), \quad x \in \mathbb{R}^d,$$

for some unimodular $c \in \mathbb{C}$ and some $(x_0, \omega_0) \in \mathbb{R}^d \times \mathbb{R}^d$, where φ is the Gaussian in (1.1).

Throughout the paper, we will frequently refer to the constant

$$(1.9) \quad \kappa_p := \frac{p-1}{p},$$

without any further reference to (1.9).

To give a flavour of our results, in dimension $d = 1$ we can state the following theorem (see Section 2 for the results in arbitrary dimension).

Theorem 1.1. *Let $p \in [1, \infty)$, $A \in (0, \infty]$, $B \in (0, \infty)$, with the proviso that $A < \infty$ when $p = 1$, and let $F \in L^p(\mathbb{R}^2)$ satisfy the constraints (1.6).*

(i) *If $p = 1$, then*

$$(1.10) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq A(1 - e^{-B/A}),$$

and equality occurs if and only if, for some $\theta \in \mathbb{R}$ and some $z_0 \in \mathbb{R}^2$,

$$(1.11) \quad F(z) = A e^{i\theta} \chi_{\mathcal{B}}(z - z_0) \quad \forall z \in \mathbb{R}^2,$$

where $\mathcal{B} \subset \mathbb{R}^2$ is the disc of area B/A , centered at the origin.

(ii) If $p > 1$ and $(B/A)^p \leq \kappa_p$ (in particular when $A = \infty$), then

$$(1.12) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \kappa_p^{\frac{1}{p}} B,$$

with equality if and only if, for some $\theta \in \mathbb{R}$ and some $z_0 \in \mathbb{R}^2$,

$$(1.13) \quad F(z) = B \kappa_p^{-\frac{1}{p}} e^{i\theta} e^{-\frac{\pi}{p-1}|z-z_0|^2}, \quad z \in \mathbb{R}^2.$$

(iii) If $p > 1$ and $(B/A)^p > \kappa_p$, then

$$(1.14) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq A \left(1 - \frac{e^{\kappa_p - (B/A)^p}}{p}\right),$$

with equality if and only if, for some $\theta \in \mathbb{R}$ and some $z_0 \in \mathbb{R}^2$,

$$(1.15) \quad F(z) = e^{i\theta} \min\{\lambda e^{-\frac{\pi}{p-1}|z-z_0|^2}, A\}, \quad z \in \mathbb{R}^2,$$

where $\lambda = A e^{(B/A)^p / (p-1) - 1/p} > A$.

Finally, if F achieves equality in (1.10), (1.12), or (1.14), then

$$\|F\|_{L^p} = B,$$

and (1.7) holds for some $f, g \in L^2(\mathbb{R})$ such that $\|f\|_{L^2} = \|g\|_{L^2} = 1$, if and only if both f and g are Gaussians of the kind (1.8) (with $d = 1$), possibly with different c 's, but with the same $(x_0, \omega_0) \in \mathbb{R}^2$ given by z_0 (the center of symmetry of $|F|$).

As a consequence, given $F \in L^p(\mathbb{R}^2)$, letting $B = \|F\|_{L^p}$ and $A = \|F\|_{L^\infty}$, one obtains the following estimates.

Corollary 1.2.

(a) Let $1 < p < \infty$. For every $F \in L^p(\mathbb{R}^2) \setminus \{0\}$ we have

$$(1.16) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \kappa_p^{\frac{1}{p}} \|F\|_{L^p},$$

with equality if and only if F is as described in (1.13), for some $B > 0$, $\theta \in \mathbb{R}$, $z_0 \in \mathbb{R}^2$.

(b) Let $1 \leq p < \infty$. For every $F \in L^\infty(\mathbb{R}^2) \cap L^p(\mathbb{R}^2) \setminus \{0\}$, satisfying $(\|F\|_{L^p} / \|F\|_{L^\infty})^p > \kappa_p$, we have

$$(1.17) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \|F\|_{L^\infty} \left(1 - \frac{e^{\kappa_p - (\|F\|_{L^p} / \|F\|_{L^\infty})^p}}{p}\right).$$

If $p > 1$, equality occurs in (1.17) if and only if F has the form in (1.15) for some $\lambda > A > 0$, $\theta \in \mathbb{R}$, $z_0 \in \mathbb{R}^2$.

If $p = 1$, equality occurs in (1.17) if and only if $F = c \chi_{\mathcal{B}}$ for some $c \in \mathbb{C} \setminus \{0\}$, where $\mathcal{B} \subset \mathbb{R}^2$ is a ball.

Some remarks are in order. When $A = \infty$ in Theorem 1.1, the extremal functions F have the form in (1.13) and $\|F\|_{L^\infty} = B\kappa_p^{1/p}$. Hence Theorem 1.1 (ii) in fact reduces to the case when only the L^p constraints is present.

As already observed in [12], the estimate (1.16) is easily seen to be *equivalent*, by duality, to Lieb's uncertainty inequality for the short-time Fourier transform (with Gaussian window) [41] (see also [15, 42]), which in dimension $d = 1$ reads as follows: for $2 \leq p < \infty$,

$$(1.18) \quad \|\mathcal{V}f\|_{L^p(\mathbb{R}^2)} \leq (2/p)^{1/p} \|f\|_{L^2(\mathbb{R})}.$$

Hence, case (b) in Corollary 1.2 is the most interesting. However, our proof does not rely on (1.18) and therefore provides an alternative proof of (1.18) as well (we address to [41] for applications of this estimate in quantum physics).

If $p = 1$, (1.17) reduces to the estimate

$$(1.19) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \|F\|_{L^\infty} (1 - e^{-\|F\|_{L^1}/\|F\|_{L^\infty}}),$$

which was proved in [29] under the stronger assumptions that F is real-valued (so that L_F is self-adjoint) and spherically symmetric, exploiting the fact that in that case L_F diagonalizes in the Hermite bases. Also, when F is the characteristic function of a measurable subset Ω of finite measure the latter estimate reduces to (1.2). Moreover, if $F \in L^1(\mathbb{R}^2)$ and $0 \leq F \leq 1$, we deduce that

$$\|L_F\|_{L^2 \rightarrow L^2} \leq 1 - e^{-\text{Tr } L_F}.$$

Indeed it is well-known that if $F \in L^1$ is non-negative then L_F is trace class and $\text{Tr } L_F = \|F\|_{L^1}$ (see e.g. [56]).

In Theorem 1.1 we excluded the case $p = 1$, $A = \infty$, because in that case the optimal constant in (1.5), which is easily seen to be $C = B$, is not achieved (see Remark 3.5).

We point out that the appearance of the Gaussians truncated above in (1.15) as extremal functions of a time-frequency concentration problem is, to our knowledge, an unprecedented fact in the literature.

The norm estimates in (1.10), (1.12), and (1.14) are obtained through a new intermediate estimate of general interest, valid under no quantitative restrictions on F , that in dimension $d = 1$ reads

$$(1.20) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \int_0^\infty (1 - e^{-\mu(t)}) dt,$$

where μ is the distribution function of $|F|$ (see Theorem 2.2 for the general statement, and Corollary 2.4 for a related symmetrization result). Then one may wonder how the right-hand side of (1.20) can be further estimated if F is subject to (1.6), and this leads in a natural way to a nonstandard variational problem (where the decreasing function μ is the unknown) described and completely solved in Section 3. Since the solution to this variational problem is (depending on the values of p , A and B) the distribution function of a

ball, of a Gaussian, or of a Gaussian truncated above, this eventually leads to Theorem 1.1 (and Theorem 2.1 in arbitrary dimension).

We also notice that the above estimates could be rephrased as uncertainty principles. As an illustration, suppose that for some $\varepsilon \in (0, 1)$, $f \in L^2(\mathbb{R}) \setminus \{0\}$ and $F \in L^\infty(\mathbb{R}^2)$, with $0 \leq F \leq 1$, we have

$$\int_{\mathbb{R}^2} F(x, \omega) |\mathcal{V}f(x, \omega)|^2 dx d\omega \geq (1 - \varepsilon) \|f\|_{L^2}^2.$$

Then it follows from (1.19) that $\|F\|_{L^1} \geq \log(1/\varepsilon)$, which generalizes [50, Theorem 1.2], where the case $F = \chi_\Omega$ was considered.

The above results could also be interpreted as optimal bounds for the norm of Toeplitz operators in the Bargmann-Fock space of entire functions [57] (cf. [29]). Also, this general scheme seems to go to the very heart of the problem and is flexible enough to encompass more general function spaces, where the norm $\|F\|_{L^p}^p$ is replaced by $\int_{\mathbb{R}^{2d}} \Phi(|F(x, \omega)|) dx d\omega$, Φ being a non-negative convex function satisfying some natural conditions, and the objective function is also of the type $\int_{\mathbb{R}^{2d}} F(x, \omega) G(|\mathcal{V}f(x, \omega)|) dx d\omega$, providing a straightforward pattern to find optimal bounds and corresponding optimal functions. Here we chose to confine ourselves to the case of the L^p setting for the sake of concreteness. However, to illustrate the scope of this approach, in Section 5 we briefly prove similar results for wavelet localization operators [21, 22]. In that case, the Cauchy wavelet plays the role of the Gaussian window (1.1), and the estimate analogous to (1.2) was recently proved in [52].

We emphasize that the concentration problem for the wavelet transform is directly related to deep issues in the theory of Bergman spaces; for example, the bound analogous of Lieb's uncertainty inequality –which we also recapture in Theorem 5.2 (ii) below (case $A = \infty$)– has recently appeared in [38] in the form of an optimal contractivity estimate for such spaces. We address the interested reader to [7, 14] for a general account of this circle of ideas, and to [44] and the references therein for the intimate connection with the Wehrl conjecture by Lieb and Solovej.

2. THE MULTIDIMENSIONAL CASE

The extension of Theorem 1.1 to arbitrary dimension $d \geq 1$ requires the introduction of the function

$$(2.1) \quad G(s) := \int_0^s e^{-(d\tau)^{\frac{1}{d}}} d\tau$$

(note that $G(s) = 1 - e^{-s}$ when $d = 1$). With this notation, we can now state the following

Theorem 2.1. *Assume $p \in [1, \infty)$, $A \in (0, \infty]$ and $B \in (0, \infty)$, with the proviso that $A < \infty$ if $p = 1$, and let $F \in L^p(\mathbb{R}^{2d})$ satisfy the constraints (1.6).*

(i) If $p = 1$, then

$$(2.2) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq A G(B/A)$$

where $G(s)$ is as in (2.1), and equality occurs if and only if, for some $\theta \in \mathbb{R}$ and some $z_0 \in \mathbb{R}^{2d}$,

$$(2.3) \quad F(z) = Ae^{i\theta} \chi_{\mathcal{B}}(z - z_0) \quad \forall z \in \mathbb{R}^{2d},$$

where $\mathcal{B} \subset \mathbb{R}^{2d}$ is the ball of volume B/A , centered at the origin.

(ii) If $p > 1$ and $(B/A)^p \leq \kappa_p^d$, then

$$(2.4) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \kappa_p^{d/p} B,$$

with equality if and only if, for some $\theta \in \mathbb{R}$ and some $z_0 \in \mathbb{R}^{2d}$,

$$(2.5) \quad F(z) = e^{i\theta} \lambda e^{-\frac{\pi}{p-1}|z-z_0|^2} \quad z \in \mathbb{R}^{2d},$$

where $\lambda = \kappa_p^{-d/p} B$.

(iii) If $p > 1$ and $(B/A)^p > \kappa_p^d$, then

$$(2.6) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \int_0^A G(u_\lambda(t)) dt,$$

where $u_\lambda(t) = (-\log((t/\lambda)^{p-1}))^d/d!$ and $\lambda > A$ is uniquely determined by the condition that $p \int_0^A t^{p-1} u_\lambda(t) dt = B^p$. Equality occurs in (2.6) if and only if, for some $\theta \in \mathbb{R}$ and some $z_0 \in \mathbb{R}^{2d}$,

$$(2.7) \quad F(z) = e^{i\theta} \min\{\lambda e^{-\frac{\pi}{p-1}|z-z_0|^2}, A\} \quad z \in \mathbb{R}^{2d}.$$

Finally, if F achieves equality in (2.2), (2.4), or (2.6), then

$$(2.8) \quad \|F\|_{L^p} = B,$$

and (1.7) holds for some $f, g \in L^2(\mathbb{R}^d)$ such that $\|f\|_{L^2} = \|g\|_{L^2} = 1$, if and only if both f and g are functions of the kind (1.8), possibly with different c 's, but with the same $(x_0, \omega_0) \in \mathbb{R}^{2d}$ that coincides with z_0 (the center of symmetry of $|F|$).

The proof of Theorem 2.1, which will be given in Section 4, partially relies on the following result, which is of standalone interest and provides an explicit bound for the norm of the operator L_F , in terms of the distribution function of $|F|$.

Theorem 2.2. Assume $F \in L^p(\mathbb{R}^{2d})$ for some $p \in [1, +\infty)$, and let

$$(2.9) \quad \mu(t) = |\{|F| > t\}|, \quad t > 0,$$

denote the distribution function of $|F|$. Then

$$(2.10) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \int_0^\infty G(\mu(t)) dt$$

where G is as in (2.1), in particular, when $d = 1$

$$\|L_F\|_{L^2 \rightarrow L^2} \leq \int_0^\infty (1 - e^{-\mu(t)}) dt.$$

Equality occurs in (2.10) if and only if $F(z) = e^{i\theta} \rho(|z - z_0|)$ for some $\theta \in \mathbb{R}$, some $z_0 \in \mathbb{R}^{2d}$, and some nonincreasing function $\rho : [0, +\infty) \rightarrow [0, +\infty)$. In this case, (1.7) holds true for some $f, g \in L^2(\mathbb{R}^d)$ such that $\|f\|_{L^2} = \|g\|_{L^2} = 1$, if (and only if, when F is not identically zero) both f and g are functions of the kind (1.8), possibly with different c 's, but with the same (x_0, ω_0) that coincides with z_0 .

To prove Theorem 2.2, we need the following result ([50, Theorem 4.1]).

Theorem 2.3 ([50]). *For every $f \in L^2(\mathbb{R}^d)$ such that $\|f\|_{L^2} = 1$, and every measurable subset $\Omega \subset \mathbb{R}^{2d}$ of measure $|\Omega| < \infty$, we have*

$$(2.11) \quad \int_{\Omega} |\mathcal{V}f(x, \omega)|^2 dx d\omega \leq G(|\Omega|)$$

where G is as in (2.1), in particular when $d = 1$

$$(2.12) \quad \int_{\Omega} |\mathcal{V}f(x, \omega)|^2 dx d\omega \leq 1 - e^{-|\Omega|}.$$

Equality occurs in (2.11) (for some set Ω such that $0 < |\Omega| < \infty$ and some f such that $\|f\|_{L^2} = 1$) if and only if f is a function of the form as in (1.8) and Ω is equivalent, in measure, to a ball of center (x_0, ω_0) .

We warn the reader that in [50] (2.11) is stated in a different way, namely with $G(s)$ defined as

$$G(s) = \frac{1}{(d-1)!} \int_0^{\pi(s/\omega_{2d})^{1/d}} t^{d-1} e^{-t} dt,$$

where ω_{2d} is the volume of the unit ball in \mathbb{R}^{2d} . Since $\omega_{2d} = \pi^d/d!$, by the change of variable $t^d = d! \tau$ one can easily check that this definition coincides with that in (2.1).

Proof of Theorem 2.2. Given f and g normalized in $L^2(\mathbb{R}^d)$, writing z to denote the variable $(x, \omega) \in \mathbb{R}^{2d}$ and abbreviating dz for $dx d\omega$, we have, by (1.4),

$$(2.13) \quad \begin{aligned} |(L_F f, g)| &\leq \int_{\mathbb{R}^{2d}} |F(z)| \cdot |\mathcal{V}f(z)| \cdot |\mathcal{V}g(z)| dz \\ &\leq \left(\int_{\mathbb{R}^{2d}} |F(z)| \cdot |\mathcal{V}f(z)|^2 dz \right)^{1/2} \left(\int_{\mathbb{R}^{2d}} |F(z)| \cdot |\mathcal{V}g(z)|^2 dz \right)^{1/2}. \end{aligned}$$

Letting $m = \text{ess sup } |F(z)|$ (and assuming $m > 0$, otherwise the claims to be proved are trivial), we apply the ‘‘layer cake’’ formula (cf. [43, Page 26])

$$(2.14) \quad |F(z)| = \int_0^m \chi_{\{|F|>t\}}(z) dt, \quad z \in \mathbb{R}^{2d},$$

and use (2.11) (with $\Omega = \{|F| > t\}$) to estimate

$$(2.15) \quad \int_{\mathbb{R}^{2d}} |F(z)| \cdot |\mathcal{V}f(z)|^2 dz = \int_0^m \left(\int_{\{|F|>t\}} |\mathcal{V}f(z)|^2 dz \right) dt \\ \leq \int_0^m G(\mu(t)) dt$$

(since $G(0) = 0$, the last integral coincides with the one in (2.10), by the definition of m). Since the same argument, with g in place of f , leads to

$$(2.16) \quad \int_{\mathbb{R}^{2d}} |F(z)| \cdot |\mathcal{V}g(z)|^2 dz \leq \int_0^m G(\mu(t)) dt,$$

from (2.13) and the arbitrariness of f and g we obtain (2.10). In fact, since $p < \infty$, L_F is a compact operator on $L^2(\mathbb{R}^d)$ (see [56]), hence there exist f, g (normalized in $L^2(\mathbb{R}^d)$) satisfying (1.7), so that having equality in (2.10) is equivalent to having equality (with these f and g) in (2.13), (2.15) and (2.16): we now characterize for what F, f and g this really occurs.

First, equality occurs in (2.15) if and only if

$$(2.17) \quad \int_{\{|F|>t\}} |\mathcal{V}f(z)|^2 dz = G(\mu(t))$$

for a.e. $t \in (0, m)$. When this happens, the validity of (2.17) for just *one* $t_0 \in (0, m)$ is enough, by Theorem 2.3, to infer that f has the form as in (1.8) (for some $(x_0, \omega_0) \in \mathbb{R}^{2d}$ and some unimodular c), and that the corresponding level set $\{|F| > t_0\}$ is (equivalent to) a ball centered at $z_0 := (x_0, \omega_0)$; then, still by Theorem 2.3, the fact that (2.17) holds for a.e. $t \in (0, m)$ implies (f being the same) that also all the other level sets $\Omega_t = \{|F| > t\}$ are equivalent to balls centered at the same z_0 (once this is known for a.e. $t \in (0, m)$, the passage to *every* $t \in (0, m)$ follows since $\Omega_t = \bigcup_{s>t} \Omega_s$). These necessary conditions on f and F are, in turn, also sufficient to guarantee (2.17) for a.e. $t \in (0, m)$. Thus, summing up, equality in (2.15) is equivalent to f being as in (1.8) and $|F(z)|$ being spherically symmetric (and radially decreasing) around z_0 , i.e. $|F(z)| = \rho(|z - z_0|)$ for some nonincreasing function $\rho : [0, +\infty) \rightarrow [0, +\infty)$.

Equality in (2.16) can be characterized similarly, so that also g must be as in (1.8) (possibly with a different c , but with the same z_0 because F is the same).

Therefore, as soon as equality occurs in (2.15) and (2.16), we have $\mathcal{V}g(z) = e^{i\alpha} \mathcal{V}f(z)$ for some $\alpha \in \mathbb{R}$, so that the second inequality in (2.13) is automatically an equality, while the first inequality, which explicitly amounts to

$$\left| \int_{\mathbb{R}^{2d}} F(z) \mathcal{V}f(z) \overline{\mathcal{V}g(z)} dz \right| \leq \int_{\mathbb{R}^{2d}} |F(z)| \cdot |\mathcal{V}f(z)| \cdot |\mathcal{V}g(z)| dz,$$

becomes an equality if and only if

$$e^{-i\theta} \int_{\mathbb{R}^2} F(z) \cdot |\mathcal{V}f(z)|^2 dz = \int_{\mathbb{R}^2} |F(z)| \cdot |\mathcal{V}f(z)|^2 dz$$

for some $\theta \in \mathbb{R}$. This, in turn, is equivalent to the condition

$$e^{-i\theta} F(z) \cdot |\mathcal{V}f(z)|^2 = |F(z)| \cdot |\mathcal{V}f(z)|^2 \quad \text{for a.e. } z \in \mathbb{R}^{2d},$$

but since $|\mathcal{V}f(z)|^2 > 0$ (note that $|\mathcal{V}f(z)|^2$ is a Gaussian when f is as in (1.8)) we see that equalities occur in (2.13), for f and g as in (1.8), if and only if $F(z) = e^{i\theta}|F(z)|$ (i.e. $F(z) = e^{i\theta}\rho(|z - z_0|)$) for a.e. $z \in \mathbb{R}^{2d}$. \square

An interesting consequence is the following symmetrization result, which shows that the norm of the localization operator L_F increases when F is replaced by its Schwarz symmetrization (we refer to [8] for a general account on symmetrization).

Corollary 2.4. *Assume $F \in L^p(\mathbb{R}^{2d})$ for some $p \in [1, +\infty)$, and let F^* denote the Schwarz symmetrization of $|F|$. Then*

$$(2.18) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \|L_{F^*}\|_{L^2 \rightarrow L^2},$$

and equality occurs if and only if $F(z) = e^{i\theta}F^*(z - z_0)$ for some $\theta \in \mathbb{R}$ and some $z_0 \in \mathbb{R}^{2d}$.

Proof. Let $\mu(t)$ be defined as in (2.9), and observe that μ is also the distribution function of F^* . Applying Theorem 2.2 to F^* , we get the bound (2.10) with F^* in place of F . But since $F^*(z)$ is, by construction, of the form $\rho(|z|)$ where $\rho : [0, +\infty) \rightarrow [0, +\infty)$ is nonincreasing, by Theorem 2.2 this bound is in fact an equality, i.e.

$$(2.19) \quad \|L_{F^*}\|_{L^2 \rightarrow L^2} = \int_0^\infty G(\mu(t)) dt,$$

which combined with (2.10) yields (2.18). Equality, by Theorem 2.2, occurs therein if and only if $F(z) = e^{i\theta}\tilde{\rho}(|z - z_0|)$, for some $\theta \in \mathbb{R}$, some $z_0 \in \mathbb{R}^{2d}$, and some nonincreasing function $\tilde{\rho} : [0, +\infty) \rightarrow [0, +\infty)$. Finally, since $\tilde{\rho}(|z - z_0|) = |F(z)|$ has the same distribution function as $\rho(|z|) = F^*(z)$, we see that $\tilde{\rho} = \rho$, hence equality in (2.18) is equivalent to $F(z) = e^{i\theta}F^*(z - z_0)$ for a.e. $z \in \mathbb{R}^{2d}$. \square

3. A NONSTANDARD VARIATIONAL PROBLEM

To prove Theorems 2.1 and 1.1 we shall build on Theorem 2.2, seeking sharp upper bounds for the right-hand side of (2.10) when F is subject to the double constraint in (1.6). Given p , A and B as in Theorem 2.1, based on the integral constraint

$$(3.1) \quad p \int_0^A t^{p-1} u(t) dt \leq B^p$$

we define the class of functions

$$(3.2) \quad \mathcal{C} = \{u : (0, A) \rightarrow [0, +\infty) \mid u \text{ is decreasing and satisfies (3.1)}\}.$$

If μ is the distribution function of $|F|$ as in (2.9), where $F \in L^p(\mathbb{R}^{2d})$ and satisfies (1.6), the L^∞ constraint can be expressed as

$$(3.3) \quad \mu(t) = 0 \quad \forall t \geq A,$$

while the L^p constraint corresponds to letting $u = \mu$ in (3.1). Hence, since μ is decreasing, we see that μ (restricted to the relevant interval $(0, A)$) belongs to the class \mathcal{C} as defined above, and to estimate the right-hand side of (2.10) it is natural to investigate the variational problem

$$(3.4) \quad \sup_{v \in \mathcal{C}} I(v) \quad \text{where} \quad I(v) := \int_0^A G(v(t)) dt$$

(integration in (3.4) can be restricted to $(0, A)$ due to (3.3), since $G(0) = 0$). We first show the existence of an extremal function.

Proposition 3.1. *The supremum in (3.4) is finite and is attained by at least one function $u \in \mathcal{C}$. Moreover, every extremal function u achieves equality in the constraint (3.1).*

Proof. Due to (3.1), for every $u \in \mathcal{C}$ and every $t \in (0, A)$ we have

$$t^p u(t) \leq p \int_0^t \tau^{p-1} u(\tau) d\tau \leq B^p,$$

which yields the pointwise bound

$$(3.5) \quad u(t) \leq B^p/t^p \quad \forall t \in (0, A), \quad \forall u \in \mathcal{C}.$$

Since G in (2.1) is increasing and bounded, from this bound we obtain

$$(3.6) \quad I(u) = \int_0^A G(u(t)) dt \leq \int_0^A G(B^p/t^p) dt < \infty \quad \forall u \in \mathcal{C}$$

(when $A = \infty$, and hence $p > 1$, the finiteness of the last integral follows from the bound $G(s) \leq s$). As a consequence, the supremum in (3.4) is finite.

Now let $u_n \in \mathcal{C}$ be a maximizing sequence for problem (3.4), i.e.

$$\lim_{n \rightarrow \infty} I(u_n) = \sup_{v \in \mathcal{C}} I(v).$$

Since each u_n satisfies a bound as in (3.5), by Helly's selection theorem we can extract a subsequence (still denoted by u_n) pointwise converging to a decreasing function $u : (0, A) \rightarrow [0, +\infty)$, which still satisfies (3.1) by Fatou's lemma, so that $u \in \mathcal{C}$ as well. As already observed, $u_n(t) \leq B^p/t^p$, and therefore by dominated convergence (arguing as in (3.6)) we have $I(u) = \lim_{n \rightarrow \infty} I(u_n)$, which proves that u is a maximizer.

Now assume that a strict inequality occurs in the constraint (3.1). Then, letting $u_\varepsilon(t) = u(t) + \varepsilon e^{-t}$ and choosing $\varepsilon > 0$ small enough, we would have $u_\varepsilon \in \mathcal{C}$: since the function $G(s)$ in (2.1) is strictly increasing, we would also have $I(u_\varepsilon) > I(u)$, which is impossible since u is a maximizer. \square

We now show that one can remove the monotonicity assumption in (3.2), while retaining the same extremal functions for the corresponding maximization problem.

Proposition 3.2. *We have*

$$(3.7) \quad \sup_{v \in \mathcal{C}} I(v) = \sup_{v \in \mathcal{C}'} I(v),$$

where

$$\mathcal{C}' = \{u : (0, A) \rightarrow [0, +\infty) \mid u \text{ is measurable and satisfies (3.1)}\}.$$

In particular, any function $u \in \mathcal{C}$ achieving the supremum on the left-hand side, also achieves the supremum on the right-hand side.

Proof. Consider an arbitrary $u \in \mathcal{C}'$, and let u^* denote the decreasing rearrangement of u (i.e. the unique decreasing, right-continuous function on $(0, A)$ which is equimeasurable with u , see [34, Section 10.12]). Since for every $s > 0$ the superlevel set $\{u > s\} \subseteq (0, A)$ has finite measure (this is trivial when $A < \infty$, and is guaranteed by (3.1) when $A = \infty$), we see that $u^* : (0, A) \rightarrow [0, +\infty)$ takes only finite values. We also claim that

$$(3.8) \quad p \int_0^A t^{p-1} u^*(t) dt \leq p \int_0^A t^{p-1} u(t) dt.$$

Indeed, let ν denote the Radon measure on $(0, A)$ with density t^{p-1} , and observe that, since t^{p-1} is an increasing function of t ,

$$\nu(E) = \int_E t^{p-1} dt \geq \int_0^{|E|} t^{p-1} dt = \nu((0, |E|))$$

for every measurable $E \subseteq (0, A)$, where $|E|$ denotes the Lebesgue measure of E . In particular, when $E = \{u > s\}$ is a superlevel set of u , we obtain

$$\nu(\{u > s\}) \geq \nu((0, |\{u > s\}|)) = \nu((0, |\{u^* > s\}|)) = \nu(\{u^* > s\}) \quad \forall s \geq 0$$

(the first equality is due to the equimeasurability of u and u^* , the second to the fact that u^* is decreasing and right-continuous). Then (3.8) follows, since

$$\begin{aligned} \int_0^A t^{p-1} u(t) dt &= \int_0^A u(t) d\nu(t) = \int_0^\infty \nu(\{u > s\}) ds \\ &\geq \int_0^\infty \nu(\{u^* > s\}) ds = \int_0^A u^*(t) d\nu(t) = \int_0^A t^{p-1} u^*(t) dt. \end{aligned}$$

Since u^* is decreasing and $u \in \mathcal{C}'$, (3.8) shows that $u^* \in \mathcal{C}$. On the other hand, u and u^* are equi-measurable, so that $I(u) = I(u^*) \leq \sup_{v \in \mathcal{C}} I(v)$. Then, from the arbitrariness of $u \in \mathcal{C}'$, we see that the inequality \geq occurs in (3.7), while the opposite inequality is trivial because $\mathcal{C} \subset \mathcal{C}'$. \square

Remark 3.3. This proof entails that if $u \in \mathcal{C}'$ is not equal a.e. to a decreasing function (i.e. if u and u^* do not coincide a.e.), then the inequality in (3.8) is strict, and hence (by the last part of Proposition 3.1) u^* cannot achieve the

supremum in (3.4). Since $I(u) = I(u^*)$, from (3.7) we infer that the second supremum in (3.7) can be achieved only by functions in \mathcal{C} . In the following, however, we will not need this fact.

Now we are in a position to completely solve problem (3.4). Let

$$(3.9) \quad \text{Log}_-(x) := \max\{-\log x, 0\}, \quad x > 0.$$

Theorem 3.4. *There exists a unique function $u \in \mathcal{C}$ attaining the supremum in (3.4), and this u achieves equality in (3.1). More precisely:*

(i) *If $p = 1$ (hence $A < \infty$), then $u(t) = B/A$ is constant on $(0, A)$, and*

$$(3.10) \quad I(u) = A G(B/A)$$

where $G(s)$ is defined as in (2.1).

(ii) *If $p > 1$ and $(B/A)^p \leq \kappa_p^d$ (in particular when $A = \infty$), then*

$$(3.11) \quad u(t) = \frac{1}{d!} (\text{Log}_-((t/\lambda)^{p-1}))^d, \quad \forall t \in (0, A)$$

where $\lambda = B\kappa_p^{-\frac{d}{p}} \leq A$, and

$$(3.12) \quad I(u) = \kappa_p^{d\kappa_p} B.$$

(iii) *If $p > 1$ and $(B/A)^p > \kappa_p^d$, then u is given by*

$$(3.13) \quad u(t) = \frac{1}{d!} (-\log((t/\lambda)^{p-1}))^d \quad \forall t \in (0, A),$$

where $\lambda > A$ is uniquely determined by the condition that u achieves equality in (3.1). Moreover, when $d = 1$, $\lambda = Ae^{(B/A)^p/(p-1)-1/p}$ and

$$(3.14) \quad I(u) = A \left(1 - \frac{e^{\kappa_p - (B/A)^p}}{p} \right) \quad (\text{when } d = 1).$$

Proof of Theorem 3.4. We already know from Proposition 3.1 that the supremum in (3.4) is achieved by some $u \in \mathcal{C}$ and that equality holds in (3.1) for any such u , i.e.

$$(3.15) \quad p \int_0^A t^{p-1} u(t) dt = B^p.$$

Moreover, by Proposition 3.2, any such u also achieves the second supremum in (3.7): we now exploit this information to determine u .

Let us first assume that $p = 1$ (case (i)). By Jensen inequality and (3.15)

$$\frac{I(u)}{A} = \frac{1}{A} \int_0^A G(u(t)) dt \leq G\left(\frac{1}{A} \int_0^A u(t) dt\right) = G(B/A)$$

with equality if and only if u is constant (i.e. $u = B/A$), since G is strictly concave. Then (3.10) follows, and case (i) is proved.

From now on, we assume that $p > 1$. Observe that u is not identically zero, due to (3.15). Thus, letting

$$(3.16) \quad M = \sup \{t \in (0, A) \mid u(t) > 0\},$$

since u is decreasing and $u \geq 0$, we have $M \in (0, A]$ and $u(t) > 0$ for $t \in (0, M)$, while $u(t) = 0$ for $t \in (M, A)$ if $M < A$.

Now fix two numbers a, b such that $0 < a < b < M$, and let $\eta \in L^\infty(0, A)$ be an arbitrary function supported in $[a, b]$ and such that

$$(3.17) \quad \int_a^b t^{p-1} \eta(t) dt = 0.$$

Since on $[a, b]$ we have $u(t) \geq u(b) > 0$, if $|\varepsilon|$ is small enough (depending on $\|\eta\|_{L^\infty}$) we have $u + \varepsilon\eta \in \mathcal{C}'$, and the function $I(u + \varepsilon\eta)$ has a maximum at $\varepsilon = 0$ since u is a maximizer of I on \mathcal{C}' . Therefore, differentiating under the integral, we obtain (without expanding G' for the moment)

$$(3.18) \quad 0 = \frac{d}{d\varepsilon} I(u + \varepsilon\eta)|_{\varepsilon=0} = \int_a^b G'(u(t)) \eta(t) dt.$$

Focusing on the interval $[a, b]$, this condition can be interpreted as follows: the restriction to $[a, b]$ of the function $G'(u(t))$ is orthogonal, in $L^2(a, b)$, to every function $\eta \in L^\infty(a, b)$ satisfying (3.17). Now, given $\eta \in L^2(a, b)$ satisfying (3.17), by a standard truncation argument (and a slight perturbation of the resulting functions) one can easily construct functions $\eta_k \in L^\infty(a, b)$, each satisfying $\int_a^b t^{p-1} \eta_k(t) dt = 0$ (hence also $\int_a^b G'(u(t)) \eta_k(t) dt = 0$), such that $\eta_k \rightarrow \eta$ in $L^2(a, b)$. Letting $k \rightarrow \infty$, we see that $G'(u(t))$ is orthogonal, in $L^2(a, b)$, to any $\eta \in L^2(a, b)$ (not necessarily L^∞) satisfying (3.17), and therefore

$$(3.19) \quad G'(u(t)) = ct^{p-1} \quad \text{for a.e. } t \in [a, b]$$

for some constant $c > 0$. In fact, by the arbitrariness of the interval $[a, b] \subset (0, M)$, letting $a \rightarrow 0^+$ and $b \rightarrow M^-$ one can see that (3.19) holds for a.e. $t \in (0, M)$, and hence for *every* $t \in (0, M)$ as $u(t)$ is decreasing. Defining a new constant λ by $\lambda^{1-p} = c$, we can write this as

$$(3.20) \quad G'(u(t)) = (t/\lambda)^{p-1} \quad \forall t \in (0, M)$$

and hence, since $G'(s) = \exp(-(s d!)^{1/d})$, we find

$$(3.21) \quad u(t) = \begin{cases} \frac{1}{d!} (-\log((t/\lambda)^{p-1}))^d & \text{if } t \in (0, M) \\ 0 & \text{if } t \in (M, A) \end{cases}$$

the second case being meaningful only if $M < A$. Observe that, since $u(t) > 0$ on $(0, M)$, we must have

$$(3.22) \quad \lambda \geq M,$$

in particular $M < \infty$. Now we prove that, if $M < A$, then actually $\lambda = M$, so that by (3.21) $u(t)$ is continuous also at $t = M$. Arguing by contradiction, i.e. assuming that $M < A$ but $\lambda > M$, we choose $\delta > 0$ small enough (so

that $M - \delta > 0$ and $M + \delta < A$) and define the function

$$(3.23) \quad \eta(t) = \begin{cases} -t^{1-p} & \text{if } t \in (M - \delta, M) \\ t^{1-p} & \text{if } t \in (M, M + \delta) \\ 0 & \text{otherwise} \end{cases}$$

in such a way that (3.17) holds true. Defining $u_\varepsilon(t) = u(t) + \varepsilon\eta(t)$, (3.23) and (3.21) (combined with our assumption that $\lambda > M$) imply that $u_\varepsilon \geq 0$ on $(0, A)$, provided $\varepsilon > 0$ is small enough. This, combined with (3.17), yields for small $\varepsilon > 0$ an admissible competitor $u_\varepsilon \in \mathcal{C}'$, so that $I(u_\varepsilon) \leq I(u)$ since u is a maximizer of I . Hence, differentiating under the integral as done for (3.18), we obtain

$$0 \geq \lim_{\varepsilon \rightarrow 0^+} \frac{I(u_\varepsilon) - I(u)}{\varepsilon} = \int_{M-\delta}^{M+\delta} G'(u(t))\eta(t) dt.$$

According to (3.23), (3.20), and the fact that $G'(u(t)) = G'(0) = 1$ when $t > M$, the last inequality reduces to

$$0 \geq - \int_{M-\delta}^M \lambda^{1-p} dt + \int_M^{M+\delta} t^{1-p} dt.$$

Dividing by δ and letting $\delta \rightarrow 0^+$ we find $M \geq \lambda$, which is the desired contradiction. Hence $\lambda = M$, if $M < A$.

This shows that, regardless of whether $M < A$ or $M = A$, (3.21) can be written, with the notation (3.9), as

$$(3.24) \quad u(t) = \frac{1}{d!} (\text{Log}_-((t/\lambda)^{p-1}))^d \quad \forall t \in (0, A).$$

Now let $h(\lambda) = p \int_0^A t^{p-1} u_\lambda(t) dt$, $\lambda > 0$, where u_λ is the function on the right-hand side of (3.24). Then $h(\lambda)$ is strictly increasing, so that there is at most one value of $\lambda > 0$ such that (3.15) (i.e. $h(\lambda) = B^p$) is satisfied, and this proves the uniqueness of u , as a maximizer of I , via (3.24). To quantify λ and make (3.24) more explicit, note that if $A < \infty$ (by the change of variable $t/A = \tau$)

$$(3.25) \quad h(A) = \frac{pA^p}{d!} \int_0^1 \tau^{p-1} (-\log \tau^{p-1})^d d\tau = A^p \frac{(p-1)^d}{p^d} = \kappa_p^d A^p.$$

In case (ii), when $B^p \leq \kappa_p^d A^p$, the strict monotonicity of $h(\lambda)$ and the condition $h(\lambda) = B^p$ imply that $\lambda \leq A$. Then, by (3.24), (3.15) becomes

$$\begin{aligned} B^p &= \frac{p}{d!} \int_0^\lambda t^{p-1} (-\log((t/\lambda)^{p-1}))^d dt \\ &= \frac{p\lambda^p}{d!} \int_0^1 \tau^{p-1} (-\log \tau^{p-1})^d d\tau = \lambda^p \frac{(p-1)^d}{p^d} = \lambda^p \kappa_p^d, \end{aligned}$$

so that $\lambda = B\kappa_p^{-d/p}$ and (3.11) follows from (3.24). Since $G(0) = 0$, using (3.24) we also have

$$I(u) = \int_0^\lambda G((-\log((t/\lambda)^{p-1}))^d/d!) dt = \lambda \int_0^1 G((-\log \tau^{p-1})^d/d!) d\tau.$$

As already observed, by (2.1) and a change of variable we have

$$G(s) = \frac{1}{(d-1)!} \int_0^{(ds)^{1/d}} t^{d-1} e^{-t} dt,$$

so that, by the Fubini Theorem,

$$\int_0^1 G((-\log \tau^{p-1})^d/d!) d\tau = \frac{1}{(d-1)!} \int_0^{+\infty} t^{d-1} e^{-\frac{p}{p-1}t} dt = \kappa_p^d.$$

Hence

$$I(u) = \lambda \kappa_p^d = B\kappa_p^{-d/p} \kappa_p^d = B\kappa_p^{d\kappa_p},$$

and (3.12) follows.

Similarly, in case (iii), when $B^p > \kappa_p^d A^p$, the condition $h(\lambda) = B^p$ and (3.25) imply that $\lambda > A$, and (3.24) simplifies to (3.13). Then (3.15) becomes

$$B^p = \frac{p}{d!} \int_0^A t^{p-1} (-\log((t/\lambda)^{p-1}))^d dt \quad (\lambda > A),$$

but now the values of λ and $I(u)$ cannot be computed explicitly, except when $d = 1$. In this case, computing the last integral, one finds $\lambda = Ae^{(B/A)^{p/(p-1)} - 1/p}$ as claimed after (3.13), and since when $d = 1$ $G(s) = 1 - e^{-s}$ and $u(t) = -\log((t/\lambda)^{p-1})$ by (3.13), we can compute

$$I(u) = \int_0^A G(u(t)) dt = \int_0^A (1 - (t/\lambda)^{p-1}) dt = A \left(1 - \frac{(A/\lambda)^{p-1}}{p} \right),$$

and (3.14) follows if one replaces λ with the value mentioned above. \square

Remark 3.5. In Theorem 3.4, when $p = 1$ we assumed $A < \infty$. Indeed, when $p = 1$ and $A = \infty$, there is no extremal function: for every $u \in \mathcal{C} \setminus \{0\}$ we have, by (3.1) (with $p = 1, A = \infty$),

$$I(u) = \int_0^{+\infty} G(u(t)) dt < \int_0^{+\infty} u(t) dt \leq B,$$

whereas if we consider the sequence of functions u_n on $(0, +\infty)$ given by $u_n(t) = B/n$ for $t \in (0, n)$, and $u_n(t) = 0$ elsewhere, we have $I(u_n) = G(B/n)/n \rightarrow B$ as $n \rightarrow \infty$.

4. PROOF OF THE MAIN RESULT IN ARBITRARY DIMENSION

This section is devoted to the proof of Theorems 2.1 and 1.1.

Proof of Theorems 2.1 and 1.1. First, observe that (i), (ii), and the last sentence of Theorem 1.1 follows, respectively, from (i), (ii) and the last sentence of Theorem 2.1, as particular cases when $d = 1$. The same is true of (iii), with the difference that (1.14) and (1.15) are more explicit than (2.6) and (2.7), since when $d = 1$ the value of λ can be found explicitly. Thus, we shall prove only Theorem 2.1, discussing how (1.14) and the value of λ in (1.15) are obtained when $d = 1$.

Let $F \in L^p(\mathbb{R}^{2d})$ satisfy (1.6), and let μ be its distribution function as in (2.9). Observing that, due to (3.3) and $G(0) = 0$, the integral in (2.10) can be restricted to $(0, A)$, recalling (3.4) and (3.2) we have from (2.10)

$$(4.1) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq \int_0^A G(\mu(t)) dt = I(\mu) \leq \sup_{v \in \mathcal{C}} I(v)$$

(note that $\mu \in \mathcal{C}$). On the other hand, this supremum is achieved by a unique function $u \in \mathcal{C}$ completely characterized in Theorem 3.4, so that

$$(4.2) \quad \|L_F\|_{L^2 \rightarrow L^2} \leq I(u) = \int_0^A G(u(t)) dt.$$

Then (2.2) and (2.4) follow, respectively, from (3.10) and (3.12), while (2.6) coincides with (4.2) because, in (2.6), u_λ is the function u in (3.13). In particular, when $d = 1$ and the value of $I(u)$ is the one in (3.14), (4.2) yields (1.14).

Now observe that equality in (2.2), (2.4) or (2.6) (i.e. equality in (4.2)) occurs if and only if the two inequalities in (4.1) reduce to equalities, which corresponds to the simultaneous validity of these two conditions:

(a) Equality occurs in (2.10). By Theorem 2.2, this means that $F(z) = e^{i\theta} \rho(|z - z_0|)$ for some $\theta \in \mathbb{R}$ and $z_0 \in \mathbb{R}^{2d}$, where $\rho : [0, +\infty) \rightarrow [0, +\infty)$ is decreasing. Therefore, for every $t \in (0, A)$, the set $\{|F| > t\}$ is a ball (centered at z_0) of measure $\mu(t)$ by (2.9), whereas $\{|F| > t\} = \emptyset$ if $t \geq A$, by (1.6).

(b) μ (restricted to $(0, A)$ and regarded as an element of \mathcal{C}) achieves the supremum in (3.4). By Theorem 3.4, this means that μ is either the constant function equal to B/A , or the function in (3.11), or the function in (3.13), depending on whether we are in case (i), (ii) or (iii) (note that these three cases match those of Theorem 2.1), and moreover

$$(4.3) \quad p \int_0^A t^{p-1} \mu(t) dt = B^p.$$

Combining (a) and (b), we see that in case (i), where $\mu = B/A$ on $(0, A)$, for every $t \in (0, A)$ the set $\{|F| > t\}$ is a ball (centered at z_0) of measure B/A , and hence F has the form in (2.3).

In case (ii), where μ is as in (3.11) with $\lambda = B\kappa_p^{-\frac{d}{p}}$, we can reconstruct the function ρ as follows. For every $t \in (0, \lambda)$, the radius r of a $2d$ -ball of

measure $\mu(t)$ (such as the set $\{|F| > t\}$) satisfies the equation

$$(4.4) \quad \frac{\pi^d r^{2d}}{d!} = \mu(t) = \frac{(-\log(t/\lambda)^{p-1})^d}{d!}$$

(note that $\pi^d/d!$ is the volume of the unit ball in \mathbb{R}^{2d}), and clearly $\rho(r) = t$. Solving for t we obtain

$$(4.5) \quad t = \rho(r) = \lambda e^{-\frac{\pi}{p-1}r^2} \quad \forall t \in (0, \lambda),$$

and this determines $\varphi(r)$ for every $r > 0$. Since $F(z) = e^{i\theta}\varphi(|z - z_0|)$, (2.5) is proved.

Finally, in case (iii), where the restriction of $\mu(t)$ to $(0, A)$ is as in (3.13), since now $\lambda > A$ we see that μ has a jump at $t = A$ (recall that $\mu(t) = 0$ when $t \geq A$). However, if $t \in (0, A)$, to reconstruct φ we can proceed as in (4.4), now obtaining

$$(4.6) \quad t = \rho(r) = \lambda e^{-\frac{\pi}{p-1}r^2} \quad \forall t \in (0, A).$$

The difference with respect to (4.5) is that, since now $A < \lambda$, (4.6) determines $\rho(r)$ only for $r > r_0$, where r_0 is the value of r that, plugged into (4.4), yields $t = A$. However, since ρ is decreasing and moreover $\rho \leq A$ by (1.6), we obtain that $\rho(r) = A$ when $r \in (0, r_0]$. In other words,

$$\rho(r) = \min \left\{ \lambda e^{-\frac{\pi}{p-1}r^2}, A \right\}, \quad r > 0,$$

and hence (2.7) is proved. When $d = 1$, the value of λ is the one mentioned after (3.13), so that also (1.15) is established.

Notice that (2.8) follows from (2.9) and the formula

$$\|F\|_{L^p}^p = p \int_0^\infty t^{p-1} \mu(t) dt = p \int_0^A t^{p-1} \mu(t) dt,$$

combined with (4.3). Finally, the characterization of those f, g satisfying (1.7) follows from Theorem 2.2, since any F achieving equalities in (4.1) achieves, in particular, also equality in (2.10). \square

5. WAVELET LOCALIZATION OPERATORS

In this section, we normalize the Fourier transform as

$$\widehat{f}(\omega) = \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} e^{-i\omega t} f(t) dt.$$

For $\beta > 0$ consider the so-called Cauchy wavelet $\psi_\beta \in L^2(\mathbb{R})$ [21, 22, 33] defined by

$$\widehat{\psi}_\beta(\omega) = \frac{1}{c_\beta} \chi_{[0, +\infty)}(\omega) \omega^\beta e^{-\omega},$$

where $c_\beta > 0$ and $c_\beta^2 = 2\pi 2^{-2\beta} \Gamma(2\beta)$, so that $\|\widehat{\psi}_\beta\|_{L^2(\mathbb{R}_+, d\omega/\omega)}^2 = 1/(2\pi)$ ($\mathbb{R}_+ = (0, +\infty)$). This normalization is chosen for the corresponding wavelet transform, defined below, to be an isometry. Observe also that $\|\psi_\beta\|_{L^2}^2 = \beta/(2\pi)$.

Let $H^2(\mathbb{R})$ be the Hardy space of functions in $L^2(\mathbb{R})$ whose Fourier transform is supported in $[0, +\infty)$, endowed with the L^2 norm (hence $\|f\|_{H^2} = \|f\|_{L^2}$). In particular, $\psi_\beta \in H^2(\mathbb{R})$. For $f \in H^2(\mathbb{R})$ we consider the corresponding *wavelet transform* $\mathcal{W}_{\psi_\beta} f$, defined by

$$\mathcal{W}_{\psi_\beta} f(x, y) = \frac{1}{\sqrt{y}} \int_{\mathbb{R}} f(t) \overline{\psi_\beta\left(\frac{t-x}{y}\right)} dt$$

for $(x, y) \in \mathbb{R} \times \mathbb{R}_+$.

Consider now the measure $d\nu = y^{-2} dx dy$, that is the left Haar measure on $\mathbb{R} \times \mathbb{R}_+ \simeq \mathbb{C}_+$ regarded as the “ $ax + b$ ” group. Then, as anticipated, $\mathcal{W}_{\psi_\beta} : H^2(\mathbb{R}) \rightarrow L^2(\mathbb{R} \times \mathbb{R}_+, d\nu)$ is an isometry (not onto); cf. [33].

For $F \in L^p(\mathbb{R} \times \mathbb{R}_+, d\nu)$, $1 \leq p \leq \infty$ and $f, g \in H^2(\mathbb{R})$ we define the *wavelet localization operator* $L_{F, \beta}$ by

$$\langle L_{F, \beta} f, g \rangle_{L^2} = \int_{\mathbb{R} \times \mathbb{R}_+} F \mathcal{W}_{\psi_\beta} f \overline{\mathcal{W}_{\psi_\beta} g} d\nu.$$

Spectral properties of such operator have been studied by many authors; see e.g. [21, 22, 56] and the references therein. Our goal is to find optimal estimates

$$(5.1) \quad \|L_{F, \beta}\|_{H^2 \rightarrow H^2} \leq C = C(p, A, B, \beta),$$

where $F \in L^p(\mathbb{C}_+, d\nu)$ is subject to the double constraint

$$(5.2) \quad \|F\|_{L^\infty(\mathbb{C}_+)} \leq A \quad \text{and} \quad \|F\|_{L^p(\mathbb{C}_+, d\nu)} \leq B$$

(cfr. (1.5) and (1.6)), where p , A , and B are as in Theorem 2.1.

We will need the following result from [52], which is the analogue of Theorem 2.3 for the wavelet transform:

Theorem 5.1. *Let $\beta > 0$. For every ν -measurable set $\Delta \subset \mathbb{R} \times \mathbb{R}_+$ such that $\nu(\Delta) < \infty$, and every $f \in H^2(\mathbb{R})$ such that $\|f\|_{L^2} = 1$, we have*

$$\int_{\Delta} |\mathcal{W}_{\psi_\beta} f|^2 d\nu \leq G_H(\nu(\Delta)),$$

where

$$(5.3) \quad G_H(s) = 1 - \left(1 + \frac{s}{4\pi}\right)^{-2\beta}, \quad s > 0.$$

Moreover, if $\nu(\Delta) > 0$, equality occurs if and only if $f(t)$ has the form

$$(5.4) \quad t \mapsto \frac{c}{\sqrt{y_0}} \psi_\beta\left(\frac{t - x_0}{y_0}\right)$$

for some $c \in \mathbb{C}$ such that $|c|^2 = 2\pi/\beta$ and some $(x_0, y_0) \in \mathbb{R} \times \mathbb{R}_+$, and Δ is ν -equivalent to a hyperbolic disc of center $z_0 = x_0 + iy_0$.

We recall that in the Poincaré upper half-plane \mathbb{C}_+ , endowed with the hyperbolic metric $y^{-2}(dx^2 + dy^2)$, the open hyperbolic disc of center z_0 and ν -measure $s > 0$ is given by

$$(5.5) \quad \left| \frac{z - z_0}{z - \bar{z}_0} \right|^2 < 1 - \left(1 + \frac{s}{4\pi} \right)^{-1}.$$

This is most easily checked in the Poincaré unit disc $|w| < 1$, with the metric $4(1 - |w|^2)^{-2}|dw|^2$ and corresponding measure $4(1 - |w|^2)^{-2}dA(w)$, which is isometric to the above upper half-plane via the map $w = (z - i)/(z + i)$ (by a Möbius transformation one can reduce to the case $z_0 = i$ and the hyperbolic disc $|z - i|/|z + i| < r$ is mapped to $|w| < r$).

To state the main result (the analogue of Theorem 1.1 for the wavelet transform), it is convenient to define the constants

$$(5.6) \quad \sigma = \frac{p - 1}{2\beta p + 1}, \quad \alpha = \frac{p - 1}{2\beta + 1}.$$

Theorem 5.2. *Assume $\beta > 0$, let A , B and p be as in Theorem 2.1, and consider $F \in L^p(\mathbb{C}_+, d\nu)$ satisfying (5.2).*

(i) *If $p = 1$ (and hence $A < \infty$), then*

$$(5.7) \quad \|L_{F,\beta}\|_{H^2 \rightarrow H^2} \leq A G_H(B/A),$$

and equality occurs if and only if $F = e^{i\theta} A \chi_{\mathcal{B}}$ for some $\theta \in \mathbb{R}$, where $\mathcal{B} \subset \mathbb{C}_+$ is a hyperbolic disc of measure $\nu(\mathcal{B}) = B/A$, centered at some $z_0 \in \mathbb{C}_+$.

(ii) *If $p > 1$ and $(B/A)^p \leq 4\pi\sigma$ (in particular if $A = \infty$), then*

$$(5.8) \quad \|L_{F,\beta}\|_{H^2 \rightarrow H^2} \leq \frac{2\beta}{(4\pi)^{1/p}} \sigma^{\kappa_p} B,$$

with equality if and only if

$$(5.9) \quad F(z) = e^{i\theta} \lambda \left(1 - \left| \frac{z - z_0}{z - \bar{z}_0} \right|^2 \right)^{-\alpha}, \quad z \in \mathbb{C}_+,$$

for some $\theta \in \mathbb{R}$ and $z_0 \in \mathbb{C}_+$, with $\lambda = B(4\pi\sigma)^{-1/p}$.

(iii) *If $p > 1$ and $(B/A)^p > 4\pi\sigma$, then*

$$(5.10) \quad \|L_{F,\beta}\|_{H^2 \rightarrow H^2} \leq A \left(1 - p^{2\beta} \left(\frac{\sigma}{\alpha} \right)^{2\beta+1} \left(1 + \frac{(B/A)^p}{4\pi} \right)^{-2\beta} \right),$$

with equality if and only if

$$(5.11) \quad F(z) = e^{i\theta} \min \left\{ \lambda \left(1 - \left| \frac{z - z_0}{z - \bar{z}_0} \right|^2 \right)^{-\alpha}, A \right\}, \quad z \in \mathbb{C}_+,$$

for some $\theta \in \mathbb{R}$ and $z_0 \in \mathbb{C}_+$, with λ as in (5.20).

In the above cases where equality occurs, we have $\|F\|_{L^p} = B$, and there holds $\|L_{F,\beta}\|_{H^2 \rightarrow H^2} = |\langle L_{F,\beta} f, g \rangle|$ for some $f, g \in H^2(\mathbb{R})$ such that $\|f\|_{L^2} = \|g\|_{L^2} = 1$, if and only if both f and g have the form as in (5.4), possibly with different c 's, but with the same $(x_0, y_0) \in \mathbb{R} \times \mathbb{R}_+$ defined by $z_0 = x_0 + iy_0$.

In particular, for $p = 1$ we, therefore, obtain the estimate

$$\|L_{F,\beta}\|_{H^2 \rightarrow H^2} \leq \|F\|_{L^\infty} \left(1 - \left(1 + \frac{\|F\|_{L^\infty}/\|F\|_{L^1}}{4\pi}\right)^{-2\beta}\right),$$

which is the wavelet counterpart of (1.19).

As for the short-time Fourier transform, an intermediate step towards the proof of this theorem is the following result (the analogue of Theorem 2.2 for the wavelet transform):

Theorem 5.3. *Assume $F \in L^p(\mathbb{C}_+, d\nu)$ for some $p \in [1, \infty)$, and let*

$$(5.12) \quad \mu(t) = \nu(\{|F(z)| > t\}), \quad t > 0,$$

denote the distribution function of $|F|$. Then

$$(5.13) \quad \|L_{F,\beta}\|_{H^2 \rightarrow H^2} \leq \int_0^\infty G_H(\mu(t)) dt$$

where G_H is as in (5.3), with equality if and only if $F(z) = e^{i\theta}|F(z)|$ and $|F(z)| = \rho(|z - z_0|^2/|z - \bar{z}_0|^2)$, for some $\theta \in \mathbb{R}$, some $z_0 \in \mathbb{C}_+$ and some decreasing function $\rho : [0, 1) \rightarrow [0, +\infty)$. In this case, there holds

$$(5.14) \quad |\langle L_F f, g \rangle| = \|L_F\|_{L^2 \rightarrow L^2}$$

for some $f, g \in H^2(\mathbb{R}^d)$ such that $\|f\|_{L^2} = \|g\|_{L^2} = 1$, if (and only if, when F is not identically zero) both f and g are functions of the kind (5.4), possibly with different c 's, but with the same (x_0, y_0) defined by $x_0 + iy_0 = z_0$.

The proof of the last two theorems follows the same pattern as the proof of Theorems 1.1 and 2.2, through the study of a variational problem similar to (3.4) and results analogous to those of Section 3. We omit the details, limiting ourselves to describing how the proofs can be adapted to the hyperbolic setting.

Apart from obvious changes (such as replacing \mathbb{R}^{2d} with \mathbb{C}_+ and dz with $d\nu(z)$ in the integrals etc.), Theorem 5.3 is proved by the same slicing argument as Theorem 2.2, now using Theorem 5.1 in place of Theorem 2.3, so that F achieves equality in (5.13) if and only if $F(z) = e^{i\theta}|F(z)|$ and all the superlevel sets $\{|F| > t\}$ are hyperbolic discs centered at some $z_0 \in \mathbb{C}_+$. Since every such disc has the form as in (5.5), reconstructing (for a.e. z) $|F(z)|$ as in (2.14), we see that $|F(z)|$ only depends on $|z - z_0|^2/|z - \bar{z}_0|^2$, and hence $|F(z)| = \rho(|z - z_0|^2/|z - \bar{z}_0|^2)$ for some (necessarily decreasing) function $\rho : [0, 1) \rightarrow [0, +\infty)$, as claimed.

Subsequently, to estimate the right-hand side of (5.13), one studies the variational problem

$$(5.15) \quad \sup_{v \in \mathcal{C}} I(v) \quad \text{where} \quad I(v) := \int_0^A G_H(v(t)) dt$$

(the analogue of (3.4), with the same \mathcal{C} as in (3.2)), and proceeding as in Section 3 one proves a result similar to Theorem 3.4, i.e. the existence of a

unique maximizer $u \in \mathcal{C}$, that we shall identify distinguishing three different cases as (i), (ii) and (iii) of Theorem 5.2. Then, by (5.13),

$$(5.16) \quad \|L_{F,\beta}\|_{H^2 \rightarrow H^2} \leq I(\mu) \leq \sup_{v \in \mathcal{C}} I(v) = I(u).$$

When $p = 1$ (case (i)) $u = B/A$ is constant, again by Jensen inequality (note that $G_H(s)$ is strictly convex), and this eventually leads to (5.7).

When $p > 1$, arguing exactly as for (3.19) one proves that u satisfies $G'_H(u(t)) = ct^{p-1}$ on the interval $(0, M)$ where $u > 0$, which can be rewritten as $G'_H(u(t)) = 2\beta(t/\lambda)^{p-1}$ on $(0, M)$ for some $\lambda > 0$ (cfr. (3.20)). Then, reasoning as after (3.22), one proves the continuity of u on $(0, A)$ also when $M < A$, so that computing $G'_H(s)$ from (5.3) and solving for $u(t)$, recalling (5.6) one eventually obtains

$$(5.17) \quad u(t) = u_\lambda(t) = 4\pi \max\{(t/\lambda)^{-\alpha} - 1, 0\}, \quad t \in (0, A),$$

cfr. (3.24). The value of $\lambda > 0$ can be uniquely determined by the condition that the constraint (3.1) is saturated, i.e. by solving the equation

$$(5.18) \quad B^p = h(\lambda) := p \int_0^A t^{p-1} u_\lambda(t) dt$$

(again, h is strictly increasing). Since now (cfr. (3.25)), with the notation as in (5.6), for $A < \infty$ we have

$$h(A) = 4\pi p \int_0^A t^{p-1} \left((t/A)^{-\alpha} - 1 \right) dt = \frac{4\pi(p-1)}{2\beta p + 1} A^p = 4\pi\sigma A^p,$$

we see from (5.18) that $\lambda \leq A$ when $(B/A)^p \leq 4\pi\sigma$ (case (ii)), while $\lambda > A$ when $(B/A)^p > 4\pi\sigma$ (case (iii)). In case (ii), by (5.17), (5.18) becomes

$$B^p = h(\lambda) = 4\pi p \int_0^\lambda t^{p-1} \left((t/\lambda)^{-\alpha} - 1 \right) dt = 4\pi\sigma\lambda^p,$$

whence

$$(5.19) \quad \lambda = B(4\pi\sigma)^{-1/p}.$$

Moreover, since $G_H(0) = 0$, an explicit computation based on (5.15) and (5.17) gives

$$\begin{aligned} I(u_\lambda(t)) &= \int_0^\lambda \left(1 - ((t/\lambda)^{-\alpha})^{-2\beta} \right) dt = \lambda \int_0^1 \left(1 - \tau^{2\alpha\beta} \right) d\tau \\ &= \lambda \left(1 - \frac{1}{2\alpha\beta + 1} \right) = \frac{2\beta}{(4\pi)^{1/p}} \sigma^{1-1/p} B, \end{aligned}$$

and (5.8) follows from (5.16).

Similarly, in case (iii) where $(B/A)^p > \frac{4\pi(p-1)}{2\beta p + 1}$ and $\lambda > A$, (5.17) simplifies to $u(t) = 4\pi((t/\lambda)^{-\alpha} - 1)$, so that (5.18) becomes

$$B^p = 4\pi p \int_0^A t^{p-1} \left((t/\lambda)^{-\alpha} - 1 \right) dt = 4\pi A^p \left(\frac{p\sigma}{\alpha} \left(\frac{A}{\lambda} \right)^{-\alpha} - 1 \right)$$

whence

$$(5.20) \quad \lambda = A \left(\frac{4\pi A^p p \sigma}{\alpha(B^p + 4A^p \pi)} \right)^{-\frac{1}{\alpha}}$$

and, by an elementary computation,

$$I(u_\lambda) = \int_0^A \left(1 - ((t/\lambda)^{-\alpha})^{-2\beta} \right) dt = A \left(1 - p^{2\beta} \left(\frac{\sigma}{\alpha} \right)^{2\beta+1} \left(1 + \frac{(B/A)^p}{4\pi} \right)^{-2\beta} \right).$$

Then (5.10) follows from (5.16).

Finally, F attains equality in (5.7), (5.8), or (5.10) (i.e. equalities in (5.16)), if and only if equality occurs in (5.13) and $\mu(t)$ coincides with $u(t)$ (extended to $u(t) = 0$ for $t \geq A$, if $A < \infty$). This, by Theorem 5.3, means that $F = e^{i\theta}|F|$ and $|F(z)| = \rho(|z - z_0|^2/|z - \bar{z}_0|^2)$ (where $z_0 \in \mathbb{C}_+$ and $\rho : [0, 1) \rightarrow [0, +\infty)$ is decreasing), thus we can reconstruct $|F(z)|$ by knowledge of its distribution function $\mu(t) = u(t)$. Indeed, since for every $t \in (0, A)$ the superlevel sets $\{|F| > t\}$ (being hyperbolic discs centered at z_0) have the form as in (5.5) with $s = u(t)$, i.e.

$$\{|F(z)| > t\} = \left\{ z \in \mathbb{C}_+ \quad \text{such that} \quad \left| \frac{z - z_0}{z - \bar{z}_0} \right|^2 < 1 - \left(1 + \frac{u(t)}{4\pi} \right)^{-1} \right\},$$

using the layer cake representation (2.14) (with $m = A$ and $z \in \mathbb{C}_+$) one obtains, for a.e. $z \in \mathbb{C}_+$,

$$|F(z)| = \sup \left\{ t \in (0, A) \quad \text{such that} \quad \left| \frac{z - z_0}{z - \bar{z}_0} \right|^2 < 1 - \left(1 + \frac{u(t)}{4\pi} \right)^{-1} \right\}$$

In case (i), when $p = 1$ and $u(t) = B/A$ is constant on $(0, A)$, we see that $F(z)$ is as claimed after (5.7). On the other hand, when $p > 1$, an elementary computation based on (5.17) reveals that

$$|F(z)| = \min \left\{ \lambda \left(1 - \left| \frac{z - z_0}{z - \bar{z}_0} \right|^2 \right)^{\frac{1}{\alpha}}, A \right\}.$$

In case (ii), where $\lambda \leq A$ (as given in (5.19)), since $F = e^{i\theta}|F|$ one obtains (5.9). Similarly, in case (iii) where $\lambda > A$ (as given in (5.20)), one obtains (5.11).

REFERENCES

- [1] L. D. Abreu and M. Dörfler. An inverse problem for localization operators. *Inverse Problems*, 28(11):115001, 16, 2012.
- [2] L. D. Abreu, K. Gröchenig, and J. L. Romero. On accumulated spectrograms. *Trans. Amer. Math. Soc.*, 368(5):3629–3649, 2016.
- [3] L. D. Abreu, Z. Mouayn, and F. Voigtlaender. A fractal uncertainty principle for Bergman spaces and analytic wavelets. *J. Math. Anal. Appl.*, 519(1):Paper No. 126699, 17, 2023.
- [4] L. D. Abreu, J. M. Pereira, and J. L. Romero. Sharp rates of convergence for accumulated spectrograms. *Inverse Problems*, 33(11):115008, 12, 2017.
- [5] L. D. Abreu and M. Speckbacher. Deterministic guarantees for L^1 -reconstruction: A large sieve approach with geometric flexibility. *IEEE Proceedings SampTA*, 2019.

- [6] L. D. Abreu and M. Speckbacher. Donoho-Logan large sieve principles for modulation and polyanalytic Fock spaces. *Bull. Sci. Math.*, 171:Paper No. 103032, 25, 2021.
- [7] L. D. Abreu and M. Speckbacher. Donoho-Logan large sieve principles for the continuous wavelet transform, arXiv:2210.13056, 2022.
- [8] A. Baernstein, II. *Symmetrization in analysis*, volume 36 of *New Mathematical Monographs*. Cambridge University Press, Cambridge, 2019. With David Drasin and Richard S. Laugesen, With a foreword by Walter Hayman.
- [9] D. Bayer and K. Gröchenig. Time-frequency localization operators and a Berezin transform. *Integral Equations Operator Theory*, 82(1):95–117, 2015.
- [10] F. A. Berezin. Wick and anti-Wick symbols of operators. *Mat. Sb. (N.S.)*, 86(128):578–610, 1971.
- [11] F. A. Berezin and M. A. Shubin. *The Schrödinger equation*, volume 66 of *Mathematics and its Applications (Soviet Series)*. Kluwer Academic Publishers Group, Dordrecht, 1991. Translated from the 1983 Russian edition by Yu. Rajabov, D. A. Leites and N. A. Sakharova and revised by Shubin. With contributions by G. L. Litvinov and Leites.
- [12] P. Boggiatto, E. Carypis, and A. Oliaro. Two aspects of the Donoho-Stark uncertainty principle. *J. Math. Anal. Appl.*, 434(2):1489–1503, 2016.
- [13] A. Bonami and B. Demange. A survey on uncertainty principles related to quadratic forms. *Collect. Math.*, (Vol. Extra):1–36, 2006.
- [14] O. F. Brevig, J. Ortega-Cerdà, K. Seip, and J. Zhao. Contractive inequalities for Hardy spaces. *Funct. Approx. Comment. Math.*, 59(1):41–56, 2018.
- [15] E. A. Carlen. Some integral identities and inequalities for entire functions and their application to the coherent state transform. *J. Funct. Anal.*, 97(1):231–249, 1991.
- [16] E. Cordero and K. Gröchenig. Time-frequency analysis of localization operators. *J. Funct. Anal.*, 205(1):107–131, 2003.
- [17] E. Cordero and F. Nicola. Sharp continuity results for the short-time Fourier transform and for localization operators. *Monatsh. Math.*, 162(3):251–276, 2011.
- [18] M. G. Cowling and J. F. Price. Bandwidth versus time concentration: the Heisenberg-Pauli-Weyl inequality. *SIAM J. Math. Anal.*, 15(1):151–165, 1984.
- [19] I. Daubechies. Time-frequency localization operators: a geometric phase space approach. *IEEE Trans. Inform. Theory*, 34(4):605–612, 1988.
- [20] I. Daubechies. The wavelet transform, time-frequency localization and signal analysis. *IEEE Trans. Inform. Theory*, 36(5):961–1005, 1990.
- [21] I. Daubechies. *Ten lectures on wavelets*, volume 61 of *CBMS-NSF Regional Conference Series in Applied Mathematics*. Society for Industrial and Applied Mathematics (SIAM), Philadelphia, PA, 1992.
- [22] I. Daubechies and T. Paul. Time-frequency localisation operators—a geometric phase space approach. II. The use of dilations. *Inverse Problems*, 4(3):661–680, 1988.
- [23] M. A. de Gosson. *Symplectic methods in harmonic analysis and in mathematical physics*, volume 7 of *Pseudo-Differential Operators. Theory and Applications*. Birkhäuser/Springer Basel AG, Basel, 2011.
- [24] N. C. Dias, F. Luef, and J. a. N. Prata. Uncertainty principle via variational calculus on modulation spaces. *J. Funct. Anal.*, 283(8):Paper No. 109605, 30, 2022.
- [25] C. L. Fefferman. The uncertainty principle. *Bull. Amer. Math. Soc. (N.S.)*, 9(2):129–206, 1983.
- [26] C. Fernández and A. Galbis. Compactness of time-frequency localization operators on $L^2(\mathbb{R}^d)$. *J. Funct. Anal.*, 233(2):335–350, 2006.
- [27] C. Fernández and A. Galbis. Annihilating sets for the short time Fourier transform. *Adv. Math.*, 224(5):1904–1926, 2010.
- [28] G. B. Folland and A. Sitaram. The uncertainty principle: a mathematical survey. *J. Fourier Anal. Appl.*, 3(3):207–238, 1997.

- [29] A. Galbis. Norm estimates for selfadjoint Toeplitz operators on the Fock space. *Complex Anal. Oper. Theory*, 16(1):Paper No. 15, 13, 2022.
- [30] Y. V. Galperin and K. Gröchenig. Uncertainty principles as embeddings of modulation spaces. *J. Math. Anal. Appl.*, 274(1):181–202, 2002.
- [31] K. Gröchenig. *Foundations of time-frequency analysis*. Applied and Numerical Harmonic Analysis. Birkhäuser Boston, Inc., Boston, MA, 2001.
- [32] K. Gröchenig. Uncertainty principles for time-frequency representations. In *Advances in Gabor analysis*, Appl. Numer. Harmon. Anal., pages 11–30. Birkhäuser Boston, Boston, MA, 2003.
- [33] A. Grossmann, J. Morlet, and T. Paul. Transforms associated to square integrable group representations. II. Examples. *Ann. Inst. H. Poincaré Phys. Théor.*, 45(3):293–309, 1986.
- [34] G. H. Hardy, J. E. Littlewood, and G. Pólya. *Inequalities*. Cambridge, at the University Press, 1952. 2d ed.
- [35] H. Knutsen. Daubechies’ time-frequency localization operator on Cantor type sets I. *J. Fourier Anal. Appl.*, 26(3):Paper No. 47, 19, 2020.
- [36] H. Knutsen. Daubechies’ time-frequency localization operator on Cantor type sets II. *J. Funct. Anal.*, 282(9):Paper No. 109412, 28, 2022.
- [37] H. Knutsen. A fractal uncertainty principle for the short-time Fourier transform and Gabor multipliers. *Appl. Comput. Harmon. Anal.*, 62:365–389, 2023.
- [38] A. Kulikov. Functionals with extrema at reproducing kernels. *Geom. Funct. Anal.*, 32(4):938–949, 2022.
- [39] H. J. Landau. An overview of time and frequency limiting. In *Fourier techniques and applications (Kensington, 1983)*, pages 201–220. Plenum, New York, 1985.
- [40] H. J. Landau and H. O. Pollak. Prolate spheroidal wave functions, Fourier analysis and uncertainty. II. *Bell System Tech. J.*, 40:65–84, 1961.
- [41] E. H. Lieb. Proof of an entropy conjecture of Wehrl. *Comm. Math. Phys.*, 62(1):35–41, 1978.
- [42] E. H. Lieb. Integral bounds for radar ambiguity functions and Wigner distributions. *J. Math. Phys.*, 31(3):594–599, 1990.
- [43] E. H. Lieb and M. Loss. *Analysis*, volume 14 of *Graduate Studies in Mathematics*. American Mathematical Society, Providence, RI, second edition, 2001.
- [44] E. H. Lieb and J. P. Solovej. Wehrl-type coherent state entropy inequalities for $SU(1, 1)$ and its $AX + B$ subgroup. In *Partial differential equations, spectral theory, and mathematical physics—the Ari Laptev anniversary volume*, EMS Ser. Congr. Rep., pages 301–314. EMS Press, Berlin, 2021.
- [45] F. Luef and E. Skrettingland. Convolutions for localization operators. *J. Math. Pures Appl. (9)*, 118:288–316, 2018.
- [46] F. Luef and E. Skrettingland. On accumulated Cohen’s class distributions and mixed-state localization operators. *Constr. Approx.*, 52(1):31–64, 2020.
- [47] S. Mallat. *A wavelet tour of signal processing. The sparse way*. Elsevier/Academic Press, Amsterdam, third edition, 2009. With contributions from Gabriel Peyré.
- [48] C. Muscalu and W. Schlag. *Classical and multilinear harmonic analysis. Vol. I*, volume 137 of *Cambridge Studies in Advanced Mathematics*. Cambridge University Press, Cambridge, 2013.
- [49] F. Nicola, J. L. Romero, and S. I. Trapasso. On the existence of optimizers for time-frequency concentration problems. *Calc. Var. Partial Differential Equations*, 62(1):Paper No. 21, 21, 2023.
- [50] F. Nicola and P. Tilli. The Faber-Krahn inequality for the short-time Fourier transform. *Invent. Math.*, 230(1):1–30, 2022.
- [51] F. Nicola and S. I. Trapasso. A note on the HRT conjecture and a new uncertainty principle for the short-time Fourier transform. *J. Fourier Anal. Appl.*, 26(4):Paper No. 68, 7, 2020.

- [52] J. Ramos and P. Tilli. A Faber-Krahn inequality for wavelet transforms, *Bull. Lond. Math. Soc.*, to appear.
- [53] K. Seip. Reproducing formulas and double orthogonality in Bargmann and Bergman spaces. *SIAM J. Math. Anal.*, 22(3):856–876, 1991.
- [54] D. Slepian. Some comments on Fourier analysis, uncertainty and modeling. *SIAM Rev.*, 25(3):379–393, 1983.
- [55] A. Wigderson and Y. Wigderson. The uncertainty principle: variations on a theme. *Bull. Amer. Math. Soc. (N.S.)*, 58(2):225–261, 2021.
- [56] M. W. Wong. *Wavelet transforms and localization operators*, volume 136 of *Operator Theory: Advances and Applications*. Birkhäuser Verlag, Basel, 2002.
- [57] K. Zhu. *Analysis on Fock spaces*, volume 263 of *Graduate Texts in Mathematics*. Springer, New York, 2012.

DIPARTIMENTO DI SCIENZE MATEMATICHE, POLITECNICO DI TORINO, CORSO DUCA
DEGLI ABRUZZI 24, 10129 TORINO, ITALY.

Email address: `fabio.nicola@polito.it`

DIPARTIMENTO DI SCIENZE MATEMATICHE, POLITECNICO DI TORINO, CORSO DUCA
DEGLI ABRUZZI 24, 10129 TORINO, ITALY.

Email address: `paolo.tilli@polito.it`