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# Pseudo-differential operators with isotropic symbols, Wick and anti-Wick operators, and hypoellipticity 

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#### Abstract

We study the link between $\Psi$ dos and Wick operators via the Bargmann transform. We deduce a formula for the symbol of the Wick operator in terms of the short-time Fourier transform of the Weyl symbol. This gives characterizations of Wick symbols of $\Psi$ dos of Shubin type and of infinite order, and results on composition. We prove a series expansion of Wick operators in terms of anti-Wick operators which leads to a sharp Gårding inequality and transition of hypoellipticity between Wick and Shubin symbols. Finally we show continuity results for anti-Wick operators, and estimates for the Wick symbols of anti-Wick operators. © 2022 The Author(s). Published by Elsevier Masson SAS. This is an open access article under the CC BY license (http://creativecommons.org/licenses/by/4.0/).


R É S U M É
Nous étudions le lien entre $\Psi$ dos et opérateurs Wick passant par la transformée de Bargmann. Nous déduisons une formule pour le symbole de l'opérateur de Wick en termes de la transformée de Fourier à court terme du symbole de Weyl. Ceci donne des caractérisations des symboles Wick de $\Psi$ dos de type Shubin et d'ordre infini, et des résultats sur composition. Nous prouvons un développemet en série des opérateurs Wick dans opérateurs anti-Wick qui conduit à une forte inégalité Gårding et transition de hypoellipticité entre les symboles Wick et Shubin. Enfin, nous montrons des résultats de continuité pour opérateurs anti-Wick et des estimations pour les symboles Wick des opérateurs anti-Wick.
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## 0. Introduction

In the paper we investigate conjugation with the Bargmann transformation of pseudo-differential and Toeplitz operators on $\mathbf{R}^{d}$ with isotropic symbols, and we explore relations between Wick and anti-Wick operators. Particularly we consider Shubin operators and operators of infinite order. This gives rise to analytic type pseudo-differential operators on $\mathbf{C}^{d}$ that are called Wick or Berezin operators because of the fundamental contributions by F. Berezin [5,6], which in turns goes back to some ideas in [33] by G. C. Wick.

Let $a$ be a suitable locally bounded function on $\mathbf{C}^{2 d}$ such that $z \mapsto a(z, w)$ is analytic, $z, w \in \mathbf{C}^{d}$. Then the Wick operator $\mathrm{Op}_{\mathfrak{V}}(a)$ with symbol $a$ is the operator which takes an appropriate entire function $F$ on $\mathbf{C}^{d}$ into the entire function

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}(a) F(z)=\pi^{-d} \int_{\mathbf{C}^{d}} a(z, w) F(w) e^{(z-w, w)} d \lambda(w), \tag{0.1}
\end{equation*}
$$

where $d \lambda$ is the Lebesgue measure and $(\cdot, \cdot)$ is the scalar product on $\mathbf{C}^{d}$. (See [17] and Section 1 for notation.) Wick operators appear naturally in several problems in analysis and its applications, e.g. in quantum mechanics. For example, the harmonic oscillator, the creation and annihilation operators take the simple forms

$$
F \mapsto\left\langle z, \nabla_{z}\right\rangle F+c F, \quad F \mapsto z_{j} F \quad \text { and } \quad F \mapsto \partial_{z_{j}} F
$$

respectively, for some constant $c$, in the Wick formulation (see [3]).
An advantage of the Wick calculus compared to corresponding operators on functions and distributions defined on $\mathbf{R}^{d}$ is that in almost all situations, the involved functions are entire, which admits the use of the powerful techniques of complex analysis. (A more general approach is studied in [29], where the Wick calculus is formulated in terms of spaces of formal power series expansions instead of spaces of entire functions.) The possible lack of analyticity of $a(z, w)$ in (0.1) with respect to the $w$ variable is removable in the sense that for any suitable Wick symbol $a$, there is a unique $a_{0}$ such that $(z, w) \mapsto a_{0}(z, \bar{w})$ is entire, and $\mathrm{Op}_{\mathfrak{V}}(a)=\mathrm{Op}_{\mathfrak{V}}\left(a_{0}\right)$ (see e.g. [29,30]). Consequently it is no restriction to assume that $a(z, w)$ in (0.1) is analytic in $z$ and conjugate analytic in $w$, which we do in the introduction henceforth. Any linear and continuous operator from the Schwartz space, a Fourier invariant Gelfand-Shilov space or Pilipović space, to the corresponding distribution spaces, respectively, is in a unique way transformed into a Wick operator by the Bargmann transform (see [29]).

Several operators in quantum mechanics are so-called Shubin operators, i.e. pseudo-differential operators

$$
\operatorname{Op}(\mathfrak{a}) f(x)=(2 \pi)^{-\frac{d}{2}} \int_{\mathbf{R}^{d}} \mathfrak{a}(x, \xi) \hat{f}(\xi) e^{i\langle x, \xi\rangle} d \xi, \quad f \in \mathscr{S}\left(\mathbf{R}^{d}\right)
$$

where the symbol $\mathfrak{a}$ belongs to the Shubin class $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$, the set of all $\mathfrak{a} \in C^{\infty}\left(\mathbf{R}^{2 d}\right)$ such that

$$
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} \mathfrak{a}(x, \xi)\right| \lesssim \omega(x, \xi)(1+|x|+|\xi|)^{-\rho|\alpha+\beta|}, \quad \alpha, \beta \in \mathbf{N}^{d}
$$

(See e.g. [17,26].) Here $\omega$ is a suitable weight function on $\mathbf{R}^{2 d}$ and $0 \leqslant \rho \leqslant 1$. Partial differential operators with polynomial coefficients, e.g. the creation and annihilation operators or the harmonic oscillator mentioned above, are examples of Shubin operators. In Section 2 we prove that the Bargmann image of Shubin operators with symbols in $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ is the set of all Wick operators in (0.1) such that $a$ belongs to $\mathcal{A}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$. This means that $\mathbf{C}^{2 d} \ni(z, w) \mapsto a(z, \bar{w})$ is an entire function that satisfies

$$
\begin{equation*}
\left|\partial_{z}^{\beta} \bar{\partial}_{w}^{\gamma} a(z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z+w\rangle^{-\rho|\beta+\gamma|}\langle z-w\rangle^{-N} \tag{0.2}
\end{equation*}
$$

for every $N \geqslant 0$.
An important subclass of Wick operators is the anti-Wick operators, which are Wick operators where the symbol $a(z, w)$ does not depend on $z$. That is, for an appropriate measurable function $a_{0}$ on $\mathbf{C}^{d}$, its anti-Wick operator is given by

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{Z}}^{\mathrm{aw}}\left(a_{0}\right) F(z)=\pi^{-d} \int_{\mathbf{C}^{d}} a_{0}(w) F(w) e^{(z-w, w)} d \lambda(w) . \tag{0.1}
\end{equation*}
$$

Again $F$ is a suitable entire function on $\mathbf{C}^{d}$. The anti-Wick operators can also be described as the Bargmann images of Toeplitz operators on $\mathbf{R}^{d}$. (See e.g. $[20,26,30]$ for the definition of Toeplitz operators.)

A feature of Toeplitz operators and anti-Wick operators, useful for energy estimates in quantum mechanics and time-frequency analysis, is that non-negative symbols give rise to non-negative operators. (Cf. e.g. [20,18,19].) An operator $T=\operatorname{Op}_{\mathfrak{V}}(a)$ with $a$ satisfying ( 0.2 ) for every $N \geqslant 0$, is called positive (nonnegative), if there is a constant $C>0(C \geqslant 0)$ such that

$$
(T F, F)_{A^{2}} \geqslant C\|F\|_{A^{2}}^{2}
$$

for every analytic polynomial $F$ on $\mathbf{C}^{d}$, where $(\cdot, \cdot)_{A^{2}}$ is the scalar product induced by the Hilbert norm

$$
\|F\|_{A^{2}}=\pi^{-\frac{d}{2}}\left(\int_{\mathbf{C}^{d}}|F(z)|^{2} e^{-|z|^{2}} d \lambda(z)\right)^{\frac{1}{2}}
$$

The implication from non-negative symbols to non-negative operators is not relevant for Wick operators in (0.1) when $a(z, w)$ is not constant with respect to $z$, since the analyticity of the map $z \mapsto a(z, w)$ implies that $a(z, w)$ is non-real almost everywhere. For such symbols it is instead natural to check whether positivity of the map $w \mapsto a(w, w)$ leads to positive operators (see e.g. [5,6,12]). By choosing

$$
d=1, \quad a(z, w)=1-2 z \bar{w}+2 z^{2} \bar{w}^{2} \quad \text { and } \quad F(z)=z
$$

we obtain

$$
a(w, w)=\left(1-|w|^{2}\right)^{2}+|w|^{4}>0 \quad \text { but } \quad\left(\mathrm{Op}_{\mathfrak{V}}(a) F, F\right)_{A^{2}}=-1<0 .
$$

Consequently $\mathrm{Op}_{\mathfrak{V}}(a)$ may fail to be a non-negative operator even though $a(w, w)$ is positive.
On the other hand, for certain conditions on $a$, we deduce in Section 3 a weaker positivity result for Wick operators, which is equivalent to the sharp Gårding inequality in isotropic pseudo-differential calculus on $\mathbf{R}^{d}$ (see Theorem 18.6.7 and the proof of Theorem 18.6.8 in [17]). That is, for $a \in \widehat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ with $\omega(z)=\langle z\rangle^{2 \rho}$ and $\rho>0$ we prove

$$
\begin{equation*}
\operatorname{Re}\left(\mathrm{Op}_{\mathfrak{V}}(a) F, F\right)_{A^{2}} \geqslant-C\|F\|_{A^{2}}^{2} \tag{0.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\operatorname{Im}\left(\mathrm{Op}_{\mathfrak{V}}(a) F, F\right)_{A^{2}}\right| \leqslant C\|F\|_{A^{2}}^{2}, \quad \text { when } \quad a(w, w) \geqslant 0 \tag{0.4}
\end{equation*}
$$

(cf. Theorem 4.2). In particular we obtain energy estimates also for Wick operators with symbols that are non-negative on the diagonal.

The latter result is obtained by approximating Wick operators by anti-Wick operators, using for the Wick operator (0.1) with $a \in \mathcal{A}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ the remarkable identity

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}(a)=\sum_{|\alpha|<M} \frac{(-1)^{|\alpha|}}{\alpha!} \mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(b_{\alpha}\right)+\mathrm{Op}_{\mathfrak{V}}\left(c_{M}\right) \quad \text { where } \quad b_{\alpha}(w)=\partial_{z}^{\alpha} \bar{\partial}_{w}^{\alpha} a(w, w), \tag{0.5}
\end{equation*}
$$

for some $c_{M} \in \mathcal{A}_{\mathrm{Sh}, \rho}^{\left(\omega_{M}\right)}\left(\mathbf{C}^{2 d}\right)$ with $\omega_{M}(z)=\omega(z)\langle z\rangle^{-2 M \rho}$. Here we again assume $\rho>0$. The decay conditions on $b_{\alpha}$ and $c_{M}$ are, respectively,

$$
\begin{equation*}
\left|\partial_{w}^{\beta} \bar{\partial}_{w}^{\gamma} b_{\alpha}(w)\right| \lesssim \omega(\sqrt{2} \bar{w})\langle w\rangle^{-\rho|2 \alpha+\beta+\gamma|}, \quad \alpha, \beta, \gamma \in \mathbf{N}^{d} \tag{0.6}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\partial_{z}^{\beta} \bar{\partial}_{w}^{\gamma} c_{M}(z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z\rangle^{-2 M \rho}\langle z+w\rangle^{-\rho|\beta+\gamma|}\langle z-w\rangle^{-N}, \tag{0.7}
\end{equation*}
$$

for every $N \geqslant 0$. Consequently, several Wick operators can essentially be expressed as linear combinations of anti-Wick operators. The expansion (0.5) is deduced in Section 3 using Taylor expansion and integration by parts, see Proposition 3.1 and Remark 3.2. The estimates (0.6) and (0.7) are consequences of (2.5)" and Proposition 3.3.

The conditions on $b_{\alpha}$ are the same as the conditions on $a(0.2)$, restricted to the diagonal $z=w$, and with improved decay. On the diagonal, the growth term $e^{\frac{1}{2}|z-w|^{2}}$ disappears, which dominates in (0.2) when $|z-w| \gtrsim|z|$ or $|z-w| \gtrsim|w|$. The right-hand side of (0.6) becomes as large as possible when $\alpha=\beta=\gamma=0$, that is $b_{0}$ is the dominating term in the sum (0.5).

The conditions on $c_{M}$ are the same as the estimates (0.2) again with improved decay due to the factor $\langle z\rangle^{-2 M \rho}$.

For polynomial symbols, (0.5) agrees with the integral formula [5, Theorem 3] due to Berezin which carry over Wick operators into anti-Wick operators. For the general case, (0.5) is analogous to the approximation technique of pseudo-differential operators on $\mathbf{R}^{d}$ in terms of Toeplitz operators given in [26, Theorem 24.1] and its proof, by Shubin.

The anti-Wick symbols in (0.5) $b_{\alpha}(w)=\partial_{z}^{\alpha} \bar{\partial}_{w}^{\alpha} a(w, w)$ extend to have the property that $\partial_{z}^{\alpha} \bar{\partial}_{w}^{\alpha} a(z, w)$ is entire in $z$ and conjugate entire in $w$. Note that restriction to the diagonal also appears in the positivity condition (0.3) on Wick symbols.

The sharp Gårding inequality (0.3) is reached by using the fact that $\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(b_{0}\right)$ is non-negative, and that if $T$ is either $\mathrm{Op}_{\mathfrak{V}}^{\text {aw }}\left(b_{\alpha}\right)$ or $\mathrm{Op}_{\mathfrak{V}}\left(c_{N}\right)$ for $\alpha \neq 0$, then $\|T F\|_{A^{2}} \lesssim\|F\|_{A^{2}}$ when $F \in A\left(\mathbf{C}^{d}\right)$ is a polynomial.

In Section 5 we deduce links concerning ellipticity, hypoellipticity (in Shubin's sense) and weak ellipticity between Shubin and Wick symbols. The notion of hypoelliptic symbol resembles hypoelliptic symbols in Shubin's sense (see [26]). More specifically, we say that the symbol $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ is hypoelliptic of order $\rho_{0} \geqslant 0$, whenever there is an $R>0$ such that

$$
|\mathfrak{a}(x, \xi)| \gtrsim \omega(x, \xi)\langle(x, \xi)\rangle^{-\rho_{0}} \quad \text { and } \quad\left|\partial^{\alpha} \mathfrak{a}(x, \xi)\right| \lesssim|\mathfrak{a}(x, \xi)|\langle(x, \xi)\rangle^{-\rho|\alpha|}
$$

when $|(x, \xi)| \geqslant R$.
A linear operator $T$ from $\mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$ to $\mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$ is called globally hypoelliptic if

$$
T f=g, f \in \mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right), g \in \mathscr{S}\left(\mathbf{R}^{d}\right) \quad \Rightarrow \quad f \in \mathscr{S}\left(\mathbf{R}^{d}\right) .
$$

(See e.g. [11].) It can be proved that a pseudo-differential operator with hypoelliptic symbol in Shubin's sense is globally hypoelliptic as operator (see e.g. [26, Corollary 25.1]).

We show, similarly to our investigations of the sharp Gårding inequality and for expansion (0.5), that ellipticity, hypoellipticity and to some degree weak ellipticity for the Shubin symbol $\mathfrak{a}$ can be characterized by certain conditions for the corresponding Wick symbol $a(z, w)$ along the diagonal $z=w$. For example, let $\mathfrak{a}$ be a polynomial on $\mathbf{R}^{d}$ with principal symbol $\mathfrak{a}_{p}$, and let $a(z, w)$ be a polynomial in $z, \bar{w} \in \mathbf{C}^{d}$ with
principal part $a_{p}$. Then $\mathfrak{a}$ is elliptic means that $\mathfrak{a}_{p}(x, \xi) \neq 0$ when $(x, \xi) \neq(0,0)$, and $a$ is elliptic means that $a_{p}(z, z) \neq 0$ when $z \neq 0$. For such $\mathfrak{a}$ we prove

$$
\mathfrak{a} \text { is elliptic } \Leftrightarrow a \text { is elliptic, }
$$

when $a(z, w)$ is the Wick symbol corresponding to $\mathfrak{a}$ (which must be a polynomial in $z$ and $\bar{w}$ ).
Our investigations include the Bargmann transform of certain operators of infinite order, i.e. pseudodifferential operators with ultra-differentiable symbols that are permitted to grow faster than any polynomial at infinity together with their derivatives. Particularly we consider Wick operators of infinite order, i.e. the Bargmann images $\mathrm{Op}_{\mathfrak{V}}(a)$ of operators $\operatorname{Op}(\mathfrak{a})$ of infinite order in [1], and characterize their images under the Bargmann transform (see Theorem 2.6). Then we deduce in Subsections 3.2 and 3.3 continuity results for anti-Wick operators which holds for the symbols $b_{\alpha}$ in (0.5) when $\mathrm{Op}_{\mathfrak{V}}(a)$ is the Bargmann image of an operator of infinite order.

In fact, in Subsection 3.2 we show that $\mathrm{Op}_{\mathfrak{V} \mathfrak{j}}^{\text {aw }}\left(b_{\alpha}\right)$ possess several other continuity properties than what is valid for $\mathrm{Op}_{\mathfrak{V}}(a)$ in the expansion (0.5) (see Propositions 3.6 and 3.9). In Subsection 3.3 we deduce estimates of the Wick symbol $b_{\alpha}^{\text {aw }}$ to the anti-Wick operator $\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(b_{\alpha}\right)$, i.e. the unique element $b_{\alpha}^{\text {aw }} \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ such that $\mathrm{Op}_{\mathfrak{V}}\left(b_{\alpha}^{\text {aw }}\right)=\mathrm{Op}_{\mathfrak{V} \mathfrak{J}}^{\text {aw }}\left(b_{\alpha}\right)$. We show that usually, $b_{\alpha}^{\text {aw }}$ satisfies stronger conditions than $a$ when $\mathrm{Op}_{\mathfrak{V}}(a)$ is a Wick operator of infinite order (see Theorems 3.11, 3.14 and 3.13).

The paper is organized as follows. In Section 1 we set the stage by providing necessary background notions and fixing the notation. It contains useful properties for weight functions, Gelfand-Shilov spaces, the Bargmann transform, pseudo-differential operators, Wick and anti-Wick operators. Thereafter we characterize in Section 2 Shubin operators and operators of infinite order in terms of appropriate classes of Wick operators on the Bargmann side. These considerations are based on a formula for the Wick symbol expressed in terms of a short-time Fourier transform of the Weyl symbol, and admits characterization of the Wick symbols corresponding to Shubin Weyl symbols and symbols for operators of infinite order (see Proposition 2.3).

In Section 2 we also study composition and show for example that the well-known closure under composition of Shubin operators and operators of infinite orders have simple and natural proofs on the Wick symbol side.

In Section 3 we deduce series expansions of Wick operators in terms of anti-Wick operators, and between Wick symbols and symbols to corresponding Shubin operators. We also consider anti-Wick operators, and show continuity results for them. We show that the upper bounds for the Wick symbols of anti-Wick operators are stricter than for general Wick symbols.

In Section 4 we discuss lower bounds for Wick operators and deduce the sharp Gårding's inequality. Section 5 concerns ellipticity, hypoellipticity and weak ellipticity.

Finally we observe in Section 6 that a polynomial bound of a Wick symbol implies that the symbol is a polynomial. For pseudo-differential operators this corresponds to partial differential operators with polynomial coefficients. This gives a characterization of such operators as those having polynomially bounded Wick symbols.

Various types of function spaces, distribution spaces, their Bargmann images, and symbol classes for pseudo-differential, Wick and anti-Wick operators appear frequently in the paper. For the reader's convenience we summarize several of these items in Appendix A.

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## 1. Preliminaries

In this section we recall some facts on function and distribution spaces as well as on pseudo-differential operators, Wick and anti-Wick operators. Subsection 1.1 concerns weight functions and Subsection 1.2 treats Gelfand-Shilov spaces. In Subsection 1.3 we introduce the Bargmann transform and topological spaces of entire functions on $\mathbf{C}^{d}$, and in Subsection 1.4 we recall the definitions and some facts on pseudo-differential operators on $\mathbf{R}^{d}$ as well as Wick and anti-Wick operators on $\mathbf{C}^{d}$. Subsection 1.5 defines certain symbol classes for pseudo-differential operators on $\mathbf{R}^{d}$.

### 1.1. Weight functions

A weight on $\mathbf{R}^{d}$ is a positive function $\omega \in L_{l o c}^{\infty}\left(\mathbf{R}^{d}\right)$ such that $1 / \omega \in L_{l o c}^{\infty}\left(\mathbf{R}^{d}\right)$. The weight $\omega$ is called moderate if there is a positive locally bounded function $v$ such that

$$
\begin{equation*}
\omega(x+y) \leqslant C \omega(x) v(y), \quad x, y \in \mathbf{R}^{d}, \tag{1.1}
\end{equation*}
$$

for some constant $C \geqslant 1$. If $\omega$ and $v$ are weights such that (1.1) holds, then $\omega$ is also called $v$-moderate. The set of all moderate weights on $\mathbf{R}^{d}$ is denoted by $\mathscr{P}_{E}\left(\mathbf{R}^{d}\right)$. The set $\mathscr{P}\left(\mathbf{R}^{d}\right)$ consists of weights that are $v$ moderate for a polynomially bounded weight, that is a weight of the form $v(x)=\langle x\rangle^{s}$ where $\langle x\rangle=\left(1+|x|^{2}\right)^{\frac{1}{2}}$ and $s \geqslant 0$. The bracket notation is also used for complex arguments as $\langle z\rangle=\left(1+|z|^{2}\right)^{\frac{1}{2}}$ when $z \in \mathbf{C}^{d}$. In particular, $\omega \in \mathscr{P}\left(\mathbf{R}^{d}\right)$, if and only if

$$
\begin{equation*}
\omega(x+y) \leqslant C \omega(x)\langle y\rangle^{r}, \quad x, y \in \mathbf{R}^{d}, \tag{1.1}
\end{equation*}
$$

for some $r \geqslant 0$. If $s \in \mathbf{R}$ then $x \mapsto\langle x\rangle^{s}$ belongs to $\mathscr{P}\left(\mathbf{R}^{d}\right)$, due to Peetre's inequality (with optimal constant, see [25, Lemma 2.1])

$$
\begin{equation*}
\langle x+y\rangle^{s} \leqslant\left(\frac{2}{\sqrt{3}}\right)^{|s|}\langle x\rangle^{s}\langle y\rangle^{|s|} \quad x, y \in \mathbf{R}^{d}, \quad s \in \mathbf{R} . \tag{1.2}
\end{equation*}
$$

The weight $v$ is called submultiplicative if it is even and (1.1) holds for $\omega=v$. If (1.1) holds and $v$ is submultiplicative then

$$
\begin{gather*}
\frac{\omega(x)}{v(y)} \lesssim \omega(x+y) \lesssim \omega(x) v(y),  \tag{1.3}\\
v(x+y) \lesssim v(x) v(y) \quad \text { and } \quad v(x)=v(-x), \quad x, y \in \mathbf{R}^{d} .
\end{gather*}
$$

The notation $A(\theta) \lesssim B(\theta), \theta \in \Omega$, means that there is a constant $c>0$ such that $A(\theta) \leqslant c B(\theta)$ for all $\theta \in \Omega$.
If $\omega$ is a moderate weight then by [30] there is a submultiplicative weight $v$ such that (1.1) and (1.3) hold. If $v$ is submultiplicative then

$$
\begin{equation*}
1 \lesssim v(x) \lesssim e^{r|x|} \tag{1.4}
\end{equation*}
$$

for some constant $r>0$ (cf. [14]). In particular, if $\omega$ is moderate, then

$$
\begin{equation*}
\omega(x+y) \lesssim \omega(x) e^{r|y|} \quad \text { and } \quad e^{-r|x|} \lesssim \omega(x) \lesssim e^{r|x|}, \quad x, y \in \mathbf{R}^{d} \tag{1.5}
\end{equation*}
$$

for some $r>0$. If not otherwise specified the symbol $v$ always denote a submultiplicative weight.

### 1.2. Gelfand-Shilov spaces

Let $s, \sigma>0$. The Gelfand-Shilov space $\mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right)\left(\Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right)\right)$ of Roumieu (Beurling) type consists of all $f \in C^{\infty}\left(\mathbf{R}^{d}\right)$ such that

$$
\begin{equation*}
\|f\|_{\mathcal{S}_{s, h}^{\sigma}} \equiv \sup \left(\frac{\left|x^{\alpha} \partial^{\beta} f(x)\right|}{h^{|\alpha+\beta|} \alpha!^{s} \beta!^{\sigma}}\right) \tag{1.6}
\end{equation*}
$$

is finite for some (every) $h>0$ (see [13]). The supremum refers to all $\alpha, \beta \in \mathbf{N}^{d}$ and $x \in \mathbf{R}^{d}$. The seminorms $\|\cdot\|_{\mathcal{S}_{s, h}^{\sigma}}$ induce an inductive limit topology for the space $\mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right)$ and a projective limit topology for $\Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right)$. The latter space is a Fréchet space under this topology. The space $\mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right) \neq\{0\}$, if and only if $s+\sigma \geqslant 1$, and $\Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right) \neq\{0\}$, if and only if $s+\sigma>1$, see [13,22]. We write $\mathcal{S}_{s}\left(\mathbf{R}^{d}\right)=\mathcal{S}_{s}^{s}\left(\mathbf{R}^{d}\right)$ and $\Sigma_{s}\left(\mathbf{R}^{d}\right)=\Sigma_{s}^{s}\left(\mathbf{R}^{d}\right)$.

The Gelfand-Shilov distribution spaces $\left(\mathcal{S}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)$ and $\left(\Sigma_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)$ are the (strong) dual spaces of $\mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right)$ and $\Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right)$, respectively.

The embeddings

$$
\begin{align*}
\mathcal{S}_{s_{1}}^{\sigma_{1}}\left(\mathbf{R}^{d}\right) \hookrightarrow \Sigma_{s_{2}}^{\sigma_{2}}\left(\mathbf{R}^{d}\right) \hookrightarrow & \mathcal{S}_{s_{2}}^{\sigma_{2}}\left(\mathbf{R}^{d}\right) \hookrightarrow \mathscr{S}\left(\mathbf{R}^{d}\right) \\
& \hookrightarrow \mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right) \hookrightarrow\left(\mathcal{S}_{s_{2}}^{\sigma_{2}}\right)^{\prime}\left(\mathbf{R}^{d}\right) \hookrightarrow\left(\Sigma_{s_{2}}^{\sigma_{2}}\right)^{\prime}\left(\mathbf{R}^{d}\right) \hookrightarrow\left(\mathcal{S}_{s_{1}}^{\sigma_{1}}\right)^{\prime}\left(\mathbf{R}^{d}\right), \\
& s_{1}+\sigma_{1} \geqslant 1, s_{1}<s_{2}, \sigma_{1}<\sigma_{2}, \tag{1.7}
\end{align*}
$$

are dense. For topological spaces $A$ and $B, A \hookrightarrow B$ means that the inclusion $A \subseteq B$ is continuous.
The spaces $\mathcal{S}_{s}$ and $\Sigma_{s}$, and their dual spaces, admit characterizations in terms of coefficients with respect to expansions with respect to the Hermite functions

$$
h_{\alpha}(x)=\pi^{-\frac{d}{4}}(-1)^{|\alpha|}\left(2^{|\alpha|} \alpha!\right)^{-\frac{1}{2}} e^{\frac{|x|^{2}}{2}}\left(\partial^{\alpha} e^{-|x|^{2}}\right), \quad \alpha \in \mathbf{N}^{d} .
$$

The set of Hermite functions on $\mathbf{R}^{d}$ is an orthonormal basis for $L^{2}\left(\mathbf{R}^{d}\right)$. We use $\mathcal{H}_{0}\left(\mathbf{R}^{d}\right)$ to denote the space of finite linear combinations of Hermite functions. Then $\mathcal{H}_{0}\left(\mathbf{R}^{d}\right)$ is dense in the Schwartz space $\mathscr{S}\left(\mathbf{R}^{d}\right)$, as well as in $\mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$, with respect to its weak* topology. The same conclusion is true for $\Sigma_{s}\left(\mathbf{R}^{d}\right)$ when $s>\frac{1}{2}$, $\mathcal{S}_{s}\left(\mathbf{R}^{d}\right)$ when $s \geqslant \frac{1}{2}$ and their distribution dual spaces $\Sigma_{s}^{\prime}\left(\mathbf{R}^{d}\right)$ and $\mathcal{S}_{s}^{\prime}\left(\mathbf{R}^{d}\right)$. A function or (ultra-)distribution $f$ in any of these spaces possess an expansion of the form

$$
\begin{equation*}
f=\sum_{\alpha \in \mathbf{N}^{d}} c(f, \alpha) h_{\alpha}, \quad c(f, \alpha)=\left(f, h_{\alpha}\right), \quad \alpha \in \mathbf{N}^{d} . \tag{1.8}
\end{equation*}
$$

Here $(\cdot, \cdot)$ denotes the unique extensions of the $L^{2}$ form, which is linear in the first variable and conjugate linear in the second variable, from $\mathcal{H}_{0}\left(\mathbf{R}^{d}\right) \times \mathcal{H}_{0}\left(\mathbf{R}^{d}\right)$ to $\mathcal{S}_{s}^{\prime}\left(\mathbf{R}^{d}\right) \times \mathcal{S}_{s}\left(\mathbf{R}^{d}\right)$ or $\Sigma_{s}^{\prime}\left(\mathbf{R}^{d}\right) \times \Sigma_{s}\left(\mathbf{R}^{d}\right)$. We recall that (cf. [24, Chapter V. 3 ])

$$
\begin{align*}
f \in \mathscr{S}\left(\mathbf{R}^{d}\right) & \Leftrightarrow|c(f, \alpha)| \lesssim\langle\alpha\rangle^{-N} \text { for every } N \geqslant 0,  \tag{1.9}\\
f \in \mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right) & \Leftrightarrow|c(f, \alpha)| \lesssim\langle\alpha\rangle^{N} \quad \text { for some } N \geqslant 0 .
\end{align*}
$$

The topology on $\mathscr{S}\left(\mathbf{R}^{d}\right)$ is equivalent to the Fréchet space topology defined by the sequence space seminorms

$$
\mathscr{S}\left(\mathbf{R}^{d}\right) \ni f \mapsto \sum_{\alpha \in \mathbf{N}^{d}}\langle\alpha\rangle^{2 N}|c(f, \alpha)|^{2}, \quad N \geqslant 0 .
$$

For $f \in \mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$ the sum in (1.8) converges in the weak* topology.

The Hermite functions are eigenfunctions to the harmonic oscillator $H=H_{d} \equiv|x|^{2}-\Delta$ and to the Fourier transform $\mathscr{F}$, given by

$$
\mathscr{F} f(\xi)=\widehat{f}(\xi) \equiv(2 \pi)^{-\frac{d}{2}} \int_{\mathbf{R}^{d}} f(x) e^{-i\langle x, \xi\rangle} d x, \quad \xi \in \mathbf{R}^{d}
$$

when $f \in L^{1}\left(\mathbf{R}^{d}\right)$. Here $\langle\cdot, \cdot\rangle$ denotes the scalar product on $\mathbf{R}^{d}$. In fact

$$
H_{d} h_{\alpha}=(2|\alpha|+d) h_{\alpha}, \quad \alpha \in \mathbf{N}^{d} .
$$

The Fourier transform $\mathscr{F}$ extends uniquely to homeomorphisms on $\mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$, from $\left(\mathcal{S}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)$ to $\left(\mathcal{S}_{\sigma}^{s}\right)^{\prime}\left(\mathbf{R}^{d}\right)$ and from $\left(\Sigma_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)$ to $\left(\Sigma_{\sigma}^{s}\right)^{\prime}\left(\mathbf{R}^{d}\right)$. It also restricts to homeomorphisms on $\mathscr{S}\left(\mathbf{R}^{d}\right)$, from $\mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right)$ to $\mathcal{S}_{\sigma}^{s}\left(\mathbf{R}^{d}\right)$, from $\Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right)$ to $\Sigma_{\sigma}^{s}\left(\mathbf{R}^{d}\right)$, and to a unitary operator on $L^{2}\left(\mathbf{R}^{d}\right)$. Similar facts hold true when the Fourier transform is replaced by a partial Fourier transform.

Let $\phi \in \mathscr{S}\left(\mathbf{R}^{d}\right) \backslash 0$ be fixed. We use the transform

$$
\begin{align*}
\mathcal{T}_{\phi} f(x, \xi) & =(2 \pi)^{-\frac{d}{2}} e^{i\langle x, \xi\rangle}\left(f, e^{i<\cdot, \xi\rangle} \phi(\cdot-x)\right) \\
& =e^{i\langle x, \xi\rangle} \mathscr{F}(f \cdot \overline{\phi(\cdot-x)})(\xi)  \tag{1.10}\\
& =\mathscr{F}(f(\cdot+x) \bar{\phi})(\xi), \quad x, \xi \in \mathbf{R}^{d},
\end{align*}
$$

where $f \in \mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$ and $\phi \in \mathscr{S}\left(\mathbf{R}^{d}\right) \backslash 0$ (cf. [8]). If $f, \phi \in \mathscr{S}\left(\mathbf{R}^{d}\right)$ then

$$
\begin{align*}
\mathcal{T}_{\phi} f(x, \xi) & =(2 \pi)^{-\frac{d}{2}} e^{i\langle x, \xi\rangle} \int_{\mathbf{R}^{d}} f(y) \overline{\phi(y-x)} e^{-i\langle y, \xi\rangle} d y  \tag{1.10}\\
& =(2 \pi)^{-\frac{d}{2}} \int_{\mathbf{R}^{d}} f(y+x) \overline{\phi(y)} e^{-i\langle y, \xi\rangle} d y, \quad x, \xi \in \mathbf{R}^{d} .
\end{align*}
$$

We notice that the short-time Fourier transform $V_{\phi} f$ of $f$ is given by

$$
\begin{equation*}
V_{\phi} f(x, \xi)=e^{-i\langle x, \xi\rangle} \mathcal{T}_{\phi} f(x, \xi) \tag{1.11}
\end{equation*}
$$

That is, $\mathcal{T}_{\phi}$ is a modulated short-time Fourier transform. Thus by [30, Theorem 2.3] it follows that the definition of the map $(f, \phi) \mapsto \mathcal{T}_{\phi} f$ from $\mathscr{S}\left(\mathbf{R}^{d}\right) \times \mathscr{S}\left(\mathbf{R}^{d}\right)$ to $\mathscr{S}\left(\mathbf{R}^{2 d}\right)$ is uniquely extendable to a continuous map from $\mathcal{S}_{s}^{\prime}\left(\mathbf{R}^{d}\right) \times \mathcal{S}_{s}^{\prime}\left(\mathbf{R}^{d}\right)$ to $\mathcal{S}_{s}^{\prime}\left(\mathbf{R}^{2 d}\right)$, and restricts to a continuous map from $\mathcal{S}_{s}\left(\mathbf{R}^{d}\right) \times \mathcal{S}_{s}\left(\mathbf{R}^{d}\right)$ to $\mathcal{S}_{s}\left(\mathbf{R}^{2 d}\right)$. The same conclusion holds with $\Sigma_{s}$ in place of $\mathcal{S}_{s}$, at each place.

The adjoint $\mathcal{T}_{\phi}^{*}$ is given by

$$
\left(\mathcal{T}_{\phi}^{*} F, g\right)_{L^{2}\left(\mathbf{R}^{d}\right)}=\left(F, \mathcal{T}_{\phi} g\right)_{L^{2}\left(\mathbf{R}^{2 d}\right)}
$$

for $F \in \mathcal{S}_{s}^{\prime}\left(\mathbf{R}^{2 d}\right)$ and $g \in \mathcal{S}_{s}\left(\mathbf{R}^{d}\right)$, and similarly with $\Sigma_{s}$ or with $\mathscr{S}$ in place of $\mathcal{S}_{s}$ at each occurrence. When $F$ is a polynomially bounded measurable function we write

$$
\begin{equation*}
\mathcal{T}_{\phi}^{*} F(y)=(2 \pi)^{-\frac{d}{2}} \iint_{\mathbf{R}^{2 d}} F(x, \xi) e^{i\langle y-x, \xi\rangle} \phi(y-x) d x d \xi, \tag{1.12}
\end{equation*}
$$

where the integral is defined weakly so that $\left(\mathcal{T}_{\phi}^{*} F, g\right)_{L^{2}\left(\mathbf{R}^{d}\right)}=\left(F, \mathcal{T}_{\phi} g\right)_{L^{2}\left(\mathbf{R}^{2 d}\right)}$ for $g \in \mathscr{S}\left(\mathbf{R}^{d}\right)$. The identity (1.12) is called Moyal's formula.

We have

$$
\begin{equation*}
\left(\mathcal{T}_{\psi}^{*} \circ \mathcal{T}_{\phi}\right) f=(\psi, \phi) f, \quad f \in \mathcal{S}_{s}^{\prime}\left(\mathbf{R}^{d}\right), \phi, \psi \in \mathcal{S}_{s}\left(\mathbf{R}^{d}\right) \tag{1.13}
\end{equation*}
$$

and similarly with $\Sigma_{s}$ or with $\mathscr{S}$ in place of $\mathcal{S}_{s}$ at each occurrence.
Two important features of $\mathcal{T}_{\phi}$ which distinguish it from the short-time Fourier transform are the differential identities

$$
\begin{equation*}
\partial_{x}^{\alpha} \mathcal{T}_{\phi} f(x, \xi)=\mathcal{T}_{\phi}\left(\partial^{\alpha} f\right)(x, \xi), \quad \alpha \in \mathbf{N}^{d} \tag{1.14}
\end{equation*}
$$

and

$$
\begin{equation*}
D_{\xi}^{\beta} \mathcal{T}_{\phi} f(x, \xi)=\mathcal{T}_{g_{\beta}} f(x, \xi), \quad \beta \in \mathbf{N}^{d}, \quad \phi_{\beta}(x)=(-x)^{\beta} \phi(x) . \tag{1.15}
\end{equation*}
$$

By (1.11) it follows that characterizations of Gelfand-Shilov spaces and their distribution spaces in terms of estimates of their short-time Fourier transforms carry over to estimates on $\mathcal{T}_{\phi}$ in place of $V_{\phi}$. For example we have the following (see e.g. [15,27] for the proof of (1) and [31] for the proof of (2)). See also [10] for related results.

Proposition 1.1. Let $s, \sigma>0, \phi \in \mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right) \backslash 0\left(\phi \in \Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right) \backslash 0\right)$ and let $f \in\left(\mathcal{S}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)\left(f \in\left(\Sigma_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)\right)$. Then the following is true:
(1) $f \in \mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right)\left(f \in \Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right)\right)$ if and only if

$$
\begin{equation*}
\left|\mathcal{T}_{\phi} f(x, \xi)\right| \lesssim e^{-r\left(|x|^{\frac{1}{s}}+|\xi|^{\frac{1}{\sigma}}\right)}, \quad x, \xi \in \mathbf{R}^{d} \tag{1.16}
\end{equation*}
$$

for some (every) $r>0$.
(2) $f \in\left(\mathcal{S}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)\left(f \in\left(\Sigma_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)\right)$ if and only if

$$
\begin{equation*}
\left|\mathcal{T}_{\phi} f(x, \xi)\right| \lesssim e^{r\left(|x|^{\frac{1}{s}}+|\xi|^{\frac{1}{\sigma}}\right)}, \quad x, \xi \in \mathbf{R}^{d}, \tag{1.17}
\end{equation*}
$$

for every (some) $r>0$.

### 1.3. The Bargmann transform and spaces of analytic functions

If $\Omega \subseteq \mathbf{C}^{d}$ is open then $A(\Omega)$ consists of all (complex-valued) analytic functions on $\Omega$. Complex derivatives are denoted, with $z=x+i y \in \Omega$,

$$
\partial_{z_{j}}=\frac{1}{2}\left(\partial_{x_{j}}-i \partial_{y_{j}}\right), \quad \bar{\partial}_{z_{j}}=\frac{1}{2}\left(\partial_{x_{j}}+i \partial_{y_{j}}\right)
$$

for $1 \leqslant j \leqslant d$, which admits the Cauchy-Riemann equations to be written as $\bar{\partial}_{z_{j}} f=0,1 \leqslant j \leqslant d$.
The Bargmann kernel is defined by

$$
\mathfrak{A}_{d}(z, y)=\pi^{-\frac{d}{4}} \exp \left(-\frac{1}{2}\left(\langle z, z\rangle+|y|^{2}\right)+2^{1 / 2}\langle z, y\rangle\right), \quad z \in \mathbf{C}^{d}, \quad y \in \mathbf{R}^{d},
$$

where

$$
\langle z, w\rangle=\sum_{j=1}^{d} z_{j} w_{j} \quad \text { and } \quad(z, w)=\langle z, \bar{w}\rangle
$$

when

$$
z=\left(z_{1}, \ldots, z_{d}\right) \in \mathbf{C}^{d} \quad \text { and } \quad w=\left(w_{1}, \ldots, w_{d}\right) \in \mathbf{C}^{d}
$$

Sometimes $\langle\cdot, \cdot\rangle$ denotes the duality between a test function space and its dual. The context precludes confusion between its double use. The Bargmann transform $\mathfrak{V}_{d} f$ of $f \in \mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{d}\right)$ is the entire function

$$
\begin{equation*}
\mathfrak{V}_{d} f(z)=\left\langle f, \mathfrak{A}_{d}(z, \cdot)\right\rangle, \quad z \in \mathbf{C}^{d} . \tag{1.18}
\end{equation*}
$$

The right-hand side is a well defined element in $A\left(\mathbf{C}^{d}\right)$, since $y \mapsto \mathfrak{A}_{d}(z, y)$ belongs to $\mathcal{S}_{1 / 2}\left(\mathbf{R}^{d}\right)$ for $z \in \mathbf{C}^{d}$ fixed, and $\mathfrak{A}_{d}(\cdot, y)$ is entire for all $y \in \mathbf{R}^{d}$. Let $p \in[1, \infty]$ and $\omega \in \mathscr{P}_{E}\left(\mathbf{R}^{d}\right)$. Then $L_{(\omega)}^{p}\left(\mathbf{R}^{d}\right)$ consists of all $f \in L_{l o c}^{1}\left(\mathbf{R}^{d}\right)$ such that $\|f\|_{L_{(\omega)}^{p}} \equiv\|f \cdot \omega\|_{L^{p}}$ is finite. If $f \in L_{(\omega)}^{p}\left(\mathbf{R}^{d}\right)$, then

$$
\begin{equation*}
\mathfrak{V}_{d} f(z)=\int_{\mathbf{R}^{d}} \mathfrak{A}_{d}(z, y) f(y) d y=\pi^{-\frac{d}{4}} \int_{\mathbf{R}^{d}} \exp \left(-\frac{1}{2}\left(\langle z, z\rangle+|y|^{2}\right)+2^{1 / 2}\langle z, y\rangle\right) f(y) d y, \quad z \in \mathbf{C}^{d} . \tag{1.19}
\end{equation*}
$$

(Cf. [3,4,30,31].)
For $p \in(0, \infty], \omega \in \mathscr{P}_{E}\left(\mathbf{C}^{d}\right)$ and $\omega_{0}(z)=\omega(\sqrt{2} \bar{z})$, let $A_{(\omega)}^{p}\left(\mathbf{C}^{d}\right)$ be the set of all $F \in A\left(\mathbf{C}^{d}\right)$ such that

$$
\|F\|_{A_{(\omega)}^{p}} \equiv \pi^{-\frac{d}{p}}\left\|F \cdot e^{-\frac{1}{2}|\cdot|^{2}} \cdot \omega_{0}\right\|_{L^{p}}
$$

is finite, and set $A^{p}=A_{(\omega)}^{p}$ when $\omega=1$. It was proved by Bargmann [3] that

$$
\begin{equation*}
\mathfrak{V}_{d}: L^{2}\left(\mathbf{R}^{d}\right) \rightarrow A^{2}\left(\mathbf{C}^{d}\right) \tag{1.20}
\end{equation*}
$$

is bijective and isometric. The space $A^{2}\left(\mathbf{C}^{d}\right)$ is the Hilbert space of entire functions with scalar product

$$
(F, G)_{A^{2}} \equiv \int_{\mathbf{C}^{d}} F(z) \overline{G(z)} d \mu(z), \quad F, G \in A^{2}\left(\mathbf{C}^{d}\right)
$$

where $d \mu(z)=\pi^{-d} e^{-|z|^{2}} d \lambda(z)$ and $d \lambda(z)$ is the Lebesgue measure on $\mathbf{C}^{d}$. The space $A^{2}\left(\mathbf{C}^{d}\right)$ is known as the Fock or Segal-Bargmann space in quantum mechanics (see [12,16]).

In [3] it was proved that the Bargmann transform maps the Hermite functions to monomials as

$$
\begin{equation*}
\mathfrak{V}_{d} h_{\alpha}=e_{\alpha}, \quad e_{\alpha}(z)=\frac{z^{\alpha}}{\alpha!^{\frac{1}{2}}}, \quad z \in \mathbf{C}^{d}, \quad \alpha \in \mathbf{N}^{d} . \tag{1.21}
\end{equation*}
$$

The orthonormal basis $\left\{h_{\alpha}\right\}_{\alpha \in \mathbf{N}^{d}} \subseteq L^{2}\left(\mathbf{R}^{d}\right)$ is thus mapped to the orthonormal basis $\left\{e_{\alpha}\right\}_{\alpha \in \mathbf{N}^{d}} \subseteq A^{2}\left(\mathbf{C}^{d}\right)$. Bargmann also proved that there is a reproducing formula for $A^{2}\left(\mathbf{C}^{d}\right)$. Let $\Pi_{A}$ be the operator from $L^{2}(d \mu)$ to $A\left(\mathbf{C}^{d}\right)$, given by

$$
\begin{equation*}
\Pi_{A} F(z)=\int_{\mathbf{C}^{d}} F(w) e^{(z, w)} d \mu(w), \quad z \in \mathbf{C}^{d} \tag{1.22}
\end{equation*}
$$

Then $\Pi_{A}$ is the orthogonal projection from $L^{2}(d \mu)$ to $A^{2}\left(\mathbf{C}^{d}\right)$ (cf. [3]).
When we discuss extensions and restrictions of the Bargmann transform to Gelfand-Shilov spaces and their distribution spaces, we use

$$
\begin{equation*}
|z|_{s, \sigma}=|\operatorname{Re} z|^{\frac{1}{s}}+|\operatorname{Im} z|^{\frac{1}{\sigma}}, \quad z \in \mathbf{C}^{d}, \tag{1.23}
\end{equation*}
$$

and consider the seminorms

$$
\|F\|_{\mathcal{A}_{\mathscr{S} ; r}} \equiv \| F \cdot e^{-\frac{1}{2}|\cdot|^{2}\langle\cdot\rangle^{r}\left\|_{L^{\infty}}, \quad\right\| F\left\|_{\mathcal{A}_{\mathscr{S} ; r}^{\prime}} \equiv\right\| F \cdot e^{-\frac{1}{2}|\cdot|^{2}}\langle\cdot\rangle^{-r} \|_{L^{\infty}} .}
$$

and

$$
\|F\|_{\mathcal{A}_{s ; r}^{\sigma}} \equiv\left\|F \cdot e^{-\frac{1}{2}|\cdot|^{2}+r|\cdot|_{s, \sigma}}\right\|_{L^{\infty}}, \quad\|F\|_{\mathcal{A}_{s}^{\prime} ; ; r} \equiv\left\|F \cdot e^{-\frac{1}{2}|\cdot|^{2}-r|\cdot|_{s, \sigma}}\right\|_{L^{\infty}}
$$

when $F \in A\left(\mathbf{C}^{d}\right), r>0$ and $s, \sigma \geqslant \frac{1}{2}$. Then $\mathcal{A}_{0, s}^{\sigma}\left(\mathbf{C}^{d}\right)$ for $s, \sigma>\frac{1}{2}, \mathcal{A}_{\mathscr{S}}\left(\mathbf{C}^{d}\right)$ and $\left(\mathcal{A}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right)$ for $s, \sigma \geqslant \frac{1}{2}$ are the sets of all $F \in A\left(\mathbf{C}^{d}\right)$ such that

$$
\begin{equation*}
\|F\|_{\mathcal{A}_{S_{s ; r}^{\sigma}}}<\infty, \quad\|F\|_{\mathcal{A}_{\mathscr{s} ; r}}<\infty \quad \text { and } \quad\|F\|_{\mathcal{A}_{S_{s ; r}^{\prime}}^{\prime}}<\infty \tag{1.24}
\end{equation*}
$$

respectively, for every $r>0$. The spaces are equipped with the projective limit topology with respect to $r>0$, defined by each class of seminorms, respectively.

In the same way we let $\mathcal{A}_{s}^{\sigma}\left(\mathbf{C}^{d}\right)$ for $s, \sigma \geqslant \frac{1}{2}, \mathcal{A}_{\mathscr{S}}^{\prime}\left(\mathbf{C}^{d}\right)$ and $\left(\mathcal{A}_{0, s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right)$ for $s, \sigma>\frac{1}{2}$ be the sets of all $F \in A\left(\mathbf{C}^{d}\right)$ such that

$$
\begin{equation*}
\|F\|_{\mathcal{A}_{\mathcal{S}_{\xi ; r}^{g}}}<\infty, \quad\|F\|_{\mathcal{A}_{\mathscr{s} ; r}^{\prime}}<\infty \quad \text { and } \quad\|F\|_{\mathcal{A}_{\mathcal{S}_{s ; r}^{\prime}}^{\prime}}<\infty, \tag{1.25}
\end{equation*}
$$

respectively, for some $r>0$. Their topologies are the inductive limit topologies with respect to $r>0$, defined by each class of seminorms, respectively. We also set

$$
\mathcal{A}_{0, s}=\mathcal{A}_{0, s}^{s} \quad \text { and } \quad \mathcal{A}_{s}=\mathcal{A}_{s}^{s} .
$$

Then

$$
\begin{array}{ll}
\mathfrak{V}_{d}: \mathscr{S}\left(\mathbf{R}^{d}\right) \rightarrow \mathcal{A}_{\mathscr{S}}\left(\mathbf{C}^{d}\right), & \mathfrak{V}_{d}: \quad \mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right) \rightarrow \mathcal{A}_{\mathscr{S}}^{\prime}\left(\mathbf{C}^{d}\right), \\
\mathfrak{V}_{d}: \mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right) \rightarrow \mathcal{A}_{s}^{\sigma}\left(\mathbf{C}^{d}\right), & \mathfrak{V}_{d}:\left(\mathcal{S}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right) \rightarrow\left(\mathcal{A}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right) \quad s, \sigma \geqslant \frac{1}{2}
\end{array}
$$

and

$$
\mathfrak{V}_{d}: \Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right) \rightarrow \mathcal{A}_{0, s}^{\sigma}\left(\mathbf{C}^{d}\right), \quad \mathfrak{V}_{d}:\left(\Sigma_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right) \rightarrow\left(\mathcal{A}_{0, s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right), \quad s, \sigma>\frac{1}{2}
$$

are homeomorphisms [31].
From these homeomorphisms, the fact that the map (1.20) is a homeomorphism and duality properties for Gelfand-Shilov spaces, it follows that $(\cdot, \cdot)_{A^{2}}$ on $\mathcal{A}_{1 / 2}\left(\mathbf{C}^{d}\right) \times \mathcal{A}_{1 / 2}\left(\mathbf{C}^{d}\right)$ is uniquely extendable to a continuous sesqui-linear form on $\left(\mathcal{A}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right) \times \mathcal{A}_{s}^{\sigma}\left(\mathbf{C}^{d}\right)$. The (strong) dual of $\mathcal{A}_{s}^{\sigma}\left(\mathbf{C}^{d}\right)$ can be identified with $\left(\mathcal{A}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right)$ through this form. Similar facts hold for $\mathcal{A}_{0, s}^{\sigma}$ in place of $\mathcal{A}_{s}^{\sigma}$ at each occurrence. (Cf. e.g. [30,31].)

Finally let $\mathcal{A}_{b_{1} ; r}\left(\mathbf{C}^{d}\right)$ and $\mathcal{A}_{b_{\infty} ; r}\left(\mathbf{C}^{d}\right)$ for $r>0$ be the Banach spaces which consist of all $F \in A\left(\mathbf{C}^{d}\right)$ such that

$$
\|F\|_{\mathcal{A}_{b_{1} ; r}} \equiv\left\|F \cdot e^{-r|\cdot|}\right\|_{L^{\infty}} \quad \text { respectively } \quad\|F\|_{\mathcal{A}_{b_{\infty} ; r}} \equiv\left\|F \cdot e^{-r|\cdot|^{2}}\right\|_{L^{\infty}}
$$

is finite, and let $\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$ be the inductive limit of $\mathcal{A}_{b_{1} ; r}\left(\mathbf{C}^{d}\right)$ with respect to $r>0$. Also let $\mathcal{A}_{0, b_{\infty}}\left(\mathbf{C}^{d}\right)$ and $\mathcal{A}_{0, b_{\infty}}^{\prime}\left(\mathbf{C}^{d}\right)$ be the projective respectively inductive limit topologies of $\mathcal{A}_{b_{\infty} ; r}\left(\mathbf{C}^{d}\right)$ with respect to $r>0$.

It is evident that $\mathcal{A}_{\mathfrak{b}_{1}}\left(\mathbf{C}^{d}\right)$ is densely embedded in $\mathcal{A}_{s}^{\sigma}\left(\mathbf{C}^{d}\right)$ for every $s, \sigma \geqslant \frac{1}{2}$, as well as in $\mathcal{A}_{0, s}^{\sigma}\left(\mathbf{C}^{d}\right)$ for every $s, \sigma>\frac{1}{2}$. The form $(\cdot, \cdot)_{A^{2}}$ on $\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right) \times \mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$ is uniquely extendable to a continuous sesqui-linear
form on $A\left(\mathbf{C}^{d}\right) \times \mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$ and the dual of $\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$ can be identified with $A\left(\mathbf{C}^{d}\right)$. The Fréchet space topology of $A\left(\mathbf{C}^{d}\right)$ can be defined by the seminorms

$$
F \mapsto \sup _{|z| \leqslant N}|F(z)|, \quad N=1,2, \ldots
$$

(Cf. [31].)
Remark 1.2. The spaces $\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$ and $\mathcal{A}_{0, b_{\infty}}\left(\mathbf{C}^{d}\right)$ are examples of Bargmann images of special Pilipović spaces, a family of Fourier invariant topological vector spaces which are smaller than any Fourier invariant Gelfand-Shilov space, and which were introduced and investigated in [31]. For any $\sigma>0$, the Bargmann image of the Pilipovic spaces $\mathcal{H}_{b_{\sigma}}\left(\mathbf{R}^{d}\right)$ and $\mathcal{H}_{0, b_{\sigma}}\left(\mathbf{R}^{d}\right)$ are given by

$$
\mathcal{A}_{b_{\sigma}}\left(\mathbf{C}^{d}\right) \equiv\left\{F \in A\left(\mathbf{C}^{d}\right) ;|F(z)| \lesssim e^{r|z|^{\frac{2 \sigma}{\sigma+1}}} \text { for some } r>0\right\}
$$

respectively

$$
\mathcal{A}_{0, b_{\sigma}}\left(\mathbf{C}^{d}\right) \equiv\left\{F \in A\left(\mathbf{C}^{d}\right) ;|F(z)| \lesssim e^{r|z|^{\frac{2 \sigma}{\sigma+1}}} \text { for every } r>0\right\}
$$

If $\sigma>1$, then the (strong) duals of $\mathcal{A}_{b_{\sigma}}\left(\mathbf{C}^{d}\right)$ and $\mathcal{A}_{0, b_{\sigma}}\left(\mathbf{C}^{d}\right)$ are given by

$$
\mathcal{A}_{b_{\sigma}}^{\prime}\left(\mathbf{C}^{d}\right) \equiv\left\{F \in A\left(\mathbf{C}^{d}\right) ;|F(z)| \lesssim e^{r|z|^{\frac{2 \sigma}{\sigma-1}}} \text { for every } r>0\right\}
$$

respectively

$$
\mathcal{A}_{0, b_{\sigma}}^{\prime}\left(\mathbf{C}^{d}\right) \equiv\left\{F \in A\left(\mathbf{C}^{d}\right) ;|F(z)| \lesssim e^{r|z|^{\frac{2 \sigma}{\sigma-1}}} \text { for some } r>0\right\}
$$

through a unique extension of the $A^{2}$ scalar product on $\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right) \times \mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$. In particular, if $\sigma$ tends to $\infty$, it follows that some of these conditions tend to

$$
\mathcal{A}_{0, b_{\infty}}\left(\mathbf{C}^{d}\right) \equiv\left\{F \in A\left(\mathbf{C}^{d}\right) ;|F(z)| \lesssim e^{r|z|^{2}} \text { for every } r>0\right\}
$$

respectively

$$
\mathcal{A}_{0, b_{\infty}}^{\prime}\left(\mathbf{C}^{d}\right) \equiv\left\{F \in A\left(\mathbf{C}^{d}\right) ;|F(z)| \lesssim e^{r|z|^{2}} \text { for some } r>0\right\}
$$

Note that in $[31,29]$, the set $\mathcal{A}_{0, b_{\infty}}\left(\mathbf{C}^{d}\right)$ is denoted by $\mathcal{A}_{0, \frac{1}{2}}\left(\mathbf{C}^{d}\right)$, and its dual $\mathcal{A}_{0, b_{\infty}}^{\prime}\left(\mathbf{C}^{d}\right)$ is denoted by $\mathcal{A}_{0, \frac{1}{2}}^{\prime}\left(\mathbf{C}^{d}\right)$.

At many places it will be crucial to use the Gaussian window

$$
\begin{equation*}
\phi(x)=\pi^{-\frac{d}{4}} e^{-\frac{1}{2}|x|^{2}}, \quad x \in \mathbf{R}^{d} \tag{1.26}
\end{equation*}
$$

in the transform $\mathcal{T}_{\phi}$. For this $\phi$ the relationship between the Bargmann transform and $\mathcal{T}_{\phi}$ is

$$
\begin{equation*}
\mathfrak{V}_{d}=U_{\mathfrak{V}} \circ \mathcal{T}_{\phi}, \quad \text { and } \quad U_{\mathfrak{V}}^{-1} \circ \mathfrak{V}_{d}=\mathcal{T}_{\phi} \tag{1.27}
\end{equation*}
$$

where $U_{\mathfrak{V}}$ is the linear, continuous and bijective operator on $\mathscr{D}^{\prime}\left(\mathbf{R}^{2 d}\right) \simeq \mathscr{D}^{\prime}\left(\mathbf{C}^{d}\right)$, given by

$$
\begin{equation*}
U_{\mathfrak{V}} F(x+i \xi)=(2 \pi)^{\frac{d}{2}} e^{\frac{1}{2}\left(|x|^{2}+|\xi|^{2}\right)} e^{i\langle x, \xi\rangle} F(\sqrt{2} x,-\sqrt{2} \xi), \quad x, \xi \in \mathbf{R}^{d} \tag{1.28}
\end{equation*}
$$

cf. [30] in combination with (1.11).
In analytic operator theory we need subspaces of

$$
\widehat{A}\left(\mathbf{C}^{2 d}\right) \equiv\left\{\Theta K ; K \in A\left(\mathbf{C}^{2 d}\right)\right\}
$$

where the semi-conjugation operator is

$$
\begin{equation*}
(\Theta K)(z, w)=K(z, \bar{w}), \quad z, w \in \mathbf{C}^{d} \tag{1.29}
\end{equation*}
$$

If $T$ is a linear and continuous operator from $\mathcal{S}_{1 / 2}\left(\mathbf{R}^{d}\right)$ to $\mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{d}\right)$, then there is a unique $K \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ such that $\Theta K \in \mathcal{A}_{1 / 2}^{\prime}\left(\mathbf{C}^{2 d}\right)$ and $\mathfrak{V}_{d} \circ T \circ \mathfrak{V}_{d}^{-1}$ is given by

$$
\begin{equation*}
F(z) \mapsto \int_{\mathbf{C}^{d}} K(z, w) F(w) d \mu(w) . \tag{1.30}
\end{equation*}
$$

(See e.g. [29].) For these reasons we let

$$
\hat{\mathcal{A}}_{0, s}\left(\mathbf{C}^{2 d}\right), \quad \hat{\mathcal{A}}_{s}\left(\mathbf{C}^{2 d}\right), \quad \hat{\mathcal{A}}_{\mathscr{S}}\left(\mathbf{C}^{2 d}\right), \quad \hat{\mathcal{A}}_{\mathscr{A}}^{\prime}\left(\mathbf{C}^{2 d}\right), \quad \hat{\mathcal{A}}_{s}^{\prime}\left(\mathbf{C}^{2 d}\right) \quad \text { and } \quad \hat{\mathcal{A}}_{0, s}^{\prime}\left(\mathbf{C}^{2 d}\right)
$$

be the images of

$$
\mathcal{A}_{0, s}\left(\mathbf{C}^{2 d}\right), \quad \mathcal{A}_{s}\left(\mathbf{C}^{2 d}\right), \quad \mathcal{A}_{\mathscr{S}}\left(\mathbf{C}^{2 d}\right), \quad \mathcal{A}_{\mathscr{S}}^{\prime}\left(\mathbf{C}^{2 d}\right), \quad \mathcal{A}_{s}^{\prime}\left(\mathbf{C}^{2 d}\right) \quad \text { and } \quad \mathcal{A}_{0, s}^{\prime}\left(\mathbf{C}^{2 d}\right)
$$

respectively, under the map $\Theta$. We also let $\widehat{A^{p}}\left(\mathbf{C}^{2 d}\right)$ and $\widehat{\mathcal{A}}_{b_{1}}\left(\mathbf{C}^{2 d}\right)$ be the images of $A^{p}\left(\mathbf{C}^{2 d}\right)$ and $\mathcal{A}_{b_{1}}\left(\mathbf{C}^{2 d}\right)$, respectively, under the map $\Theta$. The topologies of the former spaces are inherited from the corresponding latter spaces.

The semi-conjugated Bargmann (SCB) transform is defined as

$$
\mathfrak{V}_{\Theta, d}=\Theta \circ \mathfrak{V}_{2 d} .
$$

All properties of the Bargmann transform carry over naturally to analogous properties for the SCB transform.

### 1.4. Pseudo-differential operators

Let $A$ be a real $d \times d$ matrix. The pseudo-differential operator $\mathrm{Op}_{A}(\mathfrak{a})$ with symbol $\mathfrak{a} \in \mathcal{S}_{1 / 2}\left(\mathbf{R}^{2 d}\right)$ is the linear and continuous operator on $\mathcal{S}_{1 / 2}\left(\mathbf{R}^{d}\right)$ given by

$$
\begin{equation*}
\mathrm{Op}_{A}(\mathfrak{a}) f(x)=(2 \pi)^{-d} \iint_{\mathbf{R}^{2 d}} \mathfrak{a}(x-A(x-y), \xi) f(y) e^{i\langle x-y, \xi\rangle} d y d \xi, \quad x \in \mathbf{R}^{d} . \tag{1.31}
\end{equation*}
$$

For $\mathfrak{a} \in \mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$ the pseudo-differential operator $\operatorname{Op}_{A}(\mathfrak{a})$ is defined as the continuous operator from $\mathcal{S}_{1 / 2}\left(\mathbf{R}^{d}\right)$ to $\mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{d}\right)$ with distribution kernel

$$
\begin{equation*}
K_{\mathfrak{a}, A}(x, y)=(2 \pi)^{-\frac{d}{2}} \mathscr{F}_{2}^{-1} \mathfrak{a}(x-A(x-y), x-y), \quad x, y \in \mathbf{R}^{d}, \tag{1.32}
\end{equation*}
$$

where $\mathscr{F}_{2} F$ is the partial Fourier transform of $F(x, y) \in \mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$ with respect to the $y$ variable. This definition makes sense since the mappings

$$
\begin{equation*}
\mathscr{F}_{2} \quad \text { and } \quad F(x, y) \mapsto F(x-A(x-y), x-y) \tag{1.33}
\end{equation*}
$$

are homeomorphisms on $\mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$. The map $\mathfrak{a} \mapsto K_{\mathfrak{a}, A}$ is hence a homeomorphism on $\mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$.
If $A$ and $B$ are real $d \times d$ matrices and $\mathfrak{a} \in \mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$, then there is a unique $\mathfrak{b} \in \mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$ such that $\mathrm{Op}_{A}(\mathfrak{a})=\mathrm{Op}_{B}(\mathfrak{b})$, and $\mathfrak{b}$ can be obtained by

$$
\begin{equation*}
\mathrm{Op}_{A}(\mathfrak{a})=\mathrm{Op}_{B}(\mathfrak{b}) \quad \Leftrightarrow \quad e^{i\left\langle A D_{\xi}, D_{x}\right\rangle} \mathfrak{a}(x, \xi)=e^{i\left\langle B D_{\xi}, D_{x}\right\rangle} \mathfrak{b}(x, \xi) \tag{1.34}
\end{equation*}
$$

(see [9,17]).
Remark 1.3. By Fourier's inversion formula, (1.32) and the kernel theorem [21, Theorem 2.2], [28, Theorem 2.5] for operators from Gelfand-Shilov spaces to their duals, it follows that the map $\mathfrak{a} \mapsto \mathrm{Op}_{A}(\mathfrak{a})$ is bijective from $\mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$ to the set of all linear and continuous operators from $\mathcal{S}_{1 / 2}\left(\mathbf{R}^{d}\right)$ to $\mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$.

If $A=0$ then $\operatorname{Op}_{A}(\mathfrak{a})=\operatorname{Op}_{0}(\mathfrak{a})=\operatorname{Op}(\mathfrak{a})=\mathfrak{a}(x, D)$ is the Kohn-Nirenberg or standard representation. If $A=\frac{1}{2} I_{d}$ where $I_{d}$ is the $d \times d$ identity matrix then $\operatorname{Op}_{A}(\mathfrak{a})=\mathrm{Op}^{w}(\mathfrak{a})$ is the Weyl quantization. In this paper we use mainly the Weyl quantization and we put

$$
K_{\mathfrak{a}}^{w}=K_{\mathfrak{a}, I_{d} / 2} .
$$

The Weyl product $\mathfrak{a} \# \mathfrak{b}$ of two Weyl symbols $\mathfrak{a}, \mathfrak{b} \in \mathcal{S}_{1 / 2}\left(\mathbf{R}^{2 d}\right)$ is defined as the product of symbols corresponding to operator composition. Thus

$$
\mathrm{Op}^{w}(\mathfrak{a} \# \mathfrak{b})=\mathrm{Op}^{w}(\mathfrak{a}) \circ \mathrm{Op}^{w}(\mathfrak{b})
$$

and the Weyl product can be extended to larger spaces as long as composition is well defined.
Next we recall the definition of Wick operators. Suppose that $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ satisfies

$$
\begin{equation*}
w \mapsto a(z, w) e^{r|w|-|w|^{2}} \in L^{1}\left(\mathbf{C}^{d}\right) \tag{1.35}
\end{equation*}
$$

locally uniformly with respect to $z \in \mathbf{C}^{d}$ for every $r>0$. Then the analytic pseudo-differential operator, or Wick operator $\mathrm{Op}_{\mathfrak{V}}(a)$ with symbol $a$ and acting on $F \in \mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$, is defined by

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}(a) F(z)=\int_{\mathbf{C}^{d}} a(z, w) F(w) e^{(z, w)} d \mu(w), \quad z \in \mathbf{C}^{d} \tag{1.36}
\end{equation*}
$$

(Cf. e.g. [5,12,29-31].) The condition (1.35) and $F \in \mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$ imply that the integrand on the right-hand side of (1.36) is well defined. The locally uniform condition (1.35) with respect to $z \in \mathbf{C}^{d}$ implies that $\mathrm{Op}_{\mathfrak{V}}(a) F \in A\left(\mathbf{C}^{d}\right)$.

In [29] several extensions and restrictions of $\mathrm{Op}_{\mathfrak{V}}(a)$ are given. The following result follows from [29, Theorems 2.7 and 2.8]. Here $\mathcal{L}\left(\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right), A\left(\mathbf{C}^{d}\right)\right)$ is the space of all linear and continuous operators from $\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$ to $A\left(\mathbf{C}^{d}\right)$.

Proposition 1.4. The map $a \mapsto \operatorname{Op}_{\mathfrak{V}}(a)$ from $\widehat{\mathcal{A}}_{b_{1}}\left(\mathbf{C}^{2 d}\right)$ to $\mathcal{L}\left(\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right), A\left(\mathbf{C}^{d}\right)\right)$ is uniquely extendable to $a$ bijective map from $\hat{A}\left(\mathbf{C}^{2 d}\right)$ to $\mathcal{L}\left(\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right), A\left(\mathbf{C}^{d}\right)\right)$.

Let $L_{A}\left(\mathbf{C}^{2 d}\right)$ be the set of all $a \in L_{\text {loc }}^{1}\left(\mathbf{C}^{2 d}\right)$ such that $z \mapsto a(z, w)$ is entire for almost every $w \in \mathbf{C}^{d}$ and

$$
\begin{equation*}
w \mapsto \sup _{\alpha \in \mathbf{N}^{d}}\left|\frac{\partial_{z}^{\alpha} a(z, w) \cdot e^{r|w|-|w|^{2}}}{h^{|\alpha|} \alpha!}\right| \in L^{1}\left(\mathbf{C}^{d}\right) \tag{1.37}
\end{equation*}
$$

for every $h, r>0$ and $z \in \mathbf{C}^{d}$. If $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ satisfies (1.35) then $a \in L_{A}\left(\mathbf{C}^{2 d}\right)$ as a consequence of Cauchy's integral formula. Thus $L_{A}\left(\mathbf{C}^{2 d}\right)$ is a relaxation of the former condition.

If $a \in L_{A}\left(\mathbf{C}^{2 d}\right)$ then $\mathrm{Op}_{\mathfrak{V}}(a): \mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right) \rightarrow \mathcal{A}_{b_{1}}^{\prime}\left(\mathbf{C}^{d}\right)=A\left(\mathbf{C}^{d}\right)$ is continuous. Hence the following result is a straight-forward consequence of Proposition 1.4 and the fact that $\widehat{\mathcal{A}}_{b_{1}}^{\prime}\left(\mathbf{C}^{2 d}\right)=\widehat{A}\left(\mathbf{C}^{2 d}\right)$.

Proposition 1.5. Let $a \in L_{A}\left(\mathbf{C}^{2 d}\right)$. Then there is a unique $a_{0} \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ such that $\mathrm{Op}_{\mathfrak{V}}(a)=\mathrm{Op}_{\mathfrak{V}}\left(a_{0}\right)$ as mappings from $\mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$ to $\mathcal{A}_{b_{1}}^{\prime}\left(\mathbf{C}^{d}\right)$. It holds

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}(a)=\mathrm{Op}_{\mathfrak{V}}\left(a_{0}\right) \quad \text { where } \quad a_{0}(z, w)=\pi^{-d} \int_{\mathbf{C}^{d}} a\left(z, w_{1}\right) e^{-\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right) . \tag{1.38}
\end{equation*}
$$

Proof. The operator $\Pi_{A}$ defined in (1.22) is the orthogonal projection from $L^{2}(d \mu)$ to $A^{2}\left(\mathbf{C}^{d}\right)$ which is uniquely extendable to a continuous map from

$$
\begin{equation*}
L_{0, A}\left(\mathbf{C}^{d}\right) \equiv\left\{a_{0} \in L_{\mathrm{loc}}^{1}\left(\mathbf{C}^{d}\right) ; w \mapsto a_{0}(w) e^{r|w|-|w|^{2}} \in L^{1}\left(\mathbf{C}^{d}\right) \text { for every } r>0\right\} \tag{1.39}
\end{equation*}
$$

to $A\left(\mathbf{C}^{d}\right)$ (see e.g. [30]). Hence, if $F, G \in \mathcal{A}_{b_{1}}\left(\mathbf{C}^{d}\right)$ and $a_{0}$ is given by (1.38) then

$$
\begin{aligned}
& \left(\mathrm{Op}_{\mathfrak{V}}(a) F, G\right)_{A^{2}}=\left(\left(\mathrm{Op}_{\mathfrak{V}}(a) \circ \Pi_{A}\right) F, G\right)_{A^{2}} \\
& \qquad \begin{aligned}
=\left(\int_{\mathbf{C}^{d}}\left(\int_{\mathbf{C}^{d}} a\left(\cdot, w_{1}\right) e^{\left(\cdot, w_{1}\right)} e^{\left(w_{1}, w\right)} d \mu\left(w_{1}\right)\right) F(w) d \mu(w), G\right)_{A^{2}}
\end{aligned} \\
& =\left(\int_{\mathbf{C}^{d}} a_{0}(\cdot, w) e^{(\cdot, w)} F(w) d \mu(w), G\right)_{A^{2}}=\left(\mathrm{Op}_{\mathfrak{V}}\left(a_{0}\right) F, G\right)_{A^{2}},
\end{aligned}
$$

and thus $\mathrm{Op}_{\mathfrak{V}}(a)=\mathrm{Op}_{\mathfrak{V}}\left(a_{0}\right)$ follows. The assertion now follows from Proposition 1.4 and the fact that $a_{0}$ in the integral formula of $(1.38)$ defines an element in $\widehat{A}\left(\mathbf{C}^{2 d}\right)$.

We will also consider anti-Wick operators [12,5,6] defined by

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right) F(z)=\int_{\mathbf{C}^{d}} a_{0}(w) F(w) e^{(z, w)} d \mu(w), \quad z \in \mathbf{C}^{d}, \tag{1.40}
\end{equation*}
$$

when $a_{0} \in L_{0, A}\left(\mathbf{C}^{d}\right)$ and $F$ belongs to $\mathcal{A}_{0}\left(\mathbf{C}^{d}\right)$, the space of analytic polynomials on $\mathbf{C}^{d}$. Then $a_{0} \in L_{0, A}\left(\mathbf{C}^{d}\right)$ if and only if $a(z, w) \equiv a_{0}(w)$ belongs to $L_{A}\left(\mathbf{C}^{2 d}\right)$, and then $\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right)=\mathrm{Op}_{\mathfrak{V}}(a)$. Consequently, all results for Wick operators with symbols in $L_{A}\left(\mathbf{C}^{2 d}\right)$ hold for anti-Wick operators. In particular, if $a_{0} \in L_{0, A}\left(\mathbf{C}^{d}\right)$, then $\mathrm{Op}_{\mathfrak{Z}}^{\mathrm{aw}}\left(a_{0}\right): \mathcal{A}_{\mathfrak{b}_{1}}\left(\mathbf{C}^{d}\right) \rightarrow A\left(\mathbf{C}^{d}\right)$ is continuous. We denote the Wick symbol of the anti-Wick operator $\mathrm{Op}_{\mathfrak{Y}}^{\text {aw }}\left(a_{0}\right)$ by $a_{0}^{\text {aw }}$. Then (1.38) takes the form

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right)=\mathrm{Op}_{\mathfrak{V}}\left(a_{0}^{\mathrm{aw}}\right) \quad \text { where } \quad a_{0}^{\mathrm{aw}}(z, w)=\pi^{-d} \int_{\mathbf{C}^{d}} a_{0}\left(w_{1}\right) e^{-\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right) . \tag{1.38}
\end{equation*}
$$

Pseudo-differential operators on $\mathbf{R}^{d}$ may be transferred to Wick operators on $\mathbf{C}^{d}$ by means of the Bargmann transform.

Definition 1.6. Let $\mathfrak{a} \in \mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$.
(1) the Bargmann assignment $\mathrm{S}_{\mathfrak{V} \mathfrak{a}}$ of $\mathfrak{a}$ is the unique element $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ which fulfills

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}(a)=\mathfrak{V}_{d} \circ \mathrm{Op}^{w}(\mathfrak{a}) \circ \mathfrak{V}_{d}^{*} \quad \Leftrightarrow \quad a=\mathrm{S}_{\mathfrak{V}} \mathfrak{a} ; \tag{1.41}
\end{equation*}
$$

(2) the Bargmann kernel assignment $K_{\mathfrak{V}, \mathfrak{a}}$ of $\mathfrak{a}$ is the unique element $K \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$, which is the kernel of the map $\mathfrak{V}_{d} \circ \mathrm{Op}^{w}(\mathfrak{a}) \circ \mathfrak{V}_{d}^{*}$ with respect to the sesquilinear $A^{2}$ form.

By the definitions we have

$$
\begin{equation*}
K_{\mathfrak{V}, \mathfrak{a}}(z, w)=e^{(z, w)} \mathrm{S}_{\mathfrak{V} \mathfrak{a}}(z, w) . \tag{1.42}
\end{equation*}
$$

Example 1.7. The creation and annihilation operators

$$
2^{-\frac{1}{2}}\left(x_{j}-\partial_{x_{j}}\right) \text { and } 2^{-\frac{1}{2}}\left(x_{j}+\partial_{x_{j}}\right),
$$

are transfered to the operators

$$
\begin{equation*}
F \mapsto z_{j} F \quad \text { and } \quad F \mapsto \partial_{z_{j}} F, \tag{1.43}
\end{equation*}
$$

by the Bargmann transform (see [3]). The Wick symbols of the operators in (1.43) are $z_{j}$ and $\bar{w}_{j}$, respectively $[5,30]$. By combining these identities with the fact that the Weyl symbol of $i^{-1} \partial_{x_{j}}$ equals $\xi_{j}$ we get

$$
\begin{gather*}
\mathrm{S}_{\mathfrak{V}}\left(2^{-\frac{1}{2}}\left(x_{j}-i \xi_{j}\right)\right)=z_{j}, \quad \mathrm{~S}_{\mathfrak{V}}\left(2^{-\frac{1}{2}}\left(x_{j}+i \xi_{j}\right)\right)=\bar{w}_{j}, \\
\mathrm{~S}_{\mathfrak{V}}\left(x_{j}\right)=2^{-\frac{1}{2}}\left(z_{j}+\bar{w}_{j}\right) \quad \text { and } \quad \mathrm{S}_{\mathfrak{V}}\left(\xi_{j}\right)=2^{-\frac{1}{2}} i\left(z_{j}-\bar{w}_{j}\right) . \tag{1.44}
\end{gather*}
$$

We need to compare $K_{\mathfrak{a}}^{w}$ and $K_{\mathfrak{V}, \mathfrak{a}}$. On the one hand we have for $f, g \in \mathscr{S}\left(\mathbf{R}^{d}\right)$

$$
\left(\mathrm{Op}^{w}(\mathfrak{a}) f, g\right)_{L^{2}\left(\mathbf{R}^{d}\right)}=\left(K_{\mathfrak{a}}^{w}, g \otimes \bar{f}\right)_{L^{2}\left(\mathbf{R}^{2 d}\right)}=\left(\mathfrak{V}_{2 d} K_{\mathfrak{a}}^{w}, \mathfrak{V}_{2 d}(g \otimes \bar{f})\right)_{A^{2}\left(\mathbf{C}^{2 d}\right)}
$$

and on the other hand

$$
\begin{aligned}
\left(\mathrm{Op}^{w}(\mathfrak{a}) f, g\right)_{L^{2}\left(\mathbf{R}^{d}\right)} & =\left(\mathrm{Op}_{\mathfrak{V}}(a) \mathfrak{V}_{d} f, \mathfrak{V}_{d} g\right)_{A^{2}\left(\mathbf{C}^{d}\right)} \\
& =\left(K_{\mathfrak{V}, \mathfrak{a}}, \mathfrak{V}_{d} g \otimes \overline{\mathfrak{V}_{d} f}\right)_{A^{2}\left(\mathbf{C}^{2 d}\right)} \\
& =\left(\Theta K_{\mathfrak{V}, \mathfrak{a}}, \Theta\left(\mathfrak{V}_{d} g \otimes \overline{\mathfrak{V}_{d} f}\right)\right)_{\hat{A}^{2}\left(\mathbf{C}^{2 d}\right)} .
\end{aligned}
$$

Since

$$
\Theta\left(\mathfrak{V}_{d} g \otimes \overline{\mathfrak{V}_{d} f}\right)(z, w)=\mathfrak{V}_{d} g(z) \overline{\mathfrak{V}_{d} f(\bar{w})}=\mathfrak{V}_{2 d}(g \otimes \bar{f})(z, w)
$$

we obtain

$$
\begin{equation*}
K_{\mathfrak{V}, \mathfrak{a}}=\Theta \mathfrak{V}_{2 d} K_{\mathfrak{a}}^{w}=\mathfrak{V}_{\Theta, d} K_{\mathfrak{a}}^{w} . \tag{1.45}
\end{equation*}
$$

### 1.5. Symbol classes for pseudo-differential operators on $\mathbf{R}^{d}$

In order to define a generalized family of Shubin symbol classes [26], we need to add a restriction of the involved weights. Let $\rho \in[0,1]$, and let $\mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{d}\right)$ be the set of all $\omega \in \mathscr{P}\left(\mathbf{R}^{d}\right) \cap C^{\infty}\left(\mathbf{R}^{d}\right)$ such that for every multi-index $\alpha \in \mathbf{N}^{d}$,

$$
\begin{equation*}
\left|\partial^{\alpha} \omega(x)\right| \lesssim \omega(x)\langle x\rangle^{-\rho|\alpha|}, \quad x \in \mathbf{R}^{d} . \tag{1.46}
\end{equation*}
$$

For $\omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{d}\right)$ the Shubin symbol class $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{d}\right)$ is the set of all $f \in C^{\infty}\left(\mathbf{R}^{d}\right)$ such that for every $\alpha \in \mathbf{N}^{2 d}$,

$$
\begin{equation*}
\left|\partial^{\alpha} f(x)\right| \lesssim \omega(x)\langle x\rangle^{-\rho|\alpha|}, \quad x \in \mathbf{R}^{d} . \tag{1.47}
\end{equation*}
$$

Let $\rho \in[0,1], \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{2 d}\right)$ and $A$ be a real $d \times d$ matrix. Then it follows from [26] or [17, Section 18.5] that $e^{i\left\langle A D_{\xi}, D_{x}\right\rangle}$ is a homeomorphism on $\mathrm{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$, which implies that the set

$$
\left\{\mathrm{Op}_{A}(\mathfrak{a}) ; \mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)\right\}
$$

is independent of the choice of $A$, in view of (1.34). If $B$ is another real $d \times d$ matrix and $\mathfrak{a}, \mathfrak{b} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ satisfy (1.34), then it follows from [17, Section 18.5] that

$$
\begin{equation*}
\mathfrak{a}-\mathfrak{b} \in \operatorname{Sh}_{\rho}^{\left(\omega_{\rho}\right)}\left(\mathbf{R}^{2 d}\right), \quad \text { where } \quad \omega_{\rho}(x, \xi)=\omega(x, \xi)\langle(x, \xi)\rangle^{-2 \rho} . \tag{1.48}
\end{equation*}
$$

In particular

$$
\begin{equation*}
|\mathfrak{a}(x, \xi)-\mathfrak{b}(x, \xi)| \lesssim \omega(x, \xi)\langle(x, \xi)\rangle^{-2 \rho} . \tag{1.49}
\end{equation*}
$$

We also need the symbol classes defined in [1, Definition 1.8] with symbols satisfying estimates of the form

$$
\begin{equation*}
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} \mathfrak{a}(x, \xi)\right| \lesssim h^{|\alpha+\beta|} \alpha!^{\sigma} \beta!^{s} e^{r\left(|x|^{\frac{1}{s}}+|\xi|^{\frac{1}{\sigma}}\right)}, \quad x, \xi \in \mathbf{R}^{d} . \tag{1.50}
\end{equation*}
$$

(See also [9] for the restricted case when $s=\sigma$.)
Definition 1.8. Let $s, \sigma>0$. Then
(1) $\Gamma_{s, \sigma}^{\sigma, s ; 0}\left(\mathbf{R}^{2 d}\right)$ consists of all $\mathfrak{a} \in C^{\infty}\left(\mathbf{R}^{2 d}\right)$ such that for some $r>0$, (1.50) holds for every $h>0$;
(2) $\Gamma_{s, \sigma ; 0}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$ consists of all $\mathfrak{a} \in C^{\infty}\left(\mathbf{R}^{2 d}\right)$ such that for some $h>0$, (1.50) holds for every $r>0$;
(3) $\Gamma_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$ consists of all $\mathfrak{a} \in C^{\infty}\left(\mathbf{R}^{2 d}\right)$ such that (1.50) holds for some $h>0$ and some $r>0$.

Remark 1.9. The symbol classes $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ have isotropic behavior with respect to phase space $T^{*} \mathbf{R}^{d} \simeq$ $\mathbf{R}^{2 d}$, and the same holds for the symbol classes in Definition 1.8 when $\sigma=s$. See also [9] for the restricted case when $s=\sigma$, and [2] for a bilinear extension. Important classes similar to those given by Definition 1.8 are considered in [23].

Pseudo-differential operators with symbols in the classes in Definition 1.8 are examples of so called operators of infinite order. These operators are continuous on appropriate Gelfand-Shilov (distribution) spaces [1,9]. The next result characterizes the symbol classes in Definition 1.8 by means of estimates of form

$$
\begin{equation*}
\left|\mathcal{T}_{\psi} \mathfrak{a}(x, \xi, \eta, y)\right| \lesssim e^{r_{1}\left(|x|^{\frac{1}{s}}+|\xi|^{\frac{1}{\sigma}}\right)-r_{2}\left(|\eta|^{\frac{1}{\sigma}}+|y|^{\frac{1}{s}}\right)}, \quad x, \xi, y, \eta \in \mathbf{R}^{d} . \tag{1.51}
\end{equation*}
$$

We omit the proof since the result is a special case of [1, Proposition 2.1']. We refer to [1, Subsection 1.1] for the definition of the Gelfand-Shilov spaces $\mathcal{S}_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right), \Sigma_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$ and their distribution spaces.

Proposition 1.10. Let $s, \sigma>0$ and let $\mathfrak{a} \in C^{\infty}\left(\mathbf{R}^{2 d}\right)$. Then the following is true:
(1) if $\psi \in \mathcal{S}_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right) \backslash 0$, then $\mathfrak{a} \in \Gamma_{s, \sigma ; 0}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$ if and only if (1.51) holds for some $r_{2}>0$ and every $r_{1}>0$;
(2) if $\psi \in \Sigma_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right) \backslash 0$, then $\mathfrak{a} \in \Gamma_{s, \sigma}^{\sigma, s ; 0}\left(\mathbf{R}^{2 d}\right)$ if and only if (1.51) holds for some $r_{1}>0$ and all $r_{2}>0$;
(3) if $\psi \in \Sigma_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right) \backslash 0$, then $\mathfrak{a} \in \Gamma_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$ if and only if (1.51) holds for some $r_{1}>0$ and some $r_{2}>0$.

### 1.6. Elliptic, weakly elliptic and hypoelliptic elements in $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{d}\right)$

Let $\rho \geqslant 0$ and $\omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{d}\right)$. Then $f \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{d}\right)$ is called weakly elliptic of order $\rho_{0} \geqslant 0,\left(\operatorname{in~} \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{d}\right)\right)$, or $\rho_{0}$-weakly elliptic, if there is an $R>0$ such that

$$
|f(x)| \gtrsim\langle x\rangle^{-\rho_{0}} \omega(x), \quad|x| \geqslant R .
$$

A weakly elliptic function of order 0 is called elliptic.
Let $A$ and $B$ be real $d \times d$ matrices, $\rho>0, \rho_{0} \in[0,2 \rho), \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{2 d}\right)$ and suppose that $\mathfrak{a}, \mathfrak{b} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ satisfy (1.34). It follows from (1.48) that $\mathfrak{a}$ is weakly elliptic of order $\rho_{0}$, if and only if $\mathfrak{b}$ is weakly elliptic of order $\rho_{0}$. In particular, $\mathfrak{a}$ is elliptic, if and only if $\mathfrak{b}$ is elliptic.

Next we define Shubin hypoelliptic symbols (cf. Definitions 5.1 and 25.1 in [26]).
Definition 1.11. Let $\rho>0, \rho_{0} \geqslant 0, \omega_{0} \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{d}\right)$ and $f \in \operatorname{Sh}_{\rho}^{\left(\omega_{0}\right)}\left(\mathbf{R}^{d}\right)$. Then $f$ is called hypoelliptic (in Shubin's sense in $\left.\operatorname{Sh}_{\rho}^{\left(\omega_{0}\right)}\left(\mathbf{R}^{d}\right)\right)$ of order $\rho_{0}$, if there is an $R>0$ such that for every $\alpha \in \mathbf{N}^{d}$, it holds

$$
\left|\partial^{\alpha} f(x)\right| \lesssim|f(x)|\langle x\rangle^{-\rho|\alpha|}, \quad|x| \geqslant R,
$$

and

$$
|f(x)| \gtrsim \omega_{0}(x)\langle x\rangle^{-\rho_{0}}, \quad|x| \geqslant R .
$$

Elliptic and hypoelliptic symbols are important since they give rise to parametrices. For $\rho, \omega$ as above and $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ elliptic, there is an elliptic symbol $\mathfrak{b} \in \operatorname{Sh}_{\rho}^{(1 / \omega)}\left(\mathbf{R}^{2 d}\right)$ such that

$$
\mathrm{Op}_{A}(\mathfrak{a}) \circ \mathrm{Op}_{A}(\mathfrak{b})=I+\mathrm{Op}_{A}\left(\mathfrak{c}_{1}\right) \quad \text { and } \quad \mathrm{Op}_{A}(\mathfrak{b}) \circ \mathrm{Op}_{A}(\mathfrak{a})=I+\mathrm{Op}_{A}\left(\mathfrak{c}_{2}\right)
$$

for some $\mathfrak{c}_{1}, \mathfrak{c}_{2} \in \mathscr{S}\left(\mathbf{R}^{2 d}\right)$. An operator $\mathrm{Op}(\mathfrak{c})$ with $c \in \mathscr{S}\left(\mathbf{R}^{2 d}\right)$ is regularizing in the sense that $\mathrm{Op}(\mathfrak{c})$ is continuous from $\mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$ to $\mathscr{S}\left(\mathbf{R}^{d}\right)$. (Cf. e.g. [7,26].)

## 2. Reformulation of pseudo-differential calculus using the Bargmann transform

In this section we characterize the Bargmann assignment of pseudo-differential operator symbols from Subsection 1.5, using estimates of complex derivatives. In Subsection 2.1 we show how pseudo-differential operators on $\mathbf{R}^{d}$ with Shubin symbols are transformed to Wick operators by the Bargmann transform. In Subsection 2.3 we deduce similar links between pseudo-differential operators of infinite order, given in the second part of Subsection 1.5, and suitable classes of Wick operators. Subsection 2.4 treats composition formulae for symbols of Wick operators, which leads to algebraic properties for operators in Subsection 2.1 and 2.3. As an application we obtain short proofs of composition results for pseudo-differential operators on $\mathbf{R}^{d}$ from Subsection 1.5.

### 2.1. Wick symbols of Shubin pseudo-differential operators

The following proposition is essential in the characterization of Shubin type pseudo-differential operators on $\mathbf{R}^{d}$ by means of the corresponding Wick symbols. The Shubin classes can be characterized using the transform $\mathcal{T}_{\phi}$ by means of estimates of the form

$$
\begin{align*}
\left|\partial_{x}^{\alpha} \partial_{\xi}^{\beta} \mathcal{T}_{\phi} f(x, \xi)\right| & \lesssim \omega(x)\langle x\rangle^{-\rho|\alpha|}\langle\xi\rangle^{-N},  \tag{2.1}\\
\left|\partial_{x}^{\alpha} \mathcal{T}_{\phi} f(x, \xi)\right| & \lesssim \omega(x)\langle x\rangle^{-\rho|\alpha|}\langle\xi\rangle^{-N} \tag{2.2}
\end{align*}
$$

and

$$
\begin{equation*}
\left|\mathcal{T}_{\phi} f(x, \xi)\right| \lesssim \omega(x)\langle\xi\rangle^{-N} . \tag{2.3}
\end{equation*}
$$

The proof of the following result is similar to the proof of [8, Proposition 3.2].
Proposition 2.1. Let $0 \leqslant \rho \leqslant 1$, let $\omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{d}\right)$, and suppose $f \in \mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$ and $\phi \in \mathscr{S}\left(\mathbf{R}^{d}\right) \backslash 0$. The following conditions are equivalent:
(1) $f \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{d}\right)$,
(2) (2.1) holds true for any $N \geqslant 0$ and $\alpha, \beta \in \mathbf{N}^{d}$,
(3) (2.2) holds true for any $N \geqslant 0$ and $\alpha \in \mathbf{N}^{d}$,
and the following conditions are equivalent:
(1) $f \in \operatorname{Sh}_{0}^{(\omega)}\left(\mathbf{R}^{d}\right)$,
(2)' (2.3) holds true for any $N \geqslant 0$.

Proof. First we prove that (1) implies (2). Suppose $f \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{d}\right)$ and let $\alpha, \beta, \gamma \in \mathbf{N}^{d}$ be arbitrary. We will show

$$
\left|\xi^{\gamma} \partial_{x}^{\alpha} \partial_{\xi}^{\beta} \mathcal{T}_{\phi} f(x, \xi)\right| \lesssim \omega(x)\langle x\rangle^{-\rho|\alpha|} .
$$

By (1.14), (1.15) and integration by parts we get

$$
\begin{aligned}
\left|\xi^{\gamma} \partial_{x}^{\alpha} \partial_{\xi}^{\beta} \mathcal{T}_{\phi} f(x, \xi)\right| & =\left|\xi^{\gamma} \mathcal{T}_{\phi_{\beta}}\left(\partial^{\alpha} f\right)(x, \xi)\right| \\
& =(2 \pi)^{-\frac{d}{2}}\left|\int_{\mathbf{R}^{d}}\left(\left(i \partial_{y}\right)^{\gamma} e^{-i\langle\xi, y\rangle}\right) \overline{\phi_{\beta}(y)} \partial^{\alpha} f(x+y) d y\right| \\
& \lesssim \int_{\mathbf{R}^{d}}\left|\partial_{y}^{\gamma}\left[\overline{\phi_{\beta}(y)} \partial^{\alpha} f(x+y)\right]\right| d y \\
& =\int_{\mathbf{R}^{d}}\left|\sum_{\kappa \leqslant \gamma}\binom{\gamma}{\kappa} \partial^{\gamma-\kappa \overline{\phi_{\beta}(y)}} \partial^{\alpha+\kappa} f(x+y)\right| d y \\
& \lesssim \sum_{\kappa \leqslant \gamma}\binom{\gamma}{\kappa} \int_{\mathbf{R}^{d}}\left|\partial^{\gamma-\kappa} \phi_{\beta}(y)\right| \omega(x+y)\langle x+y\rangle^{-\rho|\alpha+\kappa|} d y .
\end{aligned}
$$

Since $\omega$ is polynomially moderate, Peetre's inequality (1.2) and the fact that $\phi \in \mathscr{S}$ give

$$
\left|\xi^{\gamma} \partial_{x}^{\alpha} \partial_{\xi}^{\beta} \mathcal{T}_{\phi} f(x, \xi)\right| \lesssim \omega(x)\langle x\rangle^{-\rho|\alpha|} \sum_{\kappa \leqslant \gamma}\binom{\gamma}{\kappa} \int_{\mathbf{R}^{d}}\left|\partial^{\gamma-\kappa} \phi_{\beta}(y)\right| \omega(y)\langle y\rangle^{|m|+\rho|\alpha+\kappa|} d y=\omega(x)\langle x\rangle^{-\rho|\alpha|} .
$$

Thus $f \in \operatorname{Sh}^{(\omega)}\left(\mathbf{R}^{d}\right)$ implies (2.1), and as a special case (2.2), and $f \in \operatorname{Sh}_{0}^{(\omega)}\left(\mathbf{R}^{d}\right)$ implies (2.3). We have proved that (1) implies (2) which in turn implies (3), and that (1)' implies (2)'.

Conversely, suppose (3), that is $f \in \mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$ and (2.2) holds for all $N \geqslant 0$ and all $\alpha \in \mathbf{N}^{d}$, which is a weaker assumption than (2). We obtain from (1.13)

$$
\begin{aligned}
f(x) & =\|\phi\|_{L^{2}}^{-2} \mathcal{T}_{\phi}^{*} \mathcal{T}_{\phi} f(x) \\
& =\|\phi\|_{L^{2}}^{-2}(2 \pi)^{-\frac{d}{2}} \iint_{\mathbf{R}^{2 d}} \mathcal{T}_{\phi} f(y, \xi) e^{i\langle\xi, x-y\rangle} \phi(x-y) d y d \xi
\end{aligned}
$$

which is an absolutely convergent integral due to (2.2) and the fact that $\phi \in \mathscr{S}\left(\mathbf{R}^{d}\right)$. We may differentiate under the integral, so integration by parts, (2.2) and Peetre's inequality give for some $N_{0} \geqslant 0$, any $\alpha \in \mathbf{N}^{d}$ and any $x \in \mathbf{R}^{d}$ that

$$
\begin{aligned}
\left|\partial^{\alpha} f(x)\right| & =\|\phi\|_{L^{2}}^{-2}(2 \pi)^{-\frac{d}{2}}\left|\iint_{\mathbf{R}^{2 d}} \mathcal{T}_{\phi} f(y, \xi) \partial_{y}^{\alpha}\left(e^{i\langle\xi, x-y\rangle} \phi(x-y)\right) d y d \xi\right| \\
& =\|\phi\|_{L^{2}}^{-2}(2 \pi)^{-\frac{d}{2}}\left|\iint_{\mathbf{R}^{2 d}} \partial_{y}^{\alpha} \mathcal{T}_{\phi} f(y, \xi) e^{i\langle\xi, x-y\rangle} \phi(x-y) d y d \xi\right| \\
& =\|\phi\|_{L^{2}}^{-2}(2 \pi)^{-\frac{d}{2}}\left|\iint_{\mathbf{R}^{2 d}} \partial_{y}^{\alpha} \mathcal{T}_{\phi} f(x-y, \xi) e^{i\langle\xi, y\rangle} \phi(y) d y d \xi\right| \\
& \lesssim \iint_{\mathbf{R}^{2 d}} \omega(x-y)\langle x-y\rangle^{-\rho|\alpha|}\langle\xi\rangle^{-d-1}|\phi(y)| d y d \xi \\
& \lesssim \omega(x)\langle x\rangle^{-\rho|\alpha|} \iint_{\mathbf{R}^{2 d}}\langle\xi\rangle^{-d-1}\langle y\rangle^{N_{0}+\rho|\alpha|}|\phi(y)| d y d \xi \approx \omega(x)\langle x\rangle^{-\rho|\alpha|} .
\end{aligned}
$$

Thus $f \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{d}\right)$ and we have proved the equivalence of (1), (2) and (3).
It remains to show that (2) implies (1) , that is (2.3) for all $N \geqslant 0$ implies $f \in \operatorname{Sh}_{0}^{(\omega)}\left(\mathbf{R}^{d}\right)$. We have for some $N_{0} \geqslant 0$, any $\alpha \in \mathbf{N}^{d}, x \in \mathbf{R}^{d}$ and $N \geqslant 0$,

$$
\begin{aligned}
\left|\partial^{\alpha} f(x)\right| & =\|\phi\|_{L^{2}}^{-2}(2 \pi)^{-\frac{d}{2}}\left|\iint_{\mathbf{R}^{2 d}} \mathcal{T}_{\phi} f(y, \xi) \partial_{x}^{\alpha}\left(e^{i\langle\xi, x-y\rangle} \phi(x-y)\right) d y d \xi\right| \\
& \lesssim \sum_{\beta \leqslant \alpha}\binom{\alpha}{\beta} \iint_{\mathbf{R}^{2 d}}\left|\mathcal{T}_{\phi} f(y, \xi)\right|\langle\xi\rangle^{|\beta|}\left|\partial^{\alpha-\beta} \phi(x-y)\right| d y d \xi \\
& \lesssim \sum_{\beta \leqslant \alpha}\binom{\alpha}{\beta} \iint_{\mathbf{R}^{2 d}} \omega(y)\langle\xi\rangle^{|\alpha|-N}\left|\partial^{\alpha-\beta} \phi(x-y)\right| d y d \xi \\
& \lesssim \omega(x) \sum_{\beta \leqslant \alpha}\binom{\alpha}{\beta} \iint_{\mathbf{R}^{2 d}}\langle\xi\rangle^{|\alpha|-N}\langle x-y\rangle^{N_{0}}\left|\partial^{\alpha-\beta} \phi(x-y)\right| d y d \xi \\
& \lesssim \omega(x)
\end{aligned}
$$

provided $N$ is sufficiently large, since $\phi \in \mathscr{S}$. This shows that $f \in \operatorname{Sh}_{0}^{(\omega)}\left(\mathbf{R}^{d}\right)$.
We may now characterize the Shubin classes $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ by estimates on their Bargmann (kernel) assignments of the forms

$$
\begin{align*}
\left|\left(\partial_{z}+\bar{\partial}_{w}\right)^{\alpha}\left(\partial_{z}-\bar{\partial}_{w}\right)^{\beta} \mathrm{S}_{\mathfrak{V}} \mathfrak{a}(z, w)\right| & \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z+w\rangle^{-\rho|\alpha+\beta|}\langle z-w\rangle^{-N},  \tag{2.4}\\
\left|\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} \mathrm{S}_{\mathfrak{V} \mathfrak{a}}(z, w)\right| & \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z+w\rangle^{-\rho|\alpha+\beta|}\langle z-w\rangle^{-N},  \tag{2.5}\\
\left|\mathrm{~S}_{\mathfrak{V} \mathfrak{a}}(z, w)\right| & \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z-w\rangle^{-N} \tag{2.6}
\end{align*}
$$

and

$$
\begin{equation*}
\left|K_{\mathfrak{V}, \mathfrak{a}}(z, w)\right| \lesssim e^{\frac{1}{2}\left(|z|^{2}+|w|^{2}\right)} \omega(\sqrt{2} \bar{z})\langle z-w\rangle^{-N} . \tag{2.7}
\end{equation*}
$$

Theorem 2.2. Let $0 \leqslant \rho \leqslant 1, \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{2 d}\right)$ and $\mathfrak{a} \in \mathscr{S}^{\prime}\left(\mathbf{R}^{2 d}\right)$. The following conditions are equivalent:
(1) $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$,
(2) (2.4) holds true for every $N \geqslant 0, z, w \in \mathbf{C}^{d}$ and $\alpha, \beta \in \mathbf{N}^{d}$,
(3) (2.5) holds true for every $N \geqslant 0, z, w \in \mathbf{C}^{d}$ and $\alpha, \beta \in \mathbf{N}^{d}$,
and the following conditions are equivalent:
$(1)^{\prime} \mathfrak{a} \in \operatorname{Sh}_{0}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$,
(2)' (2.6) holds true for any $N \in \mathbf{N}$ and $z, w \in \mathbf{C}^{d}$,
(3)' (2.7) holds true for any $N \in \mathbf{N}$ and $z, w \in \mathbf{C}^{d}$.

For the proof we need the following proposition of independent interest. Here we recall that $\mathrm{S}_{\mathfrak{V}}$ is bijective from $\mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$ to the set

$$
\begin{equation*}
\left\{a \in \widehat{A}\left(\mathbf{C}^{2 d}\right) ;|a(z, w)| \lesssim e^{\left(\frac{1}{2}+r\right)|z-w|^{2}} \text { for every } r>0\right\} . \tag{2.8}
\end{equation*}
$$

Proposition 2.3. Let $\psi(x, \xi)=\left(\frac{2}{\pi}\right)^{\frac{d}{2}} e^{-\left(|x|^{2}+|\xi|^{2}\right)}, x, \xi \in \mathbf{R}^{d}, \mathfrak{a} \in \mathcal{S}_{1 / 2}^{\prime}\left(\mathbf{R}^{2 d}\right)$ and a belongs to the set in (2.8). Then

$$
\begin{equation*}
\mathrm{S}_{\mathfrak{F} \mathfrak{a}}(z, w)=(2 \pi)^{\frac{d}{2}} e^{\frac{1}{2}|z-w|^{2}} \mathcal{T}_{\psi} \mathfrak{a}\left(\frac{x+y}{\sqrt{2}},-\frac{\xi+\eta}{\sqrt{2}}, \sqrt{2}(\eta-\xi), \sqrt{2}(y-x)\right), \tag{2.9}
\end{equation*}
$$

and

$$
\begin{equation*}
\left(\mathrm{S}_{\mathfrak{V}}^{-1} a\right)(x,-\xi)=\left(\frac{2}{\pi}\right)^{d} \int_{\mathbf{C}^{d}} a\left(\frac{z}{\sqrt{2}}-w, \frac{z}{\sqrt{2}}+w\right) e^{-2|w|^{2}} d \lambda(w), \tag{2.10}
\end{equation*}
$$

with $z=x+i \xi, w=y+i \eta$ and $x, y, \xi, \eta \in \mathbf{R}^{d}$.
Proof. Let $\phi(x, y)=\pi^{-\frac{d}{2}} e^{-\frac{1}{2}\left(|x|^{2}+|y|^{2}\right)}$ for $x, y \in \mathbf{R}^{d}$, and let $K_{\mathfrak{a}}^{w}$ be the kernel of $\mathrm{Op}^{w}(\mathfrak{a})$. Then $\psi=$ $\mathscr{F}_{2}(\phi \circ \kappa)$, where $\kappa(x, y)=(x+y / 2, x-y / 2)$. By (1.27) (or [29, Eq. (1.35)]) and [8, Lemma 4.1] we have

$$
\begin{aligned}
\mathfrak{V}_{\Theta, d} K_{\mathfrak{a}}^{w}(z, w) & =\mathfrak{V}_{2 d} K_{\mathfrak{a}}^{w}(z, \bar{w})=\mathfrak{V}_{2 d} K_{\mathfrak{a}}^{w}((x, y)+i(\xi,-\eta)) \\
& =(2 \pi)^{d} e^{\frac{1}{2}\left(|z|^{2}+|w|^{2}\right)+i(\langle x, \xi\rangle-\langle y, \eta\rangle)} \mathcal{T}_{\phi} K_{\mathfrak{a}}^{w}(\sqrt{2}(x, y),-\sqrt{2}(\xi,-\eta)) \\
& =(2 \pi)^{\frac{d}{2}} e^{\frac{1}{2}\left(|z|^{2}+|w|^{2}\right)+i(\langle y, \xi\rangle-\langle x, \eta\rangle)} \mathcal{T}_{\psi} \mathfrak{a}\left(\frac{x+y}{\sqrt{2}},-\frac{\xi+\eta}{\sqrt{2}}, \sqrt{2}(\eta-\xi), \sqrt{2}(y-x)\right),
\end{aligned}
$$

$$
=(2 \pi)^{\frac{d}{2}} e^{\frac{1}{2}\left(|z|^{2}+|w|^{2}\right)+i \operatorname{Im}(z, w)} \mathcal{T}_{\psi} \mathfrak{a}\left(\frac{x+y}{\sqrt{2}},-\frac{\xi+\eta}{\sqrt{2}}, \sqrt{2}(\eta-\xi), \sqrt{2}(y-x)\right) .
$$

Together with the identity

$$
|z|^{2}+|w|^{2}+2 i \operatorname{Im}(z, w)=|z-w|^{2}+2(z, w)
$$

this gives

$$
\begin{equation*}
\mathfrak{V}_{\Theta, d} K_{\mathfrak{a}}^{w}(z, w)=(2 \pi)^{\frac{d}{2}} e^{\frac{1}{2}|z-w|^{2}+(z, w)} \mathcal{T}_{\psi} \mathfrak{a}\left(\frac{x+y}{\sqrt{2}},-\frac{\xi+\eta}{\sqrt{2}}, \sqrt{2}(\eta-\xi), \sqrt{2}(y-x)\right) . \tag{2.11}
\end{equation*}
$$

A combination of this identity with (1.42) and (1.45) gives (2.9).
In order to prove (2.10), we use Moyal's formula (1.12), (1.13) and the fact that $\|\psi\|_{L^{2}}=1$. This implies that the inverse of $\mathcal{T}_{\psi}$ is given by

$$
\begin{aligned}
\left(\mathcal{T}_{\psi}^{-1} F\right)(x, \xi) & =\left(\mathcal{T}_{\psi}^{*} F\right)(x, \xi) \\
& =(2 \pi)^{-d} \iiint \int_{\mathbf{R}^{4 d}} F\left(x_{1}, \xi_{1}, \eta_{1}, y_{1}\right) \psi\left(x-x_{1}, \xi-\xi_{1}\right) e^{i\left\langle\left\langle x-x_{1}, \eta_{1}\right\rangle+\left\langle y_{1}, \xi-\xi_{1}\right\rangle\right)} d x_{1} d \xi_{1} d \eta_{1} d y_{1}
\end{aligned}
$$

Writing

$$
G(z, w)=F(x, \xi, \eta, y), \quad z=x+i \xi, w=y+i \eta,
$$

we obtain

$$
\begin{equation*}
\mathcal{T}_{\psi}^{*} F(x, \xi)=2^{d}(2 \pi)^{-\frac{3 d}{2}} \iint_{\mathbf{C}^{2 d}} G\left(w_{1}, w_{2}\right) e^{-\left|z-w_{1}\right|^{2}} e^{i \operatorname{Im}\left\langle z-w_{1}, w_{2}\right\rangle} d \lambda\left(w_{1}\right) d \lambda\left(w_{2}\right) \tag{2.12}
\end{equation*}
$$

If $\mathfrak{a}=\mathcal{T}_{\psi}^{*} F$ and $a=\mathrm{S}_{\mathfrak{V} \mathfrak{a}}$, then (2.9) shows that

$$
a(z, w)=(2 \pi)^{\frac{d}{2}} e^{\frac{1}{2}|z-w|^{2}} G\left(\frac{\bar{z}+\bar{w}}{\sqrt{2}}, \sqrt{2}(w-z)\right)
$$

which gives

$$
G(z, w)=(2 \pi)^{-\frac{d}{2}} e^{-\frac{1}{4}|w|^{2}} a\left(\frac{2 \bar{z}-w}{2 \sqrt{2}}, \frac{2 \bar{z}+w}{2 \sqrt{2}}\right) .
$$

Inserting this into (2.12) we get

$$
\mathcal{T}_{\psi}^{*} F(x,-\xi)=\frac{1}{2^{d} \pi^{2 d}} \iint_{\mathbf{C}^{2 d}} a\left(\frac{2 \bar{w}_{1}-w_{2}}{2 \sqrt{2}}, \frac{2 \bar{w}_{1}+w_{2}}{2 \sqrt{2}}\right) e^{-\left|\bar{z}-w_{1}\right|^{2}} e^{-\frac{1}{4}\left|w_{2}\right|^{2}} e^{i \operatorname{Im}\left\langle\bar{z}-w_{1}, w_{2}\right\rangle} d \lambda\left(w_{1}\right) d \lambda\left(w_{2}\right),
$$

and by taking

$$
\frac{2 \bar{w}_{1}-w_{2}}{2 \sqrt{2}}-\frac{z}{\sqrt{2}} \quad \text { and } \quad \frac{2 \bar{w}_{1}+w_{2}}{2 \sqrt{2}}-\frac{z}{\sqrt{2}}
$$

as new variables of integration and using (1.22), we obtain

$$
\begin{aligned}
\mathcal{T}_{\psi}^{*} F(x,-\xi) & =\frac{2^{d}}{\pi^{2 d}} \iint_{\mathbf{C}^{2 d}} a\left(w_{1}+\frac{z}{\sqrt{2}}, w_{2}+\frac{z}{\sqrt{2}}\right) e^{-\left(\left|w_{1}\right|^{2}+\left|w_{2}\right|^{2}\right)} e^{2 i \operatorname{Im}\left(w_{1}, w_{2}\right)} d \lambda\left(w_{1}\right) d \lambda\left(w_{2}\right) \\
& =2^{d} \iint_{\mathbf{C}^{2} d} a\left(w_{1}+\frac{z}{\sqrt{2}}, w_{2}+\frac{z}{\sqrt{2}}\right) e^{2 i \operatorname{Im}\left(w_{1}, w_{2}\right)} d \mu\left(w_{1}\right) d \mu\left(w_{2}\right) \\
& =2^{d} \int_{\mathbf{C}^{d}}\left(\int_{\mathbf{C}^{d}} a\left(w_{1}+\frac{z}{\sqrt{2}}, w_{2}+\frac{z}{\sqrt{2}}\right) e^{\left(w_{1}, w_{2}\right)} e^{\left(-w_{2}, w_{1}\right)} d \mu\left(w_{1}\right)\right) d \mu\left(w_{2}\right) \\
& =2^{d} \int_{\mathbf{C}^{d}} a\left(-w_{2}+\frac{z}{\sqrt{2}}, w_{2}+\frac{z}{\sqrt{2}}\right) e^{-\left|w_{2}\right|^{2}} d \mu\left(w_{2}\right) \\
& =\left(\frac{2}{\pi}\right)^{d} \int_{\mathbf{C}^{d}} a\left(\frac{z}{\sqrt{2}}-w, \frac{z}{\sqrt{2}}+w\right) e^{-2|w|^{2}} d \lambda(w) .
\end{aligned}
$$

Proof of Theorem 2.2. Combining Propositions 2.1 and 2.3, writing $z+w=2 z+w-z$, we obtain that $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ if and only if for all $\alpha, \beta \in \mathbf{N}^{d}$ and $N \in \mathbf{N}$ we have

$$
\begin{aligned}
\left|\left(\partial_{x}+\partial_{y}\right)^{\alpha}\left(\partial_{\xi}+\partial_{\eta}\right)^{\beta}\left(e^{-\frac{1}{2}|z-w|^{2}} \mathrm{~S}_{\mathfrak{V}} \mathfrak{a}(z, w)\right)\right| & \lesssim \omega\left(\frac{\overline{z+w}}{\sqrt{2}}\right)\langle z+w\rangle^{-\rho|\alpha+\beta|}\langle z-w\rangle^{-N} \\
& \lesssim \omega(\sqrt{2} \bar{z})\langle z+w\rangle^{-\rho|\alpha+\beta|}\langle z-w\rangle^{-N+k}
\end{aligned}
$$

for some $k \in \mathbf{N}$ that can be absorbed into $N$.
Note that multi-index powers of the differential operators $\partial_{x}+\partial_{y}$ and $\partial_{\xi}+\partial_{\eta}$ acting on the factor $e^{-\frac{1}{2}|z-w|^{2}}=e^{-\frac{1}{2}\left(|x-y|^{2}+|\xi-\eta|^{2}\right)}$ are zero. Thus we obtain the equivalent condition

$$
\left|\left(\partial_{x}+\partial_{y}\right)^{\alpha}\left(\partial_{\xi}+\partial_{\eta}\right)^{\beta} \mathrm{S}_{\mathfrak{V}} \mathfrak{a}(z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z+w\rangle^{-\rho|\alpha+\beta|}\langle z-w\rangle^{-N} .
$$

Using the (conjugate) analyticity of $\mathrm{S}_{\mathfrak{V} \mathfrak{a}}(z, w)$ with respect to $z \in \mathbf{C}^{d}\left(w \in \mathbf{C}^{d}\right)$ we can formulate this as (2.4). We have now shown the equivalence between (1) and (2).

The equivalence between (2) and (3) follows from the binomial formulae

$$
\begin{aligned}
\left(\partial_{z}+t \bar{\partial}_{w}\right)^{\alpha} & =\sum_{\gamma \leqslant \alpha}\binom{\alpha}{\gamma} t^{|\gamma|} \partial_{z}^{\alpha-\gamma} \bar{\partial}_{w}^{\gamma}, \quad t \in\{-1,1\}, \\
\partial_{z}^{\alpha} & =2^{-|\alpha|} \sum_{\gamma \leqslant \alpha}\binom{\alpha}{\gamma}\left(\partial_{z}+\bar{\partial}_{w}\right)^{\alpha-\gamma}\left(\partial_{z}-\bar{\partial}_{w}\right)^{\gamma}
\end{aligned}
$$

and

$$
\bar{\partial}_{w}^{\beta}=2^{-|\beta|} \sum_{\gamma \leqslant \beta}\binom{\beta}{\gamma}(-1)^{|\gamma|}\left(\partial_{z}+\bar{\partial}_{w}\right)^{\beta-\gamma}\left(\partial_{z}-\bar{\partial}_{w}\right)^{\gamma} .
$$

It remains to consider the case $\rho=0$. We obtain from Propositions 2.1 and 2.3 that $\mathfrak{a} \in \operatorname{Sh}_{0}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ if and only if for all $N \in \mathbf{N}$ we have

$$
\left|\mathrm{S}_{\mathfrak{V}} \mathfrak{a}(z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z-w\rangle^{-N}, \quad z, \zeta \in \mathbf{C}^{d} .
$$

This shows the equivalence between (1) ${ }^{\prime}$ and (2)'.

Finally the equivalence of $(2)^{\prime}$ and $(3)^{\prime}$ is an immediate consequence of (1.42) and

$$
\left|e^{\frac{1}{2}\left(|z|^{2}+|w|^{2}\right)} e^{-(z, w)}\right|=e^{\frac{1}{2}\left(|z|^{2}-2 \operatorname{Re}(z, w)+|w|^{2}\right)}=e^{\frac{1}{2}|z-w|^{2}}
$$

Let $\widehat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$, be the set of all $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ such that

$$
\begin{equation*}
\left|\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a(z, w)\right| \leqslant C e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z+w\rangle^{-\rho|\alpha+\beta|}\langle z-w\rangle^{-N}, \quad N \geqslant 0 . \tag{2.13}
\end{equation*}
$$

The smallest constant $C \geqslant 0$ defines a semi-norm parameterized by $\alpha, \beta$ and $N$, and we equip $\hat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ with the Fréchet space topology defined by these semi-norms. The following result is an immediate consequence of Theorem 2.2 and its proof.

Proposition 2.4. Let $0 \leqslant \rho \leqslant 1$ and $\omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{2 d}\right)$. Then $\mathrm{S}_{\mathfrak{V}}$ is a homeomorphism from $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ to $\widehat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$.

### 2.2. Extensions and variations

There are several extensions and variations of Theorem 2.2. First we observe that by playing with $N$ in (2.4) and (2.5) and using Peetre's inequality, it follows that $\langle z+w\rangle$ in (2.4) and (2.5) can be replaced by $\Psi$, where

$$
\begin{align*}
& \Psi(z, w)=\langle z+w\rangle, \quad \Psi(z, w)=\langle z\rangle, \quad \Psi(z, w)=\langle w\rangle \\
& \Psi(z, w)=\max (\langle z\rangle,\langle w\rangle) \quad \text { or } \quad \Psi(z, w)=\min (\langle z\rangle,\langle w\rangle) \tag{2.14}
\end{align*}
$$

In particular (2.5) in Theorem 2.2 can be replaced by

$$
\begin{equation*}
\left|\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} \mathrm{S}_{\mathfrak{V} \mathfrak{a}} \mathfrak{a}(z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z}) \Psi(z, w)^{-\rho|\alpha+\beta|}\langle z-w\rangle^{-N} \tag{2.5}
\end{equation*}
$$

where $\Psi$ is given by (2.14).
Secondly, let

$$
\Omega_{k, M}=\left\{\left(\alpha_{1}, \ldots, \alpha_{k}\right) \in \mathbf{N}^{d} \times \cdots \times \mathbf{N}^{d} \simeq \mathbf{N}^{k d} ;\left|\alpha_{1}+\cdots+\alpha_{k}\right|=M\right\}
$$

where $k \geqslant 1$ and $M \geqslant 0$ are integers.
If $a \in \hat{A}\left(\mathbf{C}^{2 d}\right)$ and $\alpha \in \Omega_{4, M}$, then

$$
\begin{align*}
\partial_{x}^{\alpha_{1}} \partial_{\xi}^{\alpha_{3}} \partial_{y}^{\alpha_{2}} \partial_{\eta}^{\alpha_{4}} a & =i^{\left|\alpha_{3}\right|-\left|\alpha_{4}\right|} \partial_{z}^{\alpha_{1}+\alpha_{3}} \bar{\partial}_{w}^{\alpha_{2}+\alpha_{4}} a  \tag{2.15}\\
z & =x+i \xi, w=y+i \eta
\end{align*}
$$

because of the analyticity with respect to $z$ and conjugate analyticity with respect to $w$ for $a(z, w)$. In particular, (2.5)' implies

$$
\begin{equation*}
\left|\partial_{x}^{\alpha_{1}} \partial_{\xi}^{\alpha_{3}} \partial_{y}^{\alpha_{2}} \partial_{\eta}^{\alpha_{4}} \mathrm{~S}_{\mathfrak{V}} \mathfrak{a}(z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z}) \Psi(z, w)^{-\rho M}\langle z-w\rangle^{-N}, \quad \alpha \in \Omega_{4, M} \tag{2.5}
\end{equation*}
$$

On the other hand, if we let $\alpha_{3}=\alpha_{4}=0$ in (2.15), (2.5) ${ }^{\prime \prime}$ implies (2.5)'. Hence (2.5 $)^{\prime}$ and (2.5) " are equivalent.

Let $M \geqslant 0$ be an integer and let $T_{M}$ be the operator

$$
\begin{equation*}
T_{M}=\sum_{\alpha \in \Omega_{2, M}} c(\alpha) \partial_{z}^{\alpha_{1}} \bar{\partial}_{w}^{\alpha_{2}}, \quad c(\alpha) \in \mathbf{C}, \alpha=\left(\alpha_{1}, \alpha_{2}\right) \in \Omega_{2, M} . \tag{2.16}
\end{equation*}
$$

Then (2.5)' implies that

$$
\begin{equation*}
\left|T_{M}\left(\mathrm{~S}_{\mathfrak{V}} \mathfrak{a}\right)(z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z}) \Psi(z, w)^{-\rho M}\langle z-w\rangle^{-N}, \tag{2.5}
\end{equation*}
$$

holds for every $M \geqslant 0$ and every operator $T_{M}$ of the form (2.16). On the other hand, the operators $\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta}$ in (2.5) ${ }^{\prime}$ are special cases of the operators $T_{M}$ in (2.5)'". This shows that $\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta}$ in (2.5)' can be replaced by operators $T_{M}$ in (2.16).

In the same way it follows that (2.5)" is equivalent to (2.5) ${ }^{\prime \prime \prime}$, after $T_{M}$ in (2.16) is replaced by

$$
\begin{align*}
T_{M}=\sum_{\alpha \in \Omega_{4, M}} C(\alpha) \partial_{x}^{\alpha_{1}} \partial_{\xi}^{\alpha_{3}} \partial_{y}^{\alpha_{2}} \partial_{\eta}^{\alpha_{4}}, & z=x+i \xi, w=y+i \eta,  \tag{2.17}\\
C(\alpha) \in \mathbf{C}, & \alpha=\left(\alpha_{1}, \alpha_{2}, \alpha_{3}, \alpha_{4}\right) \in \Omega_{4, M} .
\end{align*}
$$

Finally we observe that we may replace the set of operators in (2.16) by the set of operators

$$
\begin{align*}
T_{M} & =\sum_{\alpha \in \Omega_{4, M}} C(\alpha) \partial_{z}^{\alpha_{1}} \bar{\partial}_{z}^{\alpha_{3}} \bar{\partial}_{w}^{\alpha_{2}} \partial_{w}^{\alpha_{4}},  \tag{2.18}\\
& z=x+i \xi, w=y+i \eta, \\
C(\alpha) \in \mathbf{C}, & \alpha=\left(\alpha_{1}, \alpha_{2}, \alpha_{3}, \alpha_{4}\right) \in \Omega_{4, M},
\end{align*}
$$

in the estimate (2.5) ${ }^{\prime \prime \prime}$. In fact, obviously the operators of form (2.18) contains the operators of form (2.16). Hence if $(2.5)^{\prime \prime \prime}$ holds true for operators of form (2.18), it also holds for operators of form (2.16). On the other hand, if $\alpha_{3} \neq 0$ or $\alpha_{4} \neq 0$ in (2.18), then

$$
\partial_{z}^{\alpha_{1}} \bar{\partial}_{z}^{\alpha_{3}} \bar{\partial}_{w}^{\alpha_{2}} \partial_{w}^{\alpha_{4}} a=0
$$

because of the analyticity in $z$ and conjugate analyticity in $w$ for $a(z, w)$. Consequently, it suffices to consider operators in (2.18) where all $C(\alpha)=0$ when $\alpha_{3} \neq 0$ or $\alpha_{4} \neq 0$, when investigating the condition (2.5)"'. This set of operators is exactly the set of operators in (2.16). This implies that the set of operators in (2.16) can be replaced by the set of operators in (2.18) when checking the condition (2.5)"'.

From these observations Theorem 2.2 gives the following conclusion.

Theorem 2.5. Suppose that $\rho \geqslant 0, \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{C}^{d}\right), a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ and $\Psi$ is given by (2.14). Then the following conditions are equivalent:
(1) $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$,
(2) (2.5)' holds true for every $N \geqslant 0, \alpha, \beta \in \mathbf{N}^{d}$ and $z, w \in \mathbf{C}^{d}$;
(3) for every $M \geqslant 0$, (2.5)" holds true for every $N \geqslant 0$, and $z, w \in \mathbf{C}^{d}$;
(4) for every $M \geqslant 0$, (2.5)'"' holds true for every $N \geqslant 0, T_{M}$ in (2.16) and $z, w \in \mathbf{C}^{d}$;
(5) for every $M \geqslant 0$, (2.5)"' holds true for every $N \geqslant 0, T_{M}$ in (2.17) and $z, w \in \mathbf{C}^{d}$;
(6) for every $M \geqslant 0$, (2.5)"'" holds true for every $N \geqslant 0, T_{M}$ in (2.18) and $z, w \in \mathbf{C}^{d}$.
2.3. Wick operators corresponding to Gevrey type pseudo-differential operators

Using (2.9) and (1.23) we obtain the following theorem expressed with estimates of the form

$$
\begin{equation*}
|a(z, w)| \lesssim \exp \left(\frac{1}{2}|z-w|^{2}+r_{1}|z+w|_{s, \sigma}-r_{2}|z-w|_{s, \sigma}\right) \tag{2.19}
\end{equation*}
$$

(cf. Definition 1.8). The verification is left for the reader.
Theorem 2.6. The following is true:
(1) if $s, \sigma \geqslant \frac{1}{2}$, then $\mathrm{S}_{\mathfrak{V}}$ is homeomorphic from $\Gamma_{s, \sigma ; 0}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$ to the set of all $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ such that for some $r_{2}>0$, (2.19) holds for every $r_{1}>0$;
(2) if $s, \sigma>\frac{1}{2}$, then $\mathrm{S}_{\mathfrak{V}}$ is homeomorphic from $\Gamma_{s, \sigma}^{\sigma, s ; 0}\left(\mathbf{R}^{2 d}\right)$ to the set of all $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ such that for some $r_{1}>0$, (2.19) holds for every $r_{2}>0$;
(3) if $s, \sigma>\frac{1}{2}$, then $\mathrm{S}_{\mathfrak{V}}$ is homeomorphic from $\Gamma_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$ to the set of all $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ such that (2.19) holds for some $r_{1}>0$ and some $r_{2}>0$.

Remark 2.7. The restrictions on $s$ and $\sigma$ in Theorem 2.6 are needed since we must choose $\psi$ in (1.51) as the Gauss function in Proposition 2.3. According to the proof of Theorem 2.2 this is necessary for the use of the formula (1.27) that relates $\mathcal{T}_{\phi} K_{\mathfrak{a}}^{w}$ and the Bargmann transform $\mathfrak{V}_{2 d} K_{\mathfrak{a}}^{w}$. For this $\psi$ we have $\psi \in \mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right)$ $\left(\psi \in \Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right)\right)$, if and only if $s, \sigma \geqslant \frac{1}{2}\left(s, \sigma>\frac{1}{2}\right)$.

Theorem 2.6 can be combined with continuity results in [1] to deduce continuity of Wick operators acting on the Bargmann images of $\Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right), \mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right),\left(\mathcal{S}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)$ and $\left(\Sigma_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right)$, respectively. The following result follows by a straight-forward combination of Theorems 3.8 and 3.15 in [1], (1.41) and Theorem 2.6.

Proposition 2.8. Let $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$. Then the following is true:
(1) if $s, \sigma \geqslant \frac{1}{2}$ and some $r_{2}>0$, (2.19) holds for every $r_{1}>0$, then $\mathrm{Op}_{\mathfrak{V}}(a)$ is continuous on $\mathcal{A}_{s}^{\sigma}\left(\mathbf{C}^{d}\right)$ and on $\left(\mathcal{A}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right)$;
(2) if $s, \sigma>\frac{1}{2}$ and for some $r_{1}>0$, (2.19) holds for every $r_{2}>0$, then $\mathrm{Op}_{\mathfrak{V}}(a)$ is continuous on $\mathcal{A}_{0, s}^{\sigma}\left(\mathbf{C}^{d}\right)$ and on $\left(\mathcal{A}_{0, s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right)$.

### 2.4. Composition of Wick operators

Let $a_{1}, a_{2} \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$. If composition is well defined then the complex twisted product $a_{1} \#_{\mathfrak{V}} a_{2}$ is defined by

$$
\mathrm{Op}_{\mathfrak{V}}\left(a_{1}\right) \circ \mathrm{Op}_{\mathfrak{V}}\left(a_{2}\right)=\mathrm{Op}_{\mathfrak{V}}\left(a_{1} \# \mathfrak{V} a_{2}\right) .
$$

By straight-forward computations it follows that the product $\#_{\mathfrak{V}}$ is given by

$$
\begin{equation*}
a_{1} \#_{\mathfrak{V}} a_{2}(z, w)=\pi^{-d} \int_{\mathbf{C}^{d}} a_{1}(z, u) a_{2}(u, w) e^{-(z-u, w-u)} d \lambda(u), \quad z, w \in \mathbf{C}^{d} \tag{2.20}
\end{equation*}
$$

provided the integral is well defined. Inserting derivatives, (2.20) takes the form
$\left(\partial_{z}^{\alpha_{1}} \bar{\partial}_{w}^{\beta_{1}} a_{1}\right) \not \mathbb{V}^{\prime}\left(\partial_{z}^{\alpha_{2}} \bar{\partial}_{w}^{\beta_{2}} a_{2}\right)(z, w)$

$$
\begin{equation*}
=\pi^{-d} \int_{\mathbf{C}^{d}}\left(\partial_{z}^{\alpha_{1}} \bar{\partial}_{u}^{\beta_{1}} a_{1}\right)(z, u)\left(\partial_{u}^{\alpha_{2}} \bar{\partial}_{w}^{\beta_{2}} a_{2}\right)(u, w) e^{-(z-u, w-u)} d \lambda(u), \quad z, w \in \mathbf{C}^{d} \tag{2.20}
\end{equation*}
$$

The following lemma is a product rule for the complex twisted product.
Lemma 2.9. Let $a_{1}, a_{2} \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ and suppose the integral in $(2.20)^{\prime}$ is well defined for all $z, w \in \mathbf{C}^{d}$ and all $\alpha_{1}, \alpha_{2}, \beta_{1}, \beta_{2} \in \mathbf{N}^{d}$ such that

$$
\left|\alpha_{1}+\alpha_{2}+\beta_{1}+\beta_{2}\right| \leqslant 1 .
$$

Suppose also that the integrand in (2.20) is zero at infinity. Then

$$
\begin{equation*}
\partial_{z_{j}}\left(a_{1} \#_{\mathfrak{V}} a_{2}\right)=\left(\partial_{z_{j}} a_{1}\right) \#_{\mathfrak{V}} a_{2}+a_{1} \#_{\mathfrak{V}}\left(\partial_{z_{j}} a_{2}\right), \quad j=1, \ldots, d \tag{2.21}
\end{equation*}
$$

and

$$
\begin{equation*}
\bar{\partial}_{w_{j}}\left(a_{1} \#_{\mathfrak{V}} a_{2}\right)=\left(\bar{\partial}_{w_{j}} a_{1}\right) \#_{\mathfrak{V}} a_{2}+a_{1} \#_{\mathfrak{V}}\left(\bar{\partial}_{w_{j}} a_{2}\right), \quad j=1, \ldots, d . \tag{2.22}
\end{equation*}
$$

Proof. If

$$
F_{a_{1}, a_{2}}(z, w, u)=a_{1}(z, u) a_{2}(u, w) e^{(z, u-w)+(u, w)}
$$

then

$$
\pi^{d}\left(a_{1} \not \mathfrak{N}_{2} a_{2}\right)(z, w)=\int_{\mathbf{C}^{d}} F_{a_{1}, a_{2}}(z, w, u) e^{-|u|^{2}} d \lambda(u) .
$$

This gives

$$
\pi^{d} \partial_{z_{j}}\left(a_{1} \#_{\mathfrak{V}} a_{2}\right)(z, w)=b_{1}(z, w)+b_{2}(z, w)-b_{3}(z, w),
$$

where

$$
\begin{aligned}
& b_{1}(z, w)=\int_{\mathbf{C}^{d}} F_{\partial_{z_{j}} a_{1}, a_{2}}(z, w, u) e^{-|u|^{2}} d \lambda(u), \\
& b_{2}(z, w)=\int_{\mathbf{C}^{d}} F_{a_{1}, a_{2}}(z, w, u) \bar{u}_{j} e^{-|u|^{2}} d \lambda(u)
\end{aligned}
$$

and

$$
\begin{aligned}
b_{3}(z, w) & =\bar{w}_{j} \int_{\mathbf{C}^{d}} F_{a_{1}, a_{2}}(z, w, u) e^{-|u|^{2}} d \lambda(u) \\
& =\bar{w}_{j} \pi^{d}\left(a_{1} \#_{\mathfrak{V}} a_{2}\right)(z, w)
\end{aligned}
$$

The conjugate analyticity of $u \mapsto a_{1}(z, u)$ and $u \mapsto e^{(z, u-w)}$ implies $\partial_{u_{j}} a_{1}(z, u)=\partial_{u_{j}} e^{(z, u-w)}=0$ which gives

$$
\begin{aligned}
\partial_{u_{j}} F_{a_{1}, a_{2}}(z, w, u) & =\left(a_{1}(z, u) \partial_{u_{j}} a_{2}(u, w)+\bar{w}_{j} a_{1}(z, u) a_{2}(u, w)\right) e^{(z, u-w)+(u, w)} \\
& =F_{a_{1}, \partial_{z_{j}} a_{2}}(z, w, u)+\bar{w}_{j} F_{a_{1}, a_{2}}(z, w, u) .
\end{aligned}
$$

Consider $b_{2}(z, w)$. Integration by parts gives

$$
\begin{aligned}
b_{2}(z, w) & =\int_{\mathbf{C}^{d}} F_{a_{1}, a_{2}}(z, w, u) \bar{u}_{j} e^{-|u|^{2}} d \lambda(u) \\
& =-\int_{\mathbf{C}^{d}} F_{a_{1}, a_{2}}(z, w, u) \partial_{u_{j}} e^{-|u|^{2}} d \lambda(u) \\
& =\int_{\mathbf{C}^{d}} \partial_{u_{j}} F_{a_{1}, a_{2}}(z, w, u) e^{-|u|^{2}} d \lambda(u) \\
& =\int_{\mathbf{C}^{d}} F_{a_{1}, \partial_{z_{j}} a_{2}}(z, w, u) e^{-|u|^{2}} d \lambda(u)+\bar{w}_{j} \int_{\mathbf{C}^{d}} F_{a_{1}, a_{2}}(z, w, u) e^{-|u|^{2}} d \lambda(u) \\
& =\int_{\mathbf{C}^{d}} F_{a_{1}, \partial_{z_{j}} a_{2}}(z, w, u) e^{-|u|^{2}} d \lambda(u)+b_{3}(z, w) .
\end{aligned}
$$

A combination of these identities now gives

$$
\begin{aligned}
\pi^{d} \partial_{z_{j}}\left(a_{1} \# \mathfrak{V} a_{2}\right)(z, w) & =\int_{\mathbf{C}^{d}}\left(F_{\partial_{z_{j}} a_{1}, a_{2}}(z, w, u)+F_{a_{1}, \partial_{z_{j}} a_{2}}(z, w, u)\right) e^{-|u|^{2}} d \lambda(u) \\
& =\pi^{d}\left(\partial_{z_{j}} a_{1}\right) \#_{\mathfrak{V}} a_{2}(z, w)+\pi^{d} a_{1} \#_{\mathfrak{V}}\left(\partial_{z_{j}} a_{2}\right)(z, w),
\end{aligned}
$$

and (2.21) follows.
The assertion (2.22) is proved by similar arguments.
The characterization in Theorem 2.2 (3) can be applied to prove the following composition result, which is a generalization of [26, Theorem 23.6] to include the case when $\rho=0$.

Proposition 2.10. Let $0 \leqslant \rho \leqslant 1$ and $\omega_{j} \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{2 d}\right)$ for $j=1,2$. If $\mathfrak{a}_{j} \in \operatorname{Sh}_{\rho}^{\left(\omega_{j}\right)}\left(\mathbf{R}^{2 d}\right)$ for $j=1,2$, then $\mathfrak{a}_{1} \# \mathfrak{a}_{2} \in \operatorname{Sh}_{\rho}^{\left(\omega_{1} \omega_{2}\right)}\left(\mathbf{R}^{2 d}\right)$.

Proof. If $\mathfrak{a}_{0}=\mathfrak{a}_{1} \# \mathfrak{a}_{2}$ and $a_{j}=\mathrm{S}_{\mathfrak{V}} \mathfrak{a}_{j}, j=0,1,2$, then $a_{0}=a_{1} \# \mathfrak{N} a_{2}$. From Lemma 2.9 and (2.20) we obtain for $\alpha, \beta \in \mathbf{N}^{d}$,

$$
\begin{aligned}
\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a_{0}(z, w) & =\sum_{\gamma \leqslant \alpha} \sum_{\kappa \leqslant \beta}\binom{\alpha}{\gamma}\binom{\beta}{\kappa}\left(\left(\partial_{z}^{\alpha-\gamma} \bar{\partial}_{w}^{\beta-\kappa} a_{1}\right) \# \mathfrak{V}\left(\partial_{z}^{\gamma} \bar{\partial}_{w}^{\kappa} a_{2}\right)\right)(z, w) \\
& =\pi^{-d} \sum_{\gamma \leqslant \alpha} \sum_{\kappa \leqslant \beta}\binom{\alpha}{\gamma}\binom{\beta}{\kappa} \int_{\mathbf{C}^{d}} \partial_{z}^{\alpha-\gamma} \bar{\partial}_{u}^{\beta-\kappa} a_{1}(z, u) \partial_{u}^{\gamma} \bar{\partial}_{w}^{\kappa} a_{2}(u, w) e^{(z, u-w)+(u, w)} d \mu(u) .
\end{aligned}
$$

Since $\omega_{2} \in \mathscr{P}\left(\mathbf{R}^{2 d}\right) \simeq \mathscr{P}\left(\mathbf{C}^{d}\right)$ is moderate, Theorem 2.2 gives for some $N_{0} \geqslant 0$ and any $N_{1}, N_{2} \geqslant 0$

$$
\left|\partial_{u}^{\alpha-\gamma} \bar{\partial}_{w}^{\beta-\kappa} a_{1}(z, u)\right| \lesssim \omega_{1}(\sqrt{2} \bar{z})\langle z+u\rangle^{-\rho|\alpha+\beta-\gamma-\kappa|}\langle z-u\rangle^{-N_{1}} e^{\frac{1}{2}|z-u|^{2}}
$$

and

$$
\left|\partial_{u}^{\gamma} \bar{\partial}_{w}^{\kappa} a_{2}(u, w)\right| \lesssim \omega_{2}(\sqrt{2} \bar{z})\langle z-u\rangle^{N_{0}}\langle u+w\rangle^{-\rho|\gamma+\kappa|}\langle u-w\rangle^{-N_{2}} e^{\frac{1}{2}|u-w|^{2}} .
$$

This gives

$$
\begin{equation*}
\left|\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a_{0}(z, w)\right| \lesssim \omega_{1}(\sqrt{2} \bar{z}) \omega_{2}(\sqrt{2} \bar{z}) e^{\frac{1}{2}|z-w|^{2}} \int_{\mathbf{C}^{d}} F(z, w, u) e^{\Phi(z, w, u)} d \lambda(u) \tag{2.23}
\end{equation*}
$$

where for any $N_{1} \geqslant 0$,

$$
F(z, w, u)=\langle z+u\rangle^{-\rho|\alpha+\beta-\gamma-\kappa|}\langle z-u\rangle^{N_{0}-N_{1}}\langle u+w\rangle^{-\rho|\gamma+\kappa|}\langle u-w\rangle^{-N_{2}}
$$

and

$$
\Phi(z, w, u)=-\frac{1}{2}|z-w|^{2}+\frac{1}{2}|z-u|^{2}+\frac{1}{2}|u-w|^{2}-|u|^{2}+\operatorname{Re}(z, u-w)+\operatorname{Re}(u, w)=0 .
$$

By Peetre's inequality and the facts that $\gamma \leqslant \alpha$ and $\kappa \leqslant \beta$ we get

$$
\begin{aligned}
\langle z+u\rangle^{\rho|\gamma+\kappa|}\langle u+w\rangle^{-\rho|\gamma+\kappa|} & \lesssim\langle z-w\rangle^{\rho|\gamma+\kappa|} \\
& \lesssim\langle z-u\rangle^{\rho|\gamma+\kappa|}\langle u-w\rangle^{\rho|\gamma+\kappa|} \leqslant\langle z-u\rangle^{\rho|\alpha+\beta|}\langle u-w\rangle^{\rho|\alpha+\beta|}
\end{aligned}
$$

and

$$
\langle z+u\rangle^{-\rho|\alpha+\beta|} \lesssim\langle z+w\rangle^{-\rho|\alpha+\beta|}\langle u-w\rangle^{\rho|\alpha+\beta|}
$$

wherefrom

$$
\begin{equation*}
F(z, w, u) \leqslant\langle z+w\rangle^{-\rho|\alpha+\beta|}\langle z-u\rangle^{\rho|\alpha+\beta|+N_{0}-N_{1}}\langle u-w\rangle^{2 \rho|\alpha+\beta|-N_{2}} . \tag{2.24}
\end{equation*}
$$

Hence a combination of (2.23) and (2.24) gives for any $N \geqslant 0$

$$
\begin{aligned}
& \left(\omega_{1}(\sqrt{2} \bar{z}) \omega_{2}(\sqrt{2} \bar{z})\right)^{-1}\langle z+w\rangle^{\rho|\alpha+\beta|}\left|\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a_{0}(z, w)\right| \\
& \lesssim e^{\frac{1}{2}|z-w|^{2}} \int_{\mathbf{C}^{d}}\langle z-u\rangle^{\rho|\alpha+\beta|+N_{0}-N_{1}}\langle u-w\rangle^{2 \rho|\alpha+\beta|-N_{2}} d \lambda(u) \\
& \lesssim\langle z-w\rangle^{-N} e^{\frac{1}{2}|z-w|^{2}} \int_{\mathbf{C}^{d}}\langle z-u\rangle^{\rho|\alpha+\beta|+N_{0}+N-N_{1}}\langle u-w\rangle^{2 \rho|\alpha+\beta|+N-N_{2}} d \lambda(u) .
\end{aligned}
$$

By letting

$$
N_{1} \geqslant \rho|\alpha+\beta|+N_{0}+N \quad \text { and } \quad N_{2}>2 \rho|\alpha+\beta|+N+2 d
$$

we obtain

$$
\left|\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a_{0}(z, w)\right| \lesssim \omega_{1}(\sqrt{2} \bar{z}) \omega_{2}(\sqrt{2} \bar{z})\langle z+w\rangle^{-\rho|\alpha+\beta|}\langle z-w\rangle^{-N} e^{\frac{1}{2}|z-w|^{2}} .
$$

According to Theorem 2.2 (3) this estimate implies that $\mathfrak{a}_{0} \in \operatorname{Sh}_{\rho}^{\left(\omega_{1} \omega_{2}\right)}\left(\mathbf{R}^{2 d}\right)$.
Remark 2.11. Eq. (2.20) combined with Theorem 2.6 can be used to show composition results for pseudodifferential operators with symbols in $\Gamma_{s, \sigma ; 0}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$. In fact we may use an argument similar to the proof of Proposition 2.10, but simpler since derivatives can be avoided. We obtain

$$
\mathfrak{a}_{1} \# \mathfrak{a}_{2} \in \Gamma_{s, \sigma ; 0}^{\sigma, s}\left(\mathbf{R}^{2 d}\right) \quad \text { when } \quad \mathfrak{a}_{1}, \mathfrak{a}_{2} \in \Gamma_{s, \sigma ; 0}^{\sigma, s}\left(\mathbf{R}^{2 d}\right), \quad s, \sigma \geqslant \frac{1}{2},
$$

and similarly with $\Gamma_{s, \sigma}^{\sigma, s ; 0}$ in place of $\Gamma_{s, \sigma ; 0}^{\sigma, s}$, provided $\sigma>\frac{1}{2}$. Thereby we regain parts of [1, Theorem 3.18] for certain restrictions on $s$ and $\sigma$.

## 3. Relations and estimates for Wick and anti-Wick operators

In this section we first show how to approximate a Wick operator by means of a sum of anti-Wick operators. Then we prove continuity results for anti-Wick operators with symbols having exponential type bounds. Finally we deduce estimates for the Wick symbol of these anti-Wick operators.

### 3.1. Expansion of Shubin type Wick operators with respect to anti-Wick operators

The first result can be stated for semi-conjugate analytic symbols on $\mathbf{C}^{2 d}$.
Proposition 3.1. Suppose $s \geqslant \frac{1}{2}$, $a \in \hat{\mathcal{A}}_{s}^{\prime}\left(\mathbf{C}^{2 d}\right)$, let $M \geqslant 0$ be an integer, and let

$$
a_{\alpha}(w)=\partial_{z}^{\alpha} \bar{\partial}_{w}^{\alpha} a(w, w), \quad \alpha \in \mathbf{N}^{d}
$$

and

$$
b_{\alpha}(z, w)=|\alpha| \int_{0}^{1}(1-t)^{|\alpha|-1} \partial_{z}^{\alpha} \bar{\partial}_{w}^{\alpha} a(w+t(z-w), w) d t, \quad \alpha \in \mathbf{N}^{d} \backslash 0
$$

Then

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}(a)=\sum_{|\alpha| \leqslant M} \frac{(-1)^{|\alpha|} \mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{\alpha}\right)}{\alpha!}+\sum_{|\alpha|=M+1} \frac{(-1)^{|\alpha|} \mathrm{Op}_{\mathfrak{V}}\left(b_{\alpha}\right)}{\alpha!} \tag{3.1}
\end{equation*}
$$

Proof. Taylor expansion gives

$$
a(z, w)=\sum_{|\alpha| \leqslant M} \frac{(-1)^{|\alpha|} c_{\alpha}(z, w)}{\alpha!}+\sum_{|\alpha|=M+1} \frac{(-1)^{|\alpha|} c_{0, \alpha}(z, w)}{\alpha!}
$$

where

$$
c_{\alpha}(z, w)=(-1)^{|\alpha|}(z-w)^{\alpha} \partial_{z}^{\alpha} a(w, w)
$$

and

$$
c_{0, \alpha}(z, w)=(-1)^{|\alpha|}|\alpha|(z-w)^{\alpha} \int_{0}^{1}(1-t)^{|\alpha|-1} \partial_{z}^{\alpha} a(w+t(z-w), w) d t
$$

Hence

$$
\mathrm{Op}_{\mathfrak{V}}(a)=\sum_{|\alpha| \leqslant M} \frac{(-1)^{|\alpha|} \mathrm{Op}_{\mathfrak{V}}\left(c_{\alpha}\right)}{\alpha!}+\sum_{|\alpha|=M+1} \frac{(-1)^{|\alpha|} \mathrm{Op}_{\mathfrak{V}}\left(c_{0, \alpha}\right)}{\alpha!}
$$

and the result follows if we prove

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}\left(c_{\alpha}\right)=\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{\alpha}\right) \quad \text { and } \quad \mathrm{Op}_{\mathfrak{V}}\left(c_{0, \alpha}\right)=\mathrm{Op}_{\mathfrak{V}}\left(b_{\alpha}\right) \tag{3.2}
\end{equation*}
$$

It follows from (1.38) that

$$
\mathrm{Op}_{\mathfrak{V}}\left(b_{\alpha}\right)=\mathrm{Op}_{\mathfrak{V}}\left(c_{1, \alpha}\right) \quad \text { and } \quad \mathrm{Op}_{\mathfrak{V}}\left(c_{0, \alpha}\right)=\mathrm{Op}_{\mathfrak{V}}\left(c_{2, \alpha}\right)
$$

where

$$
\begin{equation*}
c_{j, \alpha}(z, w)=(-1)^{|\alpha|} \pi^{-d}|\alpha| \int_{0}^{1}(1-t)^{|\alpha|-1} h_{j, \alpha}(a ; t, z, w) d t, \tag{3.3}
\end{equation*}
$$

$j=1,2$, with

$$
\begin{equation*}
h_{1, \alpha}(a ; t, z, w)=(-1)^{|\alpha|} \int_{\mathbf{C}^{d}} \partial_{z}^{\alpha} \bar{\partial}_{w}^{\alpha} a\left(w_{1}+t\left(z-w_{1}\right), w_{1}\right) e^{-\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right) \tag{3.4}
\end{equation*}
$$

and

$$
h_{2, \alpha}(a ; t, z, w)=\int_{\mathbf{C}^{d}}\left(z-w_{1}\right)^{\alpha} \partial_{z}^{\alpha} a\left(w_{1}+t\left(z-w_{1}\right), w_{1}\right) e^{-\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right) .
$$

Since

$$
\left(z-w_{1}\right)^{\alpha} e^{-\left(z-w_{1}, w-w_{1}\right)}=\bar{\partial}_{w_{1}}^{\alpha} e^{-\left(z-w_{1}, w-w_{1}\right)}
$$

integration by parts yields

$$
\begin{aligned}
h_{2, \alpha}(a ; t, z, w) & =\int_{\mathbf{C}^{d}} \partial_{z}^{\alpha} a\left(w_{1}+t\left(z-w_{1}\right), w_{1}\right) \bar{\partial}_{w_{1}}^{\alpha} e^{-\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right) \\
& =(-1)^{|\alpha|} \int_{\mathbf{C}^{d}} \partial_{z}^{\alpha} \bar{\partial}_{w}^{\alpha} a\left(w_{1}+t\left(z-w_{1}\right), w_{1}\right) e^{-\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right)=h_{1, \alpha}(a ; t, z, w),
\end{aligned}
$$

and the second equality in (3.2) follows. The first equality in (3.2) follows by similar arguments. The details are left for the reader.

Remark 3.2. Proposition 3.1 and its proof show that

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}(a)=\sum_{|\alpha| \leqslant M} \frac{(-1)^{|\alpha|} \mathrm{Op}_{\mathfrak{Y}}^{\mathrm{aw}}\left(a_{\alpha}\right)}{\alpha!}+\sum_{|\alpha|=M+1} \frac{(-1)^{|\alpha|} \mathrm{Op}_{\mathfrak{V}}\left(c_{1, \alpha}\right)}{\alpha!} \tag{3.1}
\end{equation*}
$$

where $c_{1, \alpha}$ is defined by (3.3) and (3.4).
In the following result we estimate $a_{\alpha}$ in Proposition 3.1 and $c_{1, \alpha}$ in (3.3) when $a=\mathrm{S}_{\mathfrak{V} \mathfrak{a}}$ satisfies (2.5) for every $N \geqslant 0$ and $\alpha, \beta \in \mathbf{N}^{d}$. By Theorem 2.2 this means that $\operatorname{Op}_{\mathfrak{V}}(a)$ is the Bargmann transform of a Shubin type operator.

Proposition 3.3. Let $0 \leqslant \rho \leqslant 1, \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{2 d}\right)$, $a \in \mathcal{\mathcal { A }}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$, and let $a_{\alpha}$ and $b_{\alpha}$ be as in Proposition 3.1 for $\alpha \in \mathbf{N}^{d}$. Then $\mathrm{Op}_{\mathfrak{V}}\left(b_{\alpha}\right)=\mathrm{Op}_{\mathfrak{V}}\left(c_{1, \alpha}\right)$ for a unique $c_{1, \alpha} \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$,

$$
\begin{equation*}
\left|\partial_{w}^{\beta} \bar{\partial}_{w}^{\gamma} a_{\alpha}(w)\right| \lesssim \omega(\sqrt{2} \bar{w})\langle w\rangle^{-\rho(2|\alpha|+|\beta+\gamma|)}, \quad \alpha, \beta, \gamma \in \mathbf{N}^{d}, \tag{3.5}
\end{equation*}
$$

and for every $N \geqslant 0$ it holds,

$$
\begin{equation*}
\left|\partial_{z}^{\beta} \bar{\partial}_{w}^{\gamma} c_{1, \alpha}(z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z+w\rangle^{-\rho(2|\alpha|+|\beta+\gamma|)}\langle z-w\rangle^{-N}, \quad \alpha, \beta, \gamma \in \mathbf{N}^{d} . \tag{3.6}
\end{equation*}
$$

Remark 3.4. The Wick symbol $c_{1, \alpha}$ in Proposition 3.3 is uniquely defined and given by (3.3) in view of Proposition 1.5, when $h_{1, \alpha}$ is defined by (3.4). The conditions in Proposition 3.3 imply that $c_{1, \alpha} \in \hat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{\left(\omega_{\alpha}\right)}\left(\mathbf{C}^{2 d}\right)$ where $\omega_{\alpha}=\langle\cdot\rangle^{-2 \rho|\alpha|} \cdot \omega$.

Proof of Proposition 3.3. The estimate (3.5) is an immediate consequence of

$$
\partial_{w}^{\beta} \bar{\partial}_{w}^{\gamma} a_{\alpha}(w)=\partial_{w}^{\alpha+\beta} \bar{\partial}_{w}^{\alpha+\gamma} a(w, w)
$$

and (2.5).
In order to prove (3.6) we first note that the uniqueness assertion for $c_{1, \alpha}$ is a consequence of Remark 3.4. Let $h_{1, \alpha}(a ; z, w)$ be the same as in the proof of Proposition 3.1. Integration by parts gives

$$
\partial_{z}^{\beta} \bar{\partial}_{w}^{\gamma} h_{1, \alpha}(a ; t, z, w)=h_{1, \alpha}\left(\partial_{z}^{\beta} \bar{\partial}_{w}^{\gamma} a ; t, z, w\right)
$$

which reduces the problem to prove that (3.6) holds for $\beta=\gamma=0$.
The assumption $a \in \widehat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ combined with $\omega$ and $\langle\cdot\rangle^{-\rho|\alpha|}$ being moderate imply

$$
\left|\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a(z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{w})\langle w\rangle^{-\rho|\alpha+\beta|}\langle z-w\rangle^{-N}
$$

for every $N \geqslant 0$. This gives

$$
e^{\operatorname{Re}(z, w)}\left|h_{1, \alpha}(a ; t, z, w)\right| \lesssim \int_{\mathbf{C}^{d}} \omega\left(\sqrt{2} \overline{w_{1}}\right) e^{\frac{t^{2}}{2}\left|z-w_{1}\right|^{2}}\left\langle w_{1}\right\rangle^{-2 \rho|\alpha|}\left\langle t\left(z-w_{1}\right)\right\rangle^{-N} e^{\operatorname{Re}\left(z+w-w_{1}, w_{1}\right)} d \lambda\left(w_{1}\right)
$$

that is

$$
\begin{align*}
& e^{-\frac{1}{4}|z-w|^{2}}\left|h_{1, \alpha}(a ; t, z, w)\right| \lesssim \int_{\mathbf{C}^{d}} \omega\left(\sqrt{2} \overline{w_{1}}\right) e^{\frac{t^{2}}{2}\left|z-w_{1}\right|^{2}}\left\langle w_{1}\right\rangle^{-2 \rho|\alpha|}\left\langle t\left(z-w_{1}\right)\right\rangle^{-N} e^{-\left|w_{1}-z_{2}\right|^{2}} d \lambda\left(w_{1}\right) \\
&=\int_{\mathbf{C}^{d}} \omega\left(\sqrt{2}\left(\overline{z_{2}+w_{1}}\right)\right) e^{\frac{t^{2}}{2}\left|z_{1}-w_{1}\right|^{2}}\left\langle z_{2}+w_{1}\right\rangle^{-2 \rho|\alpha|}\left\langle t\left(z_{1}-w_{1}\right)\right\rangle^{-N} e^{-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \tag{3.7}
\end{align*}
$$

for every $N \geqslant 0$, where $z_{1}=\frac{1}{2}(z-w)$ and $z_{2}=\frac{1}{2}(z+w)$.
If $t \in\left[0, \frac{1}{2}\right]$, then the last estimate together with the moderateness of $\omega$ gives

$$
\begin{aligned}
e^{-\left|z_{1}\right|^{2}}\left|h_{1, \alpha}(a ; t, z, w)\right| & \lesssim \omega\left(\sqrt{2} \overline{z_{2}}\right)\left\langle z_{2}\right\rangle^{-2 \rho|\alpha|} \int_{\mathbf{C}^{d}} e^{\frac{1}{8}\left|w_{1}\right|^{2}} e^{\frac{1}{8}\left|z_{1}-w_{1}\right|^{2}} e^{-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& \lesssim \omega\left(\sqrt{2} \overline{z_{2}}\right)\left\langle z_{2}\right\rangle^{-2 \rho|\alpha|} e^{\frac{1}{4}\left|z_{1}\right|^{2}} \int_{\mathbf{C}^{d}} e^{\frac{1}{4}\left|w_{1}\right|^{2}} e^{-\frac{7}{8}\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& \lesssim \omega\left(\sqrt{2} \overline{z_{2}}\right)\left\langle z_{2}\right\rangle^{-2 \rho|\alpha|} e^{\frac{1}{2}\left|z_{1}\right|^{2}}\left\langle z_{1}\right\rangle^{-N}
\end{aligned}
$$

for every $N \geqslant 0$. The moderateness of $\omega$ again gives

$$
\begin{equation*}
\left|h_{1, \alpha}(a ; t, z, w)\right| \lesssim e^{\frac{1}{2}|z-w|^{2}} \omega(\sqrt{2} \bar{z})\langle z+w\rangle^{-2 \rho|\alpha|}\langle z-w\rangle^{-N} \tag{3.8}
\end{equation*}
$$

or every $N \geqslant 0$, when $t \in\left[0, \frac{1}{2}\right]$.
Suppose instead $t \in\left[\frac{1}{2}, 1\right]$. Then $\left\langle t\left(z_{1}-w_{1}\right)\right\rangle^{-N} \asymp\left\langle z_{1}-w_{1}\right\rangle^{-N}$. Moderateness again gives

$$
\omega\left(\sqrt{2}\left(\overline{z_{2}+w_{1}}\right)\right)\left\langle z_{2}+w_{1}\right\rangle^{-2 \rho|\alpha|}\left\langle z_{1}-w_{1}\right\rangle^{-N_{0}} \lesssim \omega(\sqrt{2} \bar{z})\langle z\rangle^{-2 \rho|\alpha|}
$$

for some $N_{0}$. Hence (3.7) gives

$$
\begin{aligned}
e^{-\left|z_{1}\right|^{2}} \omega(\sqrt{2} \bar{z})^{-1}\langle z\rangle^{2 \rho|\alpha|}\left|h_{1, \alpha}(a ; t, z, w)\right| & \lesssim \int_{\mathbf{C}^{d}} e^{\frac{1}{2}\left|z_{1}-w_{1}\right|^{2}}\left\langle z_{1}-w_{1}\right\rangle^{-N} e^{-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& =e^{\left|z_{1}\right|^{2}} \int_{\mathbf{C}^{d}}\left\langle z_{1}-w_{1}\right\rangle^{-N} e^{-\frac{1}{2}\left|w_{1}+z_{1}\right|^{2}} d \lambda\left(w_{1}\right)=e^{\left|z_{1}\right|^{2}}\left\langle z_{1}\right\rangle^{-N}
\end{aligned}
$$

for every $N \geqslant 0$. This gives (3.8) also for $t \in\left[\frac{1}{2}, 1\right]$.
The result now follows by using (3.8) when estimating $\left|c_{1, \alpha}(z, w)\right|$ in (3.3) and evaluating the arising integral.

The next result, analogous to Proposition 3.3, will be useful in Section 5 when we discuss hypoellipticity for Shubin operators in the Wick setting.

Proposition 3.5. Let $\rho \geqslant 0, \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{C}^{d}\right), \omega_{t}=\omega \cdot\langle\cdot\rangle^{-2 \rho t}$ when $t \geqslant 0, a \in \hat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right), \mathfrak{a}=\mathrm{S}_{\mathfrak{V}}^{-1} a$ and $M \geqslant 0$ be an integer. Then

$$
\begin{equation*}
\mathfrak{a}(x,-\xi)=\sum_{|\alpha| \leqslant M} \frac{(-1)^{|\alpha|}\left(\partial_{z}^{\alpha} \bar{\partial}_{w}^{\alpha} a\right)\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right)}{2^{|\alpha|} \alpha!}+c_{M}(z), \quad z=x+i \xi, \tag{3.9}
\end{equation*}
$$

where

$$
\begin{equation*}
\partial_{z}^{\alpha} \bar{\partial}_{w}^{\alpha} a \in \mathcal{A}_{\mathrm{Sh}, \rho}^{\left(\omega_{|\alpha|}\right)}\left(\mathbf{C}^{2 d}\right) \quad \text { and } \quad(x, \xi) \mapsto c_{M}(x-i \xi) \in \operatorname{Sh}_{\rho}^{\left(\omega_{M+1}\right)}\left(\mathbf{R}^{2 d}\right) . \tag{3.10}
\end{equation*}
$$

Proof. The first claim in (3.10) is an immediate consequence of the definition (2.13) and Peetre's inequality.
By Taylor expanding the right-hand side of (2.10) we obtain

$$
\begin{equation*}
\mathfrak{a}(x,-\xi)=\sum_{|\alpha+\beta| \leqslant 2 M+1} \frac{(2 / \pi)^{d} I_{\alpha, \beta} \cdot\left(\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a\right)\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right)}{\alpha!\beta!}+c_{M}(z), \tag{3.11}
\end{equation*}
$$

where

$$
I_{\alpha, \beta}=\int_{\mathbf{C}^{d}}(-w)^{\alpha} \bar{w}^{\beta} e^{-2|w|^{2}} d \lambda(w),
$$

and

$$
\begin{equation*}
c_{M}(z)=2(M+1) \sum_{|\alpha+\beta|=2 M+2} \frac{(-1)^{|\beta|}}{\alpha!\beta!} \int_{0}^{1}(1-\theta)^{2 M+1} H_{\alpha, \beta}(z, \theta) d \theta \tag{3.12}
\end{equation*}
$$

with

$$
\begin{equation*}
H_{\alpha, \beta}(z, \theta)=\left(\frac{2}{\pi}\right)^{d} \int_{\mathbf{C}^{d}}\left(\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a\right)\left(\frac{z}{\sqrt{2}}-\theta w, \frac{z}{\sqrt{2}}+\theta w\right) w^{\alpha} \bar{w}^{\beta} e^{-2|w|^{2}} d \lambda(w) . \tag{3.13}
\end{equation*}
$$

The orthonormality of $\left\{e_{\alpha}\right\}_{\alpha \in \mathbf{N}^{d}} \subseteq A^{2}\left(\mathbf{C}^{d}\right)$ (cf. (1.21)) yields $I_{\alpha, \beta}=0$ if $\alpha \neq \beta$ and

$$
\begin{aligned}
I_{\alpha, \alpha} & =\int_{\mathbf{C}^{d}}(-w)^{\alpha} \bar{w}^{\alpha} e^{-2|w|^{2}} d \lambda(w) \\
& =(-1)^{|\alpha|} 2^{-d-|\alpha|} \alpha!\pi^{d} \int_{\mathbf{C}^{d}}\left|e_{\alpha}(w)\right|^{2} d \mu(w) \\
& =(-1)^{|\alpha|} 2^{-d-|\alpha|} \alpha!\pi^{d} .
\end{aligned}
$$

By inserting these values of $I_{\alpha, \beta}$ into (3.11), we obtain (3.9).
It remains to study the remainder $c_{M}$. We need to prove that $\mathfrak{c}(x, \xi)=c_{M}(x-i \xi)$ belongs to $\operatorname{Sh}_{\rho}^{\left(\omega_{M+1}\right)}\left(\mathbf{R}^{2 d}\right)$. If

$$
h_{\alpha, \beta}(z, w, \theta)=\left(\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a\right)\left(\frac{z}{\sqrt{2}}-\theta w, \frac{z}{\sqrt{2}}+\theta w\right) w^{\alpha} \bar{w}^{\beta} e^{-2|w|^{2}}
$$

then

$$
H_{\alpha, \beta}(z, \theta)=\left(\frac{2}{\pi}\right)^{d} \int_{\mathbf{C}^{d}} h_{\alpha, \beta}(z, w, \theta) d \lambda(w) .
$$

First we notice that

$$
\begin{aligned}
\partial_{z}^{\alpha} \bar{\partial}_{z}^{\beta} c_{M}(z) & =2(M+1) \sum_{|\gamma+\delta|=2 M+2} \frac{(-1)^{|\delta|}}{\gamma!\delta!} \int_{0}^{1}(1-\theta)^{2 M+1} \partial_{z}^{\alpha} \bar{\partial}_{z}^{\beta} H_{\gamma, \delta}(z, \theta) d \theta \\
\partial_{z}^{\alpha} \bar{\partial}_{z}^{\beta} H_{\gamma, \delta}(z, \theta) & =\left(\frac{2}{\pi}\right)^{d} \int_{\mathbf{C}^{d}} \partial_{z}^{\alpha} \alpha_{z}^{\beta} h_{\gamma, \delta}(z, w, \theta) d \lambda(w)
\end{aligned}
$$

and

$$
\partial_{z}^{\alpha} \bar{\partial}_{z}^{\beta} h_{\gamma, \delta}(z, w, \theta)=2^{-\frac{|\alpha+\beta|}{2}}\left(\partial_{z}^{\alpha+\gamma} \bar{\partial}_{z}^{\beta+\delta} a\right)\left(\frac{z}{\sqrt{2}}-\theta w, \frac{z}{\sqrt{2}}+\theta w\right) w^{\gamma} \bar{w}^{\delta} e^{-2|w|^{2}} .
$$

From the definition (2.13) this implies that for every $N \geqslant 0$ and some $N_{0} \geqslant 0$ we have

$$
\begin{aligned}
\left|\partial_{z}^{\alpha} \bar{\partial}_{z}^{\beta} h_{\gamma, \delta}(z, w, \theta)\right| & \lesssim e^{-2\left(1-\theta^{2}\right)|w|^{2}} \omega(\bar{z}-\sqrt{2} \theta \bar{w})\langle z\rangle^{-\rho(|\alpha+\beta|+2 M+2)}\langle\theta w\rangle^{-N-N_{0}}|w|^{2 M+2} \\
& \lesssim e^{-2(1-\theta)|w|^{2}} \omega(\bar{z})\langle z\rangle^{-\rho(|\alpha+\beta|+2 M+2)}\langle\theta w\rangle^{-N}|w|^{2 M+2} .
\end{aligned}
$$

This gives

$$
\left|\partial_{z}^{\alpha} \bar{\partial}_{z}^{\beta} H_{\gamma, \delta}(z, \theta)\right| \lesssim \omega(\bar{z})\langle z\rangle^{-\rho(|\alpha+\beta|+2 M+2)} \cdot J(\theta),
$$

where

$$
J(\theta)=\int_{\mathbf{C}^{d}} e^{-2(1-\theta)|w|^{2}}\langle\theta w\rangle^{-N}|w|^{2 M+2} d \lambda(w) .
$$

For $\theta \in\left[0, \frac{1}{2}\right]$ we get

$$
J(\theta) \leqslant \int_{\mathbf{C}^{d}} e^{-|w|^{2}}|w|^{2 M+2} d \lambda(w),
$$

which is finite and independent of $\theta$. If instead $\theta \in\left[\frac{1}{2}, 1\right]$, and choosing $N>2 d+2 M+2$, then

$$
J(\theta) \leqslant \int_{\mathbf{C}^{d}}\langle\theta w\rangle^{-N}|w|^{2 M+2} d \lambda(w),
$$

which is again finite and uniformly bounded with respect to $\theta$.

A combination of these estimates gives

$$
\left|\partial_{z}^{\alpha} \bar{\partial}_{z}^{\beta} H_{\gamma, \delta}(z, \theta)\right| \lesssim \omega(\bar{z})\langle z\rangle^{-\rho(|\alpha+\beta|+2 M+2)},
$$

which in turn implies

$$
\left|\partial_{z}^{\alpha} \bar{\partial}_{z}^{\beta} c_{M}(z)\right| \lesssim \omega(\bar{z})\langle z\rangle^{-\rho(|\alpha+\beta|+2 M+2)} .
$$

This means that $\mathfrak{c} \in \operatorname{Sh}_{\rho}^{\left(\omega_{M+1}\right)}\left(\mathbf{R}^{2 d}\right)$.

### 3.2. Continuity of anti-Wick operators with exponentially bounded symbols

Next we consider anti-Wick symbols that satisfy exponential bounds of the form

$$
\begin{equation*}
\left|a_{0}(w)\right| \lesssim e^{-r_{0}|w|^{\frac{1}{s}}}, \tag{3.14}
\end{equation*}
$$

or

$$
\begin{equation*}
\left|a_{0}(w)\right| \lesssim e^{r_{0}|w|^{\frac{1}{s}}} . \tag{3.15}
\end{equation*}
$$

In order to formulate our results we introduce new spaces of entire functions. Let $s>\frac{1}{2}, t_{0}, r>0$, and let $\mathcal{A}_{s, t_{0}, r}\left(\mathbf{C}^{d}\right)$ be the Banach space of all $F \in A\left(\mathbf{C}^{d}\right)$ such that

$$
\|F\|_{\mathcal{A}_{s, t_{0}, r}} \equiv\left\|F \cdot e^{-t_{0}|\cdot|^{2}+r|\cdot|^{\frac{1}{s}}}\right\|_{L^{\infty}}<\infty .
$$

Set

$$
\mathcal{A}_{0,\left(s, t_{0}\right)}\left(\mathbf{C}^{d}\right)=\bigcap_{r>0} \mathcal{A}_{s, t_{0}, r}\left(\mathbf{C}^{d}\right) \quad \text { and } \quad \mathcal{A}_{\left(s, t_{0}\right)}^{\prime}\left(\mathbf{C}^{d}\right)=\bigcap_{r>0} \mathcal{A}_{s, t_{0},-r}\left(\mathbf{C}^{d}\right)
$$

equipped with the projective limit topology. Likewise we set

$$
\mathcal{A}_{\left(s, t_{0}\right)}\left(\mathbf{C}^{d}\right)=\bigcup_{r>0} \mathcal{A}_{s, t_{0}, r}\left(\mathbf{C}^{d}\right) \quad \text { and } \quad \mathcal{A}_{0,\left(s, t_{0}\right)}^{\prime}\left(\mathbf{C}^{d}\right)=\bigcup_{r>0} \mathcal{A}_{s, t_{0},-r}\left(\mathbf{C}^{d}\right)
$$

equipped with the inductive limit topology.
Referring to Section 1.3 it is clear that the spaces $\mathcal{A}_{0,\left(s, t_{0}\right)}\left(\mathbf{C}^{d}\right), \mathcal{A}_{\left(s, t_{0}\right)}\left(\mathbf{C}^{d}\right), \mathcal{A}_{\left(s, t_{0}\right)}^{\prime}\left(\mathbf{C}^{d}\right)$ and $\mathcal{A}_{0,\left(s, t_{0}\right)}^{\prime}\left(\mathbf{C}^{d}\right)$ are generalizations of

$$
\begin{gathered}
\mathcal{A}_{0,\left(s, \frac{1}{2}\right)}\left(\mathbf{C}^{d}\right)=\mathfrak{V}_{d}\left(\Sigma_{s}\left(\mathbf{R}^{d}\right)\right)=\mathcal{A}_{0, s}\left(\mathbf{C}^{d}\right) \\
\mathcal{A}_{\left(s, \frac{1}{2}\right)}\left(\mathbf{C}^{d}\right)=\mathfrak{V}_{d}\left(\mathcal{S}_{s}\left(\mathbf{R}^{d}\right)\right)=\mathcal{A}_{s}\left(\mathbf{C}^{d}\right) \\
\mathcal{A}_{\left(s, \frac{1}{2}\right)}^{\prime}\left(\mathbf{C}^{d}\right)=\mathfrak{V}_{d}\left(\mathcal{S}_{s}^{\prime}\left(\mathbf{R}^{d}\right)\right)=\mathcal{A}_{s}^{\prime}\left(\mathbf{C}^{d}\right)
\end{gathered}
$$

and

$$
\mathcal{A}_{0,\left(s, \frac{1}{2}\right)}^{\prime}\left(\mathbf{C}^{d}\right)=\mathfrak{V}_{d}\left(\Sigma_{s}^{\prime}\left(\mathbf{R}^{d}\right)\right)=\mathcal{A}_{0, s}^{\prime}\left(\mathbf{C}^{d}\right)
$$

respectively.

Proposition 3.6. Let $a_{0} \in L_{l o c}^{\infty}\left(\mathbf{C}^{d}\right), s>\frac{1}{2}, 0<t_{0}<1$ and

$$
\begin{equation*}
t_{1}=\frac{1}{4\left(1-t_{0}\right)} \tag{3.16}
\end{equation*}
$$

Then the following is true:
(1) if (3.15) holds for some $r_{0}>0$ then

$$
\begin{align*}
& \operatorname{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right): \mathcal{A}_{0,\left(s, t_{0}\right)}\left(\mathbf{C}^{d}\right) \rightarrow \mathcal{A}_{0,\left(s, t_{1}\right)}\left(\mathbf{C}^{d}\right) \\
& \operatorname{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right): \mathcal{A}_{0,\left(s, t_{0}\right)}^{\prime}\left(\mathbf{C}^{d}\right) \rightarrow \mathcal{A}_{0,\left(s, t_{1}\right)}^{\prime}\left(\mathbf{C}^{d}\right) \tag{3.17}
\end{align*}
$$

are continuous;
(2) if (3.15) holds for every $r_{0}>0$ then

$$
\begin{align*}
& \mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right): \mathcal{A}_{\left(s, t_{0}\right)}\left(\mathbf{C}^{d}\right) \rightarrow \mathcal{A}_{\left(s, t_{1}\right)}\left(\mathbf{C}^{d}\right), \\
& \mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right): \mathcal{A}_{\left(s, t_{0}\right)}^{\prime}\left(\mathbf{C}^{d}\right) \rightarrow \mathcal{A}_{\left(s, t_{1}\right)}^{\prime}\left(\mathbf{C}^{d}\right) \tag{3.18}
\end{align*}
$$

are continuous.

Proof. We only prove that the first map in (3.17) is continuous. The other continuity assertions follow by similar arguments and are left for the reader.

Let $r_{2}>0$ be given, $r_{1}>r_{0}$ and $F \in \mathcal{A}_{0,\left(s, t_{0}\right)}\left(\mathbf{C}^{d}\right)$. We have for $z \in \mathbf{C}^{d}$

$$
\begin{aligned}
\left|\mathrm{Op}_{\mathfrak{V}}^{\mathrm{T}}\left(a_{0}\right) F(z)\right| e^{-t_{1}|z|^{2}+r_{2}|z|^{\frac{1}{s}}} & \lesssim e^{-t_{1}|z|^{2}+r_{2}|z|^{\frac{1}{s}}} \int_{\mathbf{C}^{d}}\left|a_{0}(w)\right||F(w)| e^{\operatorname{Re}(z, w)-|w|^{2}} d \lambda(w) \\
& \lesssim e^{-t_{1}|z|^{2}+r_{2}|z|^{\frac{1}{s}}}\|F\|_{\mathcal{A}_{s, t_{0}, r_{1}}} \int_{\mathbf{C}^{d}} e^{r_{0}|w|^{\frac{1}{s}}+t_{0}|w|^{2}-r_{1}|w|^{\frac{1}{s}}+\operatorname{Re}(z, w)-|w|^{2}} d \lambda(w) \\
& =e^{r_{2}|z|^{\frac{1}{s}}}\|F\|_{\mathcal{A}_{s, t_{0}, r_{1}}} \int_{\mathbf{C}^{d}} e^{-\left(r_{1}-r_{0}\right)|w|^{\frac{1}{s}}-\left(1-t_{0}\right)|w|^{2}+\operatorname{Re}(z, w)-t_{1}|z|^{2}} d \lambda(w) \\
& =e^{r_{2}|z|^{\frac{1}{s}}}\|F\|_{\mathcal{A}_{s, t_{0}, r_{1}}} \int_{\mathbf{C}^{d}} e^{-\left(r_{1}-r_{0}\right)|w|^{\frac{1}{s}}-\left|\sqrt{1-t_{0}} w-\frac{1}{2 \sqrt{1-t_{0}}} z\right|^{2}} d \lambda(w) \\
& =e^{r_{2}|z|^{\frac{1}{s}}}\|F\|_{\mathcal{A}_{s, t_{0}, r_{1}}} \int_{\mathbf{C}^{d}} e^{-\left(r_{1}-r_{0}\right)\left|w+\frac{1}{2\left(1-t_{0}\right)} z\right|^{\frac{1}{s}}-\left(1-t_{0}\right)|w|^{2}} d \lambda(w) \\
& \leqslant e^{\left(r_{2}-c_{1}\left(r_{1}-r_{0}\right)\right)|z|^{\frac{1}{s}}}\|F\|_{\mathcal{A}_{s, t_{0}, r_{1}}} \int_{\mathbf{C}^{d}} e^{c_{2}\left(r_{1}-r_{0}\right)|w|^{\frac{1}{s}}-\left(1-t_{0}\right)|w|^{2}} d \lambda(w) \\
& \approx\|F\|_{\mathcal{A}_{s, t_{0}, r_{1}}} e^{\left(r_{2}-c_{1}\left(r_{1}-r_{0}\right)\right)|z|^{\frac{1}{s}}}
\end{aligned}
$$

for some constants $c_{1}, c_{2}>0$. By choosing $r_{1}$ sufficiently large we get

$$
\left\|\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right) F\right\|_{\mathcal{A}_{s, t_{1}, r_{2}}} \lesssim\|F\|_{\mathcal{A}_{s, t_{0}, r_{1}}}
$$

The estimates and (1.40) imply $\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right) F \in A\left(\mathbf{C}^{d}\right)$.
Remark 3.7. Note that (3.16) implies $t_{1}>\frac{1}{4}$ and $t_{0} \leqslant t_{1}$ with equality if and only if $t_{0}=\frac{1}{2}$. Hence $\mathcal{A}_{0,\left(s, t_{0}\right)}\left(\mathbf{C}^{d}\right) \subseteq \mathcal{A}_{0,\left(s, t_{1}\right)}\left(\mathbf{C}^{d}\right)$, and similarly for the other spaces.

The particular case $t_{0}=\frac{1}{2}$ in Proposition 3.6 is the following.
Corollary 3.8. Let $a_{0} \in L_{l o c}^{\infty}\left(\mathbf{C}^{d}\right)$ and $s>\frac{1}{2}$. If (3.15) holds for some (every) $r_{0}>0$ then $\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right)$ is continuous on $\mathcal{A}_{0, s}\left(\mathbf{C}^{d}\right)\left(\right.$ on $\mathcal{A}_{s}\left(\mathbf{C}^{d}\right)$ ).

With a technique similar to the proof of Proposition 3.6 one obtains the following result.
Proposition 3.9. Let $a_{0} \in L_{l o c}^{\infty}\left(\mathbf{C}^{d}\right), s>\frac{1}{2}, 0<t_{0}<1$ and suppose (3.16) holds. Then the following is true:
(1) if (3.14) holds for all $r_{0}>0$ then

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right): \mathcal{A}_{0,\left(s, t_{0}\right)}^{\prime}\left(\mathbf{C}^{d}\right) \rightarrow \mathcal{A}_{0,\left(s, t_{1}\right)}\left(\mathbf{C}^{d}\right) \tag{3.19}
\end{equation*}
$$

is continuous;
(2) if (3.14) holds for some $r_{0}>0$ then

$$
\begin{equation*}
\mathrm{Op}_{\mathfrak{W}}^{\mathrm{aw}}\left(a_{0}\right): \mathcal{A}_{\left(s, t_{0}\right)}^{\prime}\left(\mathbf{C}^{d}\right) \rightarrow \mathcal{A}_{\left(s, t_{1}\right)}\left(\mathbf{C}^{d}\right) \tag{3.20}
\end{equation*}
$$

is continuous.
Similarly, by letting $t_{0}=\frac{1}{2}$ in Proposition 3.9 we get the following.
Corollary 3.10. Let $a_{0} \in L_{l o c}^{\infty}\left(\mathbf{C}^{d}\right)$ and $s>\frac{1}{2}$. Then the following is true:
(1) if (3.14) holds for every $r_{0}>0$ then

$$
\mathrm{Op}_{\mathfrak{Z}}^{\mathrm{aw}}\left(a_{0}\right): \mathcal{A}_{0, s}^{\prime}\left(\mathbf{C}^{d}\right) \rightarrow \mathcal{A}_{0, s}\left(\mathbf{C}^{d}\right)
$$

is continuous;
(2) if (3.14) holds for some $r_{0}>0$ then

$$
\mathrm{Op}_{\mathfrak{Z}}^{\mathrm{aw}}\left(a_{0}\right): \mathcal{A}_{s}^{\prime}\left(\mathbf{C}^{d}\right) \rightarrow \mathcal{A}_{s}\left(\mathbf{C}^{d}\right)
$$

is continuous.
3.3. Estimates of Wick symbols of anti-Wick operators with exponentially bounded symbols

For anti-Wick operators in [12, Eq. (2.94)] we have the following result.
Theorem 3.11. If $a_{0} \in L_{l o c}^{\infty}\left(\mathbf{C}^{d}\right)$ satisfies

$$
\begin{equation*}
\left|a_{0}(w)\right| \lesssim e^{r|w|^{2}}, \quad w \in \mathbf{C}^{d}, \quad \text { for some } \quad r<1, \tag{3.21}
\end{equation*}
$$

then $a_{0} \in L_{0, A}\left(\mathbf{C}^{d}\right)$ and (1.38)' holds for some $a_{0}^{\text {aw }} \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ with

$$
\left|a_{0}^{\mathrm{aw}}(z, w)\right| \lesssim e^{r_{0}|z+w|^{2}-\operatorname{Re}(z, w)}, \quad r_{0}=4^{-1}(1-r)^{-1}
$$

Proof. The claim $a_{0} \in L_{0, A}\left(\mathbf{C}^{d}\right)$ is an immediate consequence of the assumption (3.21) and the definition (1.39). The integral in (1.38)' can be estimated as

$$
\begin{aligned}
\left|\int_{\mathbf{C}^{d}} a_{0}\left(w_{1}\right) e^{-\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right)\right| & \lesssim \int_{\mathbf{C}^{d}} e^{r\left|w_{1}\right|^{2}}\left|e^{-\left(z-w_{1}, w-w_{1}\right)}\right| d \lambda\left(w_{1}\right) \\
& =e^{-\operatorname{Re}(z, w)} \int_{\mathbf{C}^{d}} e^{-(1-r)\left|w_{1}\right|^{2}} e^{\operatorname{Re}\left(z+w, w_{1}\right)} d \lambda\left(w_{1}\right) \\
& =e^{\frac{1}{4(1-r)}|z+w|^{2}-\operatorname{Re}(z, w)} \int_{\mathbf{C}^{d}} e^{-(1-r)\left|w_{1}-(z+w) /(2(1-r))\right|^{2}} d \lambda\left(w_{1}\right) \\
& =e^{r_{0}|z+w|^{2}-\operatorname{Re}(z, w)} .
\end{aligned}
$$

Remark 3.12. The condition on $a_{0}^{\text {aw }}$ in Theorem 3.11 implies that $a_{0}^{\text {aw }}$ belongs to $\widehat{\mathcal{A}}_{0, \frac{1}{2}}^{\prime}\left(\mathbf{C}^{2 d}\right)$ (see [29]). In particular it follows that $\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right)=\mathrm{Op}_{\mathfrak{V}}\left(a_{0}^{\text {aw }}\right)$ is continuous from $\mathcal{A}_{0, \frac{1}{2}}\left(\mathbf{C}^{d}\right)$ to $\mathcal{A}_{0, \frac{1}{2}}^{\prime}\left(\mathbf{C}^{d}\right)$ (cf. [29, Theorem 2.10] and Remark 1.2).

The following result concerns exponentially moderate weight functions.
Theorem 3.13. Let $a_{0} \in L_{0, A}\left(\mathbf{C}^{d}\right), a_{0}^{\text {aw }} \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ be given by (1.38)' and $\omega \in \mathscr{P}_{E}\left(\mathbf{C}^{d}\right)$. If

$$
\left|a_{0}(w)\right| \lesssim \omega(2 w), \quad w \in \mathbf{C}^{d},
$$

then

$$
\left|a_{0}^{\mathrm{aw}}(z, w)\right| \lesssim e^{\frac{1}{4}|z-w|^{2}} \omega(z+w), \quad z, w \in \mathbf{C}^{d}
$$

Proof. Let $r \geqslant 0$ be chosen such that $\omega(z+w) \lesssim \omega(z) e^{r|w|}, z, w \in \mathbf{C}^{d}$. From (1.38)' we get

$$
\begin{aligned}
\left|a_{0}^{\mathrm{aw}}(z, w)\right| & \lesssim \int_{\mathbf{C}^{d}} \omega\left(2 w_{1}\right) e^{-\operatorname{Re}\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right) \\
& =e^{-\operatorname{Re}(z, w)} \int_{\mathbf{C}^{d}} \omega\left(2 w_{1}\right) e^{\operatorname{Re}\left(z+w, w_{1}\right)-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& =e^{-\operatorname{Re}(z, w)+\frac{1}{4}|z+w|^{2}} \int_{\mathbf{C}^{d}} \omega\left(2 w_{1}\right) e^{-\left|w_{1}-(z+w) / 2\right|^{2}} d \lambda\left(w_{1}\right) \\
& =e^{\frac{1}{4}|z-w|^{2}} \int_{\mathbf{C}^{d}} \omega\left(2 w_{1}+z+w\right) e^{-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& \lesssim e^{\frac{1}{4}|z-w|^{2}} \omega(z+w) \int_{\mathbf{C}^{d}} e^{2 r\left|w_{1}\right|-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& \approx e^{\frac{1}{4}|z-w|^{2}} \omega(z+w) .
\end{aligned}
$$

The anti-Wick operators in Propositions 3.6 and 3.9 can also be described as Wick operators with symbols that have smaller growth bounds than $\hat{\mathcal{A}}_{s}\left(\mathbf{C}^{2 d}\right)$ and its dual. The following result extends Theorem 3.13 for weights of the form $e^{c|z|^{\frac{1}{s}}}$ with $c \in \mathbf{R}$ from $s \geqslant 1$ to $s \geqslant \frac{1}{2}$.

Theorem 3.14. Let $s \geqslant \frac{1}{2}\left(s>\frac{1}{2}\right), a_{0} \in L_{0, A}\left(\mathbf{C}^{d}\right)$ and let $a_{0}^{\text {aw }}$ be given by (1.38)'. Then the following is true:
(1) if (3.14) holds for some (every) $r_{0}>0$ then

$$
\begin{equation*}
\left|a_{0}^{\mathrm{aw}}(z, w)\right| \lesssim e^{\frac{1}{4}|z-w|^{2}-r|z+w|^{\frac{1}{s}}} \tag{3.22}
\end{equation*}
$$

for some (every) $r>0$;
(2) if (3.15) holds for every (some) $r_{0}>0$ then

$$
\begin{equation*}
\left|a_{0}^{\mathrm{aw}}(z, w)\right| \lesssim e^{\frac{1}{4}|z-w|^{2}+r|z+w|^{\frac{1}{s}}} \tag{3.23}
\end{equation*}
$$

for every (some) $r>0$.
Remark 3.15. Thanks to the parameter $\frac{1}{4}$ in the factor $e^{\frac{1}{4}|z-w|^{2}}$ rather than $\frac{1}{2}$, the estimates (3.23) are much stronger than the estimates (2.19) with $\sigma=s$. Corollary 3.8 can thus be seen as a consequence of Theorems 2.6 and 3.14, and [9, Definition 2.4, and Theorems 4.10 and 4.11].

Remark 3.16. The estimates for $a_{0}^{\text {aw }}$ in Theorem 3.14 may seem weak since the dominating factor $e^{\frac{1}{4}|z-w|^{2}}$ is present in (3.22) and (3.23) but absent in the original estimates (3.14) and (3.15) for $a_{0}$.

On the other hand, Wick symbols for operators with continuity involving the spaces $\mathcal{A}_{s}\left(\mathbf{C}^{d}\right)$ and $\mathcal{A}_{s}^{\prime}\left(\mathbf{C}^{d}\right)$, as well as $\mathcal{A}_{0, s}\left(\mathbf{C}^{d}\right)$ and $\mathcal{A}_{0, s}^{\prime}\left(\mathbf{C}^{d}\right)$, usually satisfy conditions of the form

$$
|a(z, w)| \lesssim e^{\frac{1}{2}|z-w|^{2} \pm r_{1}|z+w|^{\frac{1}{2}} \pm|z-w|^{\frac{1}{s}}}
$$

in view of [29, Theorems 2.9 and 2.10], and Theorem 2.6. Here the dominating factor is $e^{\frac{1}{2}|z-w|^{2}}$, which is larger than the factor $e^{\frac{1}{4}|z-w|^{2}}$ in Theorem 3.14.

This factor has a large impact on functions on $\mathbf{R}^{d}$ that are transformed back by the inverse of the Bargmann transform. For instance, if $\varepsilon>0$, then the Bargmann image of any non-trivial Gelfand-Shilov space and its distribution space contain

$$
\left\{F \in A\left(\mathbf{C}^{d}\right) ;|F(z)| \lesssim e^{\left(\frac{1}{2}-\varepsilon\right)|z|^{2}}\right\}
$$

and are contained in

$$
\left\{F \in A\left(\mathbf{C}^{d}\right) ;|F(z)| \lesssim e^{\left(\frac{1}{2}+\varepsilon\right)|z|^{2}}\right\} .
$$

The same holds true for the Bargmann images of $\mathscr{S}\left(\mathbf{R}^{d}\right)$ and $\mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$.
Theorem 3.14 is a straight-forward consequence of the following two propositions, which give more details on the relationships between $r$ and $r_{0}$ in (3.14), (3.15), (3.22) and (3.23).

Proposition 3.17. Let $s \geqslant \frac{1}{2}$ and let $r_{0}, r \in(0, \infty)$ be such that

$$
\begin{equation*}
r_{0} \in(0, \infty) \quad \text { and } \quad r<\frac{r_{0}}{4\left(1+r_{0}\right)}, \quad \text { when } \quad s=\frac{1}{2}, \tag{3.24}
\end{equation*}
$$

and

$$
\begin{equation*}
r_{0} \in(0, \infty) \quad \text { and } \quad r \leqslant 2^{-\frac{1}{s}} r_{0}, \quad \text { when } \quad s \in\left(\frac{1}{2}, \infty\right) \text {, } \tag{3.25}
\end{equation*}
$$

with strict inequality in (3.25) when $s<1$. If $a_{0} \in L_{\text {loc }}^{\infty}\left(\mathbf{C}^{d}\right)$ satisfies (3.14) and $a_{0}^{\text {aw }} \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ is given by (1.38)', then (3.22) holds.

Proposition 3.18. Let $s \geqslant \frac{1}{2}$ and $r_{0}, r \in(0, \infty)$ be such that

$$
\begin{equation*}
r_{0} \in(0,1) \quad \text { and } \quad r>\frac{r_{0}}{4\left(1-r_{0}\right)}, \quad \text { when } \quad s=\frac{1}{2} \tag{3.24}
\end{equation*}
$$

and

$$
\begin{equation*}
r_{0} \in(0, \infty) \quad \text { and } \quad r \geqslant 2^{-\frac{1}{s}} r_{0}, \quad \text { when } \quad s \in\left(\frac{1}{2}, \infty\right) \text {, } \tag{3.25}
\end{equation*}
$$

with strict inequality in (3.25)' when $s<1$. If $a_{0} \in L_{l o c}^{\infty}\left(\mathbf{C}^{d}\right)$ satisfies (3.15) and $a_{0}^{\text {aw }} \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ is given by (1.38)', then (3.23) holds.

For the proofs of Propositions 3.17 and 3.18 we use the inequalities

$$
\begin{align*}
|z|^{\theta}-|w|^{\theta} & \leqslant|z+w|^{\theta} \leqslant|z|^{\theta}+|w|^{\theta}, & & \theta \in(0,1], z, w \in \mathbf{C}^{d}  \tag{3.26}\\
|z+w|^{\theta} & \leqslant(1+\varepsilon)|z|^{\theta}+\left(1+\varepsilon^{-1}\right)|w|^{\theta}, & & \theta \in[1,2], z, w \in \mathbf{C}^{d}, \tag{3.27}
\end{align*}
$$

and

$$
\begin{equation*}
|z+w|^{\theta} \geqslant(1-\varepsilon)|z|^{\theta}+\left(1-\varepsilon^{-1}\right)|w|^{\theta}, \quad \theta \in[1,2], z, w \in \mathbf{C}^{d}, \tag{3.28}
\end{equation*}
$$

for every $\varepsilon>0$.
Proof of Proposition 3.17. Suppose that $a_{0}$ satisfies (3.14) for some $r_{0}>0$. First we consider the case $s>\frac{1}{2}$. If $s<1$ let $\varepsilon_{1}>0$ and $\varepsilon_{2}=\varepsilon_{1}^{-1}$, and if $s \geqslant 1$ let $\varepsilon_{1}=0$ and $\varepsilon_{2}=2$, and let $c=2^{-\frac{1}{s}}$. Then (1.38)', (3.26) and (3.28) give

$$
\begin{align*}
\left|a_{0}^{\mathrm{aw}}(z, w)\right| & \lesssim \int_{\mathbf{C}^{d}} e^{-r_{0}\left|w_{1}\right|^{\frac{1}{s}}} e^{-\operatorname{Re}\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right) \\
& =e^{\frac{1}{4}|z+w|^{2}-\operatorname{Re}(z, w)} \int_{\mathbf{C}^{d}} e^{-r_{0}\left|w_{1}\right|^{\frac{1}{s}}-\left|w_{1}-(z+w) / 2\right|^{2}} d \lambda\left(w_{1}\right) \\
& =e^{\frac{1}{4}|z-w|^{2}} \int_{\mathbf{C}^{d}} e^{-r_{0}\left|w_{1}+(z+w) / 2\right|^{\frac{1}{s}}-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& \leqslant e^{\frac{1}{4}|z-w|^{2}} e^{-c r_{0}\left(1-\varepsilon_{1}\right)|z+w|^{\frac{1}{s}}} \int_{\mathbf{C}^{d}} e^{-r_{0}\left(1-\varepsilon_{2}\right)\left|w_{1}\right|^{\frac{1}{s}}-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& =e^{\frac{1}{4}|z-w|^{2}} e^{-c r_{0}\left(1-\varepsilon_{1}\right)|z+w|^{\frac{1}{s}}} . \tag{3.29}
\end{align*}
$$

If $s \geqslant 1$, then $\varepsilon_{1}=0$ and $\varepsilon_{2}=2$, and the result follows from (3.29). If instead $s<1$, then the result follows by choosing $\varepsilon_{1}>0$ small enough, and we have proved the result in the case $s>\frac{1}{2}$.

Next suppose that $s=\frac{1}{2}$. For $\varepsilon_{1}>0$ and $\varepsilon_{2}=\varepsilon_{1}^{-1}$, (3.29) gives

$$
\left|a_{0}^{\text {aw }}(z, w)\right| \lesssim e^{\frac{1}{4}|z-w|^{2}} e^{-\frac{1}{4} r_{0}\left(1-\varepsilon_{1}\right)|z+w|^{2}} \int_{\mathbf{C}^{d}} e^{-\left(r_{0}\left(1-\varepsilon_{2}\right)+1\right)\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) .
$$

For any $\varepsilon_{2}<\frac{1+r_{0}}{r_{0}}$ it follows that the integral converges, and

$$
1-\varepsilon_{1}=1-\varepsilon_{2}^{-1}<\left(1+r_{0}\right)^{-1} .
$$

By the assumptions there is $\delta>0$ such that

$$
r=\frac{r_{0}(1-\delta)}{4\left(1+r_{0}\right)} .
$$

Since

$$
1-\varepsilon_{1} \nearrow\left(1+r_{0}\right)^{-1} \quad \text { as } \quad \varepsilon_{2} \nearrow \frac{1+r_{0}}{r_{0}}
$$

we may pick $0<\varepsilon_{2}<\frac{1+r_{0}}{r_{0}}$ such that

$$
\frac{1-\delta}{1+r_{0}} \leqslant 1-\varepsilon_{1}
$$

and the result follows in the case $s=\frac{1}{2}$.
Proof of Proposition 3.18. First we consider the case when $s>\frac{1}{2}$. Suppose that $a_{0}$ satisfies (3.15) for some $r_{0}>0$, let $\varepsilon_{1}, \varepsilon_{2} \geqslant 0$ be such that $\varepsilon_{1}=\varepsilon_{2}=0$ when $s \geqslant 1$ and $\varepsilon_{1} \varepsilon_{2}=1$ when $s<1$, and let $c=2^{-\frac{1}{s}}$. Then (1.38)', (3.26) and (3.27) give

$$
\begin{align*}
\left|a_{0}^{\mathrm{aw}}(z, w)\right| & \lesssim \int_{\mathbf{C}^{d}} e^{r_{0}\left|w_{1}\right|^{\frac{1}{s}}} e^{-\operatorname{Re}\left(z-w_{1}, w-w_{1}\right)} d \lambda\left(w_{1}\right) \\
& =e^{\frac{1}{4}|z+w|^{2}-\operatorname{Re}(z, w)} \int_{\mathbf{C}^{d}} e^{r_{0}\left|w_{1}\right|^{\frac{1}{s}}-\left|w_{1}-(z+w) / 2\right|^{2}} d \lambda\left(w_{1}\right) \\
& =e^{\frac{1}{4}|z-w|^{2}} \int_{\mathbf{C}^{d}} e^{r_{0}\left|w_{1}+(z+w) / 2\right|^{\frac{1}{s}}-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& \leqslant e^{\frac{1}{4}|z-w|^{2}} e^{c r_{0}\left(1+\varepsilon_{1}\right)|z+w|^{\frac{1}{s}}} \int_{\mathbf{C}^{d}} e^{r_{0}\left(1+\varepsilon_{2}\right)\left|w_{1}\right|^{\frac{1}{s}}-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) \\
& \approx e^{\frac{1}{4}|z-w|^{2}} e^{c r_{0}\left(1+\varepsilon_{1}\right)|z+w|^{\frac{1}{s}}} . \tag{3.30}
\end{align*}
$$

If $s \geqslant 1$, then $\varepsilon_{1}=\varepsilon_{2}=0$, and the result follows from (3.30). If instead $s<1$, then the result follows by choosing $\varepsilon_{1}>0$ small enough, and the result follows in the case $s>\frac{1}{2}$.

Next suppose that $s=\frac{1}{2}$. Then (3.30) gives

$$
\left|a_{0}^{\text {aw }}(z, w)\right| \lesssim e^{\frac{1}{4}|z-w|^{2}} e^{\frac{1}{4} r_{0}\left(1+\varepsilon_{1}\right)|z+w|^{2}} \int_{\mathbf{C}^{d}} e^{r_{0}\left(1+\varepsilon_{2}\right)\left|w_{1}\right|^{2}-\left|w_{1}\right|^{2}} d \lambda\left(w_{1}\right) .
$$

For any $\varepsilon_{2}<\frac{1-r_{0}}{r_{0}}$ the integral converges, and

$$
1+\varepsilon_{1}=1+\varepsilon_{2}^{-1}>\left(1-r_{0}\right)^{-1} .
$$

Since

$$
1+\varepsilon_{1} \searrow\left(1-r_{0}\right)^{-1} \quad \text { as } \quad \varepsilon_{2} \nearrow \frac{1-r_{0}}{r_{0}}
$$

the result follows in the case $s=\frac{1}{2}$ by letting $r=\frac{r_{0}\left(1+\varepsilon_{1}\right)}{4}$.
Remark 3.19. In the previous results we have mainly deduced certain estimates on the Wick symbols of anti-Wick operators. In [32], similar estimates are presented, where estimates on the anti-Wick symbols to certain Wick operators are deduced. We also refer to $[6,32]$ for further transition properties between Wick and anti-Wick operators, and to [32] for certain fundamental continuity properties for Wick and anti-Wick operators.

## 4. A lower bound for Wick operators

In this section we apply the asymptotic expansions in the previous section for Shubin-Wick operators to deduce a sharp Gårding inequality.

First we have the following result. We put $\widehat{\mathcal{A}}_{\text {Sh }, \rho}\left(\mathbf{C}^{2 d}\right)=\widehat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ when $\omega=1$.
Proposition 4.1. Let $\omega \in \mathscr{P}\left(\mathbf{C}^{d}\right), p \in[1, \infty], a \in \hat{\mathcal{A}}_{\mathrm{Sh}, 0}\left(\mathbf{C}^{2 d}\right)$ and $a_{0} \in L^{\infty}\left(\mathbf{C}^{d}\right)$. Then $\mathrm{Op}_{\mathfrak{V}}(a)$ and $\mathrm{Op}_{\mathfrak{Y}}^{\mathrm{aw}}\left(a_{0}\right)$ are both continuous on $A_{(\omega)}^{p}\left(\mathbf{C}^{d}\right)$.

The claimed continuity of $\mathrm{Op}_{\mathfrak{V}}(a)$ is a straight-forward consequence of [29, Theorem 3.3], in combination with Proposition 2.1 and the relationship $K(z, w)=a(z, w) e^{(z, w)}$ between the kernel and symbol of a Wick operator (cf. (0.1)). In order to be self-contained we include an alternative and shorter proof.

Proof. Let $F \in A_{(\omega)}^{p}\left(\mathbf{C}^{d}\right), G(z)=e^{-\frac{1}{2}|z|^{2}}|F(z) \omega(\sqrt{2} \bar{z})|$,

$$
\begin{aligned}
& H_{1}(z)=e^{-\frac{1}{2}|z|^{2}}\left|\operatorname{Op}_{\mathfrak{V}}(a) F(z) \omega(\sqrt{2} \bar{z})\right| \text { and } \\
& H_{2}(z)=e^{-\frac{1}{2}|z|^{2}}\left|\operatorname{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(a_{0}\right) F(z) \omega(\sqrt{2} \bar{z})\right| .
\end{aligned}
$$

We have

$$
\omega(\sqrt{2} \bar{z}) \lesssim \omega(\sqrt{2} \bar{w})\langle z-w\rangle^{N_{0}}
$$

for some $N_{0} \geqslant 0$. By Theorem 2.2 and (2.6) we get

$$
H_{1}(z) \lesssim e^{-\frac{1}{2}|z|^{2}} \int_{\mathbf{C}^{d}} e^{\frac{1}{2}|z-w|^{2}}\langle z-w\rangle^{-N}|F(w) \omega(\sqrt{2} \bar{z})| e^{\operatorname{Re}(z, w)-|w|^{2}} d \lambda(w)=\left(\langle\cdot\rangle^{N_{0}-N} * G\right)(z)
$$

for every $N \geqslant 0$. By choosing $N>2 d+N_{0}$ and using Young's inequality we get $\left\|H_{1}\right\|_{L^{p}} \lesssim\|G\|_{L^{p}}$ which means $\left\|\mathrm{Op}_{\mathfrak{V}}(a) F\right\|_{A_{(\omega)}^{p}} \lesssim\|F\|_{A_{(\omega)}^{p}}$, and the asserted continuity for $\mathrm{Op}_{\mathfrak{V}}(a)$ follows.

In the same way we get

$$
H_{2}(z) \lesssim\left\|a_{0}\right\|_{L^{\infty}} e^{-\frac{1}{2}|z|^{2}} \int_{\mathbf{C}^{d}}|F(w) \omega(\sqrt{2} \bar{w})|\langle z-w\rangle^{N_{0}} e^{\operatorname{Re}(z, w)-|w|^{2}} d \lambda(w) \asymp\left(\left(\langle\cdot\rangle^{N_{0}} e^{-\frac{1}{2}|\cdot|^{2}}\right) * G\right)(z),
$$

and another application of Young's inequality shows that $\left\|H_{2}\right\|_{L_{(\omega)}^{p}} \lesssim\|G\|_{L_{(\omega)}^{p}}$, that is $\left\|\mathrm{Op}_{\mathfrak{Y}}^{\mathrm{aw}}\left(a_{0}\right) F\right\|_{A_{(\omega)}^{p}} \lesssim$ $\|F\|_{A_{(\omega)}^{p}}$.

We have finally a version of the sharp Gårding inequality.
Theorem 4.2. Let $\rho>0, \omega(z)=\langle z\rangle^{2 \rho}$ and let $a \in \mathcal{\mathcal { A }}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ be such that $a(w, w) \geqslant-C_{0}$ for all $w \in \mathbf{C}^{d}$, for some constant $C_{0} \geqslant 0$. Then

$$
\begin{equation*}
\operatorname{Re}\left(\left(\mathrm{Op}_{\mathfrak{V}}(a) F, F\right)_{A^{2}}\right) \geqslant-C\|F\|_{A^{2}}^{2}, \quad F \in \mathcal{A}_{\mathscr{I}}\left(\mathbf{C}^{d}\right) \tag{4.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\operatorname{Im}\left(\left(\mathrm{Op}_{\mathfrak{V}}(a) F, F\right)_{A^{2}}\right)\right| \leqslant C\|F\|_{A^{2}}^{2}, \quad F \in \mathcal{A}_{\mathscr{S}}\left(\mathbf{C}^{d}\right) \tag{4.2}
\end{equation*}
$$

for some constant $C \geqslant 0$.

Proof. Let $b_{0}(w)=a(w, w)$. Then $\mathrm{Op}_{\mathfrak{V}}(a)=\mathrm{Op}_{\mathfrak{V}}^{\mathrm{aw}}\left(b_{0}\right)+\mathrm{Op}_{\mathfrak{V}}\left(a_{1}\right)$ for some $a_{1} \in \hat{\mathcal{A}}_{\mathrm{Sh}, \rho}\left(\mathbf{C}^{2 d}\right) \subseteq \hat{\mathcal{A}}_{\mathrm{Sh}, 0}\left(\mathbf{C}^{2 d}\right)$, in view of Proposition 3.3. Since $\Pi_{A} F=F$ for $F \in A^{2}\left(\mathbf{C}^{d}\right)$ (cf. (1.22)), the assumption $b_{0} \geqslant-C_{0}$ implies $\left(\mathrm{Op}_{\mathfrak{V}}^{\text {aw }}\left(b_{0}\right) F, F\right)_{A^{2}} \geqslant-C_{0}\|F\|_{A^{2}}^{2}$ for every $F \in \mathcal{A}_{\mathscr{S}}\left(\mathbf{C}^{d}\right)$. The operator $\mathrm{Op}_{\mathfrak{V}}\left(a_{1}\right)$ is continuous on $A^{2}\left(\mathbf{C}^{d}\right)$ in view of Proposition 4.1. A combination of these facts gives the result.

## 5. Ellipticity and hypoellipticity for Shubin and Wick operators

In this section we show that the Bargmann assignment $\mathrm{S}_{\mathfrak{V}}$ maps the sets of hypoelliptic symbols and weakly elliptic symbols in the Shubin class $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ bijectively into the sets of hypoelliptic symbols and weakly elliptic Wick symbols in $\mathcal{A}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$, respectively. (Cf. Subsection 1.6.) Then we explain some consequences for polynomial symbols.

### 5.1. Transition of weakly elliptic symbols

For symbols in $\hat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ we define ellipticity and weak ellipticity as follows.
Definition 5.1. Let $\rho>0, \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{C}^{d}\right)$ and $a \in \hat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$. Then $a$ is called weakly elliptic of order $\rho_{0} \geqslant 0$, or $\rho_{0}$-weakly elliptic, if for some $R>0$

$$
|a(z, z)| \gtrsim\langle z\rangle^{-\rho_{0}} \omega(\sqrt{2} \bar{z}), \quad|z| \geqslant R .
$$

If $a$ is weakly elliptic of order 0 then $a$ is called elliptic.
Theorem 5.2. Let $\omega \in \mathscr{P}\left(\mathbf{R}^{2 d}\right) \simeq \mathscr{P}\left(\mathbf{C}^{d}\right), \rho>0$ and $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$. Then the following is true:
(1) if $z=x+i \xi, x, \xi \in \mathbf{R}^{d}$, then

$$
\begin{equation*}
\left|\mathrm{S}_{\mathfrak{V} \mathfrak{a}}(z, z)-\mathfrak{a}(\sqrt{2} x,-\sqrt{2} \xi)\right| \lesssim \omega(\sqrt{2} \bar{z})\langle z\rangle^{-2 \rho} ; \tag{5.1}
\end{equation*}
$$

(2) if $\rho_{0} \in[0,2 \rho)$, then $\mathrm{S}_{\mathfrak{V}}$ is bijective from the set of weakly elliptic symbols in $\mathrm{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ of order $\rho_{0}$ to the set of weakly elliptic symbols in $\hat{A}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ of order $\rho_{0}$.

As a consequence of (2) in the previous theorem we get the following. Here recall Subsection 1.6 for definition of weakly elliptic symbols in Shubin classes.

Corollary 5.3. Let $\mathfrak{a}$ be as in Theorem 5.2. Then the following is true:
(1) if $\rho_{0} \in[0,2 \rho)$, then $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ is weakly elliptic of order $\rho_{0}$, if and only if $\mathrm{S}_{\mathfrak{V} \mathfrak{a}} \in \widehat{A}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ is weakly elliptic of order $\rho_{0}$;
(2) $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ is elliptic if and only if $\mathrm{S}_{\mathfrak{V} \mathfrak{a}} \in \widehat{A}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ is elliptic.

For the proof of Theorem 5.2 we need the following proposition, related to Propositions 3.1 and 3.5. It shows that the Bargmann assignment of a Shubin symbol $\mathfrak{a}$ possess convenient expansion properties.

Proposition 5.4. Let $M \geqslant 0$ be an integer, $\rho \geqslant 0, \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{2 d}\right) \simeq \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{C}^{d}\right), \omega_{k}(x, \xi)=$ $\omega(x, \xi)\langle(x, \xi)\rangle^{-2 \rho k}$ and $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$. Then for some $\mathfrak{c}_{M} \in \operatorname{Sh}_{\rho}^{\left(\omega_{M+1}\right)}\left(\mathbf{R}^{2 d}\right)$ and constants $\left\{c_{\alpha}\right\}_{|\alpha| \leqslant 2 M}$ with $c_{0}=1$, it holds

$$
\begin{equation*}
\mathrm{S}_{\mathfrak{V} \mathfrak{a}\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right)=\sum_{k=0}^{M} \mathfrak{a}_{k}(x,-\xi)+\mathfrak{c}_{M}(x,-\xi), \quad \mathfrak{a}_{k}=\sum_{|\alpha|=2 k} c_{\alpha} \partial^{\alpha} \mathfrak{a} . . . . . . . .} \tag{5.2}
\end{equation*}
$$

Proof. Let $\psi$ be as in Proposition 2.3. If we put $z=w$, then (2.9) and Taylor's formula give

$$
\begin{align*}
(2 \pi)^{d} \mathrm{~S}_{\mathfrak{V}} \mathfrak{a}\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right) & =(2 \pi)^{\frac{3 d}{2}} \mathcal{T}_{\psi} \mathfrak{a}(x,-\xi, 0,0) \\
& =2^{d} \iint_{\mathbf{R}^{2 d}} \mathfrak{a}(t+x, \tau-\xi) e^{-\left(|t|^{2}+|\tau|^{2}\right)} d t d \tau=\sum_{k=0}^{2 M+1} \mathfrak{b}_{k}(x,-\xi)+\mathfrak{c}(x,-\xi) \tag{5.3}
\end{align*}
$$

where

$$
\mathfrak{b}_{k}(x, \xi)=\frac{2^{d}}{k!} \iint_{\mathbf{R}^{2 d}}\left\langle\mathfrak{a}^{(k)}(x, \xi) ;(t, \tau), \ldots,(t, \tau)\right\rangle e^{-\left(|t|^{2}+|\tau|^{2}\right)} d t d \tau
$$

and

$$
\mathfrak{c}(x, \xi)=\frac{1}{(2 M+1)!} \int_{0}^{1}(1-\theta)^{2 M+1} \mathfrak{c}_{\theta}(x, \xi) d \theta
$$

with

$$
\mathfrak{c}_{\theta}(x, \xi)=2^{d} \iint_{\mathbf{R}^{2 d}}\left\langle\mathfrak{a}^{(2 M+2)}(x+\theta t, \xi+\theta \tau) ;(t, \tau), \ldots,(t, \tau)\right\rangle e^{-\left(|t|^{2}+|\tau|^{2}\right)} d t d \tau
$$

If $k$ is odd, then

$$
(t, \tau) \mapsto\left\langle\mathfrak{a}^{(k)}(x, \xi) ;(t, \tau), \ldots,(t, \tau)\right\rangle e^{-\left(|t|^{2}+|\tau|^{2}\right)}
$$

is odd which implies that the integral is zero. Hence $\mathfrak{b}_{k}(x, \xi)=0$ when $k$ is odd. For $k=0$ we observe that the integral for $\mathfrak{b}_{0}$ becomes

$$
2^{d} \iint_{\mathbf{R}^{2 d}} e^{-\left(|t|^{2}+|\tau|^{2}\right)} d t d \tau=(2 \pi)^{d},
$$

and it follows from these relations that

$$
(2 \pi)^{-d} \sum_{k=0}^{2 M+1} \mathfrak{b}_{k}=\sum_{k=0}^{M} \mathfrak{a}_{k},
$$

with $\mathfrak{a}_{k}$ as in (5.2) and $c_{0}=1$. Hence the result follows if we prove that the last term in (5.3) satisfies $\mathfrak{c}_{M} \in \operatorname{Sh}_{\rho}^{\left(\omega_{M+1}\right)}\left(\mathbf{R}^{2 d}\right)$.

For $\theta \in[0,1]$ and $\alpha \in \mathbf{N}^{2 d}$ we have

$$
\begin{aligned}
\left|\partial^{\alpha} \mathfrak{c}_{\theta}(x, \xi)\right| & \lesssim \iint_{\mathbf{R}^{2 d}}\left|\partial^{\alpha} \mathfrak{a}^{(2 M+2)}(x+\theta t, \xi+\theta \tau)\right|\langle(t, \tau)\rangle^{2 M+2} e^{-\left(\left.|t|\right|^{2}+|\tau|^{2}\right)} d t d \tau \\
& \lesssim \iint_{\mathbf{R}^{2 d}} \omega(x+\theta t, \xi+\theta \tau)\langle(x+\theta t, \xi+\theta \tau)\rangle^{-(2 M+2+|\alpha|) \rho}\langle(t, \tau)\rangle^{2 M+2} e^{-\left(|t|^{2}+|\tau|^{2}\right)} d t d \tau
\end{aligned}
$$

$$
\begin{aligned}
& \lesssim \omega(x, \xi)\langle(x, \xi)\rangle^{-(2 M+2+|\alpha|) \rho} \iint_{\mathbf{R}^{2 d}}\langle(t, \tau)\rangle^{N_{0}} e^{-\left(|t|^{2}+|\tau|^{2}\right)} d t d \tau \\
& =\omega(x, \xi)\langle(x, \xi)\rangle^{-(2 M+2+|\alpha|) \rho}
\end{aligned}
$$

for some $N_{0}>0$. In the last inequality we have used the fact that $\omega$ is polynomially moderate.
This implies

$$
\left|\partial^{\alpha} \mathfrak{c}(x, \xi)\right| \lesssim \int_{0}^{1}\left|\partial^{\alpha} \mathfrak{c}_{\theta}(x, \xi)\right| d \theta \lesssim \omega(x, \xi)\langle(x, \xi)\rangle^{-(2 M+2+|\alpha|) \rho},
$$

which shows that $\mathfrak{c}, \mathfrak{c}_{M} \in \operatorname{Sh}_{\rho}^{\left(\omega_{M+1}\right)}\left(\mathbf{R}^{2 d}\right)$.
Proof of Theorem 5.2. Let $\psi$ be as in Proposition 2.3 and $N=0$ in Proposition 5.4. Then

$$
\begin{equation*}
\left|\mathrm{S}_{\mathfrak{N} \mathfrak{a}}\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right)-\mathfrak{a}(x,-\xi)\right| \lesssim \omega(x,-\xi)\langle(x,-\xi)\rangle^{-2 \rho}, \tag{5.1}
\end{equation*}
$$

and (1) follows.
Suppose $\rho_{0} \in[0,2 \rho)$. Then it follows from the latter inequality that

$$
\left|\mathrm{S}_{\mathfrak{V} \mathfrak{a}}(z, z)\right| \gtrsim\langle z\rangle^{-\rho_{0}} \omega(\sqrt{2} \bar{z}), \quad|z| \geqslant R
$$

for some $R>0$, if and only if

$$
|\mathfrak{a}(x, \xi)| \gtrsim\langle(x, \xi)\rangle^{-\rho_{0}} \omega(x, \xi), \quad|z| \geqslant R
$$

for some $R>0$, and the asserted equivalence in (2) follows.

### 5.2. Shubin hypoellipticity in a Wick setting

Definition 5.5. Let $\rho>0, \rho_{0} \geqslant 0, \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{C}^{d}\right)$ and $a \in \mathcal{\mathcal { A }}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$. Then $a$ is called hypoelliptic (in the Shubin-Wick sense in $\left.\mathcal{\mathcal { A }}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)\right)$ of order $\rho_{0}$, if there is an $R>0$ such that for every $\alpha, \beta \in \mathbf{N}^{d}$, it holds

$$
\begin{equation*}
\left|\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a(z, z)\right| \lesssim|a(z, z)|\langle z\rangle^{-\rho|\alpha+\beta|}, \quad|z| \geqslant R \tag{5.4}
\end{equation*}
$$

and

$$
\begin{equation*}
|a(z, z)| \gtrsim \omega_{0}(\sqrt{2} \bar{z})\langle z\rangle^{-\rho_{0}}, \quad|z| \geqslant R . \tag{5.5}
\end{equation*}
$$

According to Definition 1.11, if $\omega, \rho$ and $\rho_{0}$ are as in the definition, then $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ is hypoelliptic of order $\rho_{0}$ means that there is an $R>0$ such that for every $\alpha \in \mathbf{N}^{2 d}$, it holds

$$
\begin{equation*}
\left|\partial^{\alpha} \mathfrak{a}(x, \xi)\right| \lesssim|\mathfrak{a}(x, \xi)|\langle(x, \xi)\rangle^{-\rho|\alpha|}, \quad|(x, \xi)| \geqslant R \tag{5.6}
\end{equation*}
$$

and

$$
\begin{equation*}
|\mathfrak{a}(x, \xi)| \gtrsim \omega(x, \xi)\langle(x, \xi)\rangle^{-\rho_{0}}, \quad|(x, \xi)| \geqslant R . \tag{5.7}
\end{equation*}
$$

Similar to Theorem 5.2 we have the following.

Theorem 5.6. Let $\rho>0, \rho_{0} \geqslant 0, \omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{R}^{2 d}\right) \simeq \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{C}^{d}\right), \mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ and $a=\mathrm{S}_{\mathfrak{V} \mathfrak{a}}$. Then $\mathfrak{a}$ is hypoelliptic of order $\rho_{0}$ in $\mathrm{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$, if and only if a is hypoelliptic of order $\rho_{0}$ in $\widehat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$.

Proof. Suppose that $\mathfrak{a} \in \operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ is hypoelliptic of order $\rho_{0}$, and choose $M \geqslant 0$ such that $2 M \rho>\rho_{0}$. Suppose that $R>0$ is chosen such that (5.6) and (5.7) are fulfilled. Then Proposition 5.4 gives for $z=x+i \xi$ with $|z| \geqslant R$ where $R>0$ is sufficiently large

$$
\begin{aligned}
\left|a\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right)\right| & \gtrsim|\mathfrak{a}(x,-\xi)|-\sum_{k=1}^{M} \sum_{|\alpha|=2 k}\left(\left|\partial^{\alpha} \mathfrak{a}(x,-\xi)\right|+|\mathfrak{c}(x,-\xi)|\right) \\
& \gtrsim|\mathfrak{a}(x,-\xi)|-|\mathfrak{a}(x,-\xi)|\langle(x,-\xi)\rangle^{-2 \rho}-\omega(x,-\xi)\langle(x,-\xi)\rangle^{-\rho(2 M+2)} \\
& \gtrsim|\mathfrak{a}(x,-\xi)|-|\mathfrak{a}(x,-\xi)|\langle(x,-\xi)\rangle^{-2 \rho} \\
& \gtrsim|\mathfrak{a}(x,-\xi)| \gtrsim \omega(x,-\xi)\langle(x,-\xi)\rangle^{-\rho_{0}},
\end{aligned}
$$

and (5.5) follows. In particular it follows from the previous estimates that

$$
\begin{equation*}
\left|a\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right)\right| \gtrsim|\mathfrak{a}(x,-\xi)|, \quad|z| \geqslant R . \tag{5.8}
\end{equation*}
$$

For fixed $\alpha, \beta \in \mathbf{N}^{d}$, let $\Omega_{k}$ be the set of all $(\gamma, \delta) \in \mathbf{N}^{2 d} \times \mathbf{N}^{2 d}$ such that $|\gamma|=2 k$ and $|\delta|=|\alpha+\beta|$. By Proposition 5.4 and (5.8) we have for some $R$ large enough and $|z| \geqslant R$,

$$
\begin{aligned}
\left|\left(\partial_{z}^{\alpha} \bar{\partial}_{w}^{\beta} a\right)\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right)\right| & \lesssim \sum_{k=0}^{M} \sum_{(\gamma, \delta) \in \Omega_{k}}\left(\left|\partial^{\gamma+\delta} \mathfrak{a}(x,-\xi)\right|+\left|\partial^{\delta} \mathfrak{c}(x,-\xi)\right|\right) \\
& \lesssim \sum_{k=0}^{M} \sum_{(\gamma, \delta) \in \Omega_{k}}\left(|\mathfrak{a}(x,-\xi)|\langle(x,-\xi)\rangle^{-\rho(2 k+|\alpha+\beta|)}+\omega(x,-\xi)\langle(x,-\xi)\rangle^{-\rho(2 M+|\alpha+\beta|)}\right) \\
& =|\mathfrak{a}(x,-\xi)|\langle(x,-\xi)\rangle^{-\rho|\alpha+\beta|}+\omega(x,-\xi)\langle(x,-\xi)\rangle^{-\rho(2 M+|\alpha+\beta|)} \\
& \lesssim|\mathfrak{a}(x,-\xi)|\langle(x,-\xi)\rangle^{-\rho|\alpha+\beta|}+|\mathfrak{a}(x,-\xi)|\langle(x,-\xi)\rangle^{\rho_{0}-\rho(2 M+|\alpha+\beta|)} \\
& \simeq|\mathfrak{a}(x,-\xi)|\langle(x,-\xi)\rangle^{-\rho|\alpha+\beta|} \lesssim\left|a\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right)\right|\langle(x,-\xi)\rangle^{-\rho|\alpha+\beta|},
\end{aligned}
$$

which implies that (5.4) holds.
This shows that $a$ is hypoelliptic of order $\rho_{0}$ in $\mathcal{\mathcal { A }}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$ when $\mathfrak{a}$ is hypoelliptic of order $\rho_{0}$ in $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$.
Suppose instead that $a$ is hypoelliptic of order $\rho_{0}$ in $\mathcal{A}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$. By using Proposition 3.5, (3.12) and (3.13) instead of Proposition 5.4, similar computations as in the first part of the proof shows that (5.6) and (5.7) hold for some $R>0$. This shows that $\mathfrak{a}$ is hypoelliptic of order $\rho_{0}$ in $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ when $a$ is hypoelliptic of order $\rho_{0}$ in $\widehat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}\left(\mathbf{C}^{2 d}\right)$, and the result follows.

### 5.3. Ellipticity in the case of polynomial symbols

Next we discuss ellipticity for polynomial symbols, i.e.

$$
\begin{equation*}
\mathfrak{a}(x, \xi)=\sum_{|\alpha+\beta| \leqslant M} c(\alpha, \beta) x^{\alpha} \xi^{\beta}, \quad x, \xi \in \mathbf{R}^{d}, \tag{5.9}
\end{equation*}
$$

and

$$
\begin{equation*}
a(z, w)=\sum_{|\alpha+\beta| \leqslant M} c(\alpha, \beta) z^{\alpha} \bar{w}^{\beta}, \quad z, w \in \mathbf{C}^{d} \tag{5.10}
\end{equation*}
$$

The corresponding principal symbols are

$$
\begin{equation*}
\mathfrak{a}_{p}(x, \xi)=\sum_{|\alpha+\beta|=M} c(\alpha, \beta) x^{\alpha} \xi^{\beta}, \quad x, \xi \in \mathbf{R}^{d} \tag{5.11}
\end{equation*}
$$

and

$$
\begin{equation*}
a_{p}(z, w)=\sum_{|\alpha+\beta|=M} c(\alpha, \beta) z^{\alpha} \bar{w}^{\beta}, \quad z, w \in \mathbf{C}^{d}, \tag{5.12}
\end{equation*}
$$

respectively.
First we relate polynomials on $\mathbf{R}^{2 d}$ to Shubin classes.

Proposition 5.7. Let $\mathfrak{a}$ and $\mathfrak{a}_{p}$ be as in (5.9) and (5.11) for some $c(\alpha, \beta) \in \mathbf{C}, \alpha, \beta \in \mathbf{N}^{d}$ and $M \geqslant 0$, and let $\omega_{M}(x, \xi)=\langle(x, \xi)\rangle^{M}, x, \xi \in \mathbf{R}^{d}$. Then the following is true:
(1) $\mathfrak{a} \in \operatorname{Sh}_{1}^{\left(\omega_{M}\right)}\left(\mathbf{R}^{2 d}\right)$;
(2) $\mathfrak{a}$ is elliptic with respect to $\omega_{M}$, if and only if $\mathfrak{a}_{p}(x, \xi) \neq 0$ when $(x, \xi) \neq 0$.

The result can be considered folklore. In order to be self-contained we present the arguments.
Proof. First we prove (1). Let $t=\max \left(\left|x_{1}\right|, \ldots,\left|x_{d}\right|,\left|\xi_{1}\right|, \ldots,\left|\xi_{d}\right|\right)$ when $x=\left(x_{1}, \ldots, x_{d}\right) \in \mathbf{R}^{d}$ and $\xi=$ $\left(\xi_{1}, \ldots, \xi_{d}\right) \in \mathbf{R}^{d}$. Then

$$
|\mathfrak{a}(x, \xi)| \leqslant \sum_{|\alpha+\beta| \leqslant M}|c(\alpha, \beta)| t^{|\alpha+\beta|} \lesssim 1+t^{M} \leqslant\langle(x, \xi)\rangle^{M}
$$

which gives the desired bound for $|\mathfrak{a}(x, \xi)|$. Since the degree of a polynomial is lowered by at least one for every differentiation we get

$$
\left|\partial^{\alpha} \mathfrak{a}(x, \xi)\right| \lesssim\langle(x, \xi)\rangle^{M-|\alpha|}
$$

for every $\alpha \in \mathbf{N}^{2 d}$, which gives (1).
In order to prove (2) we let $\mathfrak{a}_{p}$ be as in (5.11). First suppose that $\mathfrak{a}_{p}(x, \xi) \neq 0$ when $(x, \xi) \neq(0,0)$, and let $g$ be the continuous function on $\mathbf{R}^{2 d} \backslash 0$ given by

$$
g(x, \xi)=\frac{\left|\mathfrak{a}_{p}(x, \xi)\right|}{|(x, \xi)|^{M}}, \quad(x, \xi) \neq(0,0)
$$

Since $g$ is continuous and positive, and the sphere

$$
\mathbf{S}^{2 d-1}=\left\{(x, \xi) \in \mathbf{R}^{2 d} ;|x|^{2}+|\xi|^{2}=1\right\}
$$

is compact, it follows that there are constants $c_{1}, c_{2}>0$ such that

$$
c_{1} \leqslant g(x, \xi) \leqslant c_{2}, \quad(x, \xi) \in \mathbf{S}^{2 d-1}
$$

By homogeneity it now follows

$$
c_{1}|(x, \xi)|^{M} \leqslant\left|\mathfrak{a}_{p}(x, \xi)\right| \leqslant c_{2}|(x, \xi)|^{M}, \quad x, \xi \in \mathbf{R}^{d} .
$$

Hence, if

$$
\mathfrak{b}(x, \xi)=\mathfrak{a}(x, \xi)-\mathfrak{a}_{p}(x, \xi)=\sum_{|\alpha+\beta| \leqslant M-1} c(\alpha, \beta) x^{\alpha} \xi^{\beta},
$$

then the first part of the proof implies that for some constants $C>0$ and $R>0$ we have

$$
|\mathfrak{a}(x, \xi)| \geqslant\left|\mathfrak{a}_{p}(x, \xi)\right|-|\mathfrak{b}(x, \xi)| \geqslant c_{1}|(x, \xi)|^{M}-C\langle(x, \xi)\rangle^{M-1} \gtrsim\langle(x, \xi)\rangle^{M}
$$

when $|(x, \xi)| \geqslant R$. Hence $\mathfrak{a}$ is elliptic with respect to $\omega_{M}$.
Suppose instead $\mathfrak{a}_{p}\left(x_{0}, \xi_{0}\right)=0$ for some $\left(x_{0}, \xi_{0}\right) \neq(0,0)$. For any $(x, \xi)=\left(t x_{0}, t \xi_{0}\right)$ we have

$$
|\mathfrak{a}(x, \xi)| \leqslant\left|\mathfrak{a}_{p}(x, \xi)\right|+|\mathfrak{b}(x, \xi)|=\left|t^{M} \mathfrak{a}_{p}\left(x_{0}, \xi_{0}\right)\right|+|\mathfrak{b}(x, \xi)|=|\mathfrak{b}(x, \xi)| \lesssim\langle(x, \xi)\rangle^{M-1}
$$

giving that $|\mathfrak{a}(x, \xi)| \gtrsim\langle(x, \xi)\rangle^{M},|(x, \xi)| \geqslant R$, cannot hold for any $R>0$.
By Theorems 5.2, 5.6 and Proposition 5.7 we get the following. The verification is left for the reader.
Proposition 5.8. Let $a$ and $a_{p}$ be as in (5.10) and (5.12) for some $c(\alpha, \beta) \in \mathbf{C}, \alpha, \beta \in \mathbf{N}^{d}$ and $M \geqslant 0$, and let $\omega_{M}(x, \xi)=\langle(x, \xi)\rangle^{N}, x, \xi \in \mathbf{R}^{d}$. Then the following is true:
(1) $a \in \mathcal{A}_{\mathrm{Sh}, 1}^{\left(\omega_{M}\right)}\left(\mathbf{C}^{2 d}\right)$;
(2) a is elliptic in $\mathcal{A}_{\mathrm{Sh}, 1}^{\left(\omega_{M}\right)}\left(\mathbf{C}^{2 d}\right)$ if and only if $a_{p}(z, z) \neq 0$ when $z \neq 0$.

Remark 5.9. Let $\mathfrak{a}, \mathfrak{a}_{p}, a$ and $a_{p}$ be as in (5.9)-(5.12). Then it follows from Proposition 5.7 and Proposition 5.8 that $\mathfrak{a}$ is elliptic, if and only if $\mathfrak{a}_{p}$ is elliptic, and that $a$ is elliptic, if and only if $a_{p}$ is elliptic.

We have now the following.
Theorem 5.10. Let $\mathfrak{a} \in \operatorname{Sh}_{1}^{\left(\omega_{M}\right)}\left(\mathbf{R}^{2 d}\right)$ and $\mathfrak{a}_{p}$ be as in (5.9) and (5.11) for some $c(\alpha, \beta) \in \mathbf{C}, \alpha, \beta \in \mathbf{N}^{d}$ and $M \geqslant 0$. Then the following is true:
(1) the principal symbol $a_{p}(z, w)$ of $\mathrm{S}_{\mathfrak{V} \mathfrak{a}}$ is given by

$$
\begin{equation*}
a_{p}(z, w)=2^{-\frac{M}{2}} \sum_{|\alpha+\beta|=M} c(\alpha, \beta) i^{|\beta|}(z+\bar{w})^{\alpha}(z-\bar{w})^{\beta} ; \tag{5.13}
\end{equation*}
$$

(2) $\mathfrak{a}$ is elliptic in $\mathrm{Sh}_{1}^{\left(\omega_{M}\right)}\left(\mathbf{R}^{2 d}\right)$ if and only if $a_{p}$ is elliptic in $\mathcal{A}_{\mathrm{Sh}, 1}^{\left(\omega_{M}\right)}\left(\mathbf{C}^{2 d}\right)$;
(3) $\mathfrak{a}_{p}(x, \xi)>0$ for every $(x, \xi) \neq(0,0)$, if and only if $a_{p}(z, z)>0$ for every $z \neq 0$.

Proof. Let $z=x+i \xi, x, \xi \in \mathbf{R}^{d}$, i.e. $x=\frac{1}{2}(z+\bar{z})$ and $\xi=\frac{1}{2 i}(z-\bar{z})$. By Theorem 5.2 we get

$$
\begin{equation*}
a_{p}\left(2^{-\frac{1}{2}} z, 2^{-\frac{1}{2}} z\right)=\mathfrak{a}_{p}(x,-\xi) \tag{5.14}
\end{equation*}
$$

This implies

$$
a_{p}(z, z)=2^{\frac{M}{2}} \sum_{|\alpha+\beta|=M} c(\alpha, \beta) x^{\alpha}(-\xi)^{\beta}=2^{\frac{M}{2}} \sum_{|\alpha+\beta|=M} c(\alpha, \beta) 2^{-|\alpha|}(z+\bar{z})^{\alpha}(2 i)^{-|\beta|}(-(z-\bar{z}))^{\beta},
$$

which gives

$$
\begin{equation*}
a_{p}(z, z)=2^{-\frac{M}{2}} \sum_{|\alpha+\beta|=M} c(\alpha, \beta) i^{|\beta|}(z+\bar{z})^{\alpha}(z-\bar{z})^{\beta} . \tag{5.13}
\end{equation*}
$$

The formula (5.13) now follows from (5.13)' and analytic continuation, using the fact that $a_{p}(z, w)$ is analytic in $z$ and conjugate analytic in $w$.

The assertion (2) follows by a combination of Corollary 5.3, Propositions 5.7 and 5.8, and the assertion (3) is a direct consequence of (5.14).

## 6. A necessary condition for polynomially bounded Wick symbols

In [12, Section 2.7] Folland shows that polynomial symbols for pseudo-differential operators correspond to polynomial Wick and anti-Wick symbols. Thus partial differential operators with polynomial coefficients corresponds to polynomial Wick symbols.

Here we show that a Wick symbol that is polynomially bounded must be a polynomial. This gives a characterization of Wick symbols corresponding to polynomial symbols for pseudo-differential operators.

Cauchy's integral formula implies that an entire function which is polynomially bounded must be a polynomial:

Proposition 6.1. Let $F \in A\left(\mathbf{C}^{d}\right)$ with Maclaurin series

$$
F(z)=\sum_{\alpha \in \mathbf{N}^{d}} c(\alpha) e_{\alpha}(z), \quad z \in \mathbf{C}^{d}
$$

Suppose that for some $j \in\{1, \ldots, d\}, C>0, M \geqslant 0$, and an open neighborhood $I \subseteq \mathbf{C}$ of the origin it holds

$$
|F(z)| \leqslant C\left\langle z_{j}\right\rangle^{M}, \quad z_{j} \in \mathbf{C},
$$

provided $z_{k} \in I, k \in\{1, \ldots, d\} \backslash\{j\}$. Then $c(\alpha)=0$ when $\alpha_{j}>M$.
Proof. By interchanging the variables, we may assume that $j=d$. Let $R \geqslant 1$ and $\varepsilon>0$ be chosen such that

$$
D_{\varepsilon} \equiv\left\{z_{0} \in \mathbf{C} ;\left|z_{0}\right| \leqslant \varepsilon\right\} \subseteq I .
$$

Take $\alpha \in \mathbf{N}^{d}$ such that $\alpha_{d}>M$, let $\beta=\left(\alpha_{1}+1, \ldots, \alpha_{d}+1\right) \in \mathbf{N}^{d}$ and $\gamma_{\varepsilon} \subseteq \mathbf{C}$ be the boundary circle of $D_{\varepsilon}$. Then Cauchy's integral formula gives

$$
\begin{aligned}
\frac{|c(\alpha)|}{\alpha!^{\frac{1}{2}}}=\left|\frac{\partial^{\alpha} F(0)}{\alpha!}\right| & =(2 \pi)^{-d}\left|\int_{\gamma_{\varepsilon}^{d-1}} \ldots \int\left(\int_{\left|z_{d}\right|=R} \frac{F(z)}{z^{\beta}} d z_{d}\right) d z_{1} \cdots d z_{d-1}\right| \\
& \leqslant(2 \pi)^{-d} \int_{\gamma_{\varepsilon}^{d-1}} \ldots \int\left(\int_{\left|z_{d}\right|=R} \frac{|F(z)|}{\left|z^{\beta}\right|}\left|d z_{d}\right|\right)\left|d z_{1}\right| \cdots\left|d z_{d-1}\right| \\
& \lesssim R^{-\alpha_{d}\langle R\rangle^{M} \varepsilon^{-\left(\alpha_{1}+\cdots+\alpha_{d-1}\right)} \rightarrow 0}
\end{aligned}
$$

as $R \rightarrow \infty$.

Corollary 6.2. Let $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ and suppose

$$
\begin{equation*}
|a(z, w)| \lesssim\langle(z, w)\rangle^{M} \tag{6.1}
\end{equation*}
$$

for some $M \geqslant 0$. Then $a$ is a polynomial in $z \in \mathbf{C}^{d}$ and $\bar{w} \in \mathbf{C}^{d}$ of degree at most $M$.
Proof. By Proposition 6.1 it follows that $a$ is a polynomial of degree at most $2 d M$. We need to prove that the degree is at most $M$. In order to do this we may assume that $a$ has degree at least one.

For some integer $L \geqslant 1$ we have

$$
a(z, w)=a_{L}(z, w)+a_{L-1}(z, w)
$$

where

$$
a_{L}(z, w)=\sum_{|\alpha+\beta|=L} c(\alpha, \beta) z^{\alpha} \bar{w}^{\beta}
$$

is non-trivial and

$$
a_{L-1}(z, w)=\sum_{|\alpha+\beta| \leqslant L-1} c(\alpha, \beta) z^{\alpha} \bar{w}^{\beta} .
$$

Since $a_{L}$ is non-trivial, there are $z_{0}, w_{0} \in \mathbf{C}^{d}$ such that $\left|z_{0}\right|^{2}+\left|w_{0}\right|^{2}=1$ and $\left|a_{L}\left(z_{0}, w_{0}\right)\right|=c_{0} \neq 0$. By homogeneity we get

$$
\left|a_{L}\left(t z_{0}, t w_{0}\right)\right|=c_{0}|t|^{L}, \quad t \in \mathbf{R} .
$$

In the same way we get

$$
\left|a_{L-1}\left(t z_{0}, t w_{0}\right)\right| \leqslant C(1+|t|)^{L-1}, \quad t \in \mathbf{R}
$$

for some constant $C$ which is independent of $t$.
Suppose contrary to the assertion that $L>M$. For $t \in \mathbf{R}$ with $|t| \geqslant 1$ we have

$$
\begin{aligned}
\left|\frac{a\left(t z_{0}, t w_{0}\right)}{\left\langle\left(t z_{0}, t w_{0}\right)\right\rangle^{M}}\right| & \gtrsim|t|^{-M}\left(\left|a_{L}\left(t z_{0}, t w_{0}\right)\right|-\left|a_{L-1}\left(t z_{0}, t w_{0}\right)\right|\right) \\
& \geqslant|t|^{-M}\left(c_{0}|t|^{L}-C(1+|t|)^{L-1}\right) \rightarrow \infty \quad \text { as } \quad|t| \rightarrow \infty .
\end{aligned}
$$

This contradicts (6.1), and the hence our assumption that $L>M$ must be false.

## Appendix A

In this appendix we present some tables on weights, operators, spaces of entire functions on $\mathbf{C}^{d}$ and Wick symbol classes.

In the first two tables, Tables 1 and 2 we review weight classes, transforms, and operators. The next two tables, Tables 3 and 4, deal with properties of the Bargmann images of Gelfand-Shilov function spaces, the Schwartz space, and their distribution spaces. Then, in Table 5 we explain some links between the Shubin class $\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)$ (see (1.47)), and the symbol classes $\Gamma_{s, \sigma}^{\sigma, s ; 0}\left(\mathbf{R}^{2 d}\right), \Gamma_{s, \sigma ; 0}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$ and $\Gamma_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)$ (see Definition 1.8).

Table 1
Weight classes.

| Weight class | Features | Eq. ref. |
| :--- | :--- | :--- |
| $\mathscr{P}_{E}\left(\mathbf{R}^{d}\right)$ | $\omega \in L_{l o c}^{\infty}\left(\mathbf{R}^{d} ; \mathbf{R}_{+}\right), \omega(x+y) \lesssim \omega(x) e^{r\|y\|}$ | $(1.5)$ |
| $\mathscr{P}\left(\mathbf{R}^{d}\right)$ | $\omega \in L_{l o c}^{\infty}\left(\mathbf{R}^{d} ; \mathbf{R}_{+}\right), \omega(x+y) \lesssim \omega(x)\langle y\rangle^{N}$ | $(1.1)^{\prime}$ |
| $\mathscr{P}_{\text {Sh }, \rho}\left(\mathbf{R}^{d}\right)$ | $\omega \in \mathscr{P}\left(\mathbf{R}^{d}\right),\left\|\partial^{\alpha} \omega(x)\right\| \lesssim \omega(x)\langle x\rangle^{-\rho\|\alpha\|}$ | $(1.46)$ |

Table 2
Operators and transforms.

| Operator | Notation | Features | Eq. ref. |
| :--- | :--- | :--- | ---: |
| Modul. STFT | $\mathcal{T}$ | $f \mapsto(2 \pi)^{-\frac{d}{2}} \int f(y+x) \overline{\phi(y)} e^{-i\langle y, \xi\rangle} d y$ | $(1.10)^{\prime}$ |
| Bargm. transf. | $\mathfrak{V}_{d}$ | $f \mapsto \pi^{-\frac{d}{4}} \int e^{-\frac{1}{2}\left(\langle z, z\rangle+\|y\|^{2}\right)+2^{1 / 2}\langle z, y\rangle} f(y) d y$ | $(1.19)$ |
| Semi-conj. op. | $\Theta$ | $K(z, w) \mapsto K(z, \bar{w})$ | $(1.29)$ |
| Pseudo-diff. op. | $\mathrm{Op}_{A}(\mathfrak{a})$ | $f \mapsto(2 \pi)^{-d} \iint \mathfrak{a}(x-A(x-y), \xi) f(y) e^{i\langle x-y, \xi\rangle} d y d \xi$ | $(1.31)$ |
| Wick op. | $\mathrm{Op}_{\mathfrak{V}}(a)$ | $F \mapsto \int_{\mathbf{C}^{d}} a(z, w) F(w) e^{(z, w)} d \mu(w)$ | $(1.36)$ |
| Anti-Wick op. | $\mathrm{Op}_{\mathfrak{V}}^{\text {aw }}(a)$ | $F \mapsto \int_{\mathbf{C}^{d}} a(w) F(w) e^{(z, w)} d \mu(w)$ | $(1.40)$ |
| Bargm. assignm. | $\mathrm{S}_{\mathfrak{V}}$ | $\mathrm{Op}_{\mathfrak{V}}\left(\mathrm{S}_{\mathfrak{V} \mathfrak{a}} \mathfrak{a}\right)=\mathfrak{V}_{d} \circ \mathrm{Op}_{A}(\mathfrak{a}) \circ \mathfrak{V}_{d}^{*}, A=\frac{1}{2} I$ | $(1.41)$ |

Table 3
The Bargmann images of test function spaces.

| Function space | Bargmann image | $\left\|\mathfrak{V}_{d} f(z)\right\| \lesssim$ | Eq. ref. |
| :--- | :--- | :--- | :--- |
| $\mathcal{S}_{s}^{\sigma}\left(\mathbf{R}^{d}\right), s, \sigma \geqslant \frac{1}{2}$ | $\mathcal{A}_{s}^{\sigma}\left(\mathbf{C}^{d}\right)$ | $e^{\frac{\|z\|^{2}}{2}-r\|z\|_{s, \sigma}}, \exists r>0$ | $(1.6),(1.24),(1.25)$ |
| $\Sigma_{s}^{\sigma}\left(\mathbf{R}^{d}\right), s, \sigma>\frac{1}{2}$ | $\mathcal{A}_{0, s}^{\sigma}\left(\mathbf{C}^{d}\right)$ | $e^{\frac{\|z\|^{2}}{2}-r\|z\|_{s, \sigma}}, \forall r>0$ | $(1.6),(1.24),(1.25)$ |
| $\mathscr{S}\left(\mathbf{R}^{d}\right)$ | $\mathcal{A}_{\mathscr{S}}\left(\mathbf{C}^{d}\right)$ | $e^{\frac{\|z\|^{2}}{2}}\langle z\rangle^{-N}, \forall N \geqslant 0$ | $(1.6),(1.24),(1.25)$ |

Table 4
The Bargmann images of distribution spaces.

| Distribution space | Bargmann image | $\left\|\mathfrak{V}_{d} f(z)\right\| \lesssim$ | Eq. ref. |
| :--- | :--- | :--- | :--- |
| $\left(\mathcal{S}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right), s, \sigma \geqslant \frac{1}{2}$ | $\left(\mathcal{A}_{s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right)$ | $e^{\frac{\|z\|^{2}}{2}+r\|z\|_{s, \sigma}}, \forall r>0$ | $(1.6),(1.24),(1.25)$ |
| $\left(\Sigma_{s}^{\sigma}\right)^{\prime}\left(\mathbf{R}^{d}\right), s, \sigma>\frac{1}{2}$ | $\left(\mathcal{A}_{0, s}^{\sigma}\right)^{\prime}\left(\mathbf{C}^{d}\right)$ | $e^{\frac{\|z\|^{2}}{2}+r\|z\|_{s, \sigma}, \exists r>0}$ | $(1.6),(1.24),(1.25)$ |
| $\mathscr{S}^{\prime}\left(\mathbf{R}^{d}\right)$ | $e^{\frac{\|z\|^{2}}{2}}\langle\boldsymbol{z}\rangle^{N}, \exists N \geqslant 0$ | $(1.6),(1.24),(1.25)$ |  |

Table 5
Estimates for Wick symbol classes.

| Wick class | $\left\|\partial^{\alpha} a(z, w)\right\| \lesssim$ | Ref. |
| :--- | :--- | :--- |
| $\mathrm{S}_{\mathfrak{V}}\left(\operatorname{Sh}_{0}^{(\omega)}\left(\mathbf{R}^{2 d}\right)\right)$ | $e^{\frac{1}{2}\|z-w\|^{2}} \omega(\sqrt{2} \bar{z})\langle z-w\rangle^{-N} \forall N \geqslant 0, \alpha=0$ | Theorem 2.2 |
| $\mathrm{S}_{\mathfrak{V}}\left(\operatorname{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)\right)$ | $e^{\frac{1}{2}\|z-w\|^{2}} \omega_{\rho\|\alpha\|}(\sqrt{2} \bar{z})\langle z-w\rangle^{-N}, \alpha \in \mathbf{N}^{4 d}$ | Theorem 2.5 |
| $\mathrm{S}_{\mathfrak{V}}\left(\Gamma_{s, \sigma ; 0}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)\right)$ | $e^{\frac{1}{2}\|z-w\|^{2}+r_{1}\|z+w\|_{s, \sigma}-r_{2}\|z-w\|_{s, \sigma}}, \exists r_{2}>0, \forall r_{1}>0, \alpha=0$ | Theorem 2.6 |
| $\mathrm{S}_{\mathfrak{V}}\left(\Gamma_{s, \sigma}^{\sigma, s ; 0}\left(\mathbf{R}^{2 d}\right)\right)$ | $e^{\frac{1}{2}\|z-w\|^{2}+r_{1}\|z+w\|_{s, \sigma}-r_{2}\|z-w\|_{s, \sigma}}, \exists r_{1}>0, \forall r_{2}>0, \alpha=0$ | Theorem 2.6 |
| $\mathrm{S}_{\mathfrak{V}}\left(\Gamma_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)\right)$ | $e^{\frac{1}{2}\|z-w\|^{2}+r_{1}\|z+w\|_{s, \sigma}-r_{2}\|z-w\|_{s, \sigma}}, \exists r_{1}, r_{2}>0, \alpha=0$ | Theorem 2.6 |

## Here recall

$$
\langle x\rangle=\left(1+|x|^{2}\right)^{\frac{1}{2}}, \quad d \mu(w)=\pi^{-d} e^{-|w|^{2}} d \lambda(w)
$$

and

$$
|z|_{s, \sigma}=|\operatorname{Re} z|^{\frac{1}{s}}+|\operatorname{Im} z|^{\frac{1}{\sigma}}, \quad z \in \mathbf{C}^{d}
$$

where $d \lambda(w)$ is the Lebesgue measure on $\mathbf{C}^{d}$.
In Table 5 we let

$$
\begin{gathered}
\partial^{\alpha} a=\partial_{x}^{\alpha_{1}} \partial_{\xi}^{\alpha_{3}} \partial_{y}^{\alpha_{2}} \partial_{\eta}^{\alpha_{4}} a \\
z=x+i \xi, w=y+i \eta, \alpha=\left(\alpha_{1}, \alpha_{2}, \alpha_{3}, \alpha_{4}\right) \in \mathbf{N}^{2 d}
\end{gathered}
$$

when $a \in \widehat{A}\left(\mathbf{C}^{2 d}\right)$ (see (2.15)), and $\omega_{r}(z)=\omega(z)\langle z\rangle^{-r}$ when $\omega \in \mathscr{P}_{\mathrm{Sh}, \rho}\left(\mathbf{C}^{d}\right)$ and $r \in \mathbf{R}$, and recall that $\widehat{\mathcal{A}}_{\mathrm{Sh}, \rho}^{(\omega)}=\mathrm{S}_{\mathfrak{V}}\left(\mathrm{Sh}_{\rho}^{(\omega)}\left(\mathbf{R}^{2 d}\right)\right)$. For $\mathrm{S}_{\mathfrak{V}}\left(\Gamma_{s, \sigma ; 0}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)\right)$ it is assumed that $s, \sigma \geqslant \frac{1}{2}$, while for $\mathrm{S}_{\mathfrak{V}}\left(\Gamma_{s, \sigma}^{\sigma, s ; 0}\left(\mathbf{R}^{2 d}\right)\right)$ and $\mathrm{S}_{\mathfrak{V}}\left(\Gamma_{s, \sigma}^{\sigma, s}\left(\mathbf{R}^{2 d}\right)\right)$ it is assumed that $s, \sigma>\frac{1}{2}$.

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