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# Superconductive and hybrid shielding design: a 3D-modeling study

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**Abstract**—Superconducting materials are promising for fabrication of high-performance magnetic shields. For this purpose, a computational approach is crucial to guide the shielding devices’ design. In this work, we exploit a 3D modeling procedure based on a vector potential formulation to calculate the shielding properties of cup-shaped superconductors (SC) with and without the superimposition of a ferromagnetic (FM) shield. Relying on the past validation of the model, the efficiency of different shielding solutions was investigated in both axial-field (AF) and transverse-field (TF) configurations.

**Keywords**—magnetic shielding, 3D numerical modelling, bulk superconductor

## I. INTRODUCTION

Magnetic shielding is essential for several applications. Its goal is twofold, aiming at the reduction of stray field and electromagnetic noise, and at the achievement of low magnetic field background. Superconductors are key materials for shielding quasi-static magnetic fields [1] - [3]. Recently tested solutions included active [4] and passive layouts [1], [5], [6]. Focusing on the latter, a successful approach to improve shielding ability have combined growth technique able to manufacture properly shaped products with high and homogeneous critical current density, and suitable simulation tools, guiding the whole optimization process [7], [8].

Furthermore, improvements in superconductors (SC) shielding performance have been proved by superimposing a sheet of ferromagnetic (FM) material [9] - [11]. In order to investigate the non-trivial action of adding a FM sheet around the SC [12] and to optimize the shield design under realistic working conditions [13], i.e. for different orientations of applied magnetic field, a 3D modelling is needed.

In this work, we applied the numerical procedure based on the 3D magnetic vector-potential ( $\mathbf{A}$ ) formulation described in [14] to predict the shielding properties of a cup-shaped  $\text{MgB}_2$  bulks with small height/outer radius aspect ratio, studying the effect of the superimposition of different iron shields in both axial-field (AF) and transverse-field (TF) configurations.

## II. MODELLING

To model superconductor shielding properties, a numerical procedure based on the magnetic vector potential formalism

[14] was applied by a finite element calculation using COMSOL Multiphysics® [15]. As presented in [16], ensuring the collinearity between the current density and the local electric field and considering an isotropic  $J_c$ , the following equation can be introduced to relate the current density to the electric field, and, therefore to the time derivative of  $\mathbf{A}$ .

$$\mathbf{J} = \frac{J_c}{|E|} \left( |E_x| \tanh\left(\frac{E_x}{E_0}\right) \mathbf{i} + |E_y| \tanh\left(\frac{E_y}{E_0}\right) \mathbf{j} + |E_z| \tanh\left(\frac{E_z}{E_0}\right) \mathbf{k} \right) \quad (1)$$

where  $E_0$  is the conventional criterion taken usually as  $10^{-4} \text{Vm}^{-1}$  [14].

Based on a previous validation of this model to reproduce shielding experiment outputs [17], a  $J_c$  dependence on the induction field need also to be taken into account. A previous experiment [18] evidenced that the experimental  $J_c$  curves (obtained from magnetic induction cycles measured at different temperatures) were successfully fitted by the exponential relation  $J_c(B) = J_{c,0} \exp[-(B/B_0)^\gamma]$ , where  $J_{c,0} = 5.02 \times 10^8 \text{A/m}^2$ ,  $B_0 = 0.98 \text{T}$  and  $\gamma = 3.78$  were the constant parameters at  $T = 30 \text{K}$ . The same parameters were used to model the superconducting shield in all the computations reported in the following.

The magnetic properties of the FM material were defined by the interpolation of the magnetic induction versus applied field curve measured experimentally on a small piece of the soft iron already used to fabricate magnetic shield [11].

The source term for the applied magnetic field,  $\mathbf{H}_{\text{appl}}$ , was considered through the boundary conditions: at a large distance from the shield(s),  $\mathbf{B}$  was set to be equal to  $\mu_0 \mathbf{H}_{\text{appl}}$ . The applied field was always assumed increasing monotonically.

## III. NUMERICAL RESULTS

The model validation presented in [17] allows us to implement the computational procedure in order to explore new hybrid (SC+FM) configurations. Based on previous study outcomes [12], [19], the performances of three different cup-shaped arrangements (Fig.1) were investigated in both AF and TF configurations. The data are then compared with those obtained computing the shielding properties of the SC cup

only, whose dimensions are the same as for the MgB<sub>2</sub> cup studied in [18].

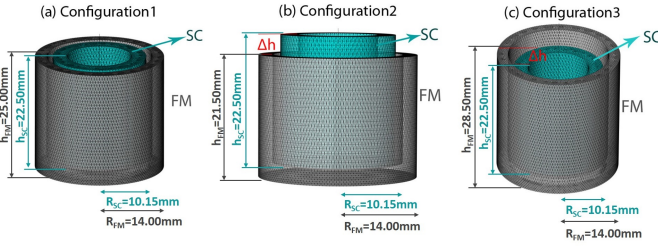


Fig. 1. Design of cup-shaped hybrid (SC+FM) configurations. SC lateral thickness: 3.15 mm, base thickness: 4.20 mm; FM edge thickness: 2.50 mm, base thickness: 2.50 mm;  $\Delta h=3.50$  mm.

For each configuration, the induction field curves along  $z$ -axis and  $y$ -axis (shown in the insets of Fig.2 and Fig.3, respectively) were computed and the shielding efficiency was estimated using the shielding factor, calculated as the ratio between the applied magnetic field and corresponding component of the local magnetic induction ( $SF = \mu_0 H_{app} / B_{z(y)}$ ). Fig.2 and Fig.3 show the SF behaviors obtained at 5 points along the cup axis in both in AF and TF mode. As one can see, the ferromagnetic sheet addition actually improves the shielding performances.

In particular, for AF configuration and with only SC solution,  $SF > 10^2$  and  $SF > 10^4$  can be reached up to  $\mu_0 H_{app} \approx 1.0$  T in the inner half of the shield and near the closed extremity, respectively (symbols in Fig.2). Adding FM sheet enlarges the high-shielding region up to a field of  $\approx 1.2$  T, even though at lower fields higher values of SF can not be gained (Fig.2).

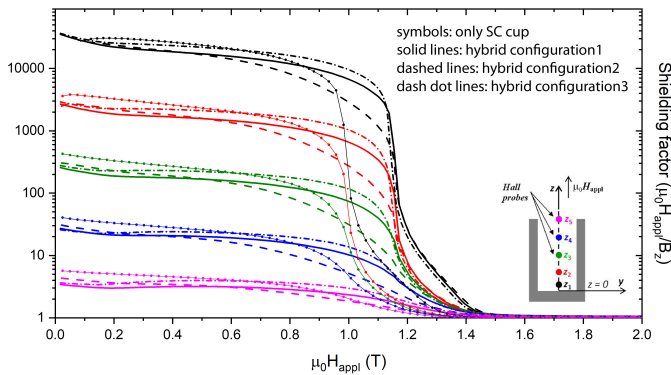


Fig. 2. Comparison between the SFs of configuration 1, 2 and 3 in AF mode considering for the superconductors the same  $J_c(B)$  dependence as that measured in [18] at  $T = 30$  K. Assuming  $(x,y,z) = (0,0,0)$  the coordinate of the shield center, the plotted curves were calculated at  $z_1 = 1$  mm,  $z_2 = 5$  mm,  $z_3 = 9.2$  mm,  $z_4 = 13.7$  mm and  $z_5 = 18.3$  mm.

On the other hand, greater enhancements are achieved in TF configuration. The improvements due to the FM addition are twofold. Firstly, likewise in AF geometry, it allows extending the zone where higher SF values are achieved up to above 1.0 T. Secondly, for low values of applied magnetic field, near the closest extremity of the shield, the screening performances

improve from a  $SF \approx 50$  with only SC (symbols Fig.3) to a  $SF \approx 120$  with configuration 2 and to a  $SF \approx 140$  with configuration 3 (Fig.3).

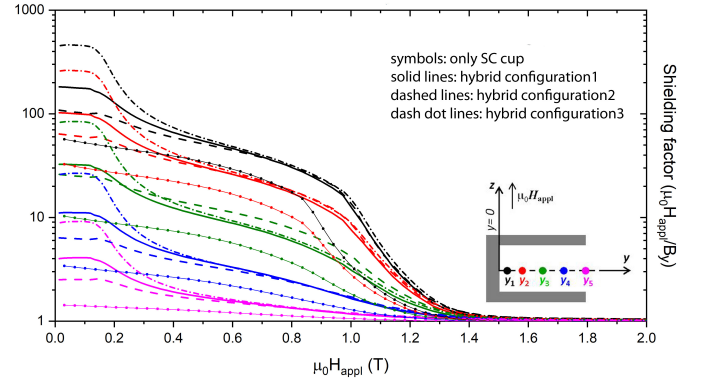


Fig. 3. Comparison between the SFs of configuration 1, 2 and 3 in TF mode considering for the superconductors the same  $J_c(B)$  dependence as that measured in [18] at  $T = 30$  K. Assuming  $(x,y,z) = (0,0,0)$  the coordinate of the shield center, the plotted curves were calculated at  $y_1 = 1$  mm,  $y_2 = 4.6$  mm,  $y_3 = 9.2$  mm,  $y_4 = 13.7$  mm and  $y_5 = 18.3$  mm.

These results show how the hybrid system can greatly improve the shielding performances obtained with only superconductor system and evidence how 3D modelling is crucial for achieving a complete information. To this end, the chosen modeling procedure was found to be a valuable tool to drive the design of future magnetic shields with small aspect ratio of height/lateral size and improved shielding performances for various orientations of the applied magnetic field.

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