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Confinement of dislocations inside a crystal with a prescribed external strain

Abstract. A system of n screw dislocations in an isotropic crystal undergoing antiplane shear is studied in the framework of linear elasticity. Imposing a suitable boundary condition for the strain, namely requesting the non-vanishing of its boundary integral, results in a confinement effect. More precisely, in the presence of an external strain with circulation equal to n times the lattice spacing, it is energetically convenient to have n distinct dislocations lying inside the crystal. The result is obtained by formulating the problem via the core radius approach and by studying the asymptotics as the core size vanishes. An iterative scheme is devised to prove the main result. This work sets the basis for studying the upscaling problem, i.e., the limit as $n \rightarrow \infty$, which is treated in [17].

Keywords. Dislocations, core radius approach, harmonic functions, divergence-measure fields

Mathematics Subject Classification (2010): 74E15 (35J25, 74B05, 49J40, 31A05).

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1 - Introduction

Starting with the pioneering work of Volterra [28], much attention has been drawn on dislocations in solids, as the ultimate cause of plasticity in crystalline materials [11, 20, 21, 26]. Dislocations are line defects in the lattice structure. The interest in dislocations became more and more evident as soon as it was understood that their presence can significantly influence the chemical and physical properties of the material. The measure of the lattice mismatch due to a dislocation is encoded in the *Burgers vector*, whose magnitude is of the order of one lattice spacing (see [16]). According to whether the Burgers vector is perpendicular or parallel to the dislocation line, ideal dislocations are classified as *edge dislocations* or *screw dislocations*, respectively. In nature, real dislocations come as a combination of these two types. For general treaties on dislocations, we refer the reader to [14, 16, 19].

In this paper we focus our attention on screw dislocations in a single isotropic crystal which occupies a cylindrical region $\Omega \times \mathbb{R}$ and which undergoes antiplane shear. According to the model proposed in [5] in the context of linearized elasticity, this allows us to study the problem in the cross section $\Omega \subset \mathbb{R}^2$. Throughout the work, we will assume that

(H1) Ω is a bounded convex open set with C^1 boundary.

We consider the lattice spacing of the material to be 2π and that all the Burgers vectors are oriented in the same direction. Therefore, every dislocation line is directed along the axis of the cylinder, is characterized by a Burgers vector of magnitude 2π along the same axis, and meets the cross section Ω at a single point. Moreover, we assume that an *external strain* acts on the crystal: we prescribe the tangential strain on $\partial\Omega$ to be of the form

$$(H2) \quad f \in L^1(\partial\Omega) \quad \text{with} \quad \int_{\partial\Omega} f(x) \, d\mathcal{H}^1(x) = 2\pi n$$

for some $n \in \mathbb{N}$. This choice of the external strain will determine at most n *distinct* dislocations inside Ω (see, e.g., [2, 23] for a comment on the topological necessity of the presence of exactly n defects; see also [9], where an evolution problem in the fractional laplacian setting is also studied), which we denote by $\mathbf{a} := (a_1, \dots, a_n) \in \Omega^n \setminus \Delta_n$, where $\Delta_n := \{\mathbf{x} = (x_1, \dots, x_n) \in \Omega^n : \text{there exist } i \neq j \text{ such that } x_i = x_j\}$.

Since the elastic energy associated with a defective material is infinite and has a logarithmic explosion in the vicinity of each dislocation a_i (see, e.g., [5, 19]), we resort to the so-called *core radius approach*, which consists in considering the energy far from the dislocations a_1, \dots, a_n . More precisely, given $\epsilon > 0$, we aim at studying the elastic energy in the perforated domain $\Omega_\epsilon(\mathbf{a}) := \Omega \setminus \bigcup_{i=1}^n \overline{B}_\epsilon(a_i)$. This approach is standard in the literature and it is employed in different contexts such as linear elasticity (see, for instance, [19, 25, 27]; also [4, 5, 22] for screw dislocations and [6] for edge dislocations), the theory of Ginzburg-Landau vortices (see, for instance, [2, 23] and the references therein), and liquid crystals (see, for instance, [12]). Since the core radius

approach eliminates the non-integrability of the strain field around the dislocations, classical variational techniques can be used. Therefore, we consider the energy

$$(1) \quad \mathcal{E}_\epsilon^{(n)}(\mathbf{a}) := \min \left\{ \frac{1}{2} \int_{\Omega_\epsilon(\mathbf{a})} |F|^2 dx : F \in \mathcal{X}_\epsilon(\mathbf{a}), F \cdot \tau = f \text{ on } \partial\Omega \setminus \bigcup_{i=1}^n \overline{B}_\epsilon(a_i) \right\},$$

where τ denotes the tangent unit vector to $\partial\Omega$ and

$$\mathcal{X}_\epsilon(\mathbf{a}) := \left\{ F \in L^2(\Omega_\epsilon(\mathbf{a}); \mathbb{R}^2) : \text{curl } F = 0 \text{ in } \mathcal{D}'(\Omega_\epsilon(\mathbf{a})), \langle F \cdot \tau, 1 \rangle_\gamma = 2\pi m \right\},$$

where $\mathcal{D}'(\Omega_\epsilon(\mathbf{a}))$ is the space of distributions on $\Omega_\epsilon(\mathbf{a})$ and γ is an arbitrary simple closed curve in $\Omega_\epsilon(\mathbf{a})$ winding once counterclockwise around m dislocations. Note that the boundary condition $F \cdot \tau$ must be intended in the sense of traces and that $\langle \cdot, \cdot \rangle_\gamma$ denotes the duality between $H^{-1/2}(\gamma)$ and $H^{1/2}(\gamma)$ (see [7]). Since these spaces are encapsulated, the energy \mathcal{E}_ϵ is monotone: if $0 < \epsilon < \eta$ then

$$(2) \quad \mathcal{E}_\epsilon^{(n)}(\mathbf{a}) \geq \mathcal{E}_\eta^{(n)}(\mathbf{a}).$$

If the a_i 's are all distinct and inside Ω , the energy (1) scales like $\pi n |\log \epsilon|$. This suggests to study the asymptotic behavior, as $\epsilon \rightarrow 0$, of the functionals $\mathcal{F}_\epsilon^{(n)} : \overline{\Omega}^n \rightarrow \mathbb{R} \cup \{+\infty\}$ defined by

$$(3) \quad \mathcal{F}_\epsilon^{(n)}(\mathbf{a}) := \mathcal{E}_\epsilon^{(n)}(\mathbf{a}) - \pi n |\log \epsilon|.$$

In this context, we say that the sequence of functionals $\mathcal{F}_\epsilon^{(n)}$ *continuously converges* in $\overline{\Omega}^n$ to $\mathcal{F}^{(n)}$ as $\epsilon \rightarrow 0$ if, for any sequence of points $\mathbf{a}^\epsilon \in \overline{\Omega}^n$ converging to $\mathbf{a} \in \overline{\Omega}^n$, the sequence (of real numbers) $\mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon)$ converges to $\mathcal{F}^{(n)}(\mathbf{a})$.

In order to write the limit functional, we introduce two objects: we take $g : \partial\Omega \rightarrow \mathbb{R}$ a primitive of f with n jump points $b_i \in \partial\Omega$ and jump amplitude 2π (see (10) for a precise definition), and for every $i \in \{1, \dots, n\}$ we set

$$(4) \quad d_i := \min_{\substack{j \in \{1, \dots, n\} \\ j \neq i}} \left\{ \frac{|a_i - a_j|}{2}, \text{dist}(a_i, \partial\Omega) \right\},$$

where $\text{dist}(\cdot, \partial\Omega)$ is the distance function from $\partial\Omega$.

Theorem 1.1. *Under the assumptions (H1) and (H2), as $\epsilon \rightarrow 0$ the functionals $\mathcal{F}_\epsilon^{(n)}$ defined by (3) continuously converge in $\overline{\Omega}^n$ to the functional $\mathcal{F}^{(n)} : \overline{\Omega}^n \rightarrow \mathbb{R} \cup \{+\infty\}$ defined as*

$$(5) \quad \begin{aligned} \mathcal{F}^{(n)}(\mathbf{a}) := & \sum_{i=1}^n \pi \log d_i + \frac{1}{2} \int_{\Omega} |\nabla v_{\mathbf{a}}|^2 dx + \sum_{i=1}^n \frac{1}{2} \int_{\Omega_{d_i}(a_i)} |K_{a_i}|^2 dx \\ & + \sum_{i=1}^n \int_{\Omega_{d_i}(a_i)} \nabla v_{\mathbf{a}} \cdot K_{a_i} dx + \sum_{i < j} \int_{\Omega} K_{a_i} \cdot K_{a_j} dx, \end{aligned}$$

if $\mathbf{a} \in \Omega^n \setminus \Delta_n$, and $\mathcal{F}^{(n)}(\mathbf{a}) := +\infty$ otherwise. Here $K_{a_i}(x) := \rho_{a_i}^{-1}(x)\hat{\theta}_{a_i}(x)$, being $(\rho_{a_i}, \theta_{a_i})$ a system of polar coordinates centered at a_i , and $v_{\mathbf{a}}$ solves

$$\begin{cases} \Delta v_{\mathbf{a}} = 0 & \text{in } \Omega, \\ v_{\mathbf{a}} = g - \sum_{i=1}^n \theta_{a_i} & \text{on } \partial\Omega. \end{cases}$$

In particular, $\mathcal{F}^{(n)}$ is continuous in $\bar{\Omega}^n$ and diverges to $+\infty$ if either at least one dislocation approaches the boundary or at least two dislocations collide, that is, $\mathcal{F}^{(n)}(\mathbf{a}) \rightarrow +\infty$ as $d_i \rightarrow 0$ for some i . Thus, $\mathcal{F}^{(n)}$ attains its minimum in $\Omega^n \setminus \Delta_n$.

We notice that, by rotating the vector fields of $\pi/2$, (5) can also be expressed as the sum of two terms: the *self energy* $\mathcal{E}_{\text{self}}$, responsible for the contribution of individual dislocations, and the *interaction energy* \mathcal{E}_{int} , depending on the mutual position of two dislocations. In formulas,

$$\mathcal{E}_{\text{self}}(a_i) := \pi \log \text{dist}(a_i, \partial\Omega) + \frac{1}{2} \int_{\Omega_{d(a_i)}(a)} |\nabla\phi_i + \nabla w_i|^2 dx + \frac{1}{2} \int_{B_{d(a)}(a)} |\nabla w_i|^2 dx,$$

and

$$\mathcal{E}_{\text{int}}(a_i, a_j) := \int_{\Omega} (\nabla\phi_i + \nabla w_i) \cdot (\nabla\phi_j + \nabla w_j) dx.$$

Here, ϕ_i and w_i are the solutions (with w_i determined up to an additive constant) to

$$\begin{cases} \Delta\phi_i = 2\pi\delta_{a_i} & \text{in } \Omega, \\ \phi_i(x) = \log|x - a_i| & \text{on } \partial\Omega, \end{cases} \quad \text{and} \quad \begin{cases} \Delta w_i = 0 & \text{in } \Omega, \\ \partial_\nu w_i(x) = \frac{1}{n}f - \partial_\nu\phi_i & \text{on } \partial\Omega. \end{cases}$$

A consequence of Theorem 1.1 is that also the energies (1) attain their minimum in Ω^n at an n -tuple of well separated points.

Corollary 1.1. *Under the assumptions (H1) and (H2), there exists $\epsilon_0 > 0$ such that, for every $\epsilon \in (0, \epsilon_0)$, the infimum problem*

$$(6) \quad \inf\{\mathcal{E}_\epsilon^{(n)}(\mathbf{a}) : \mathbf{a} \in \bar{\Omega}^n\},$$

admits a minimizer only in $\Omega^n \setminus \Delta_n$. Moreover, if $\mathbf{a}^\epsilon \in \Omega^n \setminus \Delta_n$ is a minimizer for (6), then (up to subsequences) we have $\mathbf{a}^\epsilon \rightarrow \mathbf{a}$ and $\mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon) \rightarrow \mathcal{F}^{(n)}(\mathbf{a})$, as $\epsilon \rightarrow 0$, where \mathbf{a} is a minimizer of the functional $\mathcal{F}^{(n)}$ defined in (5). In particular, for ϵ small enough, all the minimizers of problem (6) are n -tuples of distinct points that stay uniformly (with respect to ϵ) far away from the boundary and from one another.

Throughout the paper, we will always assume (H1) and (H2), even if it is not explicitly stated. We stress that the convexity and regularity assumptions

on Ω stated in (H1) provide a *uniform interior cone condition* of angle between $\pi/2$ and π , i.e.,

(7)

there exist $\pi/2 < \alpha < \pi$ and $\epsilon_\alpha > 0$ such that for every $b \in \partial\Omega$ the disk $B_{\epsilon_\alpha}(b)$ meets $\partial\Omega$ at two points b_1 and b_2 forming an angle at least α with b .

We point out that convexity and regularity play different roles: the former is conveniently assumed in order to simplify the exposition of the results (in fact, it can be removed without changing their essence); the latter, on the other hand, is fundamental in our proofs. Finally, we observe that the boundary condition of Dirichlet type $F \cdot \tau = f$, with f as in (H2), is fundamental to keep the dislocations confined inside the material. In fact, the natural boundary conditions of Neumann type imply that the dislocations migrate to the boundary and leave the domain, since, in such a case, the Dirichlet energy of the system decreases as the dislocations approach $\partial\Omega$ (see, e.g., [3, 15]).

A key feature in our analysis is the rescaling introduced in (3) (see also [1]), which is related to the so-called *Hadamard finite part* of a divergent integral (see [13]). Such type of asymptotic analysis has the advantage of keeping into account the energetic dependence on the position of the dislocations $\mathbf{a} \in \Omega^n$, whereas it is well-known that the standard rescaling obtained by dividing the energy by $|\log \epsilon|$ gives rise to an energy which only counts the number of dislocations in the bulk (see again [1]).

Some results close to those presented in this paper can be found in the literature about Ginzburg-Landau vortices (see [2] and also [23]). In this respect, the present paper can be considered as a self-contained presentation of the asymptotic results for the energy (3), presented in a language that is familiar to the dislocation community, targeted, in particular, to applied mathematicians and continuum mechanists. The statement of Theorem 1.1 is in fact similar to that of [2, Theorems I.9 and I.10], but its proof is based on an original iterative procedure (close, in spirit, to some combinatoric algorithms) which makes it quite different and useful for numerical applications. The need for such an algorithm is dictated by the fact that in the core-radius approach we allow for the cores around each dislocation to intersect with one another and with the boundary of the domain, which was avoided in [2] by introducing a safety radius. Moreover, we set here the bases and the notation for tackling the more challenging problem of the upscaling of the system of dislocations. In [17] we study the limit as $n \rightarrow \infty$ and obtain a limit energy functional defined on measures describing the distribution of dislocations in the material.

Section 2 sets the notation and presents some preparatory results. Section 3 contains a result on the properties of K and on some a priori bounds for harmonic functions. Section 4 is devoted to proving Theorem 1.1 and Corollary 1.1. Section 5 contains numerical plots of $\mathcal{F}^{(n)}$ in $\Omega = B(0, 1)$ under different boundary conditions.

2 - Preliminaries

In Subsection 2.1 we introduce the notation used throughout the paper. Then, in view of the core radius approach, in Subsection 2.2 we rewrite the energy $\mathcal{E}_\epsilon^{(n)}$ in terms of the displacement of a regular function.

2.1 - Notation

- $B_r(x)$ denotes the open disk of radius $r > 0$ centered at $x \in \mathbb{R}^2$; $\overline{B}_r(x)$ is its closure;
- for $n \geq 1$, and for $x_1, \dots, x_n \in \overline{\Omega}$, we denote $\mathbf{x} := (x_1, \dots, x_n)$ whenever there is no risk of misunderstanding; the symbol $\Omega_r(\mathbf{x})$ denotes the open set

$$\Omega_r(\mathbf{x}) := \Omega \setminus \left(\bigcup_{i=1}^n \overline{B}_r(x_i) \right);$$

$\overline{\Omega}_r(\mathbf{x})$ is its closure;

- the function $\mathbb{R}^2 \ni x \mapsto \text{dist}(x, E)$ denotes the distance of x from a set $E \subset \mathbb{R}^2$; in the particular case $E = \partial\Omega$, we define $d(x) := \text{dist}(x, \partial\Omega)$;
- we denote $\Omega^{(r)} := \{x \in \Omega : d(x) > r\}$ and by $\overline{\Omega}^{(r)}$ its closure;
- Ω^n denotes the cartesian product of n copies of Ω and $\overline{\Omega}^n$ its closure;
- for $0 < r < R$, $A_r^R(x) := B_R(x) \setminus \overline{B}_r(x)$ denotes the open annulus of internal radius r and external radius R centered at $x \in \mathbb{R}^2$;
- χ_E denotes the characteristic function of E : $\chi_E(x) = 1$ if $x \in E$; $\chi_E(x) = 0$ if $x \notin E$;
- given a set E with piecewise C^1 boundary, ν and τ denote the outer unit normal and the tangent unit vectors to ∂E , respectively;
- $\text{diam } E$ denotes the diameter of a set $E \subset \mathbb{R}^2$;
- given $x = (x_1, x_2) \in \mathbb{R}^2$, we denote by $x^\perp := (-x_2, x_1)$ the rotated vector;
- given $x \in \mathbb{R}^2$, we define (ρ_x, θ_x) as the standard polar coordinate system centered at x ; $\hat{\rho}_x$ and $\hat{\theta}_x$ denote the corresponding unit vectors;
- given $x \in \mathbb{R}^2$, we denote by K_x the vector field $K_x := \rho_x^{-1} \hat{\theta}_x$. It is easy to see that ([18])

$$(8) \quad \begin{cases} \text{div } K_x = 0 & \text{in } \mathcal{D}'(\mathbb{R}^2), \\ \text{curl } K_x = 2\pi\delta_x & \text{in } \mathcal{D}'(\mathbb{R}^2), \\ K_x \cdot \nu = 0 & \text{on } \partial B_r(x), \text{ for any } r > 0. \end{cases}$$

Notice that K_x is the absolutely continuous part of the gradient of the function θ_x , and the jump set of θ_x is the half line starting from x and passing through y with $[\theta_x] = 2\pi$ across it;

- we define the function $\omega: \bar{\Omega} \rightarrow \{1, 2\}$ by

$$(9) \quad \omega(a) = \begin{cases} 1 & \text{if } a \in \Omega, \\ 2 & \text{if } a \in \partial\Omega; \end{cases}$$

- $\text{spt } \varphi$ denotes the support of the function φ ;
- given $x \in \mathbb{R}^2$, δ_x is the Dirac measure centered at x ;
- \mathcal{H}^1 denotes the one-dimensional Hausdorff measure;
- the letter C alone represents a generic constant (possibly depending on Ω) whose value might change from line to line.

2.2 - Displacement formulation

We start by introducing the multiplicity of a dislocation.

Definition 2.1 (multiplicity of a dislocation). Let $n \in \mathbb{N}$ with $n \geq 2$ and let $a_1, \dots, a_n \in \bar{\Omega}$. We label the dislocations in such a way that the first ℓ of them ($\ell \leq n$) are all distinct, so that $a_i \neq a_k$ for all $i, k \in \{1, \dots, \ell\}$, $i \neq k$, and, for every $j \in \{\ell+1, \dots, n\}$ there exists $i(j) \in \{1, \dots, \ell\}$ such that $a_j = a_{i(j)}$. For every $i = 1, \dots, \ell$ we say that a point a_i has *multiplicity* m_i if there are $m_i - 1$ points a_j , with $j \in \{\ell+1, \dots, n\}$ that coincide with a_i . Clearly, $\sum_{i=1}^{\ell} m_i = n$.

Given n dislocations $a_1, \dots, a_n \in \bar{\Omega}$, we choose n (closed) segments $\Sigma_1, \dots, \Sigma_n$ joining the dislocations with the boundary, such that $\Omega \setminus (\Sigma_1 \cup \dots \cup \Sigma_n)$ is simply connected. One possible construction of such a family of segments is to fix a point $a^* \notin \bar{\Omega}$ and take Σ_i as the portion of the segment joining a_i with a^* lying inside $\bar{\Omega}$, for $i = 1, \dots, \ell$. With this construction, if $a_i = a_j$, then $\Sigma_i = \Sigma_j$. Moreover, we set $b_i := \Sigma_i \cap \partial\Omega$, for $i = 1, \dots, \ell$. Given $b \in \partial\Omega \setminus \{b_1, \dots, b_\ell\}$, we denote by γ_b^x the counterclockwise path in $\partial\Omega$ connecting b and x . Such parametrization induces an ordering on the points of the boundary: we say that x precedes y on $\partial\Omega$, and we write $x \preceq y$, if the support of γ_b^y contains that of γ_b^x . Without loss of generality, we may relabel the b_i 's (and the a_i 's, accordingly) so that $b \prec b_1 \preceq \dots \preceq b_\ell$. According to this notation, for $i = 1, \dots, \ell$, we define $g_{b_i}: \partial\Omega \setminus \{b_i\} \rightarrow \mathbb{R}$ as

$$g_{b_i}(x) := \begin{cases} \frac{1}{n} \int_{\gamma_b^x} f \, dt & \text{if } b \preceq x \prec b_i, \\ \frac{1}{n} \int_{\gamma_b^x} f \, dt - 2\pi & \text{if } b_i \prec x \prec b. \end{cases}$$

Moreover, we introduce

$$(10) \quad g(x) := \sum_{i=1}^{\ell} m_i g_{b_i}(x) = \int_{\gamma_b^x} f(y) dy - 2\pi \sum_{i=0}^{j-1} m_i \quad \text{if } b_{j-1} \prec x \prec b_j,$$

with $j \in \{1, \dots, \ell + 1\}$, $b_0 = b_{\ell+1} := b$, and $m_0 := 0$. Note that, since g_{b_i} is continuous except at b_i , where it has a jump of $2m_i\pi$, we have that the boundary datum g has a jump of $2\pi m_i$ at every b_i (whereas it is continuous at b).

In view of the construction above, for every $\epsilon > 0$, every connected component of $\Omega_\epsilon(\mathbf{a}) \setminus (\Sigma_1 \cup \dots \cup \Sigma_n)$ is simply connected. Therefore, the energy (1) can be expressed as

$$(11) \quad \mathcal{E}_\epsilon^{(n)}(\mathbf{a}) = \frac{1}{2} \int_{\Omega_\epsilon(\mathbf{a})} |\nabla u_\mathbf{a}^\epsilon|^2 dx,$$

where $u_\mathbf{a}^\epsilon \in H^1(\Omega_\epsilon(\mathbf{a}))$ is characterized by

$$(12) \quad \begin{cases} \Delta u_\mathbf{a}^\epsilon = 0 & \text{in } \Omega_\epsilon(\mathbf{a}) \setminus (\cup_{i=1}^n \Sigma_i), \\ [u_\mathbf{a}^\epsilon] = 2\pi m_i & \text{on } \Sigma_i \cap \Omega_\epsilon(\mathbf{a}), \\ u_\mathbf{a}^\epsilon = \sum_{i=1}^{\ell} m_i g_{b_i} = g & \text{on } \partial\Omega \setminus \cup_{i=1}^n \overline{B}_\epsilon(a_i), \\ \partial u_\mathbf{a}^\epsilon / \partial \nu = 0 & \text{on } \partial\Omega_\epsilon(\mathbf{a}) \setminus \partial\Omega, \\ \partial(u_\mathbf{a}^\epsilon)^+ / \partial \nu = \partial(u_\mathbf{a}^\epsilon)^- / \partial \nu & \text{on } (\cup_{i=1}^n \Sigma_i) \cap \Omega_\epsilon(\mathbf{a}). \end{cases}$$

We point out that the expression (11) is consistent with the definition of the energy in (1); in particular, when $a_i \equiv a$ for every $i = 1, \dots, n$, choosing $b_i \equiv b$, the datum g in (10) is given by $g(x) = \int_{\gamma_b^x} f(y) dy$, and we have $\mathcal{E}_\epsilon^{(n)}(a, \dots, a) = n^2 \mathcal{E}_\epsilon^{(1)}(a)$. Therefore, by (3), for every ϵ small enough

$$(13) \quad \mathcal{F}_\epsilon^{(n)}(\underbrace{a, \dots, a}_{n\text{-times}}) = n^2 \mathcal{F}_\epsilon^{(1)}(a) + n(n-1)\pi |\log \epsilon|.$$

The following lemma shows that the construction of the Σ_i 's described above is arbitrary.

Lemma 2.1. *The functional $\mathcal{E}_\epsilon^{(n)}$ does not depend on the choice of the discontinuity points b_i 's, nor the primitive g of f , nor the cuts Σ_i 's.*

Proof. It is enough to prove the result for $n = 1$. Let Σ be a simple smooth curve in $\overline{\Omega}$ connecting a with b , intersecting $\partial B_\epsilon(a)$ at a single point and $\partial\Omega$ only at b . Notice that each connected component of $\Omega_\epsilon(a) \setminus \Sigma$ is simply connected. The condition $\text{curl } F = 0$ in $\Omega_\epsilon(a)$ implies that there exists $u \in H^1(\Omega_\epsilon(a) \setminus \Sigma)$ such that $F = \nabla u$ in $\Omega_\epsilon(a) \setminus \Sigma$. Therefore, the energy (11) can be expressed as

$$(14) \quad \mathcal{E}_\epsilon^{(1)}(a) = \min \left\{ \frac{1}{2} \int_{\Omega_\epsilon(a)} |\nabla u|^2 dx : u \in H^1(\Omega_\epsilon(a) \setminus \Sigma), u = g \text{ on } \partial\Omega \setminus \overline{B}_\epsilon(a) \right\}.$$

To show that the energy above does not depend on the choice of the discontinuity point b , nor on the primitive g of f , nor on the curve Σ , we argue as follows.

Let g and g' be two primitives of f with the same discontinuity point. Then the two boundary data differ by a constant. The same holds true for the corresponding minimizers of (14), which have the same gradient and the same energy.

Let now assume that g and g' have two different discontinuity points, say b and b' . Denote by Σ and Σ' the segments joining a and b , and a and b' , respectively. Let ϕ and ϕ' be the angles associated with the discontinuity points b and b' , respectively, in the angular coordinate centered at a . It is not restrictive to assume that $\phi < \phi'$. Let u be the solution to (40) associated with g , b , and Σ ; define $v := u + 2\pi\chi_S$, where S is the subset of $\Omega_\epsilon(a)$ in which $\theta_a \in (\phi, \phi')$. It is easy to see that v is admissible for (14) associated with g' , b' , and Σ' . In particular, v is the solution to (40) associated with g' , b' , and Σ' ; moreover,

$$\frac{1}{2} \int_{\Omega_\epsilon(a)} |\nabla u|^2 dx = \frac{1}{2} \int_{\Omega_\epsilon(a)} |\nabla v|^2 dx.$$

This concludes the proof of the invariance with respect to the discontinuity point in the case of a straight cut.

For general curves Σ and Σ' , the region in $\Omega_\epsilon(a)$ lying between Σ and Σ' is the union of some simply connected sets S_i , $i \geq 1$, bounded by portions of Σ and Σ' . In this case, the same strategy applies, provided one adds or subtracts $2\pi\chi_{S_i}$, according to whether the portions of Σ bounding S_i precedes or follows that of Σ' in the positive orientation of the angular coordinate. This concludes the proof. \square

Notice that the functional (11) and the system (12) are well defined even in the case that there exists $i \in \{1, \dots, n\}$ such that $a_i \in \partial\Omega$. In this case, Lemma 2.1 allows us to choose the discontinuity point $b_i = a_i$ and the cut $\Sigma_i = \emptyset$.

We may write $\nabla u_{\mathbf{a}}^\epsilon$ as the sum of a singular and a regular part

$$(15) \quad \nabla u_{\mathbf{a}}^\epsilon = \sum_{i=1}^n \omega(a_i) K_{a_i} + \nabla \bar{v}_{\mathbf{a}}^\epsilon,$$

where the K_{a_i} satisfy (8) and ω is defined in (9). Here we impose that the angular coordinate θ_{a_i} centered at a_i which defines K_{a_i} has a jump of 2π on the half line containing Σ_i if $a_i \in \Omega$, and on the half line containing the outer normal to the boundary at a_i in case $a_i \in \partial\Omega$. The function $\bar{v}_{\mathbf{a}}^\epsilon \in H^1(\Omega_\epsilon(\mathbf{a}))$ solves

$$(16) \quad \begin{cases} \Delta \bar{v}_{\mathbf{a}}^\epsilon = 0 & \text{in } \Omega_\epsilon(\mathbf{a}), \\ \bar{v}_{\mathbf{a}}^\epsilon = \sum_{i=1}^n (g_{b_i} - \omega(a_i)\theta_{a_i}) & \text{on } \partial\Omega \setminus \cup_{i=1}^n \bar{B}_\epsilon(a_i), \\ \partial \bar{v}_{\mathbf{a}}^\epsilon / \partial \nu = - \sum_{i=1}^n K_{a_i} \cdot \nu & \text{on } \partial\Omega_\epsilon(\mathbf{a}) \setminus \partial\Omega. \end{cases}$$

The weight $\omega(a_i)$ is necessary in order to balance the jump of g_{b_i} at b_i . Using

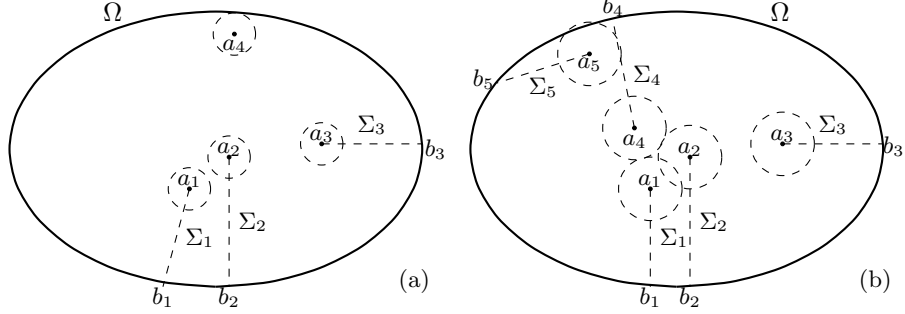


Figure 1: (a) A crystal with 4 dislocations. (b) A crystal with 5 dislocations.

(15), we can write the energy (11) as

$$\mathcal{E}_\epsilon^{(n)}(\mathbf{a}) = \frac{1}{2} \int_{\Omega_\epsilon(\mathbf{a})} \left| \sum_{i=1}^n \omega(a_i) K_{a_i} + \nabla \bar{v}_\mathbf{a}^\epsilon \right|^2 dx.$$

Remark 2.1. By linearity, the function $\bar{v}_\mathbf{a}^\epsilon$ introduced in (16) can be written as the superposition $\bar{v}_\mathbf{a}^\epsilon = \sum_{i=1}^n \bar{v}_{a_i}^\epsilon$ of the solutions $\bar{v}_{a_i}^\epsilon$ to

$$\begin{cases} \Delta \bar{v}_{a_i}^\epsilon = 0 & \text{in } \Omega_\epsilon(\mathbf{a}), \\ \bar{v}_{a_i}^\epsilon = g_{b_i} - \omega(a_i) \theta_{a_i} & \text{on } \partial\Omega \setminus \cup_{j=1}^n \bar{B}_\epsilon(a_j), \\ \partial \bar{v}_{a_i}^\epsilon / \partial \nu = -K_{a_i} \cdot \nu & \text{on } \partial\Omega_\epsilon(\mathbf{a}) \setminus \partial\Omega. \end{cases}$$

Fix $i \in \{1, \dots, n\}$. By linearity, the solution $\bar{v}_{a_i}^\epsilon$ of the system above can be written as a superposition of one function satisfying a homogeneous Neumann boundary condition and another function satisfying a homogeneous Dirichlet boundary condition, namely $\bar{v}_{a_i}^\epsilon = \hat{u}_{a_i}^\epsilon + q_{a_i}^\epsilon$, where $\hat{u}_{a_i}^\epsilon$ and $q_{a_i}^\epsilon$ solve, respectively,

$$\begin{cases} \Delta \hat{u}_{a_i}^\epsilon = 0 & \text{in } \Omega_\epsilon(\mathbf{a}), \\ \hat{u}_{a_i}^\epsilon = g_{b_i} - \omega(a_i) \theta_{a_i} & \text{on } \partial\Omega \setminus \cup_{j=1}^n \bar{B}_\epsilon(a_j), \\ \partial \hat{u}_{a_i}^\epsilon / \partial \nu = 0 & \text{on } \cup_{j=1}^n \partial B_\epsilon(a_j) \end{cases}$$

and

$$\begin{cases} \Delta q_{a_i}^\epsilon = 0 & \text{in } \Omega_\epsilon(\mathbf{a}), \\ q_{a_i}^\epsilon = 0 & \text{on } \partial\Omega \setminus \cup_{j=1}^n \bar{B}_\epsilon(a_j), \\ \partial q_{a_i}^\epsilon / \partial \nu = -K_{a_i} \cdot \nu & \text{on } \cup_{j=1}^n \partial B_\epsilon(a_j), \\ \oint_{\partial B_\epsilon(a_j)} \partial q_{a_i}^\epsilon / \partial \tau = 0 & \text{for every } j = 1, \dots, n. \end{cases}$$

3 - Auxiliary lemmas

We start by proving some useful properties of the function K satisfying (8).

Lemma 3.1 (Properties of K). *Let $y_1, y_2 \in \mathbb{R}^2$. Set $r := |y_1 - y_2|/2$ and y the midpoint $(y_1 + y_2)/2$. Then the vector fields K_{y_1} and K_{y_2} satisfy the following properties:*

(i) *the scalar product $K_{y_1} \cdot K_{y_2}$ is negative in the disk $B_r(y)$, is zero on $\partial B_r(y) \setminus \{y_1, y_2\}$, and is positive in $\mathbb{R}^2 \setminus \overline{B_r}(y)$. In particular, $\int_{B_r(y)} K_{y_1} \cdot K_{y_2} dx \geq -2\pi$;*

(ii) *the following estimate holds: $\int_{B_r(y_i)} K_{y_1} \cdot K_{y_2} dx \leq 2\pi$, for $i = 1, 2$.*

Moreover, let $\ell \in \mathbb{N}$, and let $y_1, \dots, y_\ell \in \overline{\Omega}$ be distinct points. Then,

(iii) *if $y_i \in \Omega$, for every $0 < \epsilon < d_i$ (see (4)),*
(17)

$$2\pi |\log \epsilon| + 2\pi \log d_i \leq \int_{\Omega_\epsilon(y_1, \dots, y_\ell)} |K_{y_i}|^2 dx \leq 2\pi |\log \epsilon| + 2\pi \log(\text{diam } \Omega);$$

(iv) *if $y_i \in \partial\Omega$, for every $0 < \epsilon < \min_{j \neq i} \{\epsilon_\alpha, |y_i - y_j|/2\}$,*

$$(18) \quad \alpha |\log \epsilon| + \alpha \log \epsilon_\alpha \leq \int_{\Omega_\epsilon(y_1, \dots, y_\ell)} |K_{y_i}|^2 dx \leq \pi |\log \epsilon| + \pi \log(\text{diam } \Omega).$$

Proof. To prove (i) it is enough to notice that, given a point $x \in \mathbb{R}^2 \setminus \{y_1, y_2\}$, the angle $\widehat{y_1 x y_2}$ is larger than, equal to, or smaller than $\pi/2$, according to whether $x \in B_r(y)$, $x \in \partial B_r(y)$, $x \in \mathbb{R}^2 \setminus \overline{B_r}(y)$, respectively. To prove the estimate, we consider the disk $B_r(y)$ whose diameter is the axis of the segment joining y_1 and y_2 , and we define $B_r(y)^+$ the half of $B_r(y)$ on the side of y_2 . By symmetry, we have

$$\begin{aligned} \int_{B_r(y)} K_{y_1} \cdot K_{y_2} dx &= -2 \int_{B_r(y)^+} |K_{y_1} \cdot K_{y_2}| dx \geq -\frac{2}{r} \int_{B_r(y)^+} |K_{y_2}| dx \\ &\geq -\frac{2}{r} \int_{\pi/2}^{3\pi/2} \int_0^r d\rho_2 d\theta_2 = -2\pi. \end{aligned}$$

To prove (ii), observe that $|K_{y_i}| \leq 1/r$ in $\mathbb{R}^2 \setminus \overline{B_r}(y_i)$, so that

$$\int_{B_r(y_i)} K_{y_1} \cdot K_{y_2} dx \leq \frac{1}{r} \int_{B_r(0)} |x|^{-1} dx = \pi.$$

To prove (iii) and (iv) one integrates in polar coordinates centered at y_i over the sets $A_\epsilon^{\text{diam } \Omega}(y_i) \cap \Omega$, $A_\epsilon^{d_i}(y_i) \cap \Omega$, $A_\epsilon^s(y_i) \cap \Omega$ with $s := \min_{j \neq i} \{\epsilon_\alpha, |y_i - y_j|/2\}$, and uses the convexity assumption (H1) and (7). \square

In the rest of this section, we prove some *a priori* bounds on harmonic functions that will be useful in the sequel.

Lemma 3.2. Let $a_1, \dots, a_n \in \overline{\Omega}$, let $h \in W^{1,1}(\partial\Omega)$, and let \hat{u}^ϵ be the solution to

$$(19) \quad \begin{cases} \Delta \hat{u}^\epsilon = 0 & \text{in } \Omega_\epsilon(\mathbf{a}), \\ \hat{u}^\epsilon = h & \text{on } \partial\Omega \setminus \cup_{j=1}^n \overline{B}_\epsilon(a_j), \\ \partial \hat{u}^\epsilon / \partial \nu = 0 & \text{on } \partial\Omega_\epsilon(\mathbf{a}) \cap (\cup_{j=1}^n \partial B_\epsilon(a_j)). \end{cases}$$

Then \hat{u}^ϵ is the minimizer of the Dirichlet energy in $H^1(\Omega_\epsilon(\mathbf{a}))$ with prescribed boundary datum h on $\partial\Omega \setminus \cup_{j=1}^n \overline{B}_\epsilon(a_j)$, and the following estimates hold

$$(20) \quad |\hat{u}^\epsilon(x)| \leq C \|h\|_{L^\infty(\partial\Omega)}, \quad \text{for all } x \in \overline{\Omega}_\epsilon(\mathbf{a})$$

and

$$(21) \quad \int_{\Omega_\epsilon(\mathbf{a})} |\nabla \hat{u}^\epsilon|^2 dx \leq C \|h\|_{H^{1/2}(\partial\Omega)}^2$$

for some constant $C > 0$ independent of ϵ . Moreover, in the case $a_1, \dots, a_n \in \Omega$ are distinct points and $\epsilon < \min_i d_i$, also the estimate on the gradient holds

$$(22) \quad |\nabla \hat{u}^\epsilon(x)| \leq C \|h\|_{L^\infty(\partial\Omega)}, \quad \text{for all } x \in \cup_{j=1}^n \partial B_\epsilon(a_j).$$

Proof. The minimality of \hat{u}^ϵ follows from the uniqueness of the solution to (19); estimate (20) is a consequence of the maximum principle and Hopf's Lemma, while (21) is obtained by testing \hat{u}^ϵ with the minimizer of the Dirichlet energy in $H^1(\Omega)$ with prescribed boundary datum h on $\partial\Omega \setminus \cup_{j=1}^n \overline{B}_\epsilon(a_j)$ and using the continuity of the map which associates the minimizer with the boundary datum. Finally, the proof of (22) is standard (see, e.g., [10]). \square

As a simple consequence of Lemma 3.2, we have the following estimate in the case $n = 1$.

Lemma 3.3. Let $a \in \Omega$ and let $\epsilon \in (0, d(a)/2)$. Let w^ϵ be the harmonic extension in $B_\epsilon(a)$ of \bar{u}_a^ϵ . Then there exists a constant $C > 0$ independent of ϵ such that

$$(23) \quad \int_{B_\epsilon(a)} |\nabla w^\epsilon(x)|^2 dx \leq C \epsilon^2 \|g - \theta_a\|_{L^\infty(\partial\Omega)}^2.$$

Proof. By applying Lemma 3.2 with $n = 1$, we immediately obtain that $\bar{u}_a^\epsilon \in W^{1,\infty}(\partial B_\epsilon(a))$ and

$$(24) \quad \|\bar{u}_a^\epsilon\|_{W^{1,\infty}(\partial B_\epsilon(a))} \leq C \|g - \theta_a\|_{L^\infty(\partial\Omega)}.$$

It is a known fact in the theory of harmonic functions (see [10]) that there exists a constant $C > 0$ such that for every harmonic function $\varphi \in H^1(B_1(0))$

$$\|\nabla \varphi\|_{L^2(B_1(0))}^2 \leq C \|\varphi - m(\varphi)\|_{H^{1/2}(\partial B_1(0))}^2,$$

where $m(\varphi)$ is the average of φ on $\partial B_1(0)$. Using the Poincaré inequality we have

$$\begin{aligned} & \|\varphi - m(\varphi)\|_{H^{1/2}(\partial B_1(0))}^2 \\ & \leq C \left(\|\varphi - m(\varphi)\|_{L^2(\partial B_1(0))}^2 + \int_{\partial B_1(0)} \int_{\partial B_1(0)} \frac{|\varphi(x) - \varphi(y)|^2}{|x - y|^2} dx dy \right) \\ & \leq C \left(\int_{\partial B_1(0)} \left| \frac{\partial \varphi(x)}{\partial \tau} \right|^2 dx + \int_{\partial B_1(0)} \int_{\partial B_1(0)} \frac{|\varphi(x) - \varphi(y)|^2}{|x - y|^2} dx dy \right). \end{aligned}$$

Therefore, using this and employing the change of variables $x = \epsilon x' + a$ and $y = \epsilon y' + a$ we have

$$\begin{aligned} \|\nabla w^\epsilon\|_{L^2(B_\epsilon(a))}^2 &= \int_{B_1(0)} |\nabla w^\epsilon(\epsilon x' + a)|^2 dx' \\ &\leq C \left(\int_{\partial B_1(0)} \left| \frac{\partial w^\epsilon(\epsilon x' + a)}{\partial \tau} \right|^2 dx' \right. \\ &\quad \left. + \int_{\partial B_1(0)} \int_{\partial B_1(0)} \frac{|w^\epsilon(\epsilon x' + a) - w^\epsilon(\epsilon y' + a)|^2}{|x' - y'|^2} dx' dy' \right) \\ &= C \left(\epsilon \int_{\partial B_\epsilon(a)} \left| \frac{\partial w^\epsilon(x)}{\partial \tau} \right|^2 dx + \int_{\partial B_\epsilon(a)} \int_{\partial B_\epsilon(a)} \frac{|w^\epsilon(x) - w^\epsilon(y)|^2}{|x - y|^2} dx dy \right). \end{aligned}$$

Owing to the fact that $w^\epsilon = \bar{u}^\epsilon$ on $\partial B_\epsilon(a)$ and by (24), we obtain (23). \square

The following result is similar to [2, Lemma I.5]; however, for the sake of completeness, we present here a proof.

Lemma 3.4. *Let $D_1, \dots, D_\ell \subseteq \mathbb{R}^2$ be open, bounded, simply connected sets with Lipschitz boundary, such that $\bar{D}_1, \dots, \bar{D}_\ell$ are pairwise disjoint and intersect $\bar{\Omega}$. Set $\Omega' := \Omega \setminus \cup_{i=1}^\ell \bar{D}_i$ and for every $i = 1, \dots, \ell$, let $h_i \in C^1(\bar{\Omega}')$. Consider the minimum problem*

$$(25) \quad \min_c \min_u \frac{1}{2} \int_{\Omega'} |\nabla u|^2 dx,$$

with $c = (c_1, \dots, c_\ell) \in \mathbb{R}^\ell$ and u varying in the subset of functions in $H^1(\Omega')$ satisfying

$$(26) \quad u = h_i + c_i \text{ on } \partial D_i \cap \Omega, \text{ for } i = 1, \dots, \ell.$$

Then $p \in H^1(\Omega')$ is a solution to (25) if and only if it satisfies

$$(27) \quad \begin{cases} \Delta p = 0 & \text{in } \Omega', \\ \partial p / \partial \nu = 0 & \text{on } \partial \Omega' \cap \partial \Omega, \\ \partial p / \partial \tau = \partial h_i / \partial \tau & \text{on } \partial D_i \cap \Omega, \text{ for } i = 1, \dots, \ell, \\ \int_{\partial D_i \cap \Omega} \partial p / \partial \nu d\mathcal{H}^1 = 0 & \text{for } i = 1, \dots, \ell. \end{cases}$$

Moreover, a solution p satisfies

$$(28) \quad \left| \max_{\overline{\Omega'}} p - \min_{\overline{\Omega'}} p \right| \leq \sum_{i=1}^{\ell} \mathcal{H}^1(\partial D_i \cap \Omega) \max_{\partial D_i \cap \Omega} \left| \frac{\partial h_i(x)}{\partial \tau} \right|.$$

Proof. First we remark that problem (25) is well posed, namely the class of admissible functions is not empty: indeed, the function $u := \sum_{i=1}^{\ell} \zeta_i h_i$ satisfies (26), being ζ_i suitable smooth functions in Ω' which are identically 1 in a neighborhood of D_i .

If p satisfies the minimum problem (25), then it is easy to obtain system (27) as the Euler-Lagrange conditions arising from minimality.

Conversely, assume that p satisfies (27). In order to show that p is optimal for (25), it is enough to show that the (27) has a unique solution (up to a constant). Take two solutions of (27); then their difference s is characterized by

$$(29) \quad \begin{cases} \Delta s = 0 & \text{in } \Omega', \\ \partial s / \partial \nu = 0 & \text{on } \partial \Omega' \cap \partial \Omega, \\ \partial s / \partial \tau = 0 & \text{on } \partial \Omega' \setminus \partial \Omega, \\ \int_{\partial D_i \cap \Omega} \partial s / \partial \nu \, d\mathcal{H}^1 = 0 & \text{for } i = 1, \dots, \ell. \end{cases}$$

The third condition implies that s is constant on the boundary of the holes (intersected with Ω), namely $s = s_i$ on $\partial D_i \cap \Omega$ for some $s_i \in \mathbb{R}$, for $i = 1, \dots, \ell$. In view of the maximum principle, we infer that either s attains its maximum on the boundary $\partial \Omega' \setminus \partial \Omega$ or it is constant. The former case is excluded by Hopf's Lemma combined with the second and fourth conditions in (29); therefore the latter case holds and the proof of the first statement is concluded.

We relabel the sets D_i so that the local minima of p on their boundary are ordered as follows:

$$\min_{\partial D_1 \cap \Omega} p \geq \min_{\partial D_2 \cap \Omega} p \geq \dots \geq \min_{\partial D_\ell \cap \Omega} p = \min_{\overline{\Omega'}} p.$$

Notice that, for the last equality, we have used the maximum principle, combined with the Hopf's Lemma, which prevents the minimum of p to be on $\partial \Omega$ (the same holds true for the maximum). We now define an ordered subfamily of indices as follows: we choose

$$j_1 \in \left\{ j : \max_{\partial D_j \cap \Omega} p = \max_{\overline{\Omega'}} p \right\},$$

and, by recursion, given j_i we choose

$$j_{i+1} \in \left\{ j : \max_{\partial D_j \cap \Omega} p = \max_{\cup_{k>j_i} \partial D_k \cap \Omega} p \right\}.$$

Since $j_{i+1} > j_i$, in a finite number of steps, say i^* , we find $j_{i^*} = \ell$ and the procedure stops. We clearly have

$$\max_{\overline{\Omega'}} p = \max_{\partial D_{j_1} \cap \Omega} p \geq \max_{\partial D_{j_2} \cap \Omega} p \geq \dots \geq \max_{\partial D_{j_{i^*}} \cap \Omega} p = \max_{\partial D_\ell \cap \Omega} p.$$

Moreover, we claim that, for every $i = 1, \dots, i^* - 1$

$$(30) \quad \min_{\partial D_{j_i} \cap \Omega} p \leq \max_{\partial D_{j_{i+1}} \cap \Omega} p.$$

Once proved the claim we are done, indeed

$$\begin{aligned} \left| \max_{\overline{\Omega}'} - \min_{\overline{\Omega}'} p \right| &= \max_{\partial D_{j_1} \cap \Omega} p - \min_{\partial D_{j_{i^*}} \cap \Omega} p \\ &= \sum_{i=1}^{i^*} \left(\max_{\partial D_{j_i} \cap \Omega} p - \min_{\partial D_{j_i} \cap \Omega} p \right) + \sum_{i=1}^{i^*-1} \left(\min_{\partial D_{j_i} \cap \Omega} p - \max_{\partial D_{j_{i+1}} \cap \Omega} p \right) \\ &\leq \sum_{i=1}^{\ell} \left(\max_{\partial D_i \cap \Omega} p - \min_{\partial D_i \cap \Omega} p \right) \leq \sum_{i=1}^{\ell} \mathcal{H}^1(\partial D_i \cap \Omega) \max_{\partial D_i \cap \Omega} \left| \frac{\partial h_i(x)}{\partial \tau} \right|, \end{aligned}$$

where, in the last line, we have used the Mean Value Theorem. This proves (28).

In order to prove the claim (30), we argue by contradiction: we assume that there exists $i \in \{1, \dots, i^* - 1\}$ such that

$$\underline{p} := \min_{\partial D_{j_i} \cap \Omega} p > \bar{p} := \max_{\partial D_{j_{i+1}} \cap \Omega} p.$$

Notice that, by construction, we have

$$\min_{\cup_{k=1}^{j_i} \partial D_k \cap \Omega} p > \max_{\cup_{k=j_{i+1}}^{\ell} \partial D_k \cap \Omega} p.$$

Consider the nonempty set $E := \{x \in \Omega' : \underline{p} > p(x) > \bar{p}\}$ and the function

$$\psi(x) := \begin{cases} p(x) & \text{if } p(x) \leq \bar{p}, \\ \bar{p} & \text{if } x \in E, \\ p(x) + \bar{p} - \underline{p} & \text{if } p(x) \geq \underline{p}. \end{cases}$$

The function ψ is an element of $H^1(\Omega')$, it is admissible for the minimum problem (25), and it decreases the Dirichlet energy, contradicting the optimality of p . \square

As a direct consequence of Lemma 3.4 we obtain the following corollary.

Corollary 3.1. *Let $a_1, \dots, a_\ell \in \overline{\Omega}$ be distinct points and let $(a_1^\epsilon, \dots, a_\ell^\epsilon)$ be a sequence of points in Ω^ℓ converging to (a_1, \dots, a_ℓ) as $\epsilon \rightarrow 0$, such that the disks $B_\epsilon(a_j^\epsilon)$ are pairwise disjoint. Fix $i \in \{1, \dots, \ell\}$ and let p_i^ϵ be the solution (unique up to a constant) to*

$$(31) \quad \begin{cases} \Delta p_i^\epsilon = 0 & \text{in } \Omega_\epsilon(a_1^\epsilon, \dots, a_\ell^\epsilon), \\ \partial p_i^\epsilon / \partial \nu = 0 & \text{on } \partial \Omega \setminus \cup_{j=1}^{\ell} \overline{B}_\epsilon(a_j^\epsilon), \\ \partial p_i^\epsilon / \partial \tau = -\partial \log(|x - a_i^\epsilon|) / \partial \tau & \text{on } \partial B_\epsilon(a_j^\epsilon) \cap \Omega, \text{ for every } j \neq i, \\ \partial p_i^\epsilon / \partial \tau = 0 & \text{on } \partial B_\epsilon(a_i^\epsilon) \cap \Omega, \\ \int_{\partial B_\epsilon(a_j^\epsilon) \cap \Omega} \partial p_i^\epsilon / \partial \nu = 0 & \text{for } j = 1, \dots, \ell. \end{cases}$$

Then, there exists a positive constant C independent of ϵ such that

$$(32) \quad \|p_i^\epsilon\|_{L^\infty(\Omega_\epsilon(a_1^\epsilon, \dots, a_\ell^\epsilon))} \leq C\epsilon.$$

Moreover,

$$(33) \quad \int_{\Omega_\epsilon(a_1^\epsilon, \dots, a_\ell^\epsilon)} |\nabla p_i^\epsilon|^2 dx \leq \frac{C}{|\log \epsilon|}.$$

Proof. By applying Lemma 3.4 with $D_j = B_\epsilon(a_j^\epsilon)$, and $h_j = -\log(|x - a_j^\epsilon|)$ for every $j \in \{1, \dots, \ell\} \setminus \{i\}$, $h_i = 0$, estimate (28) provides a positive constant C such that

$$\left| \max_{\Omega_\epsilon(a_1^\epsilon, \dots, a_\ell^\epsilon)} p_i^\epsilon - \min_{\Omega_\epsilon(a_1^\epsilon, \dots, a_\ell^\epsilon)} p_i^\epsilon \right| \leq 2\pi\epsilon C.$$

which implies (32).

Let ζ_ϵ be a smooth cut-off function which is identically 1 inside $B_\epsilon(0)$ and 0 outside $B_{\sqrt{\epsilon}}(0)$, and whose support vanishes as $\epsilon \rightarrow 0$. From the minimality of p_i^ϵ , comparing with $c_1 = \dots = c_\ell = 0$ and

$$u(x) = - \sum_{j \neq i} \zeta_\epsilon(x - a_j^\epsilon) \log(|x - a_j^\epsilon|)$$

in (26), we obtain

$$(34) \quad \int_{\Omega_\epsilon(a_1^\epsilon, \dots, a_\ell^\epsilon)} |\nabla p_i^\epsilon|^2 dx \leq C \left(\int_{B_{\sqrt{\epsilon}}(0)} |\nabla \zeta_\epsilon|^2 dx + |\text{spt } \zeta_\epsilon| \right),$$

for some positive constant C independent of ϵ . By infimizing with respect to $\zeta_\epsilon \in \{\zeta \in C_c^\infty(B_{\sqrt{\epsilon}}(0)), \zeta \equiv 1 \text{ in } B_\epsilon(0)\}$, the integral in the right-hand side of (34) is the capacity of a disk of radius ϵ inside a disk of radius $\sqrt{\epsilon}$, which is equal to $4\pi/|\log \epsilon|$, therefore we obtain

$$\int_{\Omega_\epsilon(a_1^\epsilon, \dots, a_\ell^\epsilon)} |\nabla p_i^\epsilon|^2 dx \leq C\pi \left(\frac{4}{|\log \epsilon|} + \epsilon \right),$$

which implies (33). The lemma is proved. \square

Lemma 3.5. *Let a_1, \dots, a_n be distinct points in Ω and let $(a_1^\epsilon, \dots, a_n^\epsilon)$ be a sequence of points in Ω^n converging to (a_1, \dots, a_n) as $\epsilon \rightarrow 0$. Then, for every $i = 1, \dots, n$, as $\epsilon \rightarrow 0$ we have*

$$(35) \quad \|\nabla \bar{v}_{a_i^\epsilon}^\epsilon - \nabla v_{a_i}\|_{L^2(\Omega_\epsilon(\mathbf{a}); \mathbb{R}^2)} \rightarrow 0,$$

where $\bar{v}_{a_i^\epsilon}^\epsilon$ is defined in Remark 2.1 and v_{a_i} satisfies

$$(36) \quad \begin{cases} \Delta v_{a_i} = 0 & \text{in } \Omega, \\ v_{a_i} = g_{b_i} - \theta_{a_i} & \text{on } \partial\Omega. \end{cases}$$

Proof. Since the limit points a_1, \dots, a_n are all distinct and far from the boundary, we may assume, without loss of generality, $\epsilon < d := \min_{i \in \{1, \dots, n\}} d_i$ (see (4)), so that, for ϵ small enough, the disks $B_\epsilon(a_i^\epsilon)$ are pairwise disjoint and do not intersect $\partial\Omega$. Let now $i \in \{1, \dots, n\}$ be fixed. By Remark 2.1, the solution $\bar{v}_{a_i^\epsilon}^\epsilon$ can be written as $\bar{v}_{a_i^\epsilon}^\epsilon = \hat{u}_i^\epsilon + q_i^\epsilon$. Then to prove (35), we will show that, for every fixed $i \in \{1, \dots, n\}$,

$$(37a) \quad \|\nabla \hat{u}_i^\epsilon - \nabla v_{a_i^\epsilon}\|_{L^2(\Omega_\epsilon(\mathbf{a}^\epsilon); \mathbb{R}^2)} \rightarrow 0,$$

$$(37b) \quad \|\nabla q_i^\epsilon\|_{L^2(\Omega_\epsilon(\mathbf{a}^\epsilon); \mathbb{R}^2)} \rightarrow 0.$$

Note that in (35) we can replace v_{a_i} with $v_{a_i^\epsilon}$ thanks to the continuity of the map $C^0(\partial\Omega) \ni h \mapsto w_h \in H^1(\Omega)$, where w_h is harmonic in Ω with boundary datum h .

By Lemma 3.2, \hat{u}_i^ϵ minimizes the Dirichlet energy in $H^1(\Omega_\epsilon(\mathbf{a}^\epsilon))$ with boundary datum $g_{b_i^\epsilon} - \theta_{a_i^\epsilon}$ on $\partial\Omega$, and by (21) we have

$$\|\nabla \hat{u}_i^\epsilon\|_{L^2(\Omega_\epsilon(\mathbf{a}^\epsilon); \mathbb{R}^2)}^2 \leq C \|g_{b_i^\epsilon} - \theta_{a_i^\epsilon}\|_{H^{1/2}(\partial\Omega)}^2.$$

Let us consider the extension of \hat{u}_i^ϵ (not relabeled), which is harmonic in every $B_\epsilon(a_j^\epsilon)$. By (20) and (22), an estimate similar to (23) holds: there exists $C > 0$ independent of ϵ such that

$$\|\nabla \hat{u}_i^\epsilon\|_{L^2(B_\epsilon(a_j^\epsilon); \mathbb{R}^2)}^2 \leq C \epsilon^2 \|g_{b_i^\epsilon} - \theta_{a_i^\epsilon}\|_{L^\infty(\partial\Omega)}^2.$$

Since all the points a_i^ϵ are in $\Omega^{(d/2)}$, by the two estimates above, we obtain a uniform bound in $H^1(\Omega)$ for \hat{u}_i^ϵ , and therefore \hat{u}_i^ϵ has a subsequence that converges weakly to a function w_i in $H^1(\Omega)$ as $\epsilon \rightarrow 0$. By the lower semicontinuity of the H^1 norm and the minimality of \hat{u}_i^ϵ and v_{a_i} , it turns out that $w_i = v_{a_i}$ and

$$\lim_{\epsilon \rightarrow 0} \|\nabla \hat{u}_i^\epsilon\|_{L^2(\Omega; \mathbb{R}^2)} = \|\nabla v_{a_i}\|_{L^2(\Omega; \mathbb{R}^2)};$$

moreover, by the triangle inequality we have that

$$\|\nabla \hat{u}_i^\epsilon - \nabla v_{a_i^\epsilon}\|_{L^2(\Omega; \mathbb{R}^2)} \leq \|\nabla \hat{u}_i^\epsilon - \nabla v_{a_i}\|_{L^2(\Omega; \mathbb{R}^2)} + \|\nabla v_{a_i} - \nabla v_{a_i^\epsilon}\|_{L^2(\Omega; \mathbb{R}^2)},$$

with v_{a_i} solution to (36). The first term in the right-hand side above vanishes as $\epsilon \rightarrow 0$. The second term converges to zero thanks to the continuity of the map that associates v_a with the boundary datum $g - \theta_a$ (observe that $g - \theta_a^\epsilon \rightarrow g - \theta_a$ in $H^{1/2}(\partial\Omega)$), so that

$$(38) \quad \|\nabla \hat{u}_i^\epsilon - \nabla v_{a_i^\epsilon}\|_{L^2(\Omega; \mathbb{R}^2)} \rightarrow 0,$$

hence (37a) follows.

To obtain an estimate for the gradient of q_i^ϵ , we define $P_i^\epsilon := (\nabla q_i^\epsilon)^\perp = (-\partial_2 q_i^\epsilon, \partial_1 q_i^\epsilon)$ and notice that P_i^ϵ is a conservative vector field in $\Omega_\epsilon(\mathbf{a})$, namely its circulation vanishes along any closed loop contained in $\Omega_\epsilon(\mathbf{a}^\epsilon)$. Therefore, there exists $p_i^\epsilon \in H^1(\Omega_\epsilon(\mathbf{a}^\epsilon))$ such that $P_i^\epsilon = \nabla p_i^\epsilon$ and, in view of Remark 2.1

with the fact that $-K_{a_\epsilon} \cdot \nu = \partial \log(|x - a_\epsilon|)/\partial \tau$, p_i^ϵ is a solution to (31). Then by (33) we deduce that the gradient of $p_i^\epsilon \rightarrow 0$ as $\epsilon \rightarrow 0$ and since the gradients of p_i^ϵ and q_i^ϵ have the same modulus, we obtain (37b). \square

Remark 3.1. From the proof of Lemma 3.5 we derive some important properties of the functions \hat{u}_i^ϵ and p_i^ϵ introduced in Remark 2.1 and (31), respectively. Let a_1, \dots, a_n be distinct points in Ω and let $(a_1^\epsilon, \dots, a_n^\epsilon)$ be a sequence of points in Ω^n converging to (a_1, \dots, a_n) as $\epsilon \rightarrow 0$. Then, for every $i = 1, \dots, n$, as $\epsilon \rightarrow 0$, we have

- (i) the function \hat{u}_i^ϵ admits an extension which converges to v_{a_i} strongly in $H^1(\Omega)$. This follows by combining the fact that $v_{a_i^\epsilon} \rightarrow v_{a_i}$ strongly in $H^1(\Omega)$, \hat{u}_i^ϵ and $v_{a_i^\epsilon}$ agree on the boundary, the Poincaré inequality, and property (38);
- (ii) the function p_i^ϵ admits a (not relabeled) extension $p_i^\epsilon \in H^1(\Omega)$. In fact, recalling that $\phi_{a_i^\epsilon}(\cdot) = \log(|\cdot - a_i^\epsilon|)$, since $p_i^\epsilon = \phi_{a_i^\epsilon} + c_j$ on $\partial B_\epsilon(a_j^\epsilon)$ for $j \neq i$ and $p_i^\epsilon = c_i$ on $\partial B_\epsilon(a_i^\epsilon)$ for some constants c_1, \dots, c_n , we define $p_i^\epsilon = \phi_{a_i^\epsilon} + c_j$ on $B_\epsilon(a_j^\epsilon)$ and $p_i^\epsilon \equiv c_i$ on $B_\epsilon(a_i^\epsilon)$. Thus, up to subtracting a constant, thanks to (32), p_i^ϵ is observed to tend to 0 strongly in $H^1(\Omega)$.

Lemma 3.6. *Let a_1, \dots, a_ℓ be distinct points in $\bar{\Omega}$ and let the functions $\bar{v}_{a_i}^\epsilon \in H^1(\Omega_\epsilon(a_1, \dots, a_\ell))$ solve the system in Remark 2.1 for some $i \in \{1, \dots, \ell\}$. Then the function $\bar{v}_{a_i}^\epsilon$ admits an extension (not relabeled) in $H^1(\Omega)$. In particular, there exist $\bar{\epsilon} > 0$ and a constant $C > 0$ such that*

$$\int_{B_\epsilon(a_j)} |\nabla \bar{v}_{a_i}^\epsilon|^2 dx \leq C,$$

for all $\epsilon < \bar{\epsilon}$ and for every $j \in \{1, \dots, \ell\}$.

Proof. We consider the domain $E := \mathbb{R}^2 \setminus (\cup_{i=1}^\ell \bar{B}_\epsilon(a_i) \cup \bar{\Omega})$ and we extend $\bar{v}_{a_i}^\epsilon$ on E by defining it as the solution to the Dirichlet problem with datum $g_{b_i} - \omega(a_i)\theta_{a_i}$ on $\partial E \cap \partial\Omega$. It is easy to check (one can use, for instance, (21)) that there exists a constant $C > 0$ such that

$$\int_E |\nabla \bar{v}_{a_i}^\epsilon|^2 dx \leq C.$$

Let $\bar{\epsilon} := \frac{1}{4} \min\{|a_i - a_j| : i, j \in \{1, \dots, \ell\}, i \neq j\}$. For $\epsilon < \bar{\epsilon}$, consider a family of functions $\zeta_\epsilon \in C^\infty(B_\epsilon(0))$ which are zero in a neighborhood of $B_{\epsilon/2}(0)$, equal to 1 in a neighborhood of $\partial B_\epsilon(0)$, and such that $\|\nabla \zeta_\epsilon\|_{L^2(B_\epsilon(0))}$ is uniformly bounded. We shall exploit the fact that $\bar{v}_{a_i}^\epsilon$ is defined in the annulus $A_\epsilon^{2\epsilon}(a_j)$ for every $j \in \{1, \dots, \ell\}$ to define its extension in $B_\epsilon(a_j)$. To this aim, consider the inversion function $\mathbb{I}_{a_j}^\epsilon : \mathbb{C} \setminus \{a_j\} \rightarrow \mathbb{C} \setminus \{a_j\}$ given by $\mathbb{I}_{a_j}^\epsilon(x) := \epsilon^2(x - a_j)/|x - a_j|^2$, and define

$$\bar{v}_{a_i}^\epsilon(x) := \begin{cases} \zeta_\epsilon(x) \bar{v}_{a_i}^\epsilon(\mathbb{I}_{a_j}^\epsilon(x)) & \text{if } x \in \mathbb{I}_{a_j}^\epsilon(A_\epsilon^{2\epsilon}(a_j)), \\ 0 & \text{elsewhere in } B_\epsilon(a_j). \end{cases}$$

Also in this case, an easy check shows that

$$\int_{B_\epsilon(a_j)} |\nabla \bar{v}_{a_i}^\epsilon|^2 dx \leq C \int_{A_\epsilon^{2\epsilon}(a_j)} |\nabla \bar{v}_{a_i}^\epsilon|^2 dx \leq C,$$

for some constant $C > 0$ independent of $\bar{\epsilon}$, and for all $\epsilon < \bar{\epsilon}$. The lemma is proved. \square

4 - The limit as $\epsilon \rightarrow 0$

This section is devoted to the proof of the main results, Theorem 1.1 and Corollary 1.1, which is presented in Subsection 4.4. Our proof strategy will rely on the results in the case of one dislocation $a \in \bar{\Omega}$, which we treat next in Subsection 4.1. To study the asymptotic behavior of the rescaled energies (3), we distinguish two different scenarios (recall the function d_i defined in (4)):

- all the limit points are in the interior of Ω and are all distinct, namely $\min_i d_i > 0$ (treated in Subsection 4.2);
- either at least two limit points coincide or one limit point is on the boundary $\partial\Omega$, namely $\min_i d_i = 0$ (treated in Subsection 4.3).

4.1 - The case $n = 1$

Given $a \in \bar{\Omega}$, the energy $\mathcal{E}_\epsilon(a) := \mathcal{E}_\epsilon^{(1)}(a)$ in (1) reads

$$\mathcal{E}_\epsilon(a) = \min \left\{ \frac{1}{2} \int_{\Omega_\epsilon(a)} |F|^2 dx : F \in L^2(\Omega_\epsilon(a); \mathbb{R}^2), \operatorname{curl} F = 0, \right. \\ \left. F \cdot \tau = f \text{ on } \partial\Omega \setminus \bar{B}_\epsilon(a) \right\},$$

and its rescaling $\mathcal{F}_\epsilon(a) := \mathcal{F}_\epsilon^{(1)}(a)$ in (3) reads

$$(39) \quad \mathcal{F}_\epsilon(a) = \frac{1}{2} \int_{\Omega_\epsilon(a)} |\nabla u_a^\epsilon|^2 dx - \pi |\log \epsilon|,$$

where u_a^ϵ solution to (12) with $n = 1$, namely

$$(40) \quad \begin{cases} \Delta u_a^\epsilon = 0 & \text{in } \Omega_\epsilon(a) \setminus \Sigma, \\ [u_a^\epsilon] = 2\pi & \text{on } \Sigma \cap \Omega_\epsilon(a), \\ u_a^\epsilon = g & \text{on } \partial\Omega \setminus \bar{B}_\epsilon(a), \\ \partial u_a^\epsilon / \partial \nu = 0 & \text{on } \partial B_\epsilon(a) \cap \Omega, \\ \partial (u_a^\epsilon)^+ / \partial \nu = \partial (u_a^\epsilon)^- / \partial \nu & \text{on } \Sigma \cap \Omega_\epsilon(a). \end{cases}$$

In the last equation above, ν is a choice of the unit normal vector to Σ . Notice that, when $\text{dist}(a, \partial\Omega) \leq \epsilon$, the choice of $b \in \partial\Omega \cap B_\epsilon(a)$ implies that $\Omega_\epsilon(a) \setminus \Sigma = \Omega_\epsilon(a)$ and the jump condition of u_a^ϵ across Σ is empty.

In this case, the decomposition (15) reads

$$(41) \quad \nabla u_a^\epsilon = \omega(a)K_a + \nabla \bar{u}_a^\epsilon,$$

where $\bar{u}_a^\epsilon \in H^1(\Omega_\epsilon(a))$ is the solution to

$$(42) \quad \begin{cases} \Delta \bar{u}_a^\epsilon = 0 & \text{in } \Omega_\epsilon(a), \\ \bar{u}_a^\epsilon = g - \omega(a)\theta_a & \text{on } \partial\Omega \setminus \bar{B}_\epsilon(a), \\ \partial \bar{u}_a^\epsilon / \partial \nu = 0 & \text{on } \partial B_\epsilon(a). \end{cases}$$

Notice that this function is the same as the \bar{v}_a^ϵ in Remark 2.1 when $n = 1$, thanks to (8).

Given $a^\epsilon \rightarrow a$ as $\epsilon \rightarrow 0$, the asymptotics of $\mathcal{F}_\epsilon(a^\epsilon)$ when $a \in \Omega$ is dealt with in Proposition 4.1, whereas the case in which $a \in \partial\Omega$ is dealt with in Proposition 4.2. These two results readily imply Theorem 1.1 in the case $n = 1$ (the continuity of the limit functional \mathcal{F} follows from the relationship between continuous convergence and Γ -convergence, see [8, Remark 4.9]).

Proposition 4.1. *Let $a \in \Omega$ and v_a satisfy (36). For every sequence $a^\epsilon \rightarrow a$ as $\epsilon \rightarrow 0$ we have*

$$(43) \quad \mathcal{F}_\epsilon(a^\epsilon) \rightarrow \pi \log d(a) + \frac{1}{2} \int_{\Omega_{d(a)}(a)} |K_a + \nabla v_a|^2 dx + \frac{1}{2} \int_{B_{d(a)}(a)} |\nabla v_a|^2 dx.$$

Proof. Since $a \in \Omega$, $\omega(a) = 1$. Since $d(a^\epsilon) \rightarrow d(a)$ and $d(a) > 0$, we can take ϵ so small so that $\epsilon < \min\{d(a^\epsilon), 1\}$. By plugging (41) into (39), $\mathcal{F}_\epsilon(a^\epsilon)$ reads

$$\mathcal{F}_\epsilon(a^\epsilon) = \frac{1}{2} \int_{\Omega_\epsilon(a^\epsilon)} |\nabla \bar{u}_{a^\epsilon}^\epsilon|^2 dx + \int_{\Omega_\epsilon(a^\epsilon)} K_{a^\epsilon} \cdot \nabla \bar{u}_{a^\epsilon}^\epsilon dx + \frac{1}{2} \int_{\Omega_\epsilon(a^\epsilon)} |K_{a^\epsilon}|^2 dx + \pi \log \epsilon.$$

Set, for brevity, $d := d(a)$, $d^\epsilon := d(a^\epsilon)$, $K := K_a$, and $K^\epsilon := K_{a^\epsilon}$. Writing

$$\begin{aligned} \frac{1}{2} \int_{\Omega_\epsilon(a^\epsilon)} |K^\epsilon|^2 dx &= \frac{1}{2} \int_{\Omega_{d^\epsilon}(a^\epsilon)} |K^\epsilon|^2 dx + \frac{1}{2} \int_{A_{d^\epsilon}^\epsilon(a^\epsilon)} |K^\epsilon|^2 dx \\ &= \frac{1}{2} \int_{\Omega_{d^\epsilon}(a^\epsilon)} |K^\epsilon|^2 dx + \pi \log \frac{d^\epsilon}{\epsilon}, \end{aligned}$$

we obtain

$$(44) \quad \begin{aligned} \mathcal{F}_\epsilon(a^\epsilon) &= \pi \log d^\epsilon + \frac{1}{2} \int_{\Omega_{d^\epsilon}(a^\epsilon)} |K^\epsilon|^2 dx + \int_{\Omega_\epsilon(a^\epsilon)} K^\epsilon \cdot \nabla \bar{u}_{a^\epsilon}^\epsilon dx \\ &\quad + \frac{1}{2} \int_{\Omega_\epsilon(a^\epsilon)} |\nabla \bar{u}_{a^\epsilon}^\epsilon|^2 dx. \end{aligned}$$

If we prove that, as $\epsilon \rightarrow 0$,

$$(45a) \quad \pi \log d^\epsilon + \frac{1}{2} \int_{\Omega_{d^\epsilon}(a^\epsilon)} |K^\epsilon|^2 dx \rightarrow \pi \log d + \frac{1}{2} \int_{\Omega_d(a)} |K|^2 dx,$$

$$(45b) \quad \int_{\Omega_\epsilon(a^\epsilon)} K^\epsilon \cdot \nabla \bar{u}_{a^\epsilon}^\epsilon dx \rightarrow \int_{\Omega_d(a)} K \cdot \nabla v_a dx,$$

$$(45c) \quad \frac{1}{2} \int_{\Omega_\epsilon(a^\epsilon)} |\nabla \bar{u}_{a^\epsilon}^\epsilon|^2 dx \rightarrow \frac{1}{2} \int_{\Omega} |\nabla v_a|^2 dx,$$

where v_a satisfies (36), then (43) follows.

Since $d^\epsilon \rightarrow d$, $K^\epsilon \chi_{B_{d^\epsilon}(a^\epsilon)}$ converges pointwise to $K \chi_{B_d(a)}$, and $K^\epsilon \chi_{B_{d^\epsilon}(a^\epsilon)}$ are uniformly bounded for ϵ small enough, (45a) follows by the Dominated Convergence Theorem.

To prove (45b), we integrate by parts to obtain

$$\int_{\Omega_\epsilon(a^\epsilon)} K^\epsilon \cdot \nabla \bar{u}_{a^\epsilon}^\epsilon dx = \int_{\partial\Omega} (K^\epsilon \cdot \nu)(g - \theta_{a^\epsilon}) dx,$$

where we have used (8) and the fact that $\bar{u}_{a^\epsilon}^\epsilon = g - \theta_{a^\epsilon}$ on $\partial\Omega$ (see (42)). Since K^ϵ and θ_{a^ϵ} are uniformly bounded in ϵ on the set $\partial\Omega$, and converge pointwise to K and θ_a , respectively, by the Dominated Convergence Theorem, we have

$$\int_{\Omega_\epsilon(a^\epsilon)} K^\epsilon \cdot \nabla \bar{u}_{a^\epsilon}^\epsilon dx \rightarrow \int_{\Omega} K \cdot \nabla v_a dx = \int_{\Omega_d(a)} K \cdot \nabla v_a dx,$$

which gives (45b). The last equality follows by the Divergence Theorem, combined with (8).

It remains to prove (45c). To do this, consider the harmonic extension w^ϵ of $\bar{u}_{a^\epsilon}^\epsilon$ inside $B_\epsilon(a^\epsilon)$. By applying Lemma 3.3 to w^ϵ with a replaced by a^ϵ , estimate (23) reads

$$(46) \quad \int_{B_\epsilon(a^\epsilon)} |\nabla w^\epsilon(x)|^2 dx \leq C\epsilon^2 \|g - \theta_{a^\epsilon}\|_{L^\infty(\partial\Omega)}^2,$$

which implies that

$$(47) \quad \|w^\epsilon\|_{H^1(B_\epsilon(a^\epsilon))} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0.$$

By combining (46) with (21), we have

$$\begin{aligned} \int_{\Omega} |\nabla w^\epsilon|^2 dx &= \int_{\Omega_\epsilon(a^\epsilon)} |\nabla \bar{u}_{a^\epsilon}^\epsilon|^2 dx + \int_{B_\epsilon(a^\epsilon)} |\nabla w^\epsilon|^2 dx \leq C \|g - \theta_{a^\epsilon}\|_{H^{1/2}(\partial\Omega)}^2 \\ &\quad + C\epsilon^2 \|g - \theta_{a^\epsilon}\|_{L^\infty(\partial\Omega)}^2. \end{aligned}$$

Therefore, letting $\epsilon \rightarrow 0$, we obtain

$$\limsup_{\epsilon \rightarrow 0} \int_{\Omega} |\nabla w^\epsilon|^2 dx \leq C \|g - \theta_a\|_{H^{1/2}(\partial\Omega)}^2,$$

which, together with Poincaré inequality, implies that w^ϵ is uniformly bounded in $H^1(\Omega)$. As a consequence there exists $w \in H^1(\Omega)$ such that (up to subsequences) $w^\epsilon \rightharpoonup w$ weakly in $H^1(\Omega)$. Since $w^\epsilon = g - \theta_{a^\epsilon}$ on $\partial\Omega$ for every ϵ , then $w = g - \theta_a$ on $\partial\Omega$. Since the \bar{u}_{a^ϵ} in (42) is the minimizer of the Dirichlet energy, we have that

$$\frac{1}{2} \int_{\Omega_\epsilon(a^\epsilon)} |\nabla \bar{u}_{a^\epsilon}^\epsilon|^2 dx \leq \frac{1}{2} \int_{\Omega_\epsilon(a^\epsilon)} |\nabla v_{a^\epsilon}|^2 dx \leq \frac{1}{2} \int_{\Omega} |\nabla v_{a^\epsilon}|^2 dx,$$

the lower semicontinuity of the H^1 norm, together with (47), gives

$$(48) \quad \begin{aligned} \frac{1}{2} \int_{\Omega} |\nabla w|^2 dx &\leq \liminf_{\epsilon \rightarrow 0} \frac{1}{2} \int_{\Omega} |\nabla w^\epsilon|^2 dx \leq \lim_{\epsilon \rightarrow 0} \frac{1}{2} \int_{\Omega} |\nabla v_{a^\epsilon}|^2 dx \\ &= \frac{1}{2} \int_{\Omega} |\nabla v_a|^2 dx, \end{aligned}$$

which implies that $w = v_a$, by the uniqueness of the minimizer v_a . Thus, all the inequalities in (48) are in fact equalities. This, together with (47), gives (45c) and completes the proof. \square

Proposition 4.2. *Let $a \in \partial\Omega$ and a^ϵ be a sequence of points in $\bar{\Omega}$ converging to a as $\epsilon \rightarrow 0$. Then there exist two constants $C_1, C_2 > 0$ independent of ϵ , a^ϵ , and a , such that*

$$(49) \quad \mathcal{F}_\epsilon(a^\epsilon) \geq C_1 |\log(\max\{\epsilon, d(a^\epsilon)\})| + C_2,$$

for every ϵ small enough. In particular, $\mathcal{F}_\epsilon(a^\epsilon) \rightarrow +\infty$ as $\epsilon \rightarrow 0$.

Proof. Since $a \in \partial\Omega$, $\omega(a) = 2$. Let α and ϵ_α be as in (7), let $\epsilon < \min\{\epsilon_\alpha, 1\}$, and $d^\epsilon < \epsilon_\alpha/2$, where we set, for brevity, $d := d(a)$ and $d^\epsilon := d(a^\epsilon)$. We distinguish two possible scenarios: the *slow collision* $\epsilon < d^\epsilon$ and the *fast collision* $\epsilon \geq d^\epsilon$. In the former case $\epsilon < d^\epsilon$, exploiting (44) we get

$$(50) \quad \begin{aligned} \mathcal{F}_\epsilon(a^\epsilon) &\geq \pi \log d^\epsilon + \frac{1}{2} \int_{\Omega_{d^\epsilon}(a^\epsilon)} |K_{a^\epsilon}|^2 dx + \int_{\Omega_{d^\epsilon}(a^\epsilon)} K_{a^\epsilon} \cdot \nabla \bar{u}_{a^\epsilon}^\epsilon dx \\ &\quad + \frac{1}{2} \int_{\Omega_{d^\epsilon}(a^\epsilon)} |\nabla \bar{u}_{a^\epsilon}^\epsilon|^2 dx \\ &= \pi \log d^\epsilon + \frac{1}{2} \int_{\Omega_{d^\epsilon}(a^\epsilon)} |K_{a^\epsilon} + \nabla \bar{u}_{a^\epsilon}^\epsilon|^2 dx. \end{aligned}$$

where we have used that $\int_{\Omega_\epsilon(a^\epsilon)} K_{a^\epsilon} \cdot \nabla \bar{u}_{a^\epsilon}^\epsilon dx = \int_{\Omega_{d^\epsilon}(a^\epsilon)} K_{a^\epsilon} \cdot \nabla \bar{u}_{a^\epsilon}^\epsilon dx$. By Lemma 2.1, we may assume that the discontinuity point b^ϵ of the boundary datum g is one of the projections of a^ϵ on $\partial\Omega$, so that $b^\epsilon \in \partial B_{d^\epsilon}(a^\epsilon) \cap \partial\Omega$. In

particular, $\Omega_{2d^\epsilon}(b^\epsilon) \subset \Omega_{d^\epsilon}(a^\epsilon)$, so that

$$\begin{aligned}
& \int_{\Omega_{d^\epsilon}(a^\epsilon)} |K_{a^\epsilon} + \nabla \bar{u}_{a^\epsilon}^\epsilon|^2 dx \geq \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |K_{a^\epsilon} + \nabla \bar{u}_{a^\epsilon}^\epsilon|^2 dx \\
& \geq \inf \left\{ \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |\nabla u|^2 dx : u \in H^1(\Omega_{2d^\epsilon}(b^\epsilon)), u = g \text{ on } \partial\Omega \setminus \bar{B}_{2d^\epsilon}(b^\epsilon) \right\} \\
(51) \quad & = \inf \left\{ \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |2K_{b^\epsilon} + \nabla u|^2 dx : u \in H^1(\Omega_{2d^\epsilon}(b^\epsilon)), \right. \\
& \qquad \qquad \qquad \left. u = g - 2\theta_{b^\epsilon} \text{ on } \partial\Omega \setminus \bar{B}_{2d^\epsilon}(b^\epsilon) \right\} \\
& = \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |2K_{b^\epsilon} + \nabla \bar{u}|^2 dx,
\end{aligned}$$

where \bar{u} solves

$$\begin{cases} \Delta \bar{u} = 0 & \text{in } \Omega_{2d^\epsilon}(b^\epsilon), \\ \bar{u} = g - 2\theta_{b^\epsilon} & \text{on } \partial\Omega \setminus \bar{B}_{2d^\epsilon}(b^\epsilon), \\ \partial \bar{u} / \partial \nu = -2K_{b^\epsilon} \cdot \nu & \text{on } \partial B_{2d^\epsilon}(b^\epsilon) \cap \Omega. \end{cases}$$

Since $K_{b^\epsilon} \cdot \nu = 0$ on $\partial B_{2d^\epsilon}(b^\epsilon)$ by (8), it follows by uniqueness that $\bar{u} = \bar{u}_b^\epsilon$, where $\bar{u}_b^\epsilon \in H^1(\Omega_\epsilon(b))$ is the solution to (42). Therefore, using Young's inequality and (21) (with $h = g - 2\theta_b$), recalling that $\epsilon < d^\epsilon$, we may bound (51) as follows:

$$\begin{aligned}
& \frac{1}{2} \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |2K_{b^\epsilon} + \nabla \bar{u}|^2 dx \\
(52) \quad & \geq 2 \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |K_{b^\epsilon}|^2 dx + 2 \int_{\Omega_{2d^\epsilon}(b^\epsilon)} K_{b^\epsilon} \cdot \nabla \bar{u} dx \\
& \geq 2 \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |K_{b^\epsilon}|^2 dx - \lambda \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |K_{b^\epsilon}|^2 dx - \frac{1}{\lambda} \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |\nabla \bar{u}|^2 dx \\
& \geq (2 - \lambda) \int_{\Omega_{2d^\epsilon}(b^\epsilon)} |K_{b^\epsilon}|^2 dx - \frac{C}{\lambda} \max_{b \in \partial\Omega} \|g - 2\theta_b\|_{H^{1/2}(\partial\Omega)}^2,
\end{aligned}$$

where $\lambda > 0$ is a constant that will be chosen later. Now, in view of the assumption on d^ϵ at the beginning of the proof, the set $\Omega_{2d^\epsilon}(b^\epsilon)$ contains a sector, which, in polar coordinates centered at b^ϵ , is the rectangle $R := (2d^\epsilon, \epsilon_\alpha) \times (\phi_0, \phi_1)$ with $\phi_1 - \phi_0 = \alpha$. Therefore, in the case $\epsilon < d^\epsilon$, by combining (50) and (51) with (52), and using the estimate from below in (18), it follows that

$$\mathcal{F}_\epsilon(a^\epsilon) \geq (\pi - (2 - \lambda)\alpha) \log d^\epsilon + (2 - \lambda)\alpha \log \frac{\epsilon_\alpha}{2} - \frac{C}{\lambda} \max_{b \in \partial\Omega} \|g - 2\theta_b\|_{H^{1/2}(\partial\Omega)}^2.$$

Recalling that $\alpha > \pi/2$ by (7), we can choose $\lambda = (\alpha - \pi/2)/\alpha$, so that the inequality above can be written as

$$(53) \quad \mathcal{F}_\epsilon(a^\epsilon) \geq C_1 |\log d^\epsilon| + C_2,$$

with

$$(54) \quad C_1 := \left(\alpha - \frac{\pi}{2}\right) \text{ and } C_2 := \left(\alpha + \frac{\pi}{2}\right) \log \frac{\epsilon_\alpha}{2} - \frac{C\alpha}{\alpha - \pi/2} \max_{b \in \partial\Omega} \|g - 2\theta_b\|_{H^{1/2}(\partial\Omega)}^2.$$

Moreover, in the case $\epsilon \geq d^\epsilon$, we consider the projection b^ϵ of a^ϵ on $\partial\Omega$, so that $b^\epsilon \in \partial B_\epsilon(a^\epsilon) \cap \partial\Omega$, and the disk $B_{2\epsilon}(b^\epsilon)$ contains $B_\epsilon(a^\epsilon)$. Arguing as in (51), (52), and using (18) with d^ϵ replaced by ϵ , we can estimate (39) as follows

$$(55) \quad \mathcal{F}_\epsilon(a^\epsilon) \geq \frac{1}{2} \int_{\Omega_{2\epsilon}(b^\epsilon)} |2K_{b^\epsilon} + \nabla \bar{u}|^2 dx - \pi |\log \epsilon| \geq C_1 |\log \epsilon| + C_2,$$

with the same constants C_1 and C_2 provided in (54).

Therefore, by (53) and (55) for every $\epsilon < \min\{\epsilon_\alpha, 1\}$, the thesis (49) follows.

□

4.2 - The case $n > 1$ with $\min_i d_i > 0$

In this case, it is convenient to write the rescaled energy (3) as the sum of the rescaled energy of each dislocation plus a remainder term accounting for interactions: recalling the expression (39) for the rescaled energy of one dislocation and the decomposition (41), we write the energy (3) as

$$(56) \quad \mathcal{F}_\epsilon^{(n)}(\mathbf{a}) = \sum_{i=1}^n \mathcal{F}_\epsilon(a_i) + \sum_{i=1}^n \mathcal{R}_\epsilon(a_i) + \sum_{i \neq j} \mathcal{G}_\epsilon(a_i, a_j),$$

where, for every $i = 1, \dots, n$, $\mathcal{F}_\epsilon(a_i)$ is given by (39),

$$(57) \quad \mathcal{R}_\epsilon(a_i) := \frac{1}{2} \int_{\Omega_\epsilon(\mathbf{a})} |K_{a_i} + \nabla \bar{v}_{a_i}^\epsilon|^2 dx - \frac{1}{2} \int_{\Omega_\epsilon(a_i)} |K_{a_i} + \nabla \bar{u}_{a_i}^\epsilon|^2 dx,$$

and, for every $i, j = 1, \dots, n$, with $i \neq j$,

$$\mathcal{G}_\epsilon(a_i, a_j) := \int_{\Omega_\epsilon(\mathbf{a})} (K_{a_i} + \nabla \bar{v}_{a_i}^\epsilon) \cdot (K_{a_j} + \nabla \bar{v}_{a_j}^\epsilon) dx,$$

$\bar{u}_{a_i}^\epsilon$ being the solution to (42) associated with a_i , and $\bar{v}_{a_i}^\epsilon$ being as in Remark 2.1.

Proposition 4.3. *Let $\mathbf{a} = (a_1, \dots, a_n) \in \Omega^n$ be an n -tuple of distinct points and let \mathbf{a}^ϵ be a sequence converging in Ω^n to \mathbf{a} as $\epsilon \rightarrow 0$. Then, for every $i = 1, \dots, n$, we have $\mathcal{R}_\epsilon(a_i^\epsilon) \rightarrow 0$ as $\epsilon \rightarrow 0$.*

Proof. For brevity, we define the d_i^ϵ 's as in (4), associated with the family $\{a_1^\epsilon, \dots, a_n^\epsilon\}$, and the d_i 's associated with the family $\{a_1, \dots, a_n\}$. By assumption $d_i^\epsilon \rightarrow d_i > 0$ for every $i = 1, \dots, n$; therefore, without loss of generality, we may take $\epsilon < \min_i d_i$. In particular, the disks $B_\epsilon(a_i^\epsilon)$ are all contained in Ω and are pairwise disjoint for every $\epsilon \in (0, \min_i d_i)$.

Fix now $i \in \{1, \dots, n\}$. The remainder (57) evaluated at a_i^ϵ can be written as $\mathcal{R}_\epsilon(a_i^\epsilon) = \mathcal{R}'_\epsilon(a_i^\epsilon) - \mathcal{R}''_\epsilon(a_i^\epsilon)$, with

$$\mathcal{R}'_\epsilon(a_i^\epsilon) := \frac{1}{2} \int_{\Omega_\epsilon(\mathbf{a}^\epsilon)} (2K_i^\epsilon + \nabla \bar{v}_i^\epsilon + \nabla \bar{u}_i^\epsilon) \cdot (\nabla \bar{v}_i^\epsilon - \nabla \bar{u}_i^\epsilon) dx,$$

and

$$\mathcal{R}''_\epsilon(a_i^\epsilon) := \frac{1}{2} \sum_{j \neq i} \int_{B_\epsilon(a_j^\epsilon)} |K_i^\epsilon + \nabla \bar{u}_i^\epsilon|^2 dx,$$

where, for brevity, we have replaced the subscript a_j^ϵ with the subscript j coupled with the superscript ϵ . By the Divergence Theorem, (42), and Remark 2.1, we have

$$\begin{aligned} \mathcal{R}'_\epsilon(a_i^\epsilon) &= \frac{1}{2} \int_{\partial\Omega_\epsilon(\mathbf{a}^\epsilon)} (-\partial \bar{v}_i^\epsilon / \partial \nu + \partial \bar{u}_i^\epsilon / \partial \nu) (\bar{v}_i^\epsilon - \bar{u}_i^\epsilon) dx \\ &= \frac{1}{2} \int_{\Omega_\epsilon(\mathbf{a}^\epsilon)} (-\nabla \bar{v}_i^\epsilon + \nabla \bar{u}_i^\epsilon) \cdot (\nabla \bar{v}_i^\epsilon - \nabla \bar{u}_i^\epsilon) dx \\ &= -\frac{1}{2} \int_{\Omega_\epsilon(\mathbf{a}^\epsilon)} |\nabla \bar{v}_i^\epsilon - \nabla \bar{u}_i^\epsilon|^2 dx. \end{aligned}$$

Then, $\mathcal{R}'_\epsilon(a_i^\epsilon) \leq \|\nabla \bar{v}_i^\epsilon - \nabla \bar{v}_i^\epsilon\|_{L^2(\Omega_\epsilon(\mathbf{a}^\epsilon); \mathbb{R}^2)} + \|\nabla \bar{v}_i^\epsilon - \nabla \bar{u}_i^\epsilon\|_{L^2(\Omega_\epsilon(\mathbf{a}^\epsilon); \mathbb{R}^2)}$, which, in view of Lemma 3.5, converges to 0, as $\epsilon \rightarrow 0$. Moreover, in view of Lemma 3.2 (with $n = 1$ and $h = g_{b_i^\epsilon} - \theta_i^\epsilon$) and the fact that $\|K_i^\epsilon\|_{L^\infty(B_\epsilon(a_j^\epsilon); \mathbb{R}^2)} \leq 1/(2d_i^\epsilon - \epsilon)$ for every $j \neq i$, we may bound $\mathcal{R}''_\epsilon(a_i^\epsilon)$ as follows:

$$\begin{aligned} \mathcal{R}''_\epsilon(a_i^\epsilon) &\leq \sum_{j \neq i} \int_{B_\epsilon(a_j^\epsilon)} (|K_i^\epsilon|^2 + |\nabla \bar{u}_i^\epsilon|^2) dx \\ &\leq \pi(n-1)\epsilon^2 \left(\frac{1}{(2d_i^\epsilon - \epsilon)^2} + C \|g_{b_i^\epsilon} - \theta_i^\epsilon\|_{L^\infty(\partial\Omega)}^2 \right), \end{aligned}$$

for some constant $C > 0$ independent of ϵ . In particular $\mathcal{R}''_\epsilon(a_i^\epsilon)$ tends to zero as $\epsilon \rightarrow 0$. This concludes the proof of the proposition. \square

Proposition 4.4. *Let $\mathbf{a} = (a_1, \dots, a_n) \in \Omega^n$ be an n -tuple of distinct points and let \mathbf{a}^ϵ be a sequence converging in Ω^n to \mathbf{a} as $\epsilon \rightarrow 0$. Then, for every $i, j = 1, \dots, n$, with $i \neq j$, we have*

$$(58) \quad \mathcal{G}_\epsilon(a_i^\epsilon, a_j^\epsilon) \rightarrow \int_{\Omega} (K_{a_i} + \nabla v_{a_i}) \cdot (K_{a_j} + \nabla v_{a_j}) dx, \quad \text{as } \epsilon \rightarrow 0,$$

v_{a_i} being the solution to (36) associated with a_i .

Proof. For brevity, we define the d_i^ϵ 's as in (4), associated with the family $\{a_1^\epsilon, \dots, a_n^\epsilon\}$, and the d_i 's associated with the family $\{a_1, \dots, a_n\}$. Fix $i, j \in \{1, \dots, n\}$, $i \neq j$. By assumption $d_i^\epsilon \rightarrow d_i > 0$ and $d_j^\epsilon \rightarrow d_j > 0$; therefore,

without loss of generality, we may take $\epsilon < \min\{d_i, d_j\}$. Since the limit points are distinct, in view of Lemma 3.5, we have

$$\chi_{\Omega_\epsilon(\mathbf{a}^\epsilon)} \nabla \bar{v}_{a_i^\epsilon}^\epsilon \rightarrow \nabla v_{a_i} \quad \text{and} \quad \chi_{\Omega_\epsilon(\mathbf{a}^\epsilon)} \nabla \bar{v}_{a_j^\epsilon}^\epsilon \rightarrow \nabla v_{a_j} \quad \text{strongly in } L^2(\Omega; \mathbb{R}^2),$$

so that

$$(59) \quad \int_{\Omega_\epsilon(\mathbf{a}^\epsilon)} \nabla \bar{v}_{a_i^\epsilon}^\epsilon \cdot \nabla \bar{v}_{a_j^\epsilon}^\epsilon \, dx \rightarrow \int_{\Omega} \nabla v_{a_i} \cdot \nabla v_{a_j} \, dx.$$

Setting for brevity $d := \min\{d_i, d_j\}$, we decompose the domain of integration as

$$\Omega_\epsilon(\mathbf{a}^\epsilon) = (\Omega_d(a_i^\epsilon, a_j^\epsilon) \cup B_d(a_i^\epsilon) \cup B_d(a_j^\epsilon)) \setminus \bigcup_{k=1}^n \bar{B}_\epsilon(a_k^\epsilon).$$

Since $K_{a_i^\epsilon} \rightarrow K_{a_i}$ and $K_{a_j^\epsilon} \rightarrow K_{a_j}$ a.e. in Ω , by the Dominated Convergence Theorem it is easy to see that, as $\epsilon \rightarrow 0$,

$$(60) \quad \int_{\Omega_d(a_i^\epsilon, a_j^\epsilon)} K_{a_i^\epsilon} \cdot K_{a_j^\epsilon} \, dx \rightarrow \int_{\Omega_d(a_i, a_j)} K_{a_i} \cdot K_{a_j} \, dx$$

and

$$(61) \quad \begin{aligned} \int_{B_d(a_i^\epsilon)} K_{a_i^\epsilon} \cdot K_{a_j^\epsilon} \, dx &= \int_{B_d(a_i)} \frac{\hat{\theta}_{a_i} \cdot \hat{\theta}_{a_j^\epsilon - (a_i^\epsilon - a_i)}}{|x - a_i| |x - a_j^\epsilon + (a_i^\epsilon - a_i)|} \, dx \\ &\rightarrow \int_{B_d(a_i)} K_{a_i} \cdot K_{a_j} \, dx. \end{aligned}$$

An analogous result holds exchanging the roles of i and j . Eventually, since the limit points a_1, \dots, a_n are distinct, we have, for $k \neq i, j$,

$$(62) \quad \int_{B_\epsilon(a_k^\epsilon)} |K_{a_i^\epsilon} \cdot K_{a_j^\epsilon}| \, dx \leq \frac{\pi \epsilon^2}{(2d_i^\epsilon - \epsilon)(2d_j^\epsilon - \epsilon)} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0,$$

and

$$(63) \quad \int_{B_\epsilon(a_i^\epsilon)} |K_{a_i^\epsilon} \cdot K_{a_j^\epsilon}| \, dx \leq \frac{2\pi\epsilon}{2d_j^\epsilon - \epsilon} \rightarrow 0, \quad \text{as } \epsilon \rightarrow 0$$

(and again, the same holds swapping the roles of i and j .) Notice that (62) and (63) are refined versions of Lemma 3.1(ii). By combining (60), (61), (62), and (63) we get

$$(64) \quad \int_{\Omega_\epsilon(\mathbf{a}^\epsilon)} K_{a_i^\epsilon} \cdot K_{a_j^\epsilon} \, dx \rightarrow \int_{\Omega} K_{a_i} \cdot K_{a_j} \, dx, \quad \text{as } \epsilon \rightarrow 0.$$

In order to study the asymptotic behavior as $\epsilon \rightarrow 0$ of the L^2 product of $K_{a_i^\epsilon}$ and $\nabla \bar{v}_{a_j^\epsilon}^\epsilon$ we use the decomposition $\bar{v}_{a_j^\epsilon}^\epsilon = \hat{u}_j^\epsilon + q_j^\epsilon$ introduced in Remark 2.1. We

recall in particular Remark 3.1(i): \hat{u}^ϵ admits an H^1 extension (not relabeled) that strongly converges to v_{a_j} ; thus, integrating by parts and exploiting again the Dominated Convergence Theorem, in the limit as $\epsilon \rightarrow 0$ we get

$$\begin{aligned}
\int_{\Omega_\epsilon(\mathbf{a}^\epsilon)} K_{a_i^\epsilon} \cdot \nabla \hat{u}_j^\epsilon \, dx &= \int_{\Omega_\epsilon(a_i^\epsilon)} K_{a_i^\epsilon} \cdot \nabla \hat{u}_j^\epsilon \, dx - \int_{\bigcup_{j \neq i} B_\epsilon(a_j^\epsilon)} K_{a_i^\epsilon} \cdot \nabla \hat{u}_j^\epsilon \, dx \\
(65) \qquad &= \int_{\partial\Omega} K_{a_i^\epsilon} \cdot \nu (g_{b_j^\epsilon} - \theta_{a_j^\epsilon}) \, dx + o(1) \\
&\rightarrow \int_{\partial\Omega} K_{a_i} \cdot \nu (g_{b_j} - \theta_{a_j}) \, dx = \int_{\Omega} K_{a_i} \cdot \nabla v_{a_j} \, dx,
\end{aligned}$$

where we have used the fact that \hat{u}_j^ϵ has vanishing L^2 norm in the disks $B_\epsilon(a_k^\epsilon)$ as $\epsilon \rightarrow 0$; while $K_{a_i^\epsilon}$ is uniformly bounded in every disk $B_\epsilon(a_k^\epsilon)$ with $k \neq i$, and satisfies $K_{a_i^\epsilon} \cdot \nu = 0$ on $\partial B_\epsilon(a_i^\epsilon)$. On the other hand, we recall that $(\nabla q_j^\epsilon)^\perp = \nabla p_j^\epsilon$ in the perforated domain, p_j^ϵ being a solution to (31). Since the solution to (31) is unique up to a constant, we choose p_j^ϵ such that $p_j^\epsilon = 0$ on $\partial B_\epsilon(a_i^\epsilon)$. Therefore, in view of Remark 3.1(ii), we infer that p_j^ϵ admits a an extension in $H^1(\Omega)$ (not relabeled) which is harmonic in every disk $B_\epsilon(a_k^\epsilon)$ and such that $\|p_j^\epsilon\|_{H^1(\Omega)} \rightarrow 0$ as $\epsilon \rightarrow 0$. Therefore, by letting $\phi_i^\epsilon(x) := \log(|x - a_i^\epsilon|)$, we have

$$\begin{aligned}
\int_{\Omega_\epsilon(\mathbf{a}^\epsilon)} K_{a_i^\epsilon} \cdot \nabla q_j^\epsilon \, dx &= \int_{\Omega_\epsilon(\mathbf{a}^\epsilon)} \nabla \phi_i^\epsilon \cdot \nabla p_j^\epsilon \, dx = \int_{\bigcup_{k=1}^n \partial B_\epsilon(a_k^\epsilon)} \phi_i^\epsilon \nabla p_j^\epsilon \cdot \nu \, dx \\
(66) \qquad &= \int_{\bigcup_{k=1}^n \partial B_\epsilon(a_k^\epsilon)} \phi_i^\epsilon \nabla (p_j^\epsilon - p_j^\epsilon(a_i^\epsilon)) \cdot \nu \, dx = \int_{\bigcup_{k \neq i} B_\epsilon(a_k^\epsilon)} \nabla \phi_i^\epsilon \cdot \nabla p_j^\epsilon \, dx,
\end{aligned}$$

and its absolute value can be estimated from above by

$$\frac{(n-1)\sqrt{\pi}\epsilon}{d_i - \epsilon} \|\nabla p_j^\epsilon\|_{L^2(\Omega)} \rightarrow 0.$$

Similarly, the same limits in (65) and (66) hold exchanging the roles of i and j . The thesis (58) follows then by putting together (59), (64), (65), and (66). \square

4.3 - The case $n > 1$ with $\min_i d_i = 0$

Lemma 4.1. *Let $\mathbf{a} = (a_1, \dots, a_n) \in \bar{\Omega}^n$ and let $0 < \epsilon < \eta$ be such that for every $a_j \in \Omega$ we have $B_\eta(a_j) \subset \Omega$ (i.e. $d_j \geq \eta$). Then there exists a positive constant C , independent of ϵ and η , such that*

$$(67) \qquad \mathcal{F}_\epsilon^{(n)}(\mathbf{a}) \geq \mathcal{F}_\eta^{(n)}(\mathbf{a}) - C.$$

Proof. We start by comparing the energies \mathcal{E}_ϵ and \mathcal{E}_η . Recalling (2), it is easy to see that

$$\begin{aligned}
(68) \qquad \mathcal{E}_\epsilon^{(n)}(\mathbf{a}) &\geq \mathcal{E}_\eta^{(n)}(\mathbf{a}) + \frac{1}{2} \int_{\bigcup_{i=1}^n A_\epsilon^\eta(a_i)} \left| \sum_{j=1}^n \omega(a_j) K_{a_j} + \nabla \bar{v}_\mathbf{a}^\epsilon \right|^2 \, dx \\
&= \mathcal{E}_\eta^{(n)}(\mathbf{a}) + J_1 + J_2 + J_3 + J_4,
\end{aligned}$$

where we have defined

$$\begin{aligned}
J_1 &:= \frac{1}{2} \sum_{i=1}^{\ell} \int_{A_{\epsilon}^{\eta}(a_i)} |\nabla \bar{v}_{\mathbf{a}}^{\epsilon}|^2 dx, & J_2 &:= \frac{1}{2} \sum_{i,j=1}^{\ell} m_j^2 \omega^2(a_j) \int_{A_{\epsilon}^{\eta}(a_i)} |K_{a_j}|^2 dx, \\
J_3 &:= \sum_{\substack{i,j,h=1 \\ h>j}}^{\ell} m_j m_h \omega(a_j) \omega(a_h) \int_{A_{\epsilon}^{\eta}(a_i)} K_{a_j} \cdot K_{a_h} dx, \\
J_4 &:= \sum_{i,j=1}^{\ell} m_j \omega(a_j) \int_{A_{\epsilon}^{\eta}(a_i)} K_{a_j} \cdot \nabla \bar{v}_{\mathbf{a}}^{\epsilon} dx.
\end{aligned}$$

The term J_1 is strictly positive, and we can neglect it. In order to bound J_2 from below, we need to distinguish two cases, according to whether $d(a_i)$ is greater or smaller than η , namely whether or not the annulus $A_{\epsilon}^{\eta}(a_i)$ is contained in Ω . In the former case, $\omega(a_i) = 1$ and a simple computation gives (see (17))

$$\begin{aligned}
(69) \quad \frac{1}{2} \sum_{j=1}^{\ell} m_j^2 \omega^2(a_j) \int_{A_{\epsilon}^{\eta}(a_i)} |K_{a_j}|^2 dx &\geq \frac{1}{2} m_i^2 \omega^2(a_i) \int_{A_{\epsilon}^{\eta}(a_i)} |K_{a_i}|^2 dx \\
&= \pi m_i^2 \log(\eta/\epsilon).
\end{aligned}$$

In the latter, $A_{\epsilon}^{\eta}(a_i)$ is not completely contained in Ω and therefore, by hypothesis, we know that $a_i \in \partial\Omega$ and in particular $\omega(a_i) = 2$, so that

$$\begin{aligned}
(70) \quad \frac{1}{2} \sum_{j=1}^{\ell} m_j^2 \omega^2(a_j) \int_{A_{\epsilon}^{\eta}(a_i)} |K_{a_j}|^2 dx &\geq 2m_i^2 \int_{A_{\epsilon}^{\eta}(a_i)} |K_{a_i}|^2 dx \\
&> \pi m_i^2 \log\left(\frac{\min(\eta, \epsilon_{\alpha})}{\epsilon}\right),
\end{aligned}$$

where ϵ_{α} is given in (7) (see (18)). Therefore, for η small enough, $\min(\eta, \epsilon_{\alpha}) = \eta$ and (70) provides the same bound as (69), hence, summing over i , we obtain

$$(71) \quad J_2 \geq \sum_{i=1}^{\ell} \pi m_i^2 \log(\eta/\epsilon).$$

Recalling the definition (3) of $\mathcal{F}_{\epsilon}^{(n)}(\mathbf{a})$, from (68) and (71) we obtain

$$(72) \quad \mathcal{F}_{\epsilon}^{(n)}(\mathbf{a}) \geq \mathcal{F}_{\eta}^{(n)}(\mathbf{a}) + \left(\sum_{i=1}^{\ell} m_i^2 - n \right) \pi \log(\eta/\epsilon) + J_3 + J_4 \geq \mathcal{F}_{\eta}^{(n)}(\mathbf{a}) + J_3 + J_4,$$

since $\eta > \epsilon$ and $\sum_{i=1}^{\ell} m_i^2 - n \geq 0$. We obtain the thesis (67) from (72), provided we bound J_3 and J_4 from below.

In view of Lemma 3.1(i), we have

$$(73) \quad J_3 = \sum_{\substack{i,j,h=1 \\ h>j}}^{\ell} m_j m_h \omega(a_j) \omega(a_h) \int_{A_{\epsilon}^{\eta}(a_i)} K_{a_j} \cdot K_{a_h} dx \geq -8\pi \ell n^2 \geq -8\pi n^3.$$

To estimate J_4 , we start by splitting $\nabla \bar{v}_{\mathbf{a}}^\epsilon = \nabla \hat{u}_{\mathbf{a}}^\epsilon + \nabla q_{\mathbf{a}}^\epsilon$, with $\hat{u}_{\mathbf{a}}^\epsilon = \sum_{k=1}^n \hat{u}_{a_k}^\epsilon$ and $q_{\mathbf{a}}^\epsilon = \sum_{k=1}^n q_{a_k}^\epsilon$, where $\hat{u}_{a_k}^\epsilon$ and $q_{a_k}^\epsilon$ are introduced in Remark 2.1. Then

$$\sum_{i=1}^{\ell} \int_{A_\epsilon^\eta(a_i)} K_{a_j} \cdot \nabla \bar{v}_{\mathbf{a}}^\epsilon dx = \sum_{i=1}^{\ell} \int_{A_\epsilon^\eta(a_i)} K_{a_j} \cdot \nabla \hat{u}_{\mathbf{a}}^\epsilon dx + \sum_{i=1}^{\ell} \int_{A_\epsilon^\eta(a_i)} K_{a_j} \cdot \nabla q_{\mathbf{a}}^\epsilon dx$$

and, by using the Divergence Theorem, (20), and the fact that $|K_{a_j}(x)| \leq |x|^{-1}$ on $\partial A_\epsilon^\eta(a_i)$, we can estimate

$$(74) \quad \begin{aligned} \sum_{i=1}^{\ell} \int_{A_\epsilon^\eta(a_i)} K_{a_j} \cdot \nabla \hat{u}_{\mathbf{a}}^\epsilon dx &= \sum_{i=1}^{\ell} \int_{\partial A_\epsilon^\eta(a_i)} K_{a_j} \cdot \nu \hat{u}_{\mathbf{a}}^\epsilon dx \\ &\geq -\bar{C} \sum_{i=1}^{\ell} \int_{\partial A_\epsilon^\eta(a_i)} |K_{a_j} \cdot \nu| dx \geq -2\pi \bar{C} n, \end{aligned}$$

where $\bar{C} := \|g - \sum_{k=1}^n \omega(a_k) \theta_{a_k}\|_{L^\infty(\partial\Omega)}$. Moreover

$$\sum_{i=1}^{\ell} \int_{A_\epsilon^\eta(a_i)} K_{a_j} \cdot \nabla q_{\mathbf{a}}^\epsilon dx = \sum_{i=1}^{\ell} \int_{A_\epsilon^\eta(a_i)} \nabla \phi_{a_j} \cdot \nabla p_{\mathbf{a}}^\epsilon dx,$$

with $p_{\mathbf{a}}^\epsilon = \sum_{i=1}^n p_{a_i}^\epsilon$, where $p_{a_i}^\epsilon$ solves (31). In particular, by using (32), we infer that the L^∞ norm of $p_{\mathbf{a}}^\epsilon$ is bounded by $Cn\epsilon$, so that we can estimate

$$(75) \quad \sum_{i=1}^{\ell} \int_{A_\epsilon^\eta(a_i)} K_{a_j} \cdot \nabla q_{\mathbf{a}}^\epsilon dx \geq -2\pi C n^2.$$

Combining (74) with (75) and summing over j , we obtain $J_4 \geq -2\pi n^2(\bar{C} + Cn)$, which, together with (72) and (73), allows us to get estimate (67), with constant $C = 2\pi n^2(\bar{C} + (C+4)n)$. The lemma is proved. \square

Proposition 4.5. *Let $\mathbf{a} \in \bar{\Omega}^n$ and let \mathbf{a}^ϵ be a sequence of points in $\bar{\Omega}^n$ converging to \mathbf{a} as $\epsilon \rightarrow 0$. If $\min_{1 \leq i \leq n} d_i = 0$ with d_i defined as in (4) then $\mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon) \rightarrow \infty$, as $\epsilon \rightarrow 0$.*

Proof. Our goal is to show that we can bound the energy $\mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon)$ from below by a quantity that explodes in the limit as $\epsilon \rightarrow 0$. This will be achieved by applying the following iterative procedure, which is performed at ϵ fixed.

Step 0 (Labeling) We start by relabeling in a more suitable way the limit dislocations and the approximating ones. According to Definition 2.1, we relabel the limit points so that the first ℓ ($1 \leq \ell \leq n$) are distinct. Moreover, we fix $\delta > 0$ such that the disks $B_\delta(a_i)$ are pairwise disjoint for $i = 1, \dots, \ell$. For every $i = 1, \dots, \ell$, there are m_i points in $\{a_1^\epsilon, \dots, a_n^\epsilon\}$ which converge to a_i . We denote these points by $a_{i,j}^\epsilon$, with $j = 1, \dots, m_i$. For ϵ small enough we clearly have $a_{i,j}^\epsilon \in B_\delta(a_i)$, for every $j = 1, \dots, m_i$ and $i = 1, \dots, \ell$.

At each iteration (from Step 1 to Step 3), we will replace the sequence (with respect to $\epsilon > 0$) of families (indexed by $i \in \{1, \dots, \ell\}$) $\{a_{i,j}^\epsilon\}_{j=1}^{m_i}$ with a sequence of singletons c_i^ϵ , still converging to a_i as $\epsilon \rightarrow 0$, with multiplicity m_i . Additionally, we will define a new core radius $\eta(\epsilon)$ and in Step 4 we will compare the energies $\mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon)$ and $\mathcal{F}_{\eta(\epsilon)}^{(n)}(\mathbf{c}^\epsilon)$.

Step 1 (Ordering) Let $i \in \{1, \dots, \ell\}$ be fixed. According to Definition 2.1, we order the family $\{a_{i,j}^\epsilon\}_{j=1}^{m_i}$ so that the first ℓ_i^ϵ are distinct, and we denote by $m_{i,j}^\epsilon$ their multiplicities. Notice that $\sum_{j=1}^{\ell_i^\epsilon} m_{i,j}^\epsilon = m_i$. With these positions, we clearly have

$$(76) \quad \mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon) = \mathcal{F}_\epsilon^{(n)}\left(\underbrace{a_{1,1}^\epsilon}_{m_{1,1}^\epsilon\text{-times}}, \dots, \underbrace{a_{1,\ell_1^\epsilon}^\epsilon}_{m_{1,\ell_1^\epsilon}^\epsilon\text{-times}}, \dots, \underbrace{a_{i,1}^\epsilon}_{m_{i,1}^\epsilon\text{-times}}, \dots, \underbrace{a_{\ell,1}^\epsilon}_{m_{\ell,1}^\epsilon\text{-times}}, \dots, \underbrace{a_{\ell,\ell_\ell^\epsilon}^\epsilon}_{m_{\ell,\ell_\ell^\epsilon}^\epsilon\text{-times}} \right).$$

In case $\ell_i^\epsilon > 1$, for every $j = 1, \dots, \ell_i^\epsilon$, we associate to the distinct points $a_{i,j}^\epsilon$ the following quantity: if $a_{i,j}^\epsilon \in \partial\Omega$, we set

$$s(a_{i,j}^\epsilon) := \min \left\{ \frac{|a_{i,j}^\epsilon - a_{i,k}^\epsilon|}{2} : k \in \{1, \dots, \ell_i^\epsilon\} \setminus \{j\} \right\},$$

whereas, if $a_{i,j}^\epsilon \in \Omega$, we set

$$s(a_{i,j}^\epsilon) := \min \left\{ d(a_{i,j}^\epsilon), \min \left\{ \frac{|a_{i,j}^\epsilon - a_{i,k}^\epsilon|}{2} : k \in \{1, \dots, \ell_i^\epsilon\} \setminus \{j\} \right\} \right\}.$$

Observe that, if the limit point $a_i \in \Omega$, for ϵ small enough $d(a_{i,j}^\epsilon)$ is always greater than any mutual semidistance $|a_{i,j}^\epsilon - a_{i,k}^\epsilon|/2$, for all $j, k \in \{1, \dots, \ell_i^\epsilon\}$. Up to reordering the different ℓ_i^ϵ points $\{a_{i,1}^\epsilon, \dots, a_{i,\ell_i^\epsilon}^\epsilon\}$ we can always suppose that

$$0 < s(a_{i,1}^\epsilon) \leq \dots \leq s(a_{i,\ell_i^\epsilon}^\epsilon).$$

In case $\ell_i^\epsilon = 1$, namely when all the $a_{i,j}^\epsilon$ coincide with $a_{i,1}^\epsilon$, we set

$$(77) \quad s(a_{i,1}^\epsilon) := \begin{cases} 0 & \text{if } a_i \in \Omega, \\ d(a_{i,1}^\epsilon) & \text{if } a_i \in \partial\Omega. \end{cases}$$

Step 2 (Stop test) If $s(a_{i,1}^\epsilon) = 0$ for every $i = 1, \dots, \ell$, then we define $c_i^\epsilon := a_{i,1}^\epsilon$, $\eta(\epsilon) := \epsilon$, and we go to Step 4. Observe that by (77) $c_i^\epsilon \in \partial\Omega$ if $a_i \in \partial\Omega$. Otherwise, we define the following quantity

$$(78) \quad \hat{s} = \hat{s}(\epsilon) := \min \{s(a_{i,1}^\epsilon) > 0 : i \in \{1, \dots, \ell\}\}$$

and we go to Step 3. Observe that the set where the minimum is taken is not empty, hence \hat{s} is finite and strictly positive.

Step 3 (Iterative step) We compare \hat{s} with ϵ .

If $\hat{s} > \epsilon$, we relabel $a_{k,j}^\epsilon$ by $\hat{a}_{k,j}^\epsilon$ and their multiplicities accordingly, set $\hat{\epsilon} := \hat{s}$, and estimate $\mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon)$ by means of (67) proved in Lemma 4.1, with $\eta = \hat{\epsilon}$, to obtain

$$(79) \quad \mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon) \geq \mathcal{F}_{\hat{\epsilon}}^{(n)}(\dots, \underbrace{\hat{a}_{k,j}^\epsilon}_{\hat{m}_{k,j}^\epsilon\text{-times}}, \dots) - C.$$

If $\hat{s} \leq \epsilon$ we distinguish two cases:

- (1) \hat{s} is equal to $|a_{i,1}^\epsilon - a_{i,2}^\epsilon|/2$, for some i : we replace the points $a_{i,1}^\epsilon$ and $a_{i,2}^\epsilon$ (and all those coinciding with either one of them) by $\hat{a}_{i,1}^\epsilon$, the mid-point between $a_{i,1}^\epsilon$ and $a_{i,2}^\epsilon$, with multiplicity $\hat{m}_{i,1}^\epsilon := m_{i,1}^\epsilon + m_{i,2}^\epsilon$. The replacement is performed simultaneously for all the indices i that realize the minimum in (78). For all the other indices, we simply relabel $a_{k,j}^\epsilon$ by $\hat{a}_{k,j}^\epsilon$ and their multiplicities accordingly.
- (2) \hat{s} is realized by $d(a_{i,1}^\epsilon)$, for some i : we replace the point $a_{i,1}^\epsilon$ (and all those coinciding with it) by $\hat{a}_{i,1}^\epsilon$, its projection to the boundary $\partial\Omega$, with multiplicity $\hat{m}_{i,1}^\epsilon := m_{i,1}^\epsilon$.

Setting $\hat{\epsilon} := \epsilon + \hat{s}$ and recalling (2), (3), and (76), we have

$$(80) \quad \mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon) \geq \mathcal{F}_{\hat{\epsilon}}^{(n)}(\dots, \underbrace{\hat{a}_{k,j}^\epsilon}_{\hat{m}_{k,j}^\epsilon\text{-times}}, \dots) - n\pi \log 2.$$

Notice that $\hat{\epsilon}$ satisfies the following bound:

$$(81) \quad \hat{\epsilon} \leq \max\{2\epsilon, \hat{s}\} \leq \max\{2\epsilon, \bar{s}\},$$

where \bar{s} is defined as the maximum value of the $s(a_{i,j}^\epsilon)$, namely

$$\bar{s} = \bar{s}(\epsilon) := \max\{s(a_{i,\ell}^\epsilon) > 0 : i \in \{1, \dots, \ell\}\}.$$

We have obtained a new family $\{\hat{a}_{i,j}^\epsilon\}$ and a new radius $\hat{\epsilon}$ and we restart the procedure by applying Step 1 to these new objects.

Notice that the procedure ends after at most n^2 iterations. Indeed, when applying Step 3, we will always fall into case (1) after at most n iterations, and then the number of distinct points will decrease when we apply (1). In conclusion, since the number of distinct points is at most n , we will reach the target situation after at most n^2 iterations of Step 3.

Step 4 (Estimates and conclusion) By combining the chain of inequalities obtained by applying Step 3 $k(\leq n^2)$ times, estimates (80) and (79) give

$$(82) \quad \mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon) \geq \mathcal{F}_{\eta(\epsilon)}^{(n)}(\underbrace{c_1^\epsilon}_{m_1\text{-times}}, \dots, \underbrace{c_\ell^\epsilon}_{m_\ell\text{-times}}) - k(C + n\pi \log 2),$$

where $\eta(\epsilon)$ is a number depending on ϵ obtained after k iterations of the procedure that defines $\hat{\epsilon}$ in Step 3. Let us estimate $\eta(\epsilon)$. At every iteration, the value \bar{s} can increase, but it is easy to see that it cannot grow larger than its double. Therefore, after k iterations of Step 3, by (81), we have

$$\eta(\epsilon) \leq 2^k \max\{\epsilon, \bar{s}\} \leq 2^{n^2} \max\{\epsilon, \bar{s}\}.$$

Since \bar{s} tends to 0 as $\epsilon \rightarrow 0$, we have $\eta(\epsilon) \rightarrow 0$ as $\epsilon \rightarrow 0$.

We now claim that the right-hand side of (82) tends to $+\infty$ as $\epsilon \rightarrow 0$. Similarly to (56), we write the energy as the sum of three contributions:

$$(83) \quad \mathcal{F}_{\eta(\epsilon)}^{(n)}\left(\underbrace{c_1^\epsilon}_{m_1\text{-times}}, \dots, \underbrace{c_\ell^\epsilon}_{m_\ell\text{-times}}\right) = \sum_{i=1}^{\ell} \mathcal{F}_{\eta(\epsilon)}^{(m_i)}\left(\underbrace{c_i^\epsilon}_{m_i\text{-times}}\right) + \sum_{i=1}^{\ell} m_i^2 \mathcal{R}_\epsilon(c_i^\epsilon) + \sum_{i \neq j} m_i m_j \mathcal{G}_\epsilon(c_i^\epsilon, c_j^\epsilon)$$

where, setting $\omega_i := \omega(a_i) = \omega(c_i^\epsilon)$ (see Step 3), and $K_i^\epsilon := K_{c_i^\epsilon}$ for brevity,

$$(84) \quad \mathcal{R}_\epsilon(c_i^\epsilon) := \frac{1}{2} \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} |\omega_i K_i^\epsilon + \nabla \bar{v}_i^\epsilon|^2 dx - \frac{1}{2} \int_{\Omega_{\eta(\epsilon)}(c_i^\epsilon)} |\omega_i K_i^\epsilon + \nabla \bar{u}_i^\epsilon|^2 dx,$$

and

$$(85) \quad \mathcal{G}_\epsilon(c_i^\epsilon, c_j^\epsilon) := \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} (\omega_i K_i^\epsilon + \nabla \bar{v}_i^\epsilon) \cdot (\omega_j K_j^\epsilon + \nabla \bar{v}_j^\epsilon) dx,$$

\bar{u}_i^ϵ being the solution to (42) associated with c_i^ϵ and \bar{v}_i^ϵ being as in Remark 2.1 for $i = 1, \dots, \ell$ and core radius $\eta(\epsilon)$.

Fix $i \in \{1, \dots, \ell\}$. If $a_i \in \Omega$, we have

$$(86) \quad \mathcal{F}_{\eta(\epsilon)}^{(m_i)}\left(\underbrace{c_i^\epsilon}_{m_i\text{-times}}\right) = m_i^2 \mathcal{F}_{\eta(\epsilon)}(c_i^\epsilon) + m_i(m_i - 1)\pi |\log \eta(\epsilon)| \geq m_i(m_i - 1)\pi |\log \eta(\epsilon)| - C,$$

where the equality follows from (13) and the inequality is a consequence of Proposition 4.1. If instead $a_i \in \partial\Omega$, by using (13) again and (49), we have

$$\mathcal{F}_{\eta(\epsilon)}^{(m_i)}\left(\underbrace{c_i^\epsilon}_{m_i\text{-times}}\right) \geq C_1 m_i^2 |\log(\max\{\eta(\epsilon), d(c_i^\epsilon)\})| + m_i(m_i - 1)\pi |\log \eta(\epsilon)| + C_2 m_i^2$$

(the constants C_1 and C_2 are those in (54)), and, since $d(c_i^\epsilon) = 0$ as noticed in Step 2, we obtain

$$(87) \quad \mathcal{F}_{\eta(\epsilon)}^{(m_i)}\left(\underbrace{c_i^\epsilon}_{m_i\text{-times}}\right) \geq (C_1 m_i^2 + m_i(m_i - 1)\pi) |\log \eta(\epsilon)| + C_2 m_i^2.$$

By Lemma 3.6 the function \bar{v}_i^ϵ can be extended inside any disk $B_{\eta(\epsilon)}(c_j^\epsilon)$ with $j \neq i$; therefore the remainder (84) can be rewritten as $\mathcal{R}_\epsilon(c_i^\epsilon) = \mathcal{R}'_\epsilon(c_i^\epsilon) - \mathcal{R}''_\epsilon(c_i^\epsilon)$ with

$$\mathcal{R}'_\epsilon(c_i^\epsilon) := \frac{1}{2} \int_{\Omega_{\eta(\epsilon)}(c_i^\epsilon)} |\omega_i K_i^\epsilon + \nabla \bar{v}_i^\epsilon|^2 dx - \frac{1}{2} \int_{\Omega_{\eta(\epsilon)}(c_i^\epsilon)} |\omega_i K_i^\epsilon + \nabla \bar{u}_i^\epsilon|^2 dx$$

and

$$\mathcal{R}''_\epsilon(c_i^\epsilon) := \frac{1}{2} \sum_{j \neq i} \int_{B_{\eta(\epsilon)}(c_j^\epsilon)} |\omega_i K_i^\epsilon + \nabla \bar{v}_i^\epsilon|^2 dx.$$

Using the minimality of \bar{u}_i^ϵ in $\Omega_{\eta(\epsilon)}(c_i^\epsilon)$ it turns out that $\mathcal{R}'_\epsilon(c_i^\epsilon) \geq 0$. Using that $|K_i^\epsilon| \leq 1/\eta(\epsilon)$ in $B_{\eta(\epsilon)}(c_j^\epsilon)$ for every $j \neq i$ and Lemma 3.6 we obtain

$$(88) \quad \begin{aligned} \mathcal{R}_\epsilon(c_i^\epsilon) &\geq -\mathcal{R}''_\epsilon(c_i^\epsilon) \\ &\geq -\frac{1}{2} \sum_{j \neq i} \int_{B_{\eta(\epsilon)}(c_j^\epsilon)} |\omega_i K_i^\epsilon|^2 dx - \frac{1}{2} \sum_{j \neq i} \int_{B_{\eta(\epsilon)}(c_j^\epsilon)} |\nabla \bar{v}_i^\epsilon|^2 dx \geq -C. \end{aligned}$$

To estimate the interaction term (85), fix also $j \in \{1, \dots, \ell\} \setminus \{i\}$. Then, we can write $\mathcal{G}_\epsilon(c_i^\epsilon, c_j^\epsilon) = \mathcal{G}'_\epsilon(c_i^\epsilon, c_j^\epsilon) + \mathcal{G}''_\epsilon(c_i^\epsilon, c_j^\epsilon) + \mathcal{G}'''_\epsilon(c_i^\epsilon, c_j^\epsilon)$ with

$$(89a) \quad \mathcal{G}'_\epsilon(c_i^\epsilon, c_j^\epsilon) := \omega_i \omega_j \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} K_i^\epsilon \cdot K_j^\epsilon dx,$$

$$(89b) \quad \mathcal{G}''_\epsilon(c_i^\epsilon, c_j^\epsilon) := \omega_j \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} \nabla \bar{v}_i^\epsilon \cdot K_j^\epsilon dx + \omega_i \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} \nabla \bar{v}_j^\epsilon \cdot K_i^\epsilon dx,$$

$$(89c) \quad \mathcal{G}'''_\epsilon(c_i^\epsilon, c_j^\epsilon) := \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} \nabla \bar{v}_i^\epsilon \cdot \nabla \bar{v}_j^\epsilon dx.$$

Thanks to Lemma 3.1(i), the functional \mathcal{G}'_ϵ in (89a) is uniformly bounded from below by a constant. To estimate \mathcal{G}''_ϵ and \mathcal{G}'''_ϵ we recall the decomposition $\bar{v}_k^\epsilon = \hat{u}_k^\epsilon + q_k^\epsilon$ and the function p_k^ϵ , solution to (31), that we introduced in Remark 2.1, where $k = i$ or $k = j$. Here the functions \hat{u}_k^ϵ and q_k^ϵ are introduced in Remark 2.1 with $(a_1^\epsilon, \dots, a_n^\epsilon)$ replaced by $(c_1^\epsilon, \dots, c_\ell^\epsilon)$, and the Dirichlet boundary condition for \hat{u}_k^ϵ given by $g_{b_k^\epsilon} - \omega_k \theta_k^\epsilon$ on $\partial\Omega$. Therefore, we can estimate (89c) as follows

$$\mathcal{G}'''_\epsilon(c_i^\epsilon, c_j^\epsilon) \leq \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} (|\nabla \hat{u}_i^\epsilon|^2 + |\nabla \hat{u}_j^\epsilon|^2 + |\nabla q_i^\epsilon|^2 + |\nabla q_j^\epsilon|^2) dx$$

and since, the gradient of q_k^ϵ coincides in modulus with the gradient of p_k^ϵ , by (21) and (33) we can bound $\mathcal{G}'''_\epsilon(c_i^\epsilon, c_j^\epsilon)$ from below. The functional in (89b) is the most delicate to treat: if both a_i and a_j are in Ω , $\mathcal{G}''_\epsilon(c_i^\epsilon, c_j^\epsilon)$ is bounded below by a constant (notice that in this case we could have applied Proposition 4.4 to $\mathcal{G}_\epsilon(c_i^\epsilon, c_j^\epsilon)$ itself and we would have concluded). In the general case, by means of the above-mentioned decomposition, (89b) reads

$$(90) \quad \begin{aligned} \mathcal{G}''_\epsilon(c_i^\epsilon, c_j^\epsilon) &= \omega_j \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} \nabla \hat{u}_i^\epsilon \cdot K_j^\epsilon dx + \omega_i \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} \nabla \hat{u}_j^\epsilon \cdot K_i^\epsilon dx \\ &\quad + \omega_j \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} \nabla q_i^\epsilon \cdot K_j^\epsilon dx + \omega_i \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} \nabla q_j^\epsilon \cdot K_i^\epsilon dx. \end{aligned}$$

We first deal with the last two terms, involving the gradients of q_i^ϵ and q_j^ϵ : from Hölder's inequality, recalling that the gradient of q_k^ϵ coincides in modulus with the gradient of p_k^ϵ , by (33) and (17) we conclude that these terms are bounded. To estimate the first two terms in (90), we use Young's inequality, (21) and (17) to obtain

$$\begin{aligned}
& \omega_j \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} \nabla \hat{u}_i^\epsilon \cdot K_j^\epsilon \, dx + \omega_i \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} \nabla \hat{u}_j^\epsilon \cdot K_i^\epsilon \, dx \\
& \leq \frac{1}{2\lambda} \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} (\omega_j |\nabla \hat{u}_i^\epsilon|^2 + \omega_i |\nabla \hat{u}_j^\epsilon|^2) \, dx \\
& \quad + \frac{\lambda}{2} \int_{\Omega_{\eta(\epsilon)}(c_1^\epsilon, \dots, c_\ell^\epsilon)} (\omega_i |K_i^\epsilon|^2 + \omega_j |K_j^\epsilon|^2) \, dx \\
& \leq \frac{C}{\lambda} \left(\|g_{b_i^\epsilon} - \omega_i \theta_i^\epsilon\|_{H^{1/2}(\partial\Omega)}^2 + \|g_{b_j^\epsilon} - \omega_j \theta_j^\epsilon\|_{H^{1/2}(\partial\Omega)}^2 \right) \\
& \quad + 2\pi\lambda |\log(\eta(\epsilon))| + 2\pi\lambda \log(\text{diam } \Omega),
\end{aligned}$$

where C is a positive constant independent of ϵ and $\lambda > 0$ is an arbitrary constant that will be chosen later. Finally, we can control the term $\mathcal{G}_\epsilon(c_i^\epsilon, c_j^\epsilon)$ as

$$(91) \quad \mathcal{G}_\epsilon(c_i^\epsilon, c_j^\epsilon) \geq -2\pi\lambda |\log(\eta(\epsilon))| - C_{ij}$$

where all the terms independent of ϵ have been included in the constant C_{ij} .

We can classify each point a_i according to whether it belongs to

1. the interior of Ω , with multiplicity $m_i = 1$;
2. the interior of Ω , with multiplicity $m_i > 1$;
3. the boundary $\partial\Omega$.

For $k = 1, 2, 3$ we denote by I_k the set of indices corresponding to those points $c_i^\epsilon \rightarrow a_i$ belonging to the k -th category. Therefore, combining (86), (87), (88), and (91), summing over i and j , we obtain the following lower bound for the energy (83):

$$\mathcal{F}_{\eta(\epsilon)}^{(n)} \left(\underbrace{c_1^\epsilon}_{m_1\text{-times}}, \dots, \underbrace{c_\ell^\epsilon}_{m_\ell\text{-times}} \right) \geq \pi(C_{\mathcal{F}} - \lambda C_{\mathcal{G}}) |\log \eta(\epsilon)| + C,$$

where

$$C_{\mathcal{F}} := \sum_{i \in I_2 \cup I_3} m_i(m_i - 1) + \sum_{i \in I_3} C_1 m_i^2, \quad C_{\mathcal{G}} := 2n^2,$$

and C is a constant independent of ϵ . The assumption $\min_{1 \leq i \leq n} d_i = 0$ guarantees that $C_{\mathcal{F}} > 0$, since in this case either $I_2 \neq \emptyset$ or $I_3 \neq \emptyset$. Therefore, choosing λ in (91) so that $0 < \lambda < C_{\mathcal{F}}/C_{\mathcal{G}}$, we obtain that the right-hand side of (82) tends to $+\infty$ as $\epsilon \rightarrow 0$. The proposition is proved. \square

4.4 - Proofs of Theorem 1.1 and of Corollary 1.1

We combine the previous results to prove the main results in the case of n dislocations. **Proof.** [Proof of Theorem 1.1] By (56), combining Propositions 4.1 and 4.2 with Propositions 4.3, 4.4, and 4.5 yields the continuous convergence of $\mathcal{F}_\epsilon^{(n)}$ to $\mathcal{F}^{(n)}$. The continuity of \mathcal{F} follows from the relationship between continuous convergence and Γ -convergence (see [8, Remark 4.9]), which implies that $\mathcal{F}^{(n)}(\mathbf{a})$ tends to $+\infty$ as either one of the a_i 's approaches the boundary or any two a_i and a_j (with $i \neq j$) become arbitrarily close. Therefore $\mathcal{F}^{(n)}$ is minimized by n -tuples $\mathbf{a} = (a_1, \dots, a_n)$ of distinct points in Ω^n . \square

Proof. [Proof of Corollary 1.1] Recalling definition (4), let $\delta > 0$ and let

$$\Omega^{\delta, n} := \{(x_1, \dots, x_n) \in \Omega^n : d_i > \delta \quad \forall i = 1, \dots, n\}.$$

The energy functional $\mathcal{E}_\epsilon^{(n)}$ is continuous over $\overline{\Omega^{\delta, n}}$, for every $\epsilon \in (0, \delta)$. Indeed, given two different configurations of dislocations \mathbf{a}' and \mathbf{a}'' , it is possible to construct a diffeomorphism Φ that maps $\overline{\Omega}_\epsilon(\mathbf{a}')$ into $\overline{\Omega}_\epsilon(\mathbf{a}'')$, keeps the boundary $\partial\Omega$ fixed, and satisfies $\|D\Phi - I\|_{L^\infty(\Omega; \mathbb{R}^{2 \times 2})}$, $\|\det D\Phi - 1\|_{L^\infty(\Omega)} = o(|\mathbf{a}' - \mathbf{a}''|)$. This implies that $\mathcal{E}_\epsilon^{(n)}(\mathbf{a}') = \mathcal{E}_\epsilon^{(n)}(\mathbf{a}'') + o(|\mathbf{a}' - \mathbf{a}''|)$.

Fix now $\epsilon > 0$ and $k \in \mathbb{N}$ and let $\mathbf{a}_{\epsilon, k}$ be such that

$$(92) \quad \inf_{\overline{\Omega}^n} \mathcal{F}_\epsilon^{(n)} \leq \mathcal{F}_\epsilon^{(n)}(\mathbf{a}_{\epsilon, k}) \leq \inf_{\overline{\Omega}^n} \mathcal{F}_\epsilon^{(n)} + 1/k.$$

Without loss of generality, we can assume that the whole sequence $\mathbf{a}_{\epsilon, k}$ converges to some $\mathbf{a}^\epsilon \in \overline{\Omega}^n$, as $k \rightarrow \infty$, and satisfies $|\mathbf{a}_{\epsilon, k} - \mathbf{a}^\epsilon| < 1/k$. We claim that there exist $\bar{\epsilon} > 0$ and $\delta > 0$ such that (defining d_i^ϵ 's as in (4), associated with the point \mathbf{a}^ϵ)

$$(93) \quad d^\epsilon := \min_i d_i^\epsilon > \delta \quad \forall \epsilon \in (0, \bar{\epsilon}).$$

Assume by contradiction that there exists a subsequence of ϵ (not relabeled) such that $d^\epsilon \rightarrow 0$ as $\epsilon \rightarrow 0$. Let k_ϵ be a sequence of natural numbers, increasing as ϵ goes to zero, and let \mathbf{a} be a cluster point of the family $\{\mathbf{a}_{\epsilon, k_\epsilon}\}_\epsilon$. In view of Theorem 1.1, thanks to [8, Corollary 7.20], we infer that \mathbf{a} is a minimizer of the functional $\mathcal{F}^{(n)}$. By the triangle inequality, we get

$$\min_i d_i \leq d^\epsilon + |\mathbf{a}^\epsilon - \mathbf{a}_{\epsilon, k_\epsilon}| + |\mathbf{a}_{\epsilon, k_\epsilon} - \mathbf{a}| < d^\epsilon + \frac{1}{k_\epsilon} + o(\epsilon) \rightarrow 0,$$

which implies that there exists an index i such that either $a_i \in \partial\Omega$ or $a_i = a_j$ for some $i \neq j$, in contradiction with Theorem 1.1.

Let $\epsilon_0 := \min\{\bar{\epsilon}, \delta\}$ and $\epsilon \in (0, \epsilon_0)$. In view of (93), the minimizing sequence $\{\mathbf{a}_{\epsilon, k}\}_k$ lies in the set $\{(x_1, \dots, x_n) \in \Omega^n : d_i > \delta \quad \forall i = 1, \dots, n\}$. Therefore, by (92), we conclude that

$$\inf_{\overline{\Omega}^n} \mathcal{F}_\epsilon^{(n)} = \lim_{k \rightarrow \infty} \mathcal{F}_\epsilon^{(n)}(\mathbf{a}_{\epsilon, k}) = \mathcal{F}_\epsilon^{(n)}(\mathbf{a}^\epsilon),$$

namely \mathbf{a}^ϵ is a minimizer of $\mathcal{F}_\epsilon^{(n)}$ made of n distinct points each of which is at distance at least δ from the boundary. Eventually, again by [8, Corollary 7.20], we infer that, up to a subsequence (not relabeled), \mathbf{a}^ϵ converges to a minimizer of $\mathcal{F}^{(n)}$ so that $\mathcal{F}_\epsilon(\mathbf{a}^\epsilon)^{(n)} \rightarrow \mathcal{F}^{(n)}(a)$ as $\epsilon \rightarrow 0$. \square

5 - Plots of the limiting energy

We plot here the limiting energy $\mathcal{F}_\epsilon^{(n)}$ of (3) in the special case $n = 1$ and Ω the unit disk centered at the origin O for two different boundary conditions. In view of Corollary 1.1, for a fixed $\epsilon > 0$ small enough, the minima of $\mathcal{F}_\epsilon^{(1)}$ are close to those of $\mathcal{F}^{(1)}$, therefore, by (3), the energy landscape provided by $\mathcal{F}^{(1)}$ is a good approximation of that of $\mathcal{E}_\epsilon^{(1)}$.

The two different boundary conditions that we consider are $f_1 = 1$ on $\partial\Omega$, and $f_2 = 2$ on $\partial\Omega \cap \{x > 0\}$ and $f_2 = 0$ on $\partial\Omega \cap \{x < 0\}$; the choices that we make for the numerics are

$$g_1 = \theta_O \quad \text{and} \quad g_2 = \begin{cases} 2\theta_O & \text{in the first quadrant,} \\ \pi & \text{in the second and third quadrants,} \\ 2\theta_O - 2\pi & \text{in the fourth quadrant.} \end{cases}$$

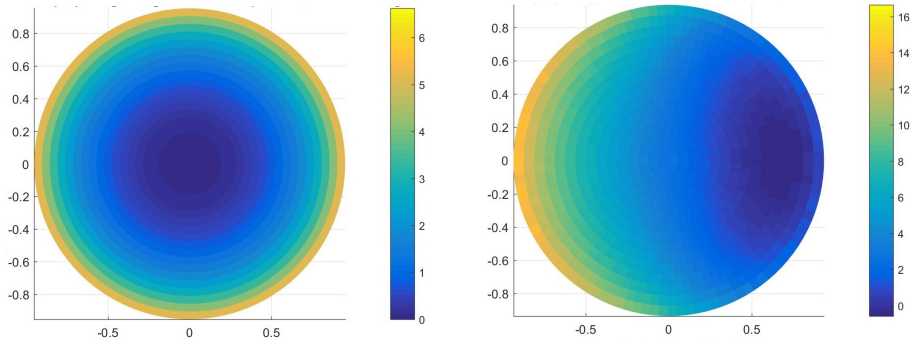


Figure 2: To the left, the functional $\mathcal{F}^{(1)}$ associated with g_1 ; to the right, the functional $\mathcal{F}^{(1)}$ associated with g_2 . Both are top views.

From Figure 2 one can deduce that the energy profile associated with f_1 is radially symmetric and has a minimum in the center. The one for f_2 is no longer radially symmetric, but symmetric with respect to the y axis and has a minimum in the interior of Ω which is located where f_2 is the largest (by symmetry, it is located along the x -axis, at about $x = 0.65$).

For $n \geq 2$, we consider Ω the unit disk centered at the origin, $f^{(n)} = n$ on $\partial\Omega$, and, for simplicity, we minimize $\mathcal{F}^{(n)}$ in a subclass of configurations: the vertices of regular n -gons centered at the origin (this particular choice is supported by

the conjecture in [24]). The optimization problem becomes (numerically) easy, as it is enough to minimize the energy with respect to the circumradius of the n -gon, which may vary from 0 to 1.

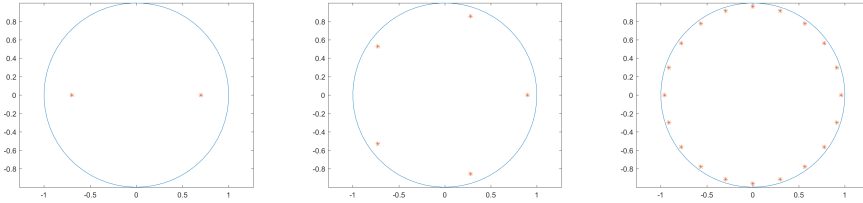


Figure 3: The optimizers of $\mathcal{F}^{(n)}$ among the vertices of regular n -gons centered at the origin for $n = 2, 5, 20$, from left to right.

Let us denote by \mathbf{p}_n an optimal n -gone of $\mathcal{F}^{(n)}$ (unique up to rotations). As it is suggested by Figure 3, as n increases we observe a decrease of the distance $\text{dist}(\mathbf{p}_n, \partial\Omega)$ and of the energy $\mathcal{F}^{(n)}(\mathbf{p}_n)$ (which is negative). More precisely, we observe the following behaviors (see Figure 4):

$$\text{dist}(\mathbf{p}_n, \partial\Omega) \sim \frac{1}{n}, \quad |\mathcal{F}^{(n)}(\mathbf{p}_n)| \sim n^2, \quad \text{as } n \rightarrow +\infty.$$

Let us mention that the investigation of minima and minimizers of $\mathcal{F}^{(n)}$ in the limit as $n \rightarrow +\infty$ is the object of our paper [17].

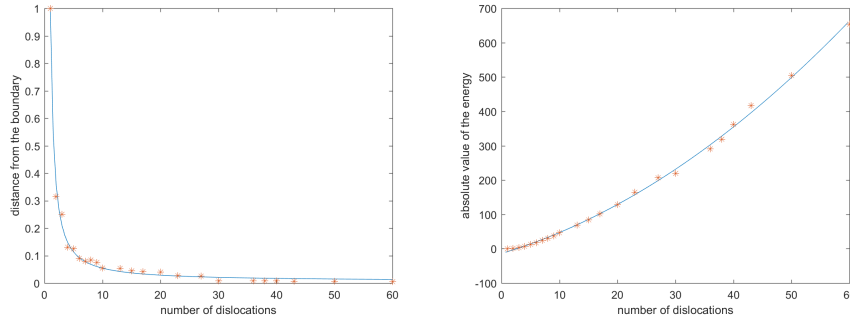


Figure 4: Plot of $\text{dist}(\mathbf{p}_n, \partial\Omega)$ (left) and plot of $|\mathcal{F}^{(n)}(\mathbf{p}_n)|$ (right) as functions of n . Test done for n dislocations, n between 1 and 60.

References

- [1] R. ALICANDRO, L. DE LUCA, A. GARRONI, and M. PONSIGLIONE, *Metastability and Dynamics of Discrete Topological Singularities in Two Dimensions: A Γ -Convergence Approach*, Arch. Rational Mech. Anal. **214** (2014), 269-330.
- [2] F. BETHUEL, H. BREZIS, and F. HÉLEIN, *Ginzburg-Landau vortices*, Progress in Nonlinear Differential Equations and their Applications, 13, Birkhäuser Boston, Inc., Boston 1994.
- [3] T. BLASS, I. FONSECA, G. LEONI, and M. MORANDOTTI, *Dynamics for systems of screw dislocations*, SIAM J. Appl. Math. **75** (2015), 393-419.
- [4] T. BLASS and M. MORANDOTTI, *Renormalized energy and Peach-Köhler forces for screw dislocations with antiplane shear*, J. Convex Anal. **24** (2017), 547-570.
- [5] P. CERMELLI and M. E. GURTIN, *The motion of screw dislocations in crystalline materials undergoing antiplane shear: glide, cross-slip, fine cross-slip*, Arch. Ration. Mech. Anal. **148** (1999), 3-52.
- [6] P. CERMELLI and G. LEONI, *Renormalized energy and forces on dislocations*, SIAM J. Math. Anal. **37** (2005), 1131-1160.
- [7] G. Q. CHEN and H. FRID, *On The Theory Of Divergence-Measure Fields And Its Applications*, Bol. Soc. Bras. Mat. **32** (2001), 401-433.
- [8] G. DAL MASO, *An introduction to Γ -convergence*, Progress in Nonlinear Differential Equations and their Applications, 8, Birkhäuser Boston Inc., Boston 1993.
- [9] S. DIPIERRO, G. PALATUCCI, and E. VALDINOCI, *Dislocation Dynamics in Crystals: A Macroscopic Theory in a Fractional Laplace Setting*, Commun. Math. Phys. **333** (2015), 1061-1105.
- [10] L. C. EVANS, *Partial differential equations*, Graduate Studies in Mathematics, 19, American Mathematical Society 1998.
- [11] A. GARRONI, G. LEONI, and M. PONSIGLIONE, *Gradient theory for plasticity via homogenization of discrete dislocations*, J. Eur. Math. Soc. **12** (2010), 1231-1266.
- [12] E. C. GARTLAND JR., A. M. SONNET, and E. G. VIRGA, *Elastic forces on nematic point defects*, Contin. Mech. Thermodyn. **14** (2002), 307-319.
- [13] J. HADAMARD, *Lectures on Cauchy's problem in linear partial differential equations*, Dover Phoenix editions, Dover Publications, New York 1923.
- [14] J. P. HIRTH and J. LOTHE, *Theory of Dislocations*, Krieger Publishing Company, Malabar 1992.
- [15] T. HUDSON and M. MORANDOTTI, *Qualitative properties of dislocation dynamics: collisions and boundary behaviour*, SIAM J. Appl. Math. **77** (2017), 1678-1705.
- [16] D. HULL and D. J. BACON, *Introduction to dislocations*, Butterworth-Heinemann, Oxford 2001.

- [17] I. LUCARDESI, M. MORANDOTTI, R. SCALA, and D. ZUCCO, *Upscaling of screw dislocations with increasing tangential strain*, Submitted, arXiv:1808.08898.
- [18] N. I. MUSKHELISHVILI, *Some Basic Problems in the Math. Theory of Elasticity*, Noordhoff, Groningen 1963.
- [19] F. R. N. NABARRO, *Theory of crystal dislocations*, International series of monographs on Physics, Clarendon, Oxford 1967.
- [20] E. OROWAN, *Zur kristallplastizität. III*, Zeitschrift für Physik **89** (1934), 634-659.
- [21] M. POLANYI, *Über eine art gitterstörung, die einen kristall plastisch machen könnte*, Zeitschrift für Physik, **89** (1934) 660-664.
- [22] M. PONSIGLIONE, *Elastic energy stored in a crystal induced by screw dislocations: from discrete to continuous*, SIAM J. Math. Anal. **39** (2007), 449-469.
- [23] E. SANDIER and S. SERFATY, *Vortices in the magnetic Ginzburg-Landau model*, Progress in Nonlinear Differential Equations and their Applications, 70. Birkhäuser Boston, Inc., Boston 2007.
- [24] E. SANDIER and M. SORET, \mathbb{S}^1 -Valued Harmonic Maps with High Topological Degree, Harmonic morphisms, harmonic maps, and related topics, 141-145, Chapman & Hall/CRC Res. Notes Math. 413, Chapman & Hall/CRC, Boca Raton 2000.
- [25] E. B. TADMOR, M. ORTIZ, and R. PHILLIPS, *Quasicontinuum analysis of defects in solids*, Philosophical Magazine A - Physics of Condensed Matter Structure Defects and Mechanical Properties **73** (1996), 1529-1563.
- [26] G. I. TAYLOR, *The Mechanism of Plastic Deformation of Crystals. Part I. Theoretical*, Proceedings of the Royal Society of London. Series A **145** (1934), 362-387.
- [27] B. VAN KOTEN, X. H. LI, M. LUSKIN, and C. ORTNER, *A computational and theoretical investigation of the accuracy of quasicontinuum methods*, in Numerical Analysis of Multiscale Problems, Lect. Notes Comput. Sci. Eng., 83, Springer, Heidelberg (2012), 67-96.
- [28] V. VOLTERRA, *Sur l'équilibre des corps élastiques multiplement connexes*, Annales scientifiques de l'École Normale Supérieure, **24** (1907), 401-517.

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